

Unravelling the Higgs nature with EFTs

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In collaboration with: Alonso, Corbett,
Éboli, Gavela, Gonzalez–Fraile,
Gonzalez–García, Merlo, Rigolin





... what kind of **Higgs**?

HEFT

SM Higgs doublet

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SM Higgs doublet

Composite Higgs

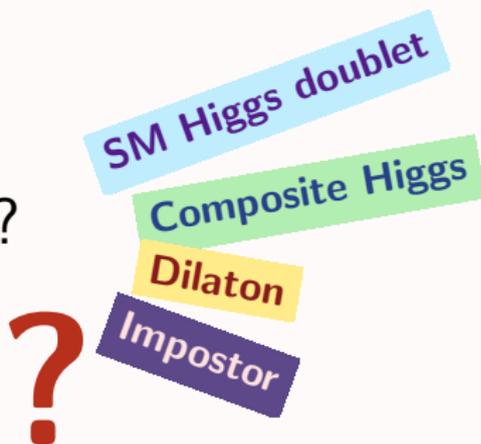


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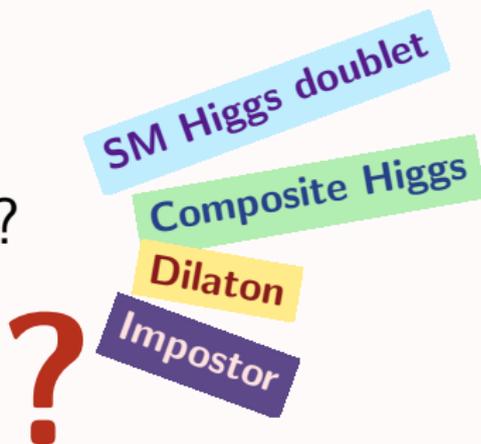
HEFT

... what kind of **Higgs**?



HEFT

... what kind of **Higgs**?



can **EFTs** help in this identification?

Strategy

So far data consistent with the SM \rightarrow linear $SU(2)_L$ doublet

Crucial to keep looking for possible departures!

*another
representation?*

dynamical origin?

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the
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is the **most generic** way to couple the Higgs
to the SM Goldstone bosons

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the **chiral effective Lagrangian**

is the **most generic** way to couple the Higgs
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linear EW doublet
[elementary Higgs]

every BSM scenario

generic singlet
[dilaton]

pGB EW doublet
[composite Higgs]

pGB EW singlet
[exotic]

corresponds to a **specific limit** of the chiral EFT description

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different BSM signals expected!

Chiral EFT: basic formalism

Appelquist, Carazzone (1980)
Longhitano (1980, 1981)

Goldstone bosons: in a bidoublet of $SU(2)_L \times U(1)_Y$

$$\mathbf{U}(x) = e^{i\pi^a(x)\sigma^a/f}, \quad \mathbf{U}(x) \mapsto L\mathbf{U}(x)R^\dagger.$$

Higgs boson: generically a gauge singlet $h(x)$.

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Important: three scales!

- Λ new resonances
- f Goldstone bosons
- v EWSB

$$\xi = v^2/f^2$$

non-linearity parameter
(not physical!)

- $\xi = 1$ technicolor
(non-linear EWSB)
- $\xi = 0$ linear EWSB

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Building blocks for the Lagrangian:

GBs $\mathbf{V}_\mu = \mathbf{D}_\mu \mathbf{U} \mathbf{U}^\dagger, \quad \mathbf{V}_\mu \mapsto L \mathbf{V}_\mu L^\dagger$

$\mathbf{T} = \mathbf{U} \sigma^3 \mathbf{U}^\dagger, \quad \mathbf{T} \mapsto L \mathbf{T} L^\dagger \rightarrow$ ~~Custodial sym.~~

Higgs $\mathcal{F}(\mathbf{h}) = 1 + 2a_i \frac{h}{v} + b_i \frac{h^2}{v^2} + \dots \quad \partial_\mu \mathcal{F}(\mathbf{h})$

Grinstein, Trott (2007)

The non-linear effective Lagrangian

With this notation, the **SM Lagrangian** reads

$$\begin{aligned}\mathcal{L}_{\text{SM}} = & (\text{kinetic terms for } \psi, W, Z, \mathcal{G}) + \\ & + \frac{1}{2} \partial_\mu h \partial^\mu h - V(h) + \\ & - \frac{(v+h)^2}{4} \text{Tr}[\mathbf{V}_\mu \mathbf{V}^\mu] + \quad \text{GB kinetic terms} \\ & \quad \quad \quad \text{gauge bosons' masses} \\ & - \frac{v+h}{\sqrt{2}} [\bar{\psi}_L \mathbf{U} Y \psi_R + \text{h.c.}] \quad \text{Yukawas}\end{aligned}$$

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BSM effects: effective operators with up to 4 derivatives

$$\Delta\mathcal{L} = \sum_i c_i \mathcal{P}_i$$

Bosonic sector
CP even

$$\mathcal{P}_W = -\frac{g^2}{4} W_{\mu\nu}^a W^{a\mu\nu} \mathcal{F}_W(h)$$

$$\mathcal{P}_B = -\frac{g'^2}{4} B_{\mu\nu} B^{\mu\nu} \mathcal{F}_B(h)$$

$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

$$\mathcal{P}_{DH} = \frac{1}{v^4} (\partial_\mu h \partial^\mu h)^2 \mathcal{F}_{DH}(h)$$

$$\mathcal{P}_{\Delta H} = \frac{1}{v^3} (\partial_\mu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$

$$\mathcal{P}_{\square H} = \frac{1}{v^2} (\square h) (\square h) \mathcal{F}_{\square H}(h)$$

$$\mathcal{P}_C = -\frac{v^2}{4} \text{Tr}(\mathbf{V}^\mu \mathbf{V}_\mu) \mathcal{F}_C(h)$$

$$\mathcal{P}_T = \frac{v^2}{4} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \mathcal{F}_T(h)$$

$$\mathcal{P}_1 = gg' B_{\mu\nu} \text{Tr}(\mathbf{T} W^{\mu\nu}) \mathcal{F}_1(h)$$

$$\mathcal{P}_2 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_2(h)$$

$$\mathcal{P}_3 = ig \text{Tr}(W_{\mu\nu} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_3(h)$$

$$\mathcal{P}_4 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_4(h)$$

$$\mathcal{P}_5 = ig \text{Tr}(W_{\mu\nu} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_5(h)$$

$$\mathcal{P}_6 = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu))^2 \mathcal{F}_6(h)$$

$$\mathcal{P}_7 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) \partial_\nu \partial^\nu \mathcal{F}_7(h)$$

$$\mathcal{P}_8 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \partial^\mu \mathcal{F}_8(h) \partial^\nu \mathcal{F}'_8(h)$$

$$\mathcal{P}_9 = \text{Tr}((\mathcal{D}_\mu \mathbf{V}^\mu)^2) \mathcal{F}_9(h)$$

$$\mathcal{P}_{10} = \text{Tr}(\mathbf{V}_\nu \mathcal{D}_\mu \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{10}(h)$$

$$\mathcal{P}_{11} = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu))^2 \mathcal{F}_{11}(h)$$

$$\mathcal{P}_{12} = g^2 \text{Tr}(\mathbf{T} W_{\mu\nu})^2 \mathcal{F}_{12}(h)$$

$$\mathcal{P}_{13} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_{13}(h)$$

$$\mathcal{P}_{14} = g \varepsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h)$$

$$\mathcal{P}_{15} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathcal{D}_\nu \mathbf{V}^\nu) \mathcal{F}_{15}(h)$$

$$\mathcal{P}_{16} = \text{Tr}([\mathbf{T}, \mathbf{V}_\nu] \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{16}(h)$$

$$\mathcal{P}_{17} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{17}(h)$$

$$\mathcal{P}_{18} = \text{Tr}(\mathbf{T} [\mathbf{V}_\mu, \mathbf{V}_\nu]) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{18}(h)$$

$$\mathcal{P}_{19} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\nu \mathcal{F}_{19}(h)$$

$$\mathcal{P}_{20} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) \partial_\nu \mathcal{F}_{20}(h) \partial^\nu \mathcal{F}'_{20}(h)$$

$$\mathcal{P}_{21} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \mathcal{F}_{21}(h) \partial^\nu \mathcal{F}'_{21}(h)$$

$$\mathcal{P}_{22} = \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\mu \mathcal{F}_{22}(h) \partial^\nu \mathcal{F}'_{22}(h)$$

$$\mathcal{P}_{23} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) (\text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{23}(h)$$

$$\mathcal{P}_{24} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{24}(h)$$

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$$\mathcal{P}_B = -\frac{g'^2}{4} B_{\mu\nu} B^{\mu\nu} \mathcal{F}_B(h)$$

$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

$$\mathcal{P}_{DH} = \frac{1}{v^4} (\partial_\mu h \partial^\mu h)^2 \mathcal{F}_{DH}(h)$$

$$\mathcal{P}_{\Delta H} = \frac{1}{v^3} (\partial_\mu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$

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$$\mathcal{P}_C = -\frac{v^2}{4} \text{Tr}(\mathbf{V}^\mu \mathbf{V}_\mu) \mathcal{F}_C(h)$$

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$$\mathcal{P}_1 = gg' B_{\mu\nu} \text{Tr}(\mathbf{TW}^{\mu\nu}) \mathcal{F}_1(h)$$

$$\mathcal{P}_2 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T}[\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_2(h)$$

$$\mathcal{P}_3 = ig \text{Tr}(W_{\mu\nu} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_3(h)$$

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Appelquist, Bernard (1980)
 Longhitano (1980,1981)
 Feruglio (1993)

No h :
ALF basis

**Custodial
symmetry
imposed**

$$\mathcal{P}_W = -\frac{g^2}{4} W_{\mu\nu}^a W^{a\mu\nu} \mathcal{F}_W(h)$$

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Massless
fermions

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$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

$$\mathcal{P}_{DH} = \frac{1}{v^4} (\partial_\mu h \partial^\mu h)^2 \mathcal{F}_{DH}(h)$$

$$\mathcal{P}_{\Delta H} = \frac{1}{v^3} (\partial_\mu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$

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$$\mathcal{P}_C = -\frac{v^2}{4} \text{Tr}(\mathbf{V}^\mu \mathbf{V}_\mu) \mathcal{F}_C(h)$$

$$\mathcal{P}_T = \frac{v^2}{4} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \mathcal{F}_T(h)$$

$$\mathcal{P}_1 = gg' B_{\mu\nu} \text{Tr}(\mathbf{T} W^{\mu\nu}) \mathcal{F}_1(h)$$

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$$\mathcal{P}_{13} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_{13}(h)$$

$$\mathcal{P}_{14} = g \varepsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h)$$

$$\mathcal{P}_{15} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathcal{D}_\nu \mathbf{V}^\nu) \mathcal{F}_{15}(h)$$

$$\mathcal{P}_{16} = \text{Tr}([\mathbf{T}, \mathbf{V}_\nu] \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{16}(h)$$

$$\mathcal{P}_{17} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{17}(h)$$

$$\mathcal{P}_{18} = \text{Tr}(\mathbf{T} [\mathbf{V}_\mu, \mathbf{V}_\nu]) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{18}(h)$$

$$\mathcal{P}_{19} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\nu \mathcal{F}_{19}(h)$$

$$\mathcal{P}_{20} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) \partial_\nu \mathcal{F}_{20}(h) \partial^\nu \mathcal{F}_{20}'(h)$$

$$\mathcal{P}_{21} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \mathcal{F}_{21}(h) \partial^\nu \mathcal{F}_{21}'(h)$$

$$\mathcal{P}_{22} = \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\mu \mathcal{F}_{22}(h) \partial^\nu \mathcal{F}_{22}'(h)$$

$$\mathcal{P}_{23} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) (\text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{23}(h)$$

$$\mathcal{P}_{24} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{24}(h)$$

$$\mathcal{P}_{25} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \partial^\nu \mathcal{F}_{25}(h)$$

$$\mathcal{P}_{26} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{26}(h)$$

$$\mathcal{P}_W = -\frac{g^2}{4} W_{\mu\nu}^a W^{a\mu\nu} \mathcal{F}_W(h)$$

$$\mathcal{P}_B = -\frac{g'^2}{4} B_{\mu\nu} B^{\mu\nu} \mathcal{F}_B(h)$$

$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

$$\mathcal{P}_{DH} = \frac{1}{v^4} (\partial_\mu h \partial^\mu h)^2 \mathcal{F}_{DH}(h)$$

$$\mathcal{P}_{\Delta H} = \frac{1}{v^3} (\partial_\mu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$

$$\mathcal{P}_{\square H} = \frac{1}{v^2} (\square h) (\square h) \mathcal{F}_{\square H}(h)$$

$$\mathcal{P}_C = -\frac{v^2}{4} \text{Tr}(\mathbf{V}^\mu \mathbf{V}_\mu) \mathcal{F}_C(h)$$

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$$\mathcal{P}_1 = gg' B_{\mu\nu} \text{Tr}(\mathbf{T} W^{\mu\nu}) \mathcal{F}_1(h)$$

$$\mathcal{P}_2 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_2(h)$$

$$\mathcal{P}_3 = ig \text{Tr}(W_{\mu\nu} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_3(h)$$

$$\mathcal{P}_4 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_4(h)$$

$$\mathcal{P}_5 = ig \text{Tr}(W_{\mu\nu} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_5(h)$$

$$\mathcal{P}_6 = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu))^2 \mathcal{F}_6(h)$$

$$\mathcal{P}_7 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) \partial_\nu \partial^\nu \mathcal{F}_7(h)$$

$$\mathcal{P}_8 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \partial^\mu \mathcal{F}_8(h) \partial^\nu \mathcal{F}'_8(h)$$

$$\mathcal{P}_9 = \text{Tr}((\mathcal{D}_\mu \mathbf{V}^\mu)^2) \mathcal{F}_9(h)$$

$$\mathcal{P}_{10} = \text{Tr}(\mathbf{V}_\nu \mathcal{D}_\mu \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{10}(h)$$

$$\mathcal{P}_{11} = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu))^2 \mathcal{F}_{11}(h)$$

$$\mathcal{P}_{12} = g^2 \text{Tr}(\mathbf{T} W_{\mu\nu})^2 \mathcal{F}_{12}(h)$$

$$\mathcal{P}_{13} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_{13}(h)$$

$$\mathcal{P}_{14} = g \varepsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h)$$

$$\mathcal{P}_{15} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathcal{D}_\nu \mathbf{V}^\nu) \mathcal{F}_{15}(h)$$

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$$\mathcal{P}_{19} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\nu \mathcal{F}_{19}(h)$$

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$$\mathcal{P}_{21} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \mathcal{F}_{21}(h) \partial^\nu \mathcal{F}'_{21}(h)$$

$$\mathcal{P}_{22} = \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\mu \mathcal{F}_{22}(h) \partial^\nu \mathcal{F}'_{22}(h)$$

$$\mathcal{P}_{23} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) (\text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{23}(h)$$

$$\mathcal{P}_{24} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{24}(h)$$

$$\mathcal{P}_{25} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \partial^\nu \mathcal{F}_{25}(h)$$

$$\mathcal{P}_{26} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{26}(h)$$

Example I

Coupling

$$A_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$A_{\mu\nu} W^{+\mu} W^{-\nu} h$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$$



generic singlet

$$2 c_2 + c_3$$

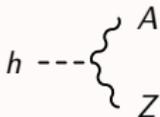
$$2 c_2 a_2 + c_3 a_3$$

$$-2 t_\theta^2 c_2 + c_3$$

$$-2 t_\theta^2 c_2 a_2 + c_3 a_3$$

Coupling

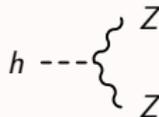
$$A_{\mu\nu} Z^\mu \partial^\nu h$$



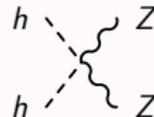
$$A_{\mu\nu} Z^\mu h \partial^\nu h$$



$$Z_{\mu\nu} Z^\mu \partial^\nu h$$



$$Z_{\mu\nu} Z^\mu h \partial^\nu h$$



generic singlet

$$2 c_4 a_4 + c_5 a_5$$

$$2 c_4 b_4 + c_5 b_5$$

$$2 t_\theta c_4 a_4 - c_5 a_5$$

$$2 t_\theta c_4 b_4 - c_5 b_5$$

Example 1

Coupling

$$A_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$A_{\mu\nu} W^{+\mu} W^{-\nu} h$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$$



generic singlet

$$2 c_2 + c_3$$

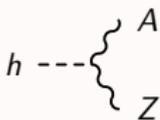
$$2 c_2 a_2 + c_3 a_3$$

$$-2t_\theta^2 c_2 + c_3$$

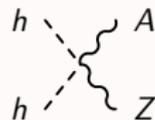
$$-2t_\theta^2 c_2 a_2 + c_3 a_3$$

Coupling

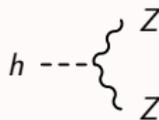
$$A_{\mu\nu} Z^\mu \partial^\nu h$$



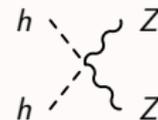
$$A_{\mu\nu} Z^\mu h \partial^\nu h$$



$$Z_{\mu\nu} Z^\mu \partial^\nu h$$



$$Z_{\mu\nu} Z^\mu h \partial^\nu h$$



generic singlet

$$2 c_4 a_4 + c_5 a_5$$

$$2 c_4 b_4 + c_5 b_5$$

$$2t_\theta c_4 a_4 - c_5 a_5$$

$$2t_\theta c_4 b_4 - c_5 b_5$$

Decorrelations:

- ▶ $\mathcal{P}_{2,3,4,5}$ are independent
- ▶ $\mathcal{F}_{2,3,4,5}(h)$ are arbitrary

If the Higgs is an elementary $SU(2)_L$ doublet

How is the linear $SU(2)$ doublet described in the chiral formalism?

- ▶ The chiral expansion shall converge to the linear one

If the Higgs is an elementary $SU(2)_L$ doublet

How is the linear $SU(2)$ doublet described in the chiral formalism?

- ▶ The chiral expansion shall converge to the linear one

For example:

HISZ linear basis

Buchmüller, Wyler (1986)

Hagiwara, Ishihara, Szalapski, Zeppenfeld (1993)

$$\mathcal{O}_{GG} = -\frac{g_s^2}{4} \Phi^\dagger \Phi G_{\mu\nu} G^{\mu\nu}$$

$$\mathcal{O}_{WW} = -\frac{g^2}{4} \Phi^\dagger W_{\mu\nu} W^{\mu\nu} \Phi$$

$$\mathcal{O}_{BB} = -\frac{g'^2}{4} \Phi^\dagger B_{\mu\nu} B^{\mu\nu} \Phi$$

$$\mathcal{O}_{BW} = -\frac{gg'}{4} \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

$$\mathcal{O}_W = \frac{ig}{2} (\mathbf{D}_\mu \Phi)^\dagger W^{\mu\nu} (\mathbf{D}_\nu \Phi)$$

$$\mathcal{O}_B = \frac{ig'}{2} (\mathbf{D}_\mu \Phi)^\dagger B^{\mu\nu} (\mathbf{D}_\nu \Phi)$$

$$\mathcal{O}_{\Phi 1} = (\mathbf{D}_\mu \Phi)^\dagger \Phi \Phi^\dagger (\mathbf{D}^\mu \Phi)$$

$$\mathcal{O}_{\Phi 2} = \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi)$$

$$\mathcal{O}_{\Phi 3} = \frac{1}{3} (\Phi^\dagger \Phi)^3$$

$$\mathcal{O}_{\Phi 4} = (\mathbf{D}_\mu \Phi)^\dagger (\mathbf{D}^\mu \Phi) (\Phi^\dagger \Phi)$$

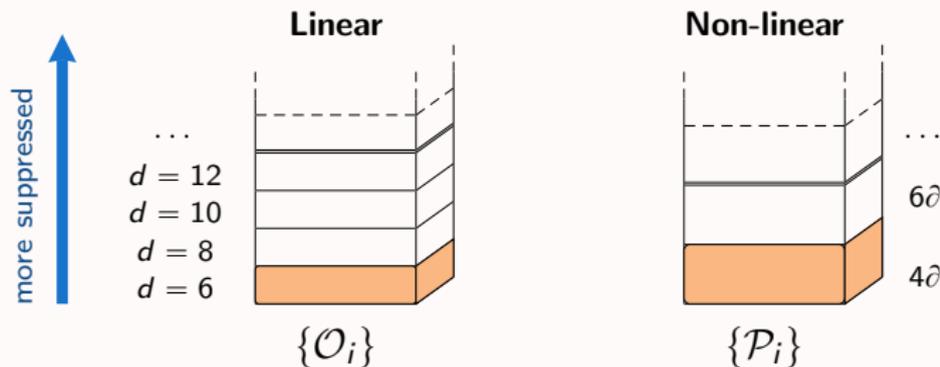
$$\mathcal{O}_{\square\Phi} = (D_\mu D^\mu \Phi)^\dagger (D_\nu D^\nu \Phi)$$

Grzadkowski, Iskrzynski, Misiak, Rosiek (2010)

 see talk by J. Gonzalez-Fraile

Linear - chiral correspondence

Two towers of operators:

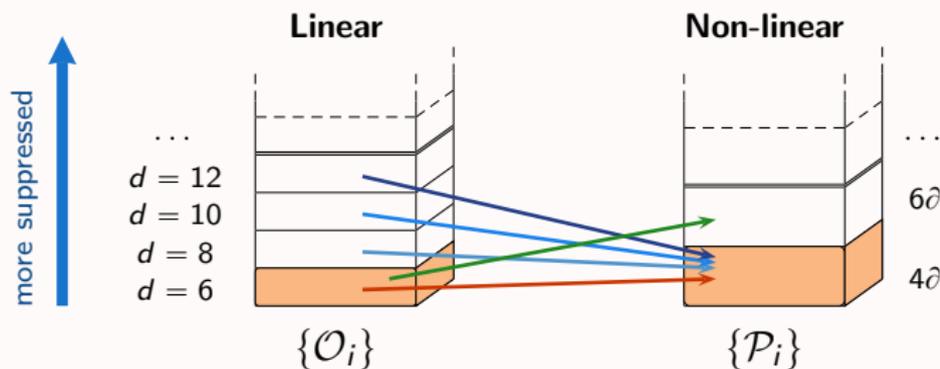


Correspondence $O_i \rightarrow P_j$

Replace in O_i :
$$\Phi \rightarrow \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Linear - chiral correspondence

Two towers of operators:



Correspondence $O_i \rightarrow P_j$

Replace in O_i :
$$\Phi \rightarrow \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

The non-linear basis

$$\mathcal{P}_W = -\frac{g^2}{4} W_{\mu\nu}^a W^{a\mu\nu} \mathcal{F}_W(h)$$

$$\mathcal{P}_B = -\frac{g'^2}{4} B_{\mu\nu} B^{\mu\nu} \mathcal{F}_B(h)$$

$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

$$\mathcal{P}_{DH} = \frac{1}{v^4} (\partial_\mu h \partial^\mu h)^2 \mathcal{F}_{DH}(h)$$

$$\mathcal{P}_{\Delta H} = \frac{1}{v^3} (\partial_\mu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$

$$\mathcal{P}_{\square H} = \frac{1}{v^2} (\square h) (\square h) \mathcal{F}_{\square H}(h)$$

$$\mathcal{P}_C = -\frac{v^2}{4} \text{Tr}(\mathbf{V}^\mu \mathbf{V}_\mu) \mathcal{F}_C(h)$$

$$\mathcal{P}_T = \frac{v^2}{4} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \mathcal{F}_T(h)$$

$$\mathcal{P}_1 = g g' B_{\mu\nu} \text{Tr}(\mathbf{T} W^{\mu\nu}) \mathcal{F}_1(h)$$

$$\mathcal{P}_2 = i g' B_{\mu\nu} \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_2(h)$$

$$\mathcal{P}_3 = i g \text{Tr}(W_{\mu\nu} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_3(h)$$

$$\mathcal{P}_4 = i g' B_{\mu\nu} \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_4(h)$$

$$\mathcal{P}_5 = i g \text{Tr}(W_{\mu\nu} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_5(h)$$

$$\mathcal{P}_6 = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu))^2 \mathcal{F}_6(h)$$

$$\mathcal{P}_7 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) \partial_\nu \partial^\nu \mathcal{F}_7(h)$$

$$\mathcal{P}_8 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \partial^\mu \mathcal{F}_8(h) \partial^\nu \mathcal{F}'_8(h)$$

$$\mathcal{P}_9 = \text{Tr}((\mathcal{D}_\mu \mathbf{V}^\mu)^2) \mathcal{F}_9(h)$$

$$\mathcal{P}_{10} = \text{Tr}(\mathbf{V}_\nu \mathcal{D}_\mu \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{10}(h)$$

$$\mathcal{P}_{11} = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu))^2 \mathcal{F}_{11}(h)$$

$$\mathcal{P}_{12} = g^2 \text{Tr}(\mathbf{T} W_{\mu\nu})^2 \mathcal{F}_{12}(h)$$

$$\mathcal{P}_{13} = i g \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_{13}(h)$$

$$\mathcal{P}_{14} = g \varepsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h)$$

$$\mathcal{P}_{15} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathcal{D}_\nu \mathbf{V}^\nu) \mathcal{F}_{15}(h)$$

$$\mathcal{P}_{16} = \text{Tr}([\mathbf{T}, \mathbf{V}_\nu] \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{16}(h)$$

$$\mathcal{P}_{17} = i g \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{17}(h)$$

$$\mathcal{P}_{18} = \text{Tr}(\mathbf{T} [\mathbf{V}_\mu, \mathbf{V}_\nu]) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{18}(h)$$

$$\mathcal{P}_{19} = \text{Tr}(\mathbf{T} \mathcal{D}_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\nu \mathcal{F}_{19}(h)$$

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$$\mathcal{P}_{25} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \partial^\nu \mathcal{F}_{25}(h)$$

$$\mathcal{P}_{26} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{26}(h)$$

Expected ξ
weights

ξ^2

ξ

ξ^4

The non-linear basis

Corresponding
to linear
 $d = 6$

$$\mathcal{P}_W = -\frac{g^2}{4} W_{\mu\nu}^a W^{a\mu\nu} \mathcal{F}_W(h)$$

$$\mathcal{P}_B = -\frac{g'^2}{4} B_{\mu\nu} B^{\mu\nu} \mathcal{F}_B(h)$$

$$\mathcal{P}_G = -\frac{g_s^2}{4} G_{\mu\nu}^A G^{A\mu\nu} \mathcal{F}_G(h)$$

$$\mathcal{P}_H = \frac{1}{2} \partial_\mu h \partial^\mu h \mathcal{F}_H(h)$$

~~$$\mathcal{P}_{\square H} = \frac{1}{v} (\partial_\nu h \partial^\mu h)^2 \mathcal{F}_{\square H}(h)$$~~

~~$$\mathcal{P}_{\Delta H} = \frac{1}{v^2} (\partial_\nu h \partial^\mu h) \square \mathcal{F}_{\Delta H}(h)$$~~

$$\mathcal{P}_{\square H} = \frac{1}{v^2} (\square h) (\square h) \mathcal{F}_{\square H}(h)$$

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$$\mathcal{P}_T = \frac{v^2}{4} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \mathcal{F}_T(h)$$

$$\mathcal{P}_1 = gg' B_{\mu\nu} \text{Tr}(\mathbf{T} W^{\mu\nu}) \mathcal{F}_1(h)$$

$$\mathcal{P}_2 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_2(h)$$

$$\mathcal{P}_3 = ig \text{Tr}(W_{\mu\nu} [\mathbf{V}^\mu, \mathbf{V}^\nu]) \mathcal{F}_3(h)$$

$$\mathcal{P}_4 = ig' B_{\mu\nu} \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_4(h)$$

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$$\mathcal{P}_6 = (\text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu))^2 \mathcal{F}_6(h)$$

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$$\mathcal{P}_8 = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \partial^\mu \mathcal{F}_8(h) \partial^\nu \mathcal{F}_8'(h)$$

$$\mathcal{P}_9 = \text{Tr}((D_\mu \mathbf{V}^\mu)^2) \mathcal{F}_9(h)$$

ξ^2

$$\mathcal{P}_{10} = \text{Tr}(\mathbf{V}_\nu D_\mu \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{10}(h)$$

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~~$$\mathcal{P}_{14} = g^2 \epsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h)$$~~

~~$$\mathcal{P}_{15} = \text{Tr}(\mathbf{T} D_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} D_\nu \mathbf{V}^\nu) \mathcal{F}_{15}(h)$$~~

~~$$\mathcal{P}_{16} = \text{Tr}([\mathbf{T}, \mathbf{V}_\nu] D_\mu \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{16}(h)$$~~

~~$$\mathcal{P}_{17} = ig \text{Tr}(\mathbf{T} W_{\mu\nu}) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{17}(h)$$~~

~~$$\mathcal{P}_{18} = \text{Tr}(\mathbf{T} [\mathbf{V}_\mu, \mathbf{V}_\nu]) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \partial^\nu \mathcal{F}_{18}(h)$$~~

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~~$$\mathcal{P}_{22} = \text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu) \partial^\mu \mathcal{F}_{22}(h) \partial^\nu \mathcal{F}_{22}'(h)$$~~

~~$$\mathcal{P}_{23} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}^\mu) (\text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{23}(h)$$~~

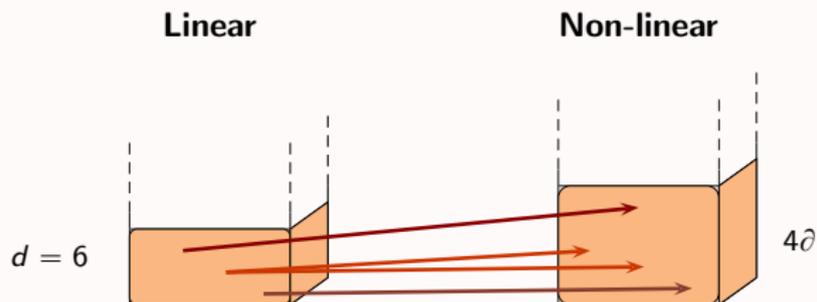
~~$$\mathcal{P}_{24} = \text{Tr}(\mathbf{V}_\mu \mathbf{V}_\nu) \text{Tr}(\mathbf{T} \mathbf{V}^\mu) \text{Tr}(\mathbf{T} \mathbf{V}^\nu) \mathcal{F}_{24}(h)$$~~

~~$$\mathcal{P}_{25} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu))^2 \partial_\nu \mathcal{F}_{25}(h)$$~~

~~$$\mathcal{P}_{26} = (\text{Tr}(\mathbf{T} \mathbf{V}_\mu) \text{Tr}(\mathbf{T} \mathbf{V}_\nu))^2 \mathcal{F}_{26}(h)$$~~

ξ^4

Correspondence between first orders



10 linear operators of $d = 6$

correspond to

17 chiral operators with 4∂

Linear doublet vs generic singlet

Already at LO the linear EFT has fewer parameters than the chiral one!

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Requirements:

1. Linearity of EWSB ▶ $\xi \rightarrow 0$

linear op. $d = 6$ → chiral op. weighted by ξ
 $d = 8$ → ξ^2
 $d = 10$ → ξ^4

▶ $\mathcal{F}(h) = (1 + h/v)^n \Rightarrow$ constraint on a_i, b_i

Linear doublet vs generic singlet

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 $d = 8 \rightarrow \xi^2$
 $d = 10 \rightarrow \xi^4$

▶ $\mathcal{F}(h) = (1 + h/v)^n \Rightarrow$ constraint on a_i, b_i

2. Gauge structure: $\Phi = \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

▶ fixes the **relative weight** of some operators to reproduce the structure of $D_\mu \Phi \propto (v+h) D_\mu \mathbf{U} + \partial_\mu h \mathbf{U} \Rightarrow$ constraint on c_i

Linear doublet vs generic singlet

For instance, consider

$$\mathcal{O}_B = \frac{ig'}{2} (\mathbf{D}^\mu \Phi)^\dagger B_{\mu\nu} (\mathbf{D}^\nu \Phi)$$

$$\mathcal{O}_W = \frac{ig}{2} (\mathbf{D}^\mu \Phi)^\dagger W_{\mu\nu} (\mathbf{D}^\nu \Phi)$$

Linear doublet vs generic singlet

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$$\begin{aligned}\mathcal{O}_B &= \frac{ig'}{2} (\mathbf{D}^\mu \Phi)^\dagger B_{\mu\nu} (\mathbf{D}^\nu \Phi) \\ \mathcal{O}_W &= \frac{ig}{2} (\mathbf{D}^\mu \Phi)^\dagger W_{\mu\nu} (\mathbf{D}^\nu \Phi)\end{aligned}$$

Replacing $\Phi \rightarrow \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ we get

$$\begin{aligned}\mathcal{O}_B &= \frac{ig'}{16} B_{\mu\nu} \text{Tr}(\mathbf{T}[\mathbf{V}^\mu, \mathbf{V}^\nu]) (v+h)^2 + \frac{ig'}{4} B_{\mu\nu} \text{Tr}(\mathbf{T}\mathbf{V}^\mu) \partial^\nu h (v+h) \\ &= v^2 \left(\frac{\mathcal{P}_2}{16} + \frac{\mathcal{P}_4}{8} \right)_{\mathcal{F}_2 = \mathcal{F}_4 = (1+h/v)^2}\end{aligned}$$

$$\begin{aligned}\mathcal{O}_W &= \frac{ig}{8} \text{Tr}(W_{\mu\nu}[\mathbf{V}^\mu, \mathbf{V}^\nu]) (v+h)^2 - \frac{ig}{2} \text{Tr}(W_{\mu\nu} \mathbf{V}^\mu) \partial^\nu h (v+h) \\ &= v^2 \left(\frac{\mathcal{P}_3}{8} - \frac{\mathcal{P}_5}{4} \right)_{\mathcal{F}_3 = \mathcal{F}_5 = (1+h/v)^2} \Rightarrow \begin{aligned} 2c_2 &= c_4 & (= c_B/8) \\ 2c_3 &= -c_5 & (= c_W/4) \\ a_i &= b_i = 1 \end{aligned}\end{aligned}$$

Back to example I

	$A_{\mu\nu} W^{+\mu} W^{-\nu}$	$A_{\mu\nu} W^{+\mu} W^{-\nu} h$	$Z_{\mu\nu} W^{+\mu} W^{-\nu}$	$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$
Coupling				
generic singlet	$2c_2 + c_3$	$2c_2 a_2 + c_3 a_3$	$-2t_\theta^2 c_2 + c_3$	$-2t_\theta^2 c_2 a_2 + c_3 a_3$
linear doublet	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$
	$A_{\mu\nu} Z^\mu \partial^\nu h$	$A_{\mu\nu} Z^\mu h \partial^\nu h$	$Z_{\mu\nu} Z^\mu \partial^\nu h$	$Z_{\mu\nu} Z^\mu h \partial^\nu h$
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generic singlet	$2c_4 a_4 + c_5 a_5$	$2c_4 b_4 + c_5 b_5$	$2t_\theta^2 c_4 a_4 - c_5 a_5$	$2t_\theta^2 c_4 b_4 - c_5 b_5$
linear doublet	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(t_\theta^2 c_B + c_W)$	$\frac{1}{4}(t_\theta^2 c_B + c_W)$

Back to example I

Coupling

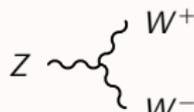
$A_{\mu\nu} W^{+\mu} W^{-\nu}$



$A_{\mu\nu} W^{+\mu} W^{-\nu} h$



$Z_{\mu\nu} W^{+\mu} W^{-\nu}$



$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$



generic singlet

$2c_2 + c_3$

$2c_2 a_2 + c_3 a_3$

$-2t_\theta^2 c_2 + c_3$

$-2t_\theta^2 c_2 a_2 + c_3 a_3$

linear doublet

$\frac{1}{8}(c_B + c_W)$

$\frac{1}{8}(c_B + c_W)$

$\frac{1}{8}(-t_\theta^2 c_B + c_W)$

$\frac{1}{8}(-t_\theta^2 c_B + c_W)$

Complete equivalence for

$2c_2 = c_4 = c_B/8$

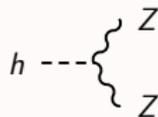
$2c_3 = -c_5 = c_W/4$

$a_i = b_i = 1$

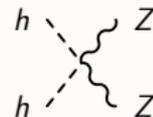
$A_{\mu\nu} Z^\mu h \partial^\nu h$



$Z_{\mu\nu} Z^\mu \partial^\nu h$



$Z_{\mu\nu} Z^\mu h \partial^\nu h$



generic singlet

$2c_4 a_4 + c_5 a_5$

$2c_4 b_4 + c_5 b_5$

$2t_\theta^2 c_4 a_4 - c_5 a_5$

$2t_\theta^2 c_4 b_4 - c_5 b_5$

linear doublet

$\frac{1}{4}(c_B - c_W)$

$\frac{1}{4}(c_B - c_W)$

$\frac{1}{4}(t_\theta^2 c_B + c_W)$

$\frac{1}{4}(t_\theta^2 c_B + c_W)$

Back to example I

Coupling

$$A_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$A_{\mu\nu} W^{+\mu} W^{-\nu} h$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu}$$



$$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$$



generic singlet

$$2c_2 + c_3$$

$$2c_2 a_2 + c_3 a_3$$

$$-2t_\theta^2 c_2 + c_3$$

$$-2t_\theta^2 c_2 a_2 + c_3 a_3$$

linear doublet

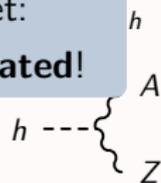
$$\frac{1}{8}(c_B + c_W)$$

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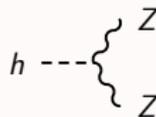
linear doublet:
couplings **correlated!**



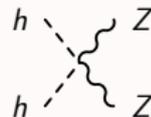
$$A_{\mu\nu} Z^\mu h \partial^\nu h$$



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$$Z_{\mu\nu} Z^\mu h \partial^\nu h$$



generic singlet

$$2c_4 a_4 + c_5 a_5$$

$$2c_4 b_4 + c_5 b_5$$

$$2t_\theta^2 c_4 a_4 - c_5 a_5$$

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linear doublet

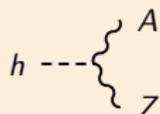
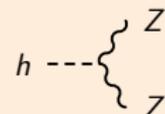
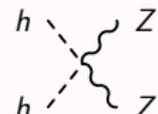
$$\frac{1}{4}(c_B - c_W)$$

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$$\frac{1}{4}(t_\theta^2 c_B + c_W)$$

$$\frac{1}{4}(t_\theta^2 c_B + c_W)$$

Back to example I

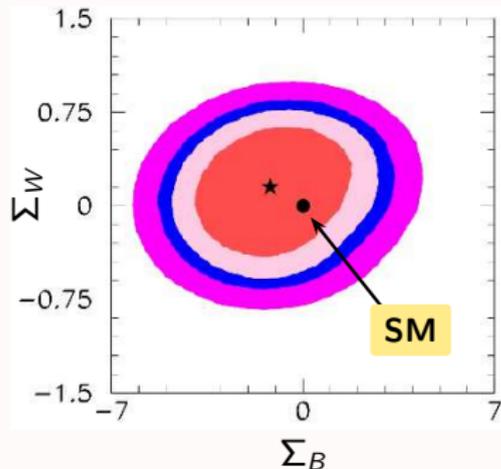
Coupling	$A_{\mu\nu} W^{+\mu} W^{-\nu}$ 	$A_{\mu\nu} W^{+\mu} W^{-\nu} h$ 	$Z_{\mu\nu} W^{+\mu} W^{-\nu}$ 	$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$ 
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generic singlet	$2c_4 a_4 + c_5 a_5$	$2c_4 b_4 + c_5 b_5$	$2t_\theta^2 c_4 a_4 - c_5 a_5$	$2t_\theta^2 c_4 b_4 - c_5 b_5$
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Combining TGV + Higgs data

A BSM sensor

$$\Sigma_B \equiv 4(2c_2 + c_2 a_4)$$

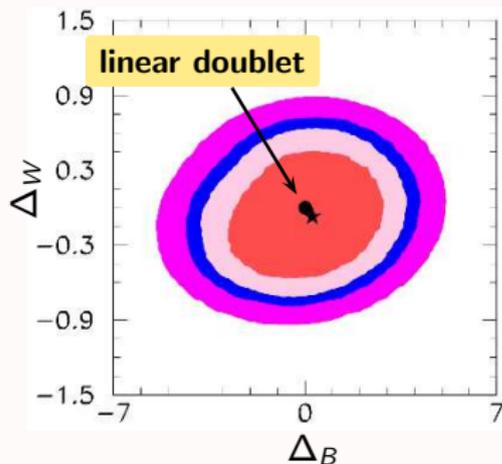
$$\Sigma_W \equiv 2(2c_3 - c_5 a_5)$$



A linear vs non-linear discriminator

$$\Delta_B \equiv 4(2c_2 - c_4 a_4)$$

$$\Delta_W \equiv 2(2c_3 + c_5 a_5)$$



χ^2 dependence after marginalizing over the other chiral parameters
Datasets: TGV (LEP) and HVV couplings (D0+CDF+LHC7+LHC8).
Colored areas: 68, 90, 95, 99% CL

Example II

In the **linear** Lagrangian

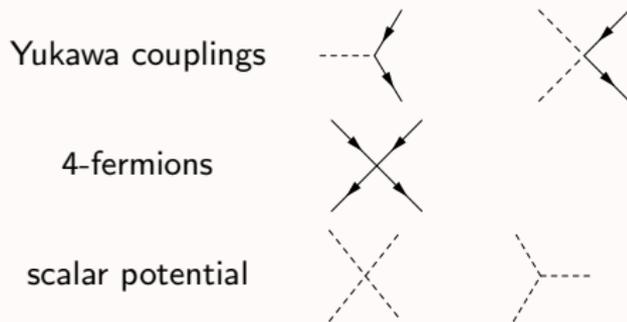
$$\mathcal{O}_{\square\Phi} = (D_\mu D^\mu \Phi)^\dagger (D_\nu D^\nu \Phi)$$

Example II

In the **linear** Lagrangian

$$\mathcal{O}_{\square\Phi} = (D_\mu D^\mu \Phi)^\dagger (D_\nu D^\nu \Phi)$$

Applying the EOM for Φ \blacktriangleright $c_{\square\Phi}$ contributes to

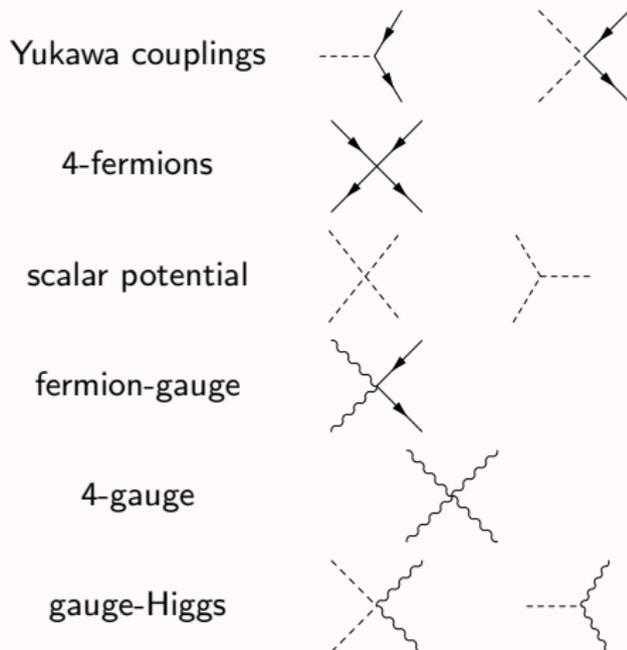


Example II

In the **chiral** Lagrangian

$$\mathcal{O}_{\square h} = \frac{1}{v^2} (\partial_\mu \partial^\mu h) (\partial_\nu \partial^\nu h)$$

Applying the EOM for h \blacktriangleright $c_{\square h}$ contributes to



Example II

In the **chiral** Lagrangian

$$\mathcal{O}_{\square h} = \frac{1}{v^2} (\partial_\mu \partial^\mu h) (\partial_\nu \partial^\nu h)$$

Applying the EOM for h \blacktriangleright $c_{\square h}$ contributes to

also from $c_{\square\phi}$

Yukawa couplings



4-fermions



scalar potential



fermion-gauge



4-gauge



gauge-Higgs



Example II

In the **chiral** Lagrangian

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Applying the EOM for h \blacktriangleright $c_{\square h}$ contributes to



also from $c_{\square\phi}$

different pattern!

Yukawa couplings



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Example II

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Yukawa couplings



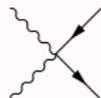
4-fermions



scalar potential



fermion-gauge



no contribution
in the linear EFT!

4-gauge



gauge-Higgs



Example II

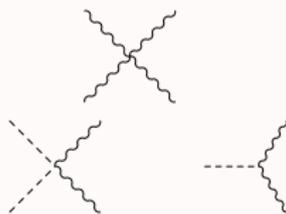
Replacing $\Phi \rightarrow \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$:

$$\mathcal{O}_{\square\Phi} \rightarrow \mathcal{P}_{\square h} + v^2 \left(\frac{1}{8} \mathcal{P}_6 + \frac{1}{4} \mathcal{P}_7 - \mathcal{P}_8 - \frac{1}{4} \mathcal{P}_9 - \frac{1}{2} \mathcal{P}_{10} \right)$$

where:

$c_6, c_9 \rightarrow$ 4-gauge

$c_7, c_8, c_9, c_{10} \rightarrow$ gauge-Higgs

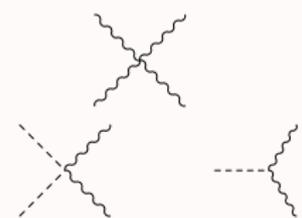


Example II

Replacing $\Phi \rightarrow \frac{v+h}{\sqrt{2}} \mathbf{U} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$:

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where:

$$\begin{array}{ll} c_6, c_9 & \rightarrow \quad 4\text{-gauge} \\ c_7, c_8, c_9, c_{10} & \rightarrow \quad \text{gauge-Higgs} \end{array}$$


The diagrams show two types of four-point interactions. The first, labeled '4-gauge', consists of four wavy lines meeting at a central point in a cross configuration. The second, labeled 'gauge-Higgs', consists of two wavy lines and two dashed lines meeting at a central point in a cross configuration.

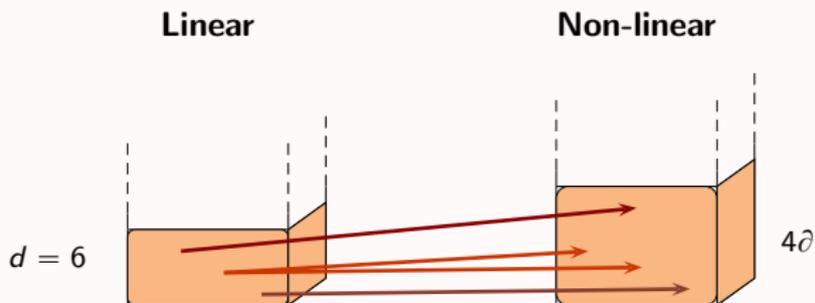
The chiral description is physically equivalent to the linear one *iff*:

$$\begin{aligned} v^2 c_{\square h} &= 8c_6 = 4c_7 = -c_8 = -4c_9 = -2c_{10} \\ a_i &= b_i = 1, \quad i = 6, \dots, 10 \end{aligned}$$

contributions to **QGV**, **gauge-Higgs** and **gauge-fermion** couplings
signal deviations from the $SU(2)_L$ doublet structure

► e.g. impact on off-shell $gg \rightarrow h^* \rightarrow VV$ from $\mathcal{P}_7 = \text{Tr}[\mathbf{V}_\mu \mathbf{V}^\mu] \partial_\nu \partial^\nu \mathcal{F}_7(h)$

Correspondence between first orders



10 linear operators of $d = 6$

correspond to

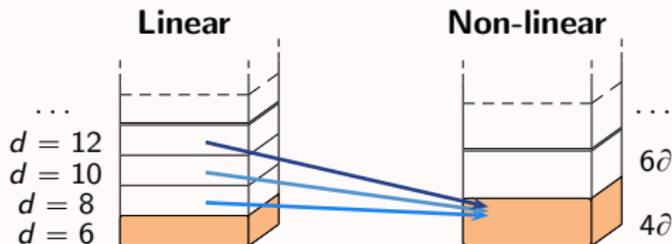
$$\mathcal{O}_B \rightarrow \frac{\mathcal{P}_2}{16} + \frac{\mathcal{P}_4}{8}$$

$$\mathcal{O}_W \rightarrow \frac{\mathcal{P}_3}{8} - \frac{\mathcal{P}_5}{4}$$

17 chiral operators with $4d$

$$\mathcal{O}_{\square\Phi} \rightarrow \mathcal{P}_{\square h} + \frac{\mathcal{P}_6}{8} + \frac{\mathcal{P}_7}{4} - \mathcal{P}_8 - \frac{\mathcal{P}_9}{4} - \frac{\mathcal{P}_{10}}{2}$$

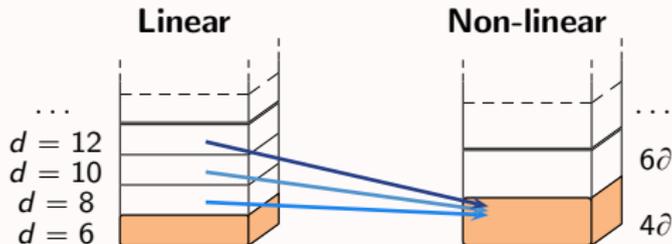
Example III



Effects that are expected to be

- ▶ leading-order corrections in the most general non-linear expansion
- ▶ higher-order corrections in the linear series

Example III



Effects that are expected to be

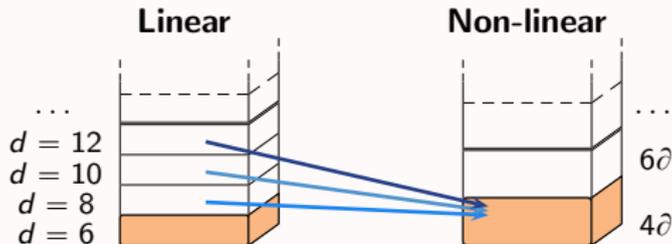
- ▶ leading-order corrections in the most general non-linear expansion
- ▶ higher-order corrections in the linear series

$$\varepsilon^{\mu\nu\rho\lambda} \left(\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\rho \Phi \right) \left(\Phi^\dagger \sigma_i \overleftrightarrow{\mathbf{D}}_\lambda \Phi \right) W_{\mu\nu}^i \quad d = 8$$



$$\mathcal{P}_{14} = g \varepsilon^{\mu\nu\rho\lambda} \text{Tr}(\mathbf{T}\mathbf{V}_\mu) \text{Tr}(\mathbf{V}_\nu W_{\rho\lambda}) \mathcal{F}_{14}(h) \quad 4\partial$$

Example III



Effects that are expected to be

- ▶ leading-order corrections in the most general non-linear expansion
- ▶ higher-order corrections in the linear series

$$\mathcal{P}_{14} \rightarrow Z_\rho \left\{ \begin{array}{l} W_\mu^+ \\ W_\nu^- \end{array} \right. - \frac{g^3 c_{14}}{2c_\theta} \varepsilon^{\mu\nu\rho\lambda} \partial_\mu W_\nu^+ W_\rho^- Z_\lambda + \text{h.c.}$$

Warning: this operator breaks custodial symmetry.

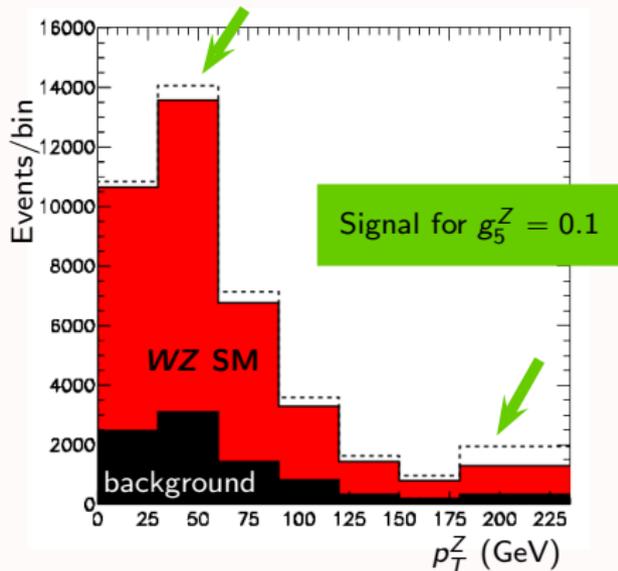
Expected LHC sensitivity

$$g_5^Z = g^2 c_{14} / 2c_\theta^2$$

Current best bound at 95% CL

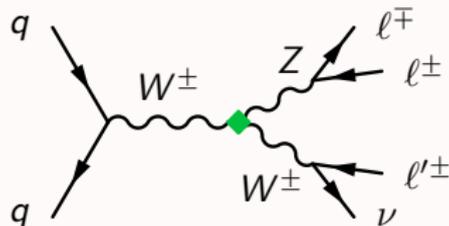
$$g_5^Z \in [-0.08, 0.04]$$

Dawson, Valencia (1994)



Simulation analysis

- ▶ WZ pair production



- ▶ binned analysis of p_T^Z distribution
- ▶ Result (95% CL)

dataset: 7+8+14 TeV
(4.7+19.6+300 fb^{-1})

$$g_5^Z \in [-0.033, 0.028]$$

If the Higgs is a pGB embedded in a doublet

So far we have compared the phenomenology of the Higgs as a

linear $SU(2)_L$ doublet

vs.

generic singlet

But what if the Higgs is a pseudo Goldstone boson ?

If the Higgs is a pGB embedded in a doublet

So far we have compared the phenomenology of the Higgs as a

linear $SU(2)_L$ doublet

vs.

generic singlet

But what if the Higgs is a pseudo Goldstone boson ?

General expectations:

1. the functions $\mathcal{F}(h)$ contain an infinite series of $\frac{h}{f} = \sqrt{\xi} \frac{h}{v}$
e.g. trigonometric

2. if h embedded in a $SU(2)_L$ doublet \Rightarrow

in concrete models, the linear expansion is recovered for $\xi \rightarrow 0$

Results from specific CH models

In specific models:

$SU(5)/SO(5)$ and $SO(5)/SO(4)$

$$c_2 \mathcal{F}_2(h) = \tilde{c}_2 \sin^2 \frac{\varphi}{2f}$$

$$c_4 \mathcal{F}_4(h) = 2 \tilde{c}_2 \sin^2 \frac{\varphi}{2f}$$

Georgi, Kaplan (1984)
Agashe, Contino, Pomarol (2004)

Alonso, IB, Gavela, Merlo, Rigolin (2014)
[hep-ph/1409.1589]

 see talk by **R. Alonso**

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$$c_4 \mathcal{F}_4(h) = 2 \tilde{c}_2 \sin^2 \frac{\varphi}{2f}$$

▶ $\mathcal{F}_2(h) = \mathcal{F}_4(h) \neq (1 + h/v)^2$

Georgi, Kaplan (1984)
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[hep-ph/1409.1589]

→ see talk by **R. Alonso**

- ▶ $\mathcal{F}_2(h) = \mathcal{F}_4(h) \neq (1 + h/v)^2$
- ▶ expanding:

$$\sin^2 \frac{\varphi}{2f} = \frac{\xi}{4} \left[1 + \frac{2h}{v} \sqrt{1 - \frac{\xi}{4}} + \frac{h^2}{v^2} \left(1 - \frac{\xi}{2} \right) \right] + \mathcal{O}(h^3) = \frac{\xi}{4} \frac{(v+h)^2}{v^2} + \mathcal{O}(\xi^2)$$

- ▶ ξ is a parameter of the model
- ▶ $\mathcal{F}(h) \rightarrow \xi(1 + h/v)^2$ for $\xi \rightarrow 0$

Results from specific CH models

In specific models:

$SU(5)/SO(5)$ and $SO(5)/SO(4)$

$$c_2 \mathcal{F}_2(h) = \tilde{c}_2 \sin^2 \frac{\varphi}{2f}$$

$$c_4 \mathcal{F}_4(h) = 2 \tilde{c}_2 \sin^2 \frac{\varphi}{2f}$$

Georgi, Kaplan (1984)
Agashe, Contino, Pomarol (2004)

Alonso, IB, Gavela, Merlo, Rigolin (2014)
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 see talk by **R. Alonso**

- ▶ $\mathcal{F}_2(h) = \mathcal{F}_4(h) \neq (1 + h/v)^2$
- ▶ expanding:

$$\sin^2 \frac{\varphi}{2f} = \frac{\xi}{4} \left[1 + \frac{2h}{v} \sqrt{1 - \frac{\xi}{4}} + \frac{h^2}{v^2} \left(1 - \frac{\xi}{2} \right) \right] + \mathcal{O}(h^3) = \frac{\xi}{4} \frac{(v+h)^2}{v^2} + \mathcal{O}(\xi^2)$$

- ▶ ξ is a parameter of the model
- ▶ $\mathcal{F}(h) \rightarrow \xi(1 + h/v)^2$ for $\xi \rightarrow 0$
- ▶ the condition $2c_2 = c_4$ is verified exactly

Back again to example I

Coupling	$A_{\mu\nu} W^{+\mu} W^{-\nu}$	$A_{\mu\nu} W^{+\mu} W^{-\nu} h$	$Z_{\mu\nu} W^{+\mu} W^{-\nu}$	$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$
generic singlet	$a 2c_2 + c_3$	$2c_2 a_2 + c_3 a_3$	$-2t_\theta^2 c_2 + c_3$	$-2t_\theta^2 c_2 a_2 + c_3 a_3$
linear doublet	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$
pGB doublet	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$

Coupling	$A_{\mu\nu} Z^\mu \partial^\nu h$	$A_{\mu\nu} Z^\mu h \partial^\nu h$	$Z_{\mu\nu} Z^\mu \partial^\nu h$	$Z_{\mu\nu} Z^\mu h \partial^\nu h$
generic singlet	$2c_4 a_4 + c_5 a_5$	$2c_4 b_4 + c_5 b_5$	$2t_\theta^2 c_4 a_4 - c_5 a_5$	$2t_\theta^2 c_4 b_4 - c_5 b_5$
linear doublet	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(t_\theta^2 c_B + c_W)$	$\frac{1}{4}(t_\theta^2 c_B + c_W)$
pGB doublet	$(2\tilde{c}_2 - \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$(2\tilde{c}_2 - \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$	$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$

Back again to example I

Coupling	$A_{\mu\nu} W^{+\mu} W^{-\nu}$	$A_{\mu\nu} W^{+\mu} W^{-\nu} h$	$Z_{\mu\nu} W^{+\mu} W^{-\nu}$	$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$
generic singlet	$a 2c_2 + c_3$	$2c_2 a_2 + c_3 a_3$	$-2t_\theta^2 c_2 + c_3$	$-2t_\theta^2 c_2 a_2 + c_3 a_3$
linear doublet	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$
pGB doublet	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$

pGB constraints

$$2c_2 = c_4 = \tilde{c}_2 \quad a_2 = a_3 = a_4 = a_5 = \sqrt{1 - \frac{\xi}{4}}$$

$$2c_3 = -c_5 = \tilde{c}_3 \quad b_4 = b_5 = 1 - \frac{\xi}{2}$$

	$Z_{\mu\nu} Z^\mu \partial^\nu h$	$Z_{\mu\nu} Z^\mu h \partial^\nu h$
generic singlet	$2c_4 a_4 + c_5 a_5$	$2c_4 b_4 + c_5 b_5$
linear doublet	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(c_B - c_W)$
pGB doublet	$(2\tilde{c}_2 - \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$(2\tilde{c}_2 - \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$
	$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$

Back again to example I

Coupling	$A_{\mu\nu} W^{+\mu} W^{-\nu}$	$A_{\mu\nu} W^{+\mu} W^{-\nu} h$	$Z_{\mu\nu} W^{+\mu} W^{-\nu}$	$Z_{\mu\nu} W^{+\mu} W^{-\nu} h$
generic singlet	$a(2c_2 + c_3)$	$2c_2 a_2 + c_3 a_3$	$-2t_\theta^2 c_2 + c_3$	$-2t_\theta^2 c_2 a_2 + c_3 a_3$
linear doublet	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$	$\frac{1}{8}(-t_\theta^2 c_B + c_W)$
pGB doublet	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(2\tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)$	$\frac{1}{2}(-2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$

pGB doublet: all couplings
correlated!

extra h legs are weighted by a
factor containing ξ

	$A_{\mu\nu} Z^\mu h \partial^\nu h$	$Z_{\mu\nu} Z^\mu \partial^\nu h$	$Z_{\mu\nu} Z^\mu h \partial^\nu h$
generic singlet	$2c_4 a_4 + c_5 a_5$	$2c_4 b_4 + c_5 b_5$	$2t_\theta^2 c_4 a_4 - c_5 a_5$
linear doublet	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(c_B - c_W)$	$\frac{1}{4}(t_\theta^2 c_B + c_W)$
pGB doublet	$(2\tilde{c}_2 - \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$	$(2\tilde{c}_2 - \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$	$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\sqrt{1 - \frac{\xi}{4}}$
			$(2t_\theta^2 \tilde{c}_2 + \tilde{c}_3)\left(1 - \frac{\xi}{2}\right)$

The chiral EFT is a powerful tool to study the Higgs nature: every scenario corresponds to a specific region of the parameter space

Disentangling a **linear $SU(2)_L$ doublet** from a **generic singlet** Higgs:

- ▶ **correlation/decorrelation** of TGV, HVV and HHVV
- ▶ the op. $\mathcal{O}_{\square\Phi}$ **does not affect** QGV, gauge-Higgs or gauge-fermion couplings, the op. $\mathcal{P}_{\square h}$ **does**
- ▶ the g_5^Z coupling is a **NNLO (d=8)/NLO (4 ∂)** effect

How to recognize a **pGB Higgs embedded as a $SU(2)_L$ doublet**:

- ▶ vertices with **different number of h legs** have characteristic relative weights that depend on ξ
- ▶ if $f \gg v$ there may be no visible difference from the linear scenario

Backup slides

Triple gauge vertices

$$\begin{aligned} \mathcal{L}_{WWW} = & -ig_{WWW} \left\{ g_1^V \left(W_{\mu\nu}^+ W^{-\mu} V^\nu - W_\mu^+ V_\nu W^{-\mu\nu} \right) + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} \right. \\ & - ig_5^V \varepsilon^{\mu\nu\rho\sigma} \left(W_\mu^+ \partial_\rho W_\nu^- - W_\nu^- \partial_\rho W_\mu^+ \right) V_\sigma + \\ & \left. + g_6^V \left(\partial_\mu W^{+\mu} W^{-\nu} - \partial_\mu W^{-\mu} W^{+\nu} \right) V_\nu \right\} \end{aligned}$$

$$g_{WWZ} = g \cos \theta, \quad g_{WW\gamma} = e$$

	Coeff. $\times e^2/s_\theta^2$	Chiral	Linear $\times v^2$
$\Delta\kappa_\gamma$	1	$-2c_1 + 2c_2 + c_3 - 4c_{12} + 2c_{13}$	$\frac{1}{8}(c_W + c_B - 2c_{BW})$
Δg_6^γ	1	$-c_9$	—
Δg_1^Z	$\frac{1}{c_\theta^2}$	$\frac{s_\theta^2}{4e^2 c_{2\theta}} c_T + \frac{2s_\theta^2}{c_{2\theta}} c_1 + c_3$	$\frac{1}{8}c_W + \frac{s_\theta^2}{4c_{2\theta}} c_{BW} - \frac{s_\theta^2}{16e^2 c_{2\theta}} c_{\Phi,1}$
$\Delta\kappa_Z$	1	$\frac{s_\theta^2}{e^2 c_{2\theta}} c_T + \frac{4s_\theta^2}{c_{2\theta}} c_1 - \frac{2s_\theta^2}{c t^2} c_2 + c_3 - 4c_{12} + 2c_{13}$	$\frac{1}{8}c_W - \frac{s_\theta^2}{8c t^2} c_B + \frac{s_\theta^2}{2c_{2\theta}} c_{BW} - \frac{s_\theta^2}{4e^2 c_{2\theta}} c_{\Phi,1}$
Δg_5^Z	$\frac{1}{c_\theta^2}$	c_{14}	—
Δg_6^Z	$\frac{1}{c_\theta^2}$	$s_\theta^2 c_9 - c_{16}$	—

$$\begin{aligned}
 \mathcal{L}_{HVV} \equiv & g_{Hgg} G_{\mu\nu}^a G^{a\mu\nu} h + g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} h + g_{HZ\gamma}^{(1)} A_{\mu\nu} Z^\mu \partial^\nu h + g_{HZ\gamma}^{(2)} A_{\mu\nu} Z^{\mu\nu} h \\
 & + g_{HZZ}^{(1)} Z_{\mu\nu} Z^\mu \partial^\nu h + g_{HZZ}^{(2)} Z_{\mu\nu} Z^{\mu\nu} h + g_{HZZ}^{(3)} Z_\mu Z^\mu h + g_{HZZ}^{(4)} Z_\mu Z^\mu \square h \\
 & + g_{HZZ}^{(5)} \partial_\mu Z^\mu Z_\nu \partial^\nu h + g_{HZZ}^{(6)} \partial_\mu Z^\mu \partial_\nu Z^\nu h \\
 & + g_{HWW}^{(1)} (W_{\mu\nu}^+ W^{-\mu} \partial^\nu h + \text{h.c.}) + g_{HWW}^{(2)} W_{\mu\nu}^+ W^{-\mu\nu} h + g_{HWW}^{(3)} W_\mu^+ W^{-\mu} h \\
 & + g_{HWW}^{(4)} W_\mu^+ W^{-\mu} \square h + g_{HWW}^{(5)} (\partial_\mu W^{+\mu} W_\nu^- \partial^\nu h + \text{h.c.}) + g_{HWW}^{(6)} \partial_\mu W^{+\mu} \partial_\nu W^{-\nu} h
 \end{aligned}$$

HVV vertices

	Coeff. $\times e^2/4v$	Chiral	Linear $\times v^2$
Δg_{Hgg}	$\frac{g_s^2}{e^2}$	$-2C_G a_G$	$-4C_{GG}$
$\Delta g_{H\gamma\gamma}$	1	$-2(c_B a_B + c_W a_W) + 8c_1 a_1 + 8c_{12} a_{12}$	$-(c_{BB} + c_{WW}) + c_{BW}$
$\Delta g_{HZ\gamma}^{(1)}$	$\frac{1}{s_{2\theta}}$	$-8(c_5 a_5 + 2c_4 a_4) - 16c_{17} a_{17}$	$2(c_W - c_B)$
$\Delta g_{HZ\gamma}^{(2)}$	$\frac{c_\theta}{s_\theta}$	$4\frac{s_\theta^2}{c_\theta^2} c_B a_B - 4c_W a_W + 8\frac{c_{2\theta}^2}{c_\theta^2} c_1 a_1 + 16c_{12} a_{12}$	$2\frac{s_\theta^2}{c_\theta^2} c_{BB} - 2c_{WW} + \frac{c_{2\theta}^2}{c_\theta^2} c_{BW}$
$\Delta g_{HZZ}^{(1)}$	$\frac{1}{c_\theta}$	$-4\frac{c_\theta^2}{s_\theta^2} c_5 a_5 + 8c_4 a_4 - 8\frac{c_\theta^2}{s_\theta^2} c_{17} a_{17}$	$\frac{c_\theta^2}{s_\theta^2} c_W + c_B$
$\Delta g_{HZZ}^{(2)}$	$-\frac{c_\theta^2}{s_\theta^2}$	$2\frac{s_\theta^4}{c_\theta^4} c_B a_B + 2c_W a_W + 8\frac{s_\theta^2}{c_\theta^2} c_1 a_1 - 8c_{12} a_{12}$	$\frac{s_\theta^4}{c_\theta^4} c_{BB} + c_{WW} + \frac{s_\theta^2}{c_\theta^2} c_{BW}$
$\Delta g_{HZZ}^{(3)}$	$\frac{m_h^2}{e^2}$	$-2c_H + 2c_C(2a_C - 1) - 8c_T(a_T - 1) - 4m_h^2 c_{Ch}$	$c_{\Phi,1} + 2c_{\Phi,4} - 2c_{\Phi,2}$
$\Delta g_{HZZ}^{(4)}$	$-\frac{1}{s_{2\theta}^2}$	$16c_7 a_7 + 32c_{25} a_{25}$	—
$\Delta g_{HZZ}^{(5)}$	$-\frac{1}{s_{2\theta}^2}$	$16c_{10} a_{10} + 32c_{19} a_{19}$	—
$\Delta g_{HZZ}^{(6)}$	$-\frac{1}{s_{2\theta}^2}$	$16c_9 a_9 + 32c_{15} a_{15}$	—
$\Delta g_{HWW}^{(1)}$	$\frac{1}{s_\theta^2}$	$-4c_5 a_5$	c_W
$\Delta g_{HWW}^{(2)}$	$\frac{1}{s_\theta^2}$	$-4c_W a_W$	$-2c_{WW}$
$\Delta g_{HWW}^{(3)}$	$\frac{m_h^2 c_\theta^2}{e^2}$	$-4c_H + 4c_C(2a_C - 1) + \frac{32e^2}{c_{2\theta}^2} c_1 + \frac{16c_\theta^2}{c_{2\theta}^2} c_T - 8m_h^2 c_{Ch} - \frac{32e^2}{s_\theta^2} c_{12}$	$\frac{-2(3c_\theta^2 - s_\theta^2)}{c_{2\theta}} c_{\Phi,1} + 4c_{\Phi,4} - 4c_{\Phi,2} + \frac{4e^2}{c_{2\theta}} c_{BW}$
$\Delta g_{HWW}^{(4)}$	$-\frac{1}{s_\theta^2}$	$8c_7 a_7$	—
$\Delta g_{HWW}^{(5)}$	$-\frac{1}{s_\theta^2}$	$4c_{10} a_{10}$	—
$\Delta g_{HWW}^{(6)}$	$-\frac{1}{s_\theta^2}$	$8c_9 a_9$	—

Quartic gauge vertices

$$\begin{aligned} \mathcal{L}_{4X} \equiv & g^2 \left\{ g_{ZZ}^{(1)} (Z_\mu Z^\mu)^2 + g_{WW}^{(1)} W_\mu^+ W^{+\mu} W_\nu^- W^{-\nu} - g_{WW}^{(2)} (W_\mu^+ W^{-\mu})^2 \right. \\ & + g_{VV'}^{(3)} W^{+\mu} W^{-\nu} (V_\mu V'_\nu + V'_\mu V_\nu) - g_{VV'}^{(4)} W_\nu^+ W^{-\nu} V^\mu V'_\mu \\ & \left. + i g_{VV'}^{(5)} e^{\mu\nu\rho\sigma} W_\mu^+ W_\nu^- V_\rho V'_\sigma \right\} \end{aligned}$$

	Coeff. $\times e^2/4s_\theta^2$	Chiral	Linear $\times v^2$
$\Delta g_{WW}^{(1)}$	1	$\frac{s_\theta^2}{e^2 c_{2\theta}} c_T + \frac{8s_\theta^2}{c_{2\theta}} c_1 + 4c_3 + 2c_{11} - 16c_{12} + 8c_{13}$	$\frac{c_W}{2} + \frac{s_\theta^2}{c_{2\theta}} c_{BW} - \frac{s_\theta^2}{4c_{2\theta} e^2} c_{\Phi 1}$
$\Delta g_{WW}^{(2)}$	1	$\frac{s_\theta^2}{e^2 c_{2\theta}} c_T + \frac{8s_\theta^2}{c_{2\theta}} c_1 + 4c_3 - 4c_6 - \frac{v^2}{2} c_{\square h} - 2c_{11} - 16c_{12} + 8c_{13}$	$\frac{c_W}{2} + \frac{s_\theta^2}{c_{2\theta}} c_{BW} - \frac{s_\theta^2}{4c_{2\theta} e^2} c_{\Phi 1}$
$\Delta g_{ZZ}^{(1)}$	$\frac{1}{c_\theta^4}$	$c_6 + \frac{v^2}{8} c_{\square h} + c_{11} + 2c_{23} + 2c_{24} + 4c_{26}$	—
$\Delta g_{ZZ}^{(3)}$	$\frac{1}{c_\theta^2}$	$\frac{s_\theta^2 c_\theta^2}{e^2 c_{2\theta}} c_T + \frac{2s_\theta^2}{c_{2\theta}} c_1 + 4c_\theta^2 c_3 - 2s_\theta^4 c_9 + 2c_{11} + 4s_\theta^2 c_{16} + 2c_{24}$	$\frac{c_W c_\theta^2}{2} + \frac{s_\theta^2}{4c_{2\theta}} c_{BW} - \frac{s_\theta^2 c_\theta^2}{4e^2 c_{2\theta}} c_{\Phi 1}$
$\Delta g_{ZZ}^{(4)}$	$\frac{1}{c_\theta}$	$\frac{2s_\theta^2 c_\theta^2}{e^2 c_{2\theta}} c_T + \frac{4s_\theta^2}{c_{2\theta}} c_1 + 8c_\theta^2 c_3 - 4c_6 - \frac{v^2}{2} c_{\square h} - 4c_{23}$	$c_W c_\theta^2 + 2 \frac{s_\theta^2}{4c_{2\theta}} c_{BW} - \frac{s_\theta^2 c_\theta^2}{2e^2 c_{2\theta}} c_{\Phi 1}$
$\Delta g_{\gamma\gamma}^{(3)}$	s_θ^2	$-2c_9$	—
$\Delta g_{\gamma Z}^{(3)}$	$\frac{s_\theta}{c_\theta}$	$\frac{s_\theta^2}{e^2 c_{2\theta}} c_T + \frac{8s_\theta^2}{c_{2\theta}} c_1 + 4c_3 + 4s_\theta^2 c_9 - 4c_{16}$	$\frac{c_W}{2} + \frac{s_\theta^2}{c_{2\theta}} c_{BW} - \frac{s_\theta^2}{4c_{2\theta} e^2} c_{\Phi 1}$
$\Delta g_{\gamma Z}^{(4)}$	$\frac{s_\theta}{c_\theta}$	$\frac{2s_\theta^2}{e^2 c_{2\theta}} c_T + \frac{16s_\theta^2}{c_{2\theta}} c_1 + 8c_3$	$c_W + 2 \frac{s_\theta^2}{c_{2\theta}} c_{BW} - \frac{s_\theta^2}{2c_{2\theta} e^2} c_{\Phi 1}$
$\Delta g_{\gamma Z}^{(5)}$	$\frac{s_\theta}{c_\theta}$	$8c_{14}$	—

Coefficients of custodial preserving operators

$c_i \mathcal{F}_i(h)$	$SU(5)/SO(5)$	$SU(3)/SU(2) \times U(1)$	linear $d = 6$
$\mathcal{F}_B(h)$	$1 - 4g'^2 \tilde{c}_{B\Sigma} \cos^2 \frac{\varphi}{2f}$	$1 - g'^2 \frac{\tilde{c}_{B\Sigma}}{6} \left(1 + 3 \cos \frac{2\varphi}{f}\right)$	$1 + \frac{(v+h)^2}{2} g'^2 c_{BB}$
$\mathcal{F}_W(h)$	$1 - 4g^2 \tilde{c}_{W\Sigma} \cos^2 \frac{\varphi}{2f}$	$1 - 2g^2 \tilde{c}_{W\Sigma} \cos \frac{\varphi}{f}$	$1 + \frac{(v+h)^2}{2} g^2 c_{WW}$
$c_{\square H} \mathcal{F}_{\square H}(h)$	$-2\tilde{c}_6 \xi$	$-2\tilde{c}_6 \xi$	$\frac{v^2}{2} c_{\square\Phi}$
$c_{\Delta H} \mathcal{F}_{\Delta H}(h)$	—	—	—
$c_{DH} \mathcal{F}_{DH}(h)$	$4(\tilde{c}_4 + \tilde{c}_5) \xi^2$	$2(2\tilde{c}_4 + 2\tilde{c}_5 + \tilde{c}_7) \xi^2$	—
$c_1 \mathcal{F}_1(h)$	$\tilde{c}_1 \sin^2 \frac{\varphi}{2f}$	$\frac{\tilde{c}_1}{4} \sin^2 \frac{\varphi}{f}$	$\frac{(v+h)^2}{4} c_{BW}$
$c_2 \mathcal{F}_2(h)$	$\tilde{c}_2 \sin^2 \frac{\varphi}{2f}$	$\frac{\tilde{c}_2}{4} \sin^2 \frac{\varphi}{f}$	$\frac{(v+h)^2}{8} c_B$
$c_3 \mathcal{F}_3(h)$	$2\tilde{c}_3 \sin^2 \frac{\varphi}{2f}$	$\frac{\tilde{c}_3}{2} \sin^2 \frac{\varphi}{f}$	$\frac{(v+h)^2}{8} c_W$
$c_4 \mathcal{F}_4(h)$	$\tilde{c}_2 \sqrt{\xi} \sin \frac{\varphi}{f}$	$\frac{\tilde{c}_2}{2} \sqrt{\xi} \sin \frac{2\varphi}{f}$	$\frac{v(v+h)}{2} c_B$
$c_5 \mathcal{F}_5(h)$	$-2\tilde{c}_3 \sqrt{\xi} \sin \frac{\varphi}{f}$	$-2\tilde{c}_3 \sqrt{\xi} \sin \frac{\varphi}{f}$	$-\frac{v(v+h)}{2} c_W$
$c_6 \mathcal{F}_6(h)$	$16\tilde{c}_4 \sin^4 \frac{\varphi}{2f} - \frac{1}{2} \tilde{c}_6 \sin^2 \frac{\varphi}{f}$	$8(2\tilde{c}_4 + \tilde{c}_7) \sin^4 \frac{\varphi}{2f} - \frac{1}{2} \tilde{c}_6 \sin^2 \frac{\varphi}{f}$	$\frac{(v+h)^2}{8} c_{\square\Phi}$
$c_7 \mathcal{F}_7(h)$	$-2\tilde{c}_6 \sqrt{\xi} \sin \frac{\varphi}{f}$	$-2\tilde{c}_6 \sqrt{\xi} \sin \frac{\varphi}{f}$	$\frac{v(v+h)}{2} c_{\square\Phi}$
$c_8 \mathcal{F}_8(h)$	$-16\tilde{c}_5 \xi \sin^2 \frac{\varphi}{2f} + 4\tilde{c}_6 \xi \cos^2 \frac{\varphi}{2f}$	$-4(4\tilde{c}_5 + \tilde{c}_7) \xi \sin^2 \frac{\varphi}{2f} + 4\tilde{c}_6 \xi \cos^2 \frac{\varphi}{2f}$	$-v^2 c_{\square\Phi}$
$c_9 \mathcal{F}_9(h)$	$4\tilde{c}_6 \sin^2 \frac{\varphi}{2f}$	$4\tilde{c}_6 \sin^2 \frac{\varphi}{2f}$	$-\frac{(v+h)^2}{4} c_{\square\Phi}$
$c_{10} \mathcal{F}_{10}(h)$	$4\tilde{c}_6 \sqrt{\xi} \sin \frac{\varphi}{f}$	$4\tilde{c}_6 \sqrt{\xi} \sin \frac{\varphi}{f}$	$-v(v+h) c_{\square\Phi}$
$c_{11} \mathcal{F}_{11}(h)$	$16\tilde{c}_5 \sin^4 \frac{\varphi}{2f}$	$16\tilde{c}_5 \sin^4 \frac{\varphi}{2f}$	—
$c_{20} \mathcal{F}_{20}(h)$	$-16\tilde{c}_4 \xi \sin^2 \frac{\varphi}{2f}$	$-4(4\tilde{c}_4 + \tilde{c}_7) \xi \sin^2 \frac{\varphi}{2f}$	—

Coefficients of custodial breaking operators

$c_i \mathcal{F}_i(h)$	$SU(3)/(SU(2) \times U(1))$	$c_i \mathcal{F}_i(h)$	$SU(3)/(SU(2) \times U(1))$
$c_{12} \mathcal{F}_{12}(h)$	$\tilde{c}_{W\Sigma} \sin^4 \frac{\varphi}{2f}$	$c_{21} \mathcal{F}_{21}(h)$	$8\tilde{c}_4 \xi \sin^4 \frac{\varphi}{2f} - 2\tilde{c}_7 \xi \cos \frac{\varphi}{f} \sin^2 \frac{\varphi}{2f},$
$c_{13} \mathcal{F}_{13}(h)$	$2\tilde{c}_3 \sin^4 \frac{\varphi}{2f}$	$c_{22} \mathcal{F}_{22}(h)$	$8\tilde{c}_5 \xi \sin^4 \frac{\varphi}{2f} + 2\xi \tilde{c}_7 \sin^2 \frac{\varphi}{2f} - 2\tilde{c}_6 \xi \sin^2 \frac{\varphi}{2f} (1 + 2 \cos \frac{\varphi}{f})$
$c_{15} \mathcal{F}_{15}(h)$	$-2\tilde{c}_6 \sin^4 \frac{\varphi}{2f}$	$c_{23} \mathcal{F}_{23}(h)$	$-16\tilde{c}_4 \sin^6 \frac{\varphi}{2f} + \tilde{c}_6 \sin^2 \frac{\varphi}{2f} \sin^2 \frac{\varphi}{f} +$ $2\tilde{c}_7 \sin^4 \frac{\varphi}{2f} (\cos \frac{\varphi}{f} - 3)$
$c_{16} \mathcal{F}_{16}(h)$	$4\tilde{c}_6 \sin^4 \frac{\varphi}{2f}$	$c_{24} \mathcal{F}_{24}(h)$	$-4(4\tilde{c}_5 + \tilde{c}_6) \sin^6 \frac{\varphi}{2f} + \tilde{c}_7 \sin^2 \frac{\varphi}{2f} \sin^2 \frac{\varphi}{f}$
$c_{17} \mathcal{F}_{17}(h)$	$2\tilde{c}_3 \sqrt{\xi} \sin^2 \frac{\varphi}{2f} \sin \frac{\varphi}{f}$	$c_{25} \mathcal{F}_{25}(h)$	$2\tilde{c}_6 \sqrt{\xi} \sin^2 \frac{\varphi}{2f} \sin \frac{\varphi}{f}$
$c_{18} \mathcal{F}_{18}(h)$	$2(\tilde{c}_6 - \tilde{c}_7) \sqrt{\xi} \sin^2 \frac{\varphi}{2f} \sin \frac{\varphi}{f}$	$c_{26} \mathcal{F}_{26}(h)$	$2(2(\tilde{c}_4 + \tilde{c}_5) + \tilde{c}_6 + \tilde{c}_7) \sin^8 \frac{\varphi}{2f}$
$c_{19} \mathcal{F}_{19}(h)$	$-4\tilde{c}_6 \sqrt{\xi} \sin^2 \frac{\varphi}{2f} \sin \frac{\varphi}{f}$		

Linear - chiral correspondence

$$\mathcal{O}_{BB} = \frac{v^2}{2} \mathcal{P}_B(h)$$

$$\mathcal{O}_{WW} = \frac{v^2}{2} \mathcal{P}_W(h)$$

$$\mathcal{O}_{GG} = -\frac{2v^2}{g_s^2} \mathcal{P}_G(h)$$

$$\mathcal{O}_{BW} = \frac{v^2}{8} \mathcal{P}_1(h)$$

$$\mathcal{O}_B = \frac{v^2}{16} \mathcal{P}_2(h) + \frac{v^2}{8} \mathcal{P}_4(h)$$

$$\mathcal{O}_W = \frac{v^2}{8} \mathcal{P}_3(h) - \frac{v^2}{4} \mathcal{P}_5(h)$$

$$\mathcal{O}_{\Phi,1} = \frac{v^2}{2} \mathcal{P}_H(h) - \frac{v^2}{4} \mathcal{F}(h) \mathcal{P}_T(h)$$

$$\mathcal{O}_{\Phi,2} = v^2 \mathcal{P}_H(h)$$

$$\mathcal{O}_{\Phi,4} = \frac{v^2}{2} \mathcal{P}_H(h) + \frac{v^2}{2} \mathcal{F}(h) \mathcal{P}_C(h)$$

$$\mathcal{O}_{\square\Phi} = \frac{v^2}{2} \mathcal{P}_{\square H}(h) + \frac{v^2}{8} \mathcal{P}_6(h) + \frac{v^2}{4} \mathcal{P}_7(h) - v^2 \mathcal{P}_8(h) - \frac{v^2}{4} \mathcal{P}_9(h) - \frac{v^2}{2} \mathcal{P}_{10}(h)$$

$$\mathcal{P}_{DH}(h), \mathcal{P}_{20}(h) \rightarrow [\mathbf{D}_\mu \Phi^\dagger \mathbf{D}^\mu \Phi]^2$$

$$\mathcal{P}_{11}(h), \mathcal{P}_{18}(h), \mathcal{P}_{21}(h), \mathcal{P}_{22}(h), \mathcal{P}_{23}(h), \mathcal{P}_{24}(h) \rightarrow [\mathbf{D}^\mu \Phi^\dagger \mathbf{D}_\nu \Phi]^2$$

$$\mathcal{P}_{12}(h) \rightarrow (\Phi^\dagger W^{\mu\nu} \Phi)^2$$

$$\mathcal{P}_{13}(h), \mathcal{P}_{17}(h) \rightarrow (\Phi^\dagger W^{\mu\nu} \Phi) \mathbf{D}_\mu \Phi^\dagger \mathbf{D}_\nu \Phi$$

$$\mathcal{P}_{14}(h) \rightarrow \varepsilon^{\mu\nu\rho\lambda} (\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\rho \Phi) (\Phi^\dagger \sigma_i \overleftrightarrow{\mathbf{D}}_\lambda \Phi) W_{\mu\nu}^i$$

$$\mathcal{P}_{15}(h), \mathcal{P}_{19}(h) \rightarrow [\Phi^\dagger \mathbf{D}_\mu \mathbf{D}^\mu \Phi - \mathbf{D}_\mu \mathbf{D}^\mu \Phi^\dagger]^2$$

$$\mathcal{P}_{16}(h), \mathcal{P}_{25}(h) \rightarrow (\mathbf{D}^\nu \Phi^\dagger \mathbf{D}_\mu \mathbf{D}^\mu \Phi - \mathbf{D}_\mu \mathbf{D}^\mu \Phi^\dagger \mathbf{D}^\nu \Phi) (\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\nu \Phi)$$

$$\mathcal{P}_{26}(h) \rightarrow \left[(\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\mu \Phi) (\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\nu \Phi) \right]^2$$

Linear - chiral correspondence (SILH)

$$\mathcal{O}_g^{\text{SILH}} = \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu}$$

$$\mathcal{O}_\gamma^{\text{SILH}} = \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi$$

$$\mathcal{O}_B^{\text{SILH}} = (\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\mu \Phi) \partial_\nu \hat{B}^{\mu\nu}$$

$$\mathcal{O}_W^{\text{SILH}} = \frac{ig}{2} (\Phi^\dagger \sigma^i \overleftrightarrow{\mathbf{D}}_\mu \Phi) \mathbf{D}_\nu W_i^{\mu\nu}$$

$$\mathcal{O}_{HB}^{\text{SILH}} = (\mathbf{D}_\mu \Phi)^\dagger (\mathbf{D}_\nu \Phi) \hat{B}^{\mu\nu}$$

$$\mathcal{O}_{HW}^{\text{SILH}} = (\mathbf{D}_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (\mathbf{D}_\nu \Phi)$$

$$\mathcal{O}_T^{\text{SILH}} = \frac{1}{2} (\Phi^\dagger \overleftrightarrow{\mathbf{D}}_\mu \Phi) (\Phi^\dagger \overleftrightarrow{\mathbf{D}}^\mu \Phi)$$

$$\mathcal{O}_H^{\text{SILH}} = \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi)$$

$$\mathcal{O}_6^{\text{SILH}} = \frac{1}{3} (\Phi^\dagger \Phi)^3$$

$$\mathcal{O}_y^{\text{SILH}} = (\Phi^\dagger \Phi) f_L \Phi \mathcal{Y} f_R + \text{h.c.}$$

$$\mathcal{O}_g^{\text{SILH}} = \frac{v^2}{2g_s^2} \partial_G(h)$$

$$\mathcal{O}_\gamma^{\text{SILH}} = \frac{v^2}{2} \partial_B(h)$$

$$\mathcal{O}_B^{\text{SILH}} = \frac{v^2}{8} (\mathcal{P}_2(h) + 2\mathcal{P}_4(h)) + \frac{v^2}{8} \mathcal{P}_1(h) + \frac{v^2}{2} \mathcal{P}_B(h)$$

$$\mathcal{O}_{HB}^{\text{SILH}} = \frac{v^2}{16} (\mathcal{P}_2(h) + 2\mathcal{P}_4(h))$$

$$\mathcal{O}_W^{\text{SILH}} = \frac{v^2}{4} (\mathcal{P}_3(h) - 2\mathcal{P}_5(h)) + \frac{v^2}{8} \mathcal{P}_1(h) + \frac{v^2}{2} \mathcal{P}_W(h)$$

$$\mathcal{O}_{HW}^{\text{SILH}} = \frac{v^2}{8} (\mathcal{P}_3(h) - 2\mathcal{P}_5(h))$$

$$\mathcal{O}_T^{\text{SILH}} = \frac{v^2}{2} \mathcal{F}(h) \mathcal{P}_T(h)$$

$$\mathcal{O}_H^{\text{SILH}} = v^2 \mathcal{P}_H(h)$$

$$\mathcal{O}_y^{\text{SILH}} = 3v^2 \mathcal{P}_H(h) + v^2 \mathcal{F}(h) \mathcal{P}_C(h) - \frac{(v+h)^3}{2} \frac{\delta V(h)}{\delta h}$$

Best bounds on the chiral coefficients

