

# Connecting UV completions with HEFTs

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HEFT2014



# Outline

- The HEFT approach, briefly
- Framework for HEFT studies
- Complete analysis of HEFT
- Limitations of HEFTs
- Benchmarks for HEFTs



The HEFT approach , briefly



# HEFT

## Bottom-up approach

operators w/ SM particles and symmetries,  
plus the **newcomer**, the **Higgs**

in this talk

linear realization, a choice of basis

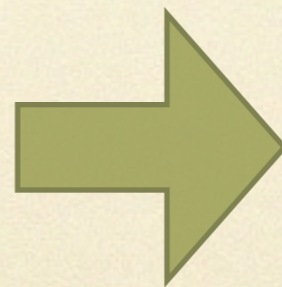
$$\begin{aligned} \mathcal{L} \supset & \frac{\bar{c}_H}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{g'^2 \bar{c}_\gamma}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_s^2 \bar{c}_g}{m_W^2} \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} \\ & + \frac{2ig \bar{c}_{HW}}{m_W^2} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + \frac{ig' \bar{c}_{HB}}{m_W^2} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & + \frac{ig \bar{c}_W}{m_W^2} [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig' \bar{c}_B}{2m_W^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{\bar{c}_t}{v^2} y_t \Phi^\dagger \Phi \Phi^\dagger \cdot \bar{Q}_L t_R + \frac{\bar{c}_b}{v^2} y_b \Phi^\dagger \Phi \Phi \cdot \bar{Q}_L b_R + \frac{\bar{c}_\tau}{v^2} y_\tau \Phi^\dagger \Phi \Phi \cdot \bar{L}_L \tau_R. \end{aligned}$$



# How do we look for New Physics using HEFTs

## Higgs anomalous couplings

HDOs generate  
HVV interactions  
with more  
derivatives



$$-\frac{1}{4}h g_{hVV}^{(1)} V_{\mu\nu} V^{\mu\nu}$$

$$-h g_{hVV}^{(2)} V_\nu \partial_\mu V^{\mu\nu}$$

$$-\frac{1}{4}h \tilde{g}_{hVV} V_{\mu\nu} \tilde{V}^{\mu\nu}$$

example.

$$g_{hww}^{(1)} = \frac{2g}{m_W} \bar{c}_{HW}$$

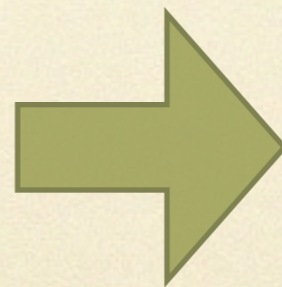
$$g_{hww}^{(2)} = \frac{g}{m_W} (\bar{c}_W + \bar{c}_{HW})$$



# How do we look for New Physics using HEFTs

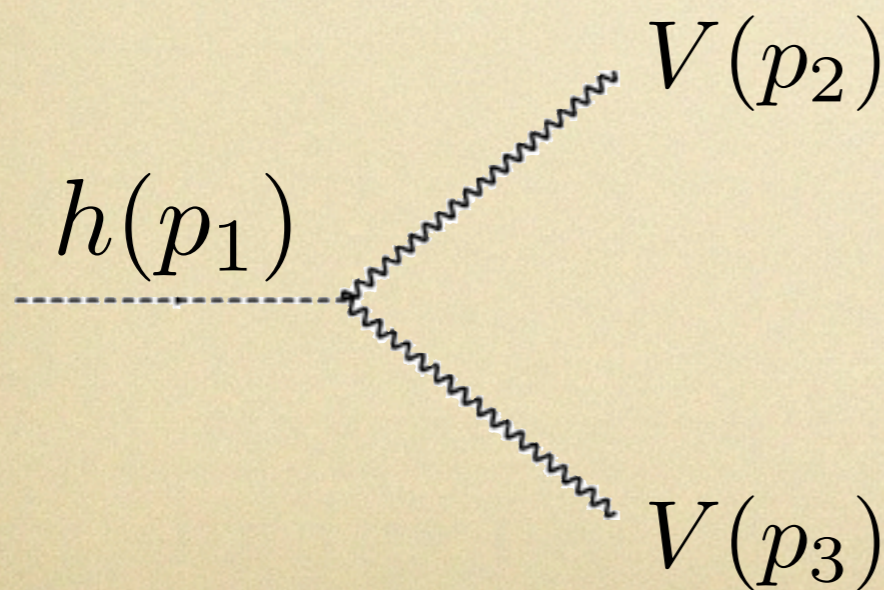
## Higgs anomalous couplings

HDOs generate  
 HVV interactions  
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$$\begin{aligned}
 & -\frac{1}{4} h \underline{g_{hVV}^{(1)}} V_{\mu\nu} V^{\mu\nu} \\
 & -h \underline{g_{hVV}^{(2)}} V_\nu \partial_\mu V^{\mu\nu} \\
 & -\frac{1}{4} h \underline{\tilde{g}_{hVV}} V_{\mu\nu} \tilde{V}^{\mu\nu}
 \end{aligned}$$

**ex.** Feynman rule if  $m_h > 2m_V$



$$i\eta_{\mu\nu} \left( \underline{g_{hVV}^{(1)}} \left( \frac{\hat{s}}{2} - m_V^2 \right) + 2 \underline{g_{hVV}^{(2)}} m_V^2 \right)$$

$$-i \underline{g_{hVV}^{(1)}} p_3^\mu p_2^\nu$$

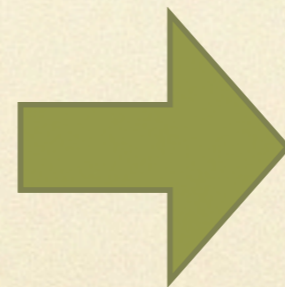
$$-i \underline{\tilde{g}_{hVV}} \epsilon^{\mu\nu\alpha\beta} p_{2,\alpha} p_{3,\beta}$$



# How do we look for New Physics using HEFTs

## Higgs anomalous couplings

HDOs generate  
HVV interactions  
with more  
derivatives

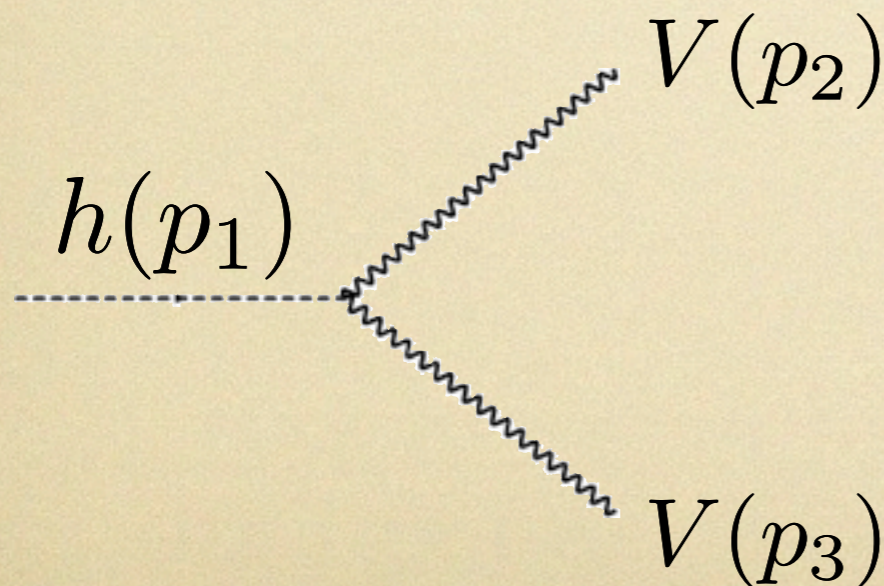


$$-\frac{1}{4}h \underline{g_{hVV}^{(1)}} V_{\mu\nu} V^{\mu\nu}$$

$$-h \underline{g_{hVV}^{(2)}} V_\nu \partial_\mu V^{\mu\nu}$$

$$-\frac{1}{4}h \underline{\tilde{g}_{hVV}} V_{\mu\nu} \tilde{V}^{\mu\nu}$$

**ex.** Feynman rule if  $m_h > 2m_V$



**total rates, COM,  
angular,  
inv mass and pT  
distributions**



# Framework for HEFT studies

Alloul, Fuks, VS. 1310.5150  
VS and Williams. In prep.



Higgs BRs

eHDECAY

Contino et al. 1303.3876



Production rates and kinematic distributions

depend on cuts  
need radiation and detector effects

Simulation tools



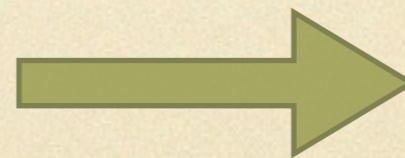
## Production rates and kinematic distributions

depend on cuts  
need radiation and detector effects

Simulation tools

coefficients

$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$



**Collider  
simulation**

observables

Limit coefficients  
= new physics



In this talk we use

1. Feynrules HDOs involving Higgs and TGCs

Alloul, Fuks, VS. 1310.5150

links to CalcHEP, LoopTools, Madgraph...

HEFT->Madgraph-> Pythia... -> FastSim / FullSim



In this talk we use

1. Feynrules HDOs involving Higgs and TGCs

Alloul, Fuks, VS. 1310.5150

links to CalcHEP, LoopTools, Madgraph...

HEFT->Madgraph->Pythia... -> FastSim / FullSim

2. QCD NLO HDOs involving Higgs and TGCs

VS and Williams. In prep.

MCFM and POWHEG

Pythia, Herwig... -> FastSim / FullSim

MC@NLO : see talk of M. Zaro

also VBF@NLO



# Complete analysis of HEFT

Ellis, VS and You. 1404.3667+work in preparation



# Number of independent operators in HEFT

In the SILH basis

$$\bar{c}_i \equiv \{ \bar{c}_H, \bar{c}_{t,b,\tau}, \bar{c}_W, \bar{c}_{HW}, \bar{c}_{HB}, \bar{c}_\gamma, \bar{c}_g \} .$$

Note that

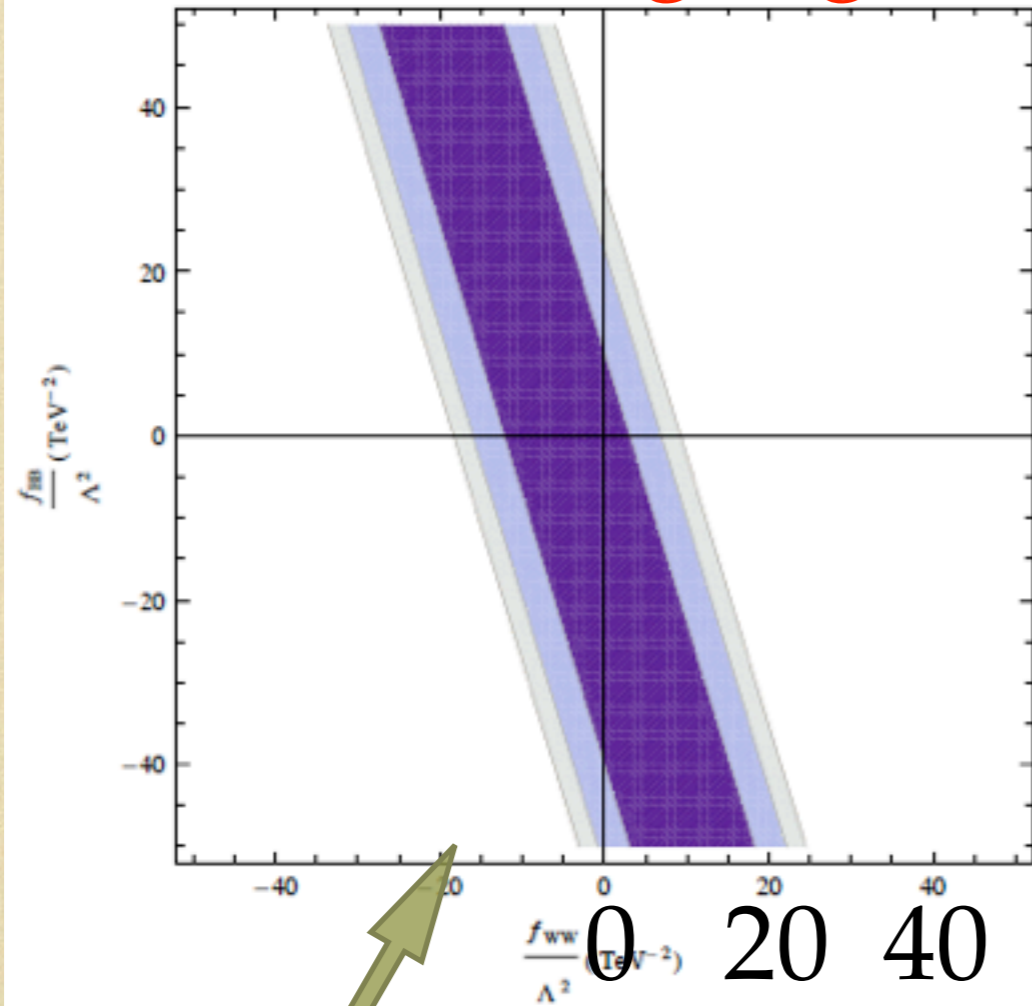
We have eliminated operators which  
contribute to STU at tree level,  
(LHC cannot compete)

but kept operators at loop order in STU...

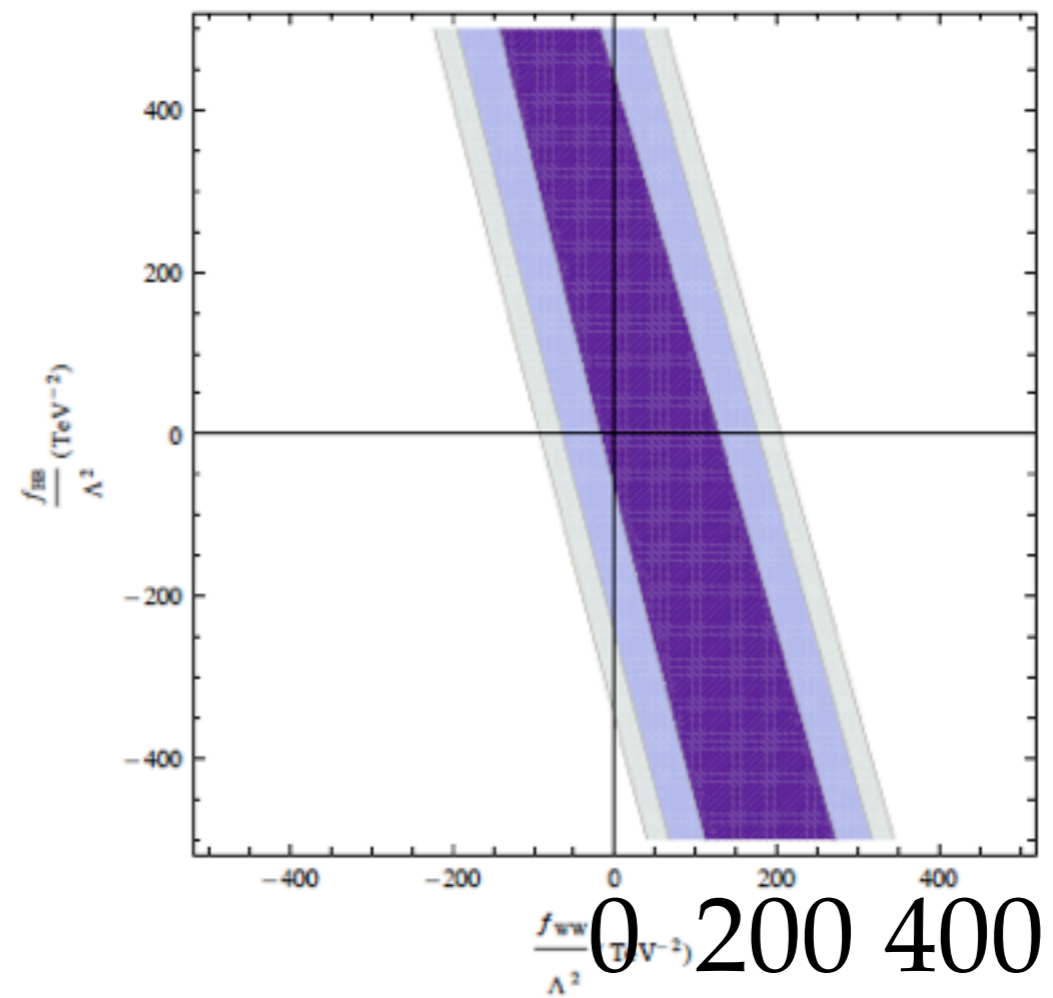


renormalization / matching **is** important

leading log



ren. scheme dep.



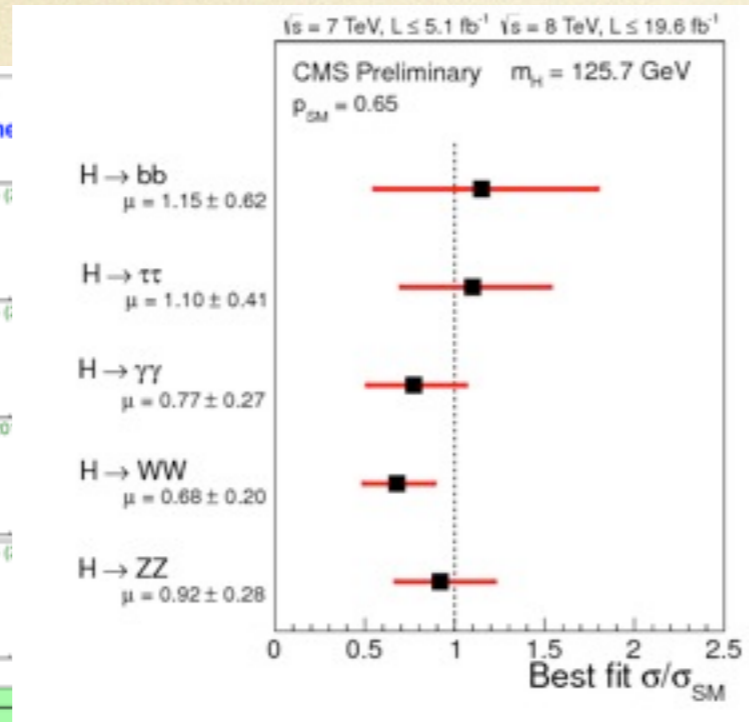
Cheng, Dawson, Zhang. 1311.3107

Masso, VS. 1211.1320

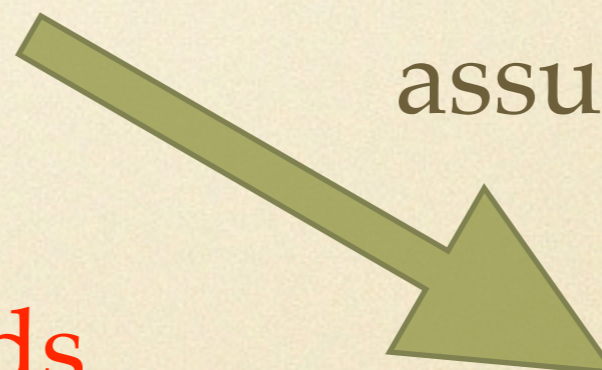
operators at loop-order in STU



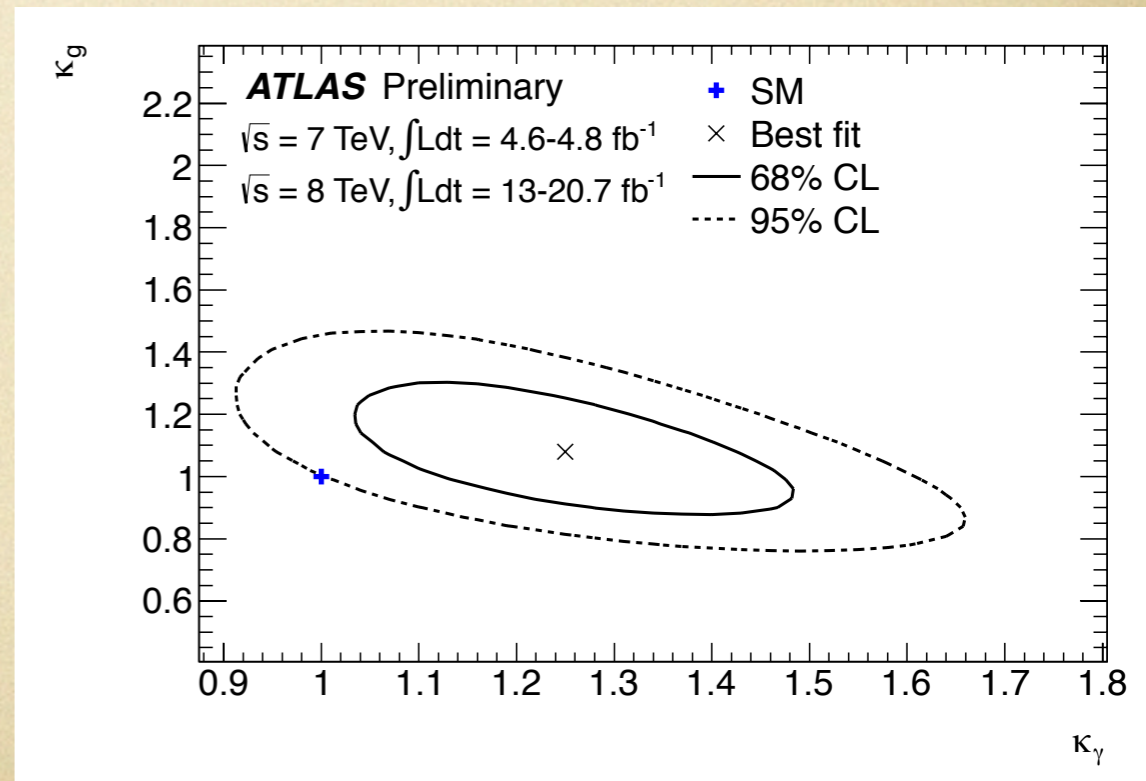
Usually, maximize information on signal strengths



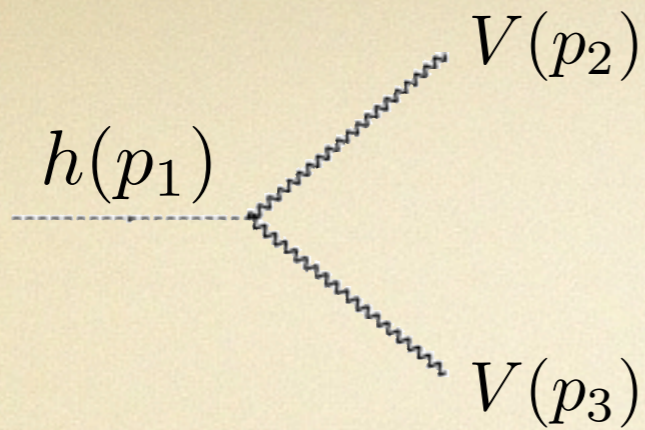
make some assumptions



+ likelihoods







$$i\eta_{\mu\nu} \left( \underline{g_{hVV}^{(1)}} \left( \frac{\hat{s}}{2} - m_V^2 \right) + \underline{2g_{hVV}^{(2)}} m_V^2 \right) - \underline{i g_{hVV}^{(1)}} p_3^\mu p_2^\nu - \underline{i \tilde{g}_{hVV}} \epsilon^{\mu\nu\alpha\beta} p_{2,\alpha} p_{3,\beta}$$

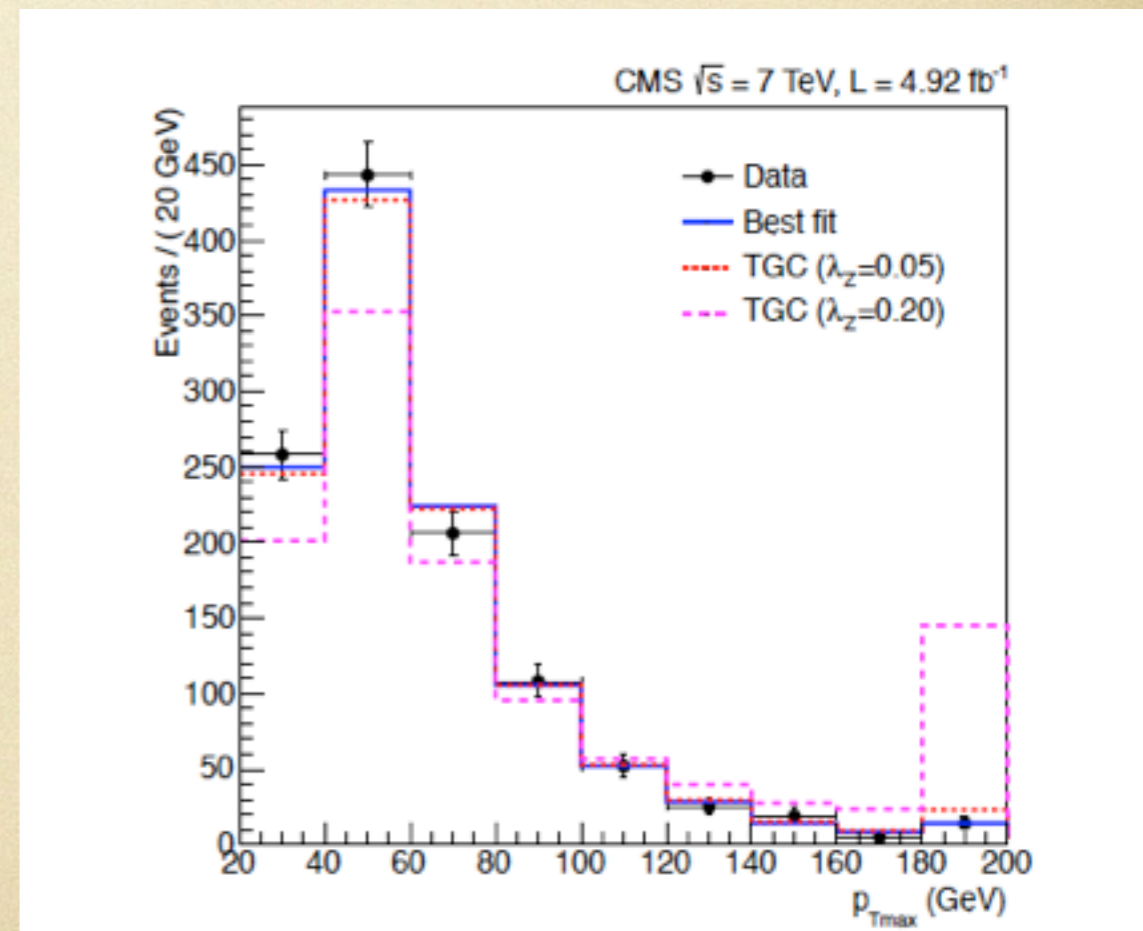
More information in kinematic distributions

For example, in EW physics

**TGCs**

instead of total rates

More Higgs data:  
total rates  $\rightarrow$  kinematics



leading lepton  $p_T$



What is the most sensitive Higgs channel to kinematics at LHC8?



# What is the most sensitive Higgs channel to kinematics at LHC8?

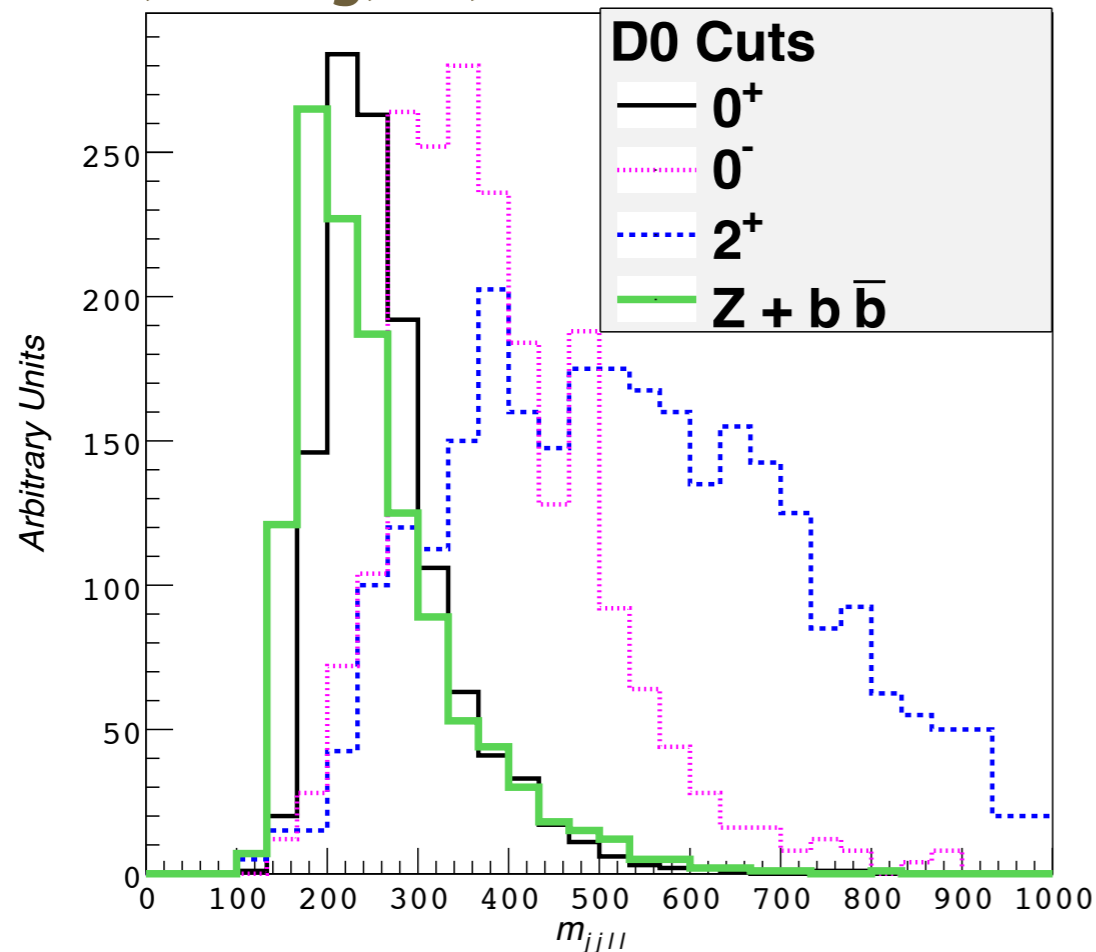
## Associated production

very sensitive to the Lorentz structure of the vertex

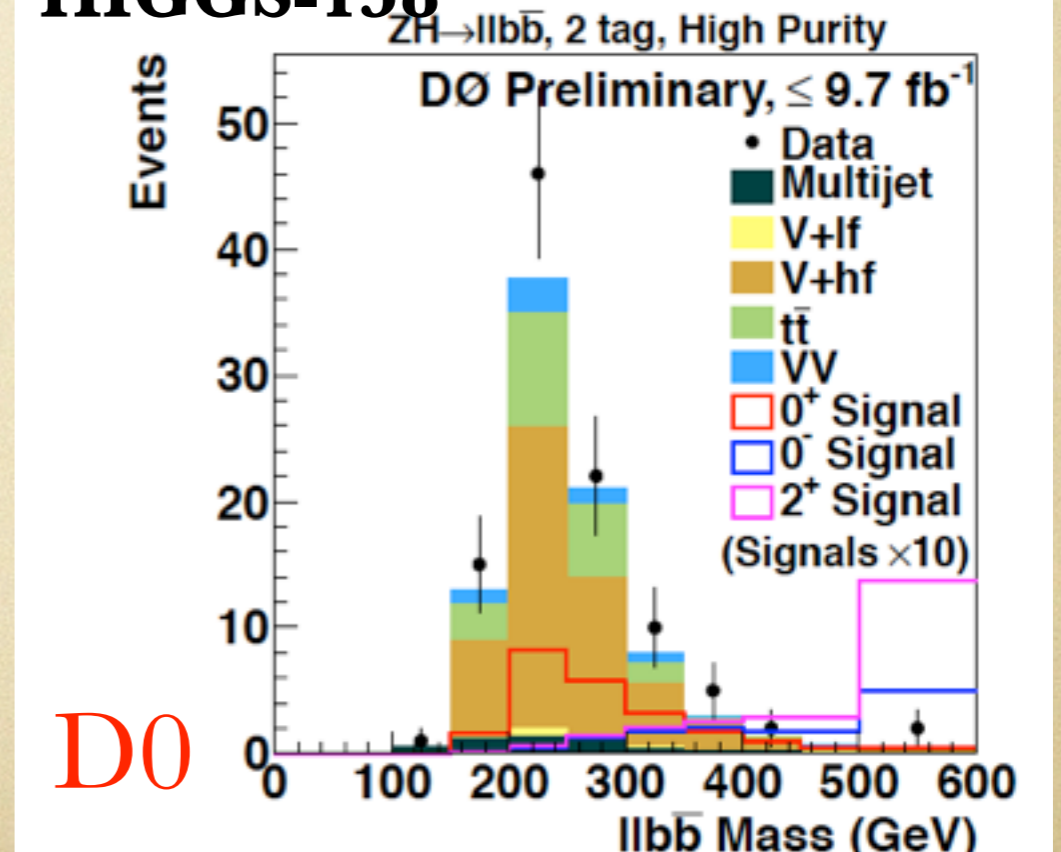
Test JCP of the Higgs

$$m_V h$$

Ellis, Hwang, VS, You. 1208.6002



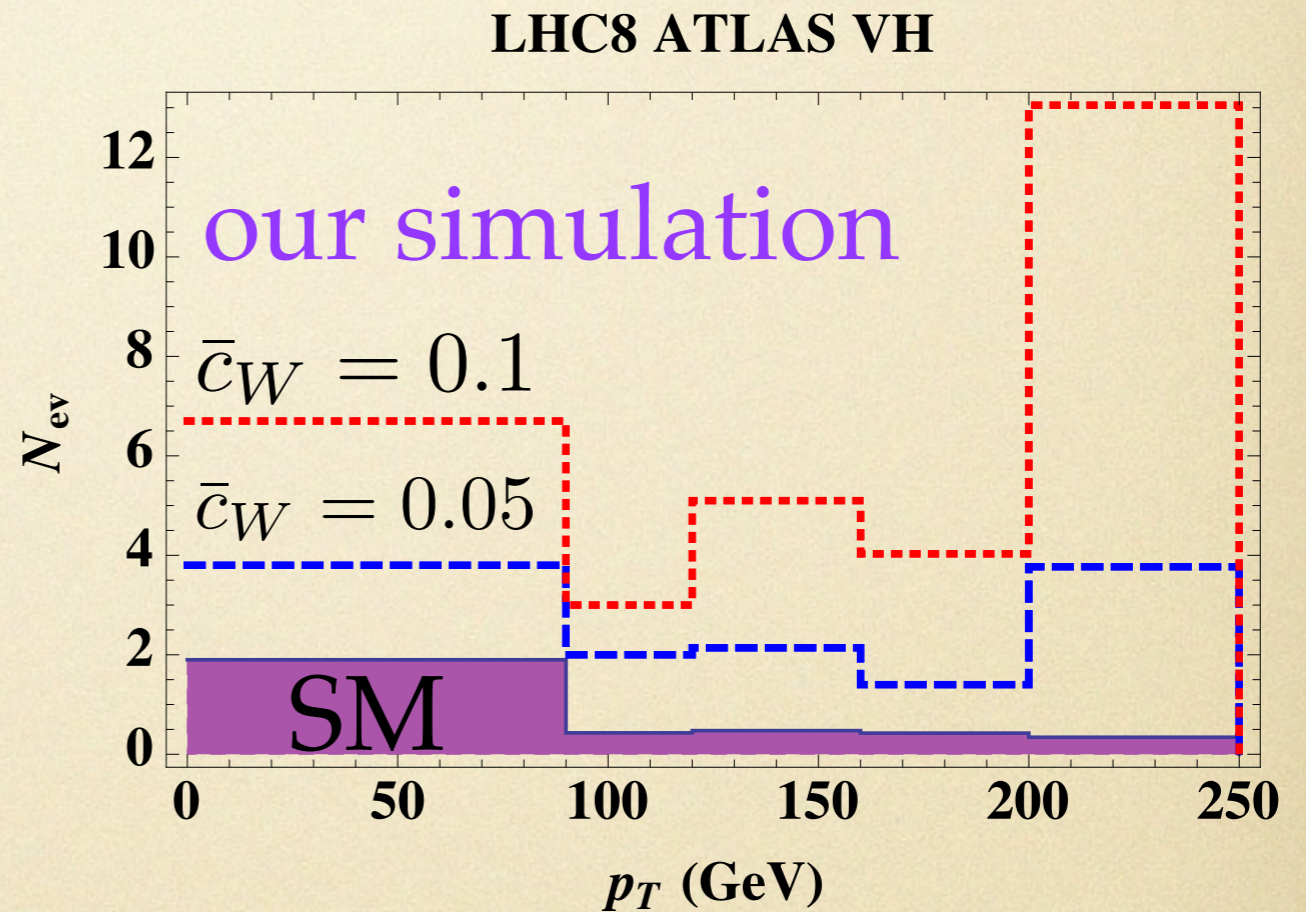
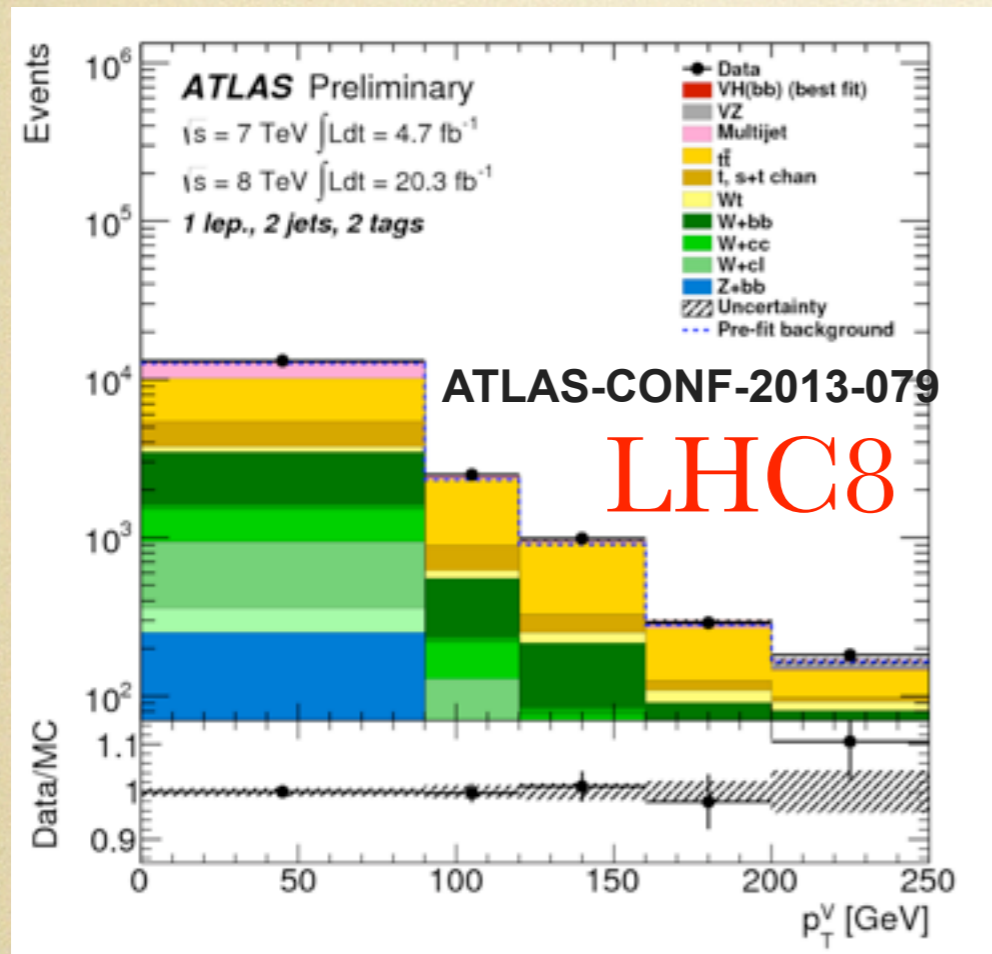
HIGGS-138





For the scalar Higgs boson

# Kinematics of associated production at LHC8



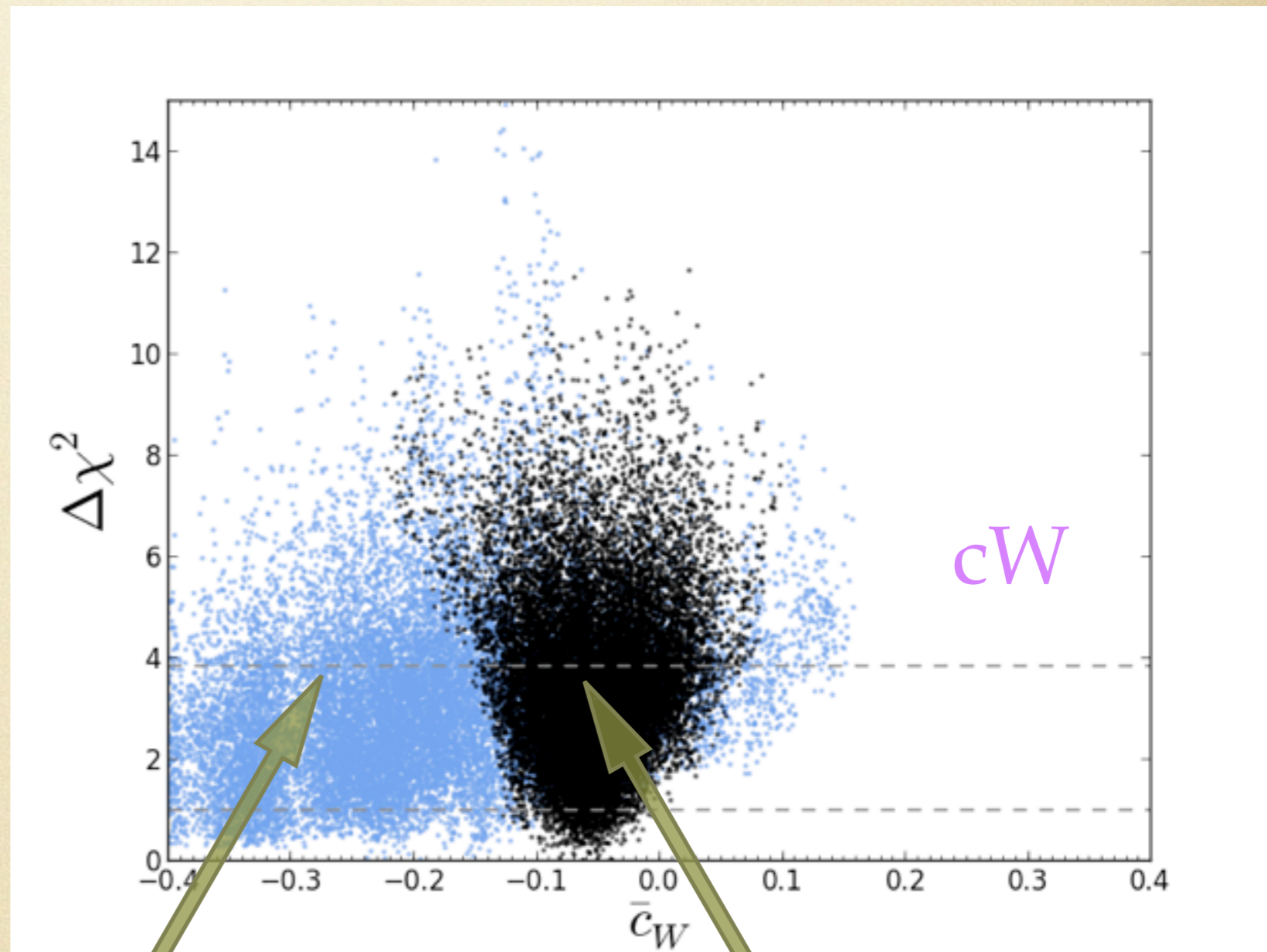
Feynrules -> MG5-> pythia->Delphes3  
verified for SM/BGs => expectation for HEFT

inclusive cross section is less  
sensitive than distribution



Besides, breaking of blind directions requires information on HV production

Global fit to 8 parameters



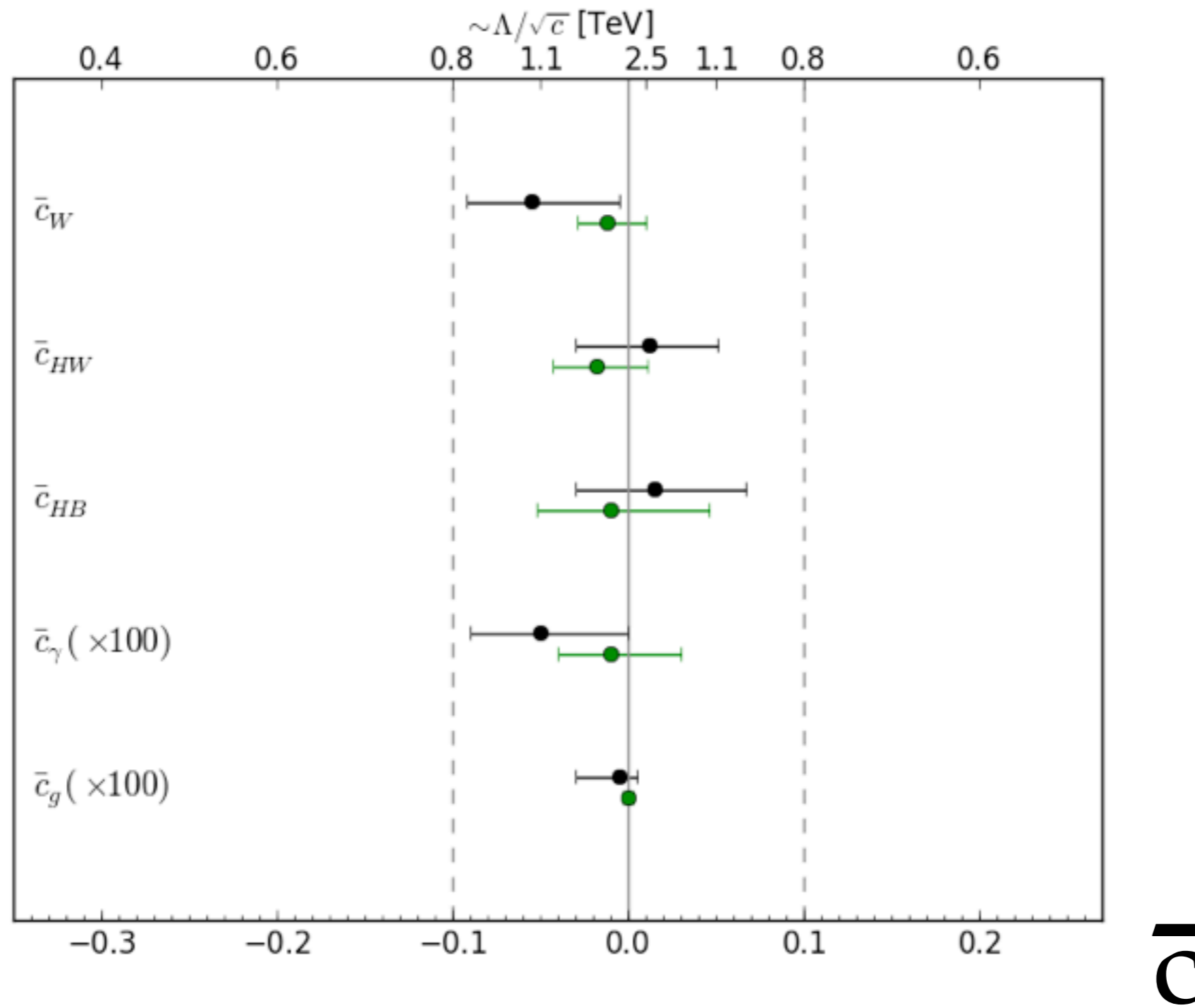
without AP

with AP



# Putting it all together

black global fit  
green one-by-one fit



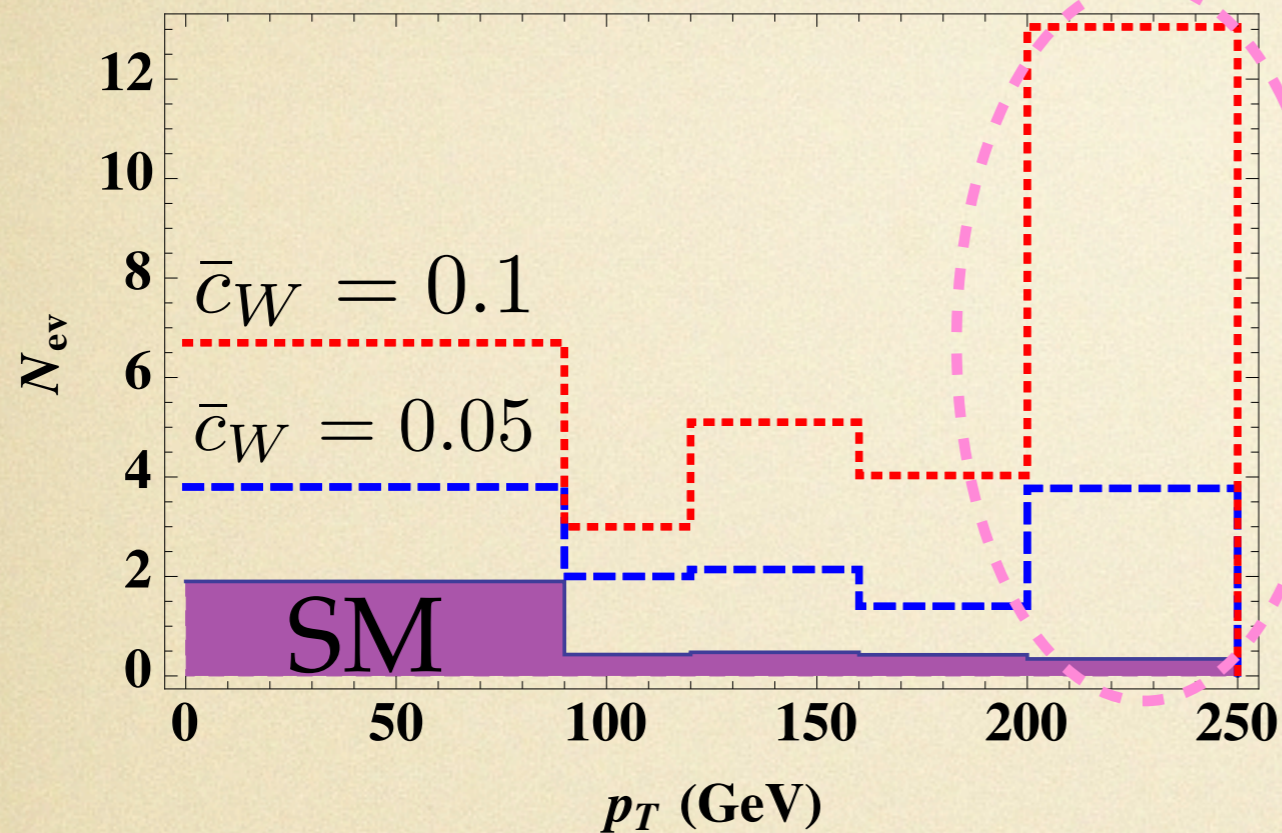
ct,cd,cH: weaker constraints



# Limitations of HEFTs



## LHC8 ATLAS VH



most sensitive bin:  
overflow (last) bin

At high- $p_T$   
sensitive to dynamics of new physics

breakdown of EFT

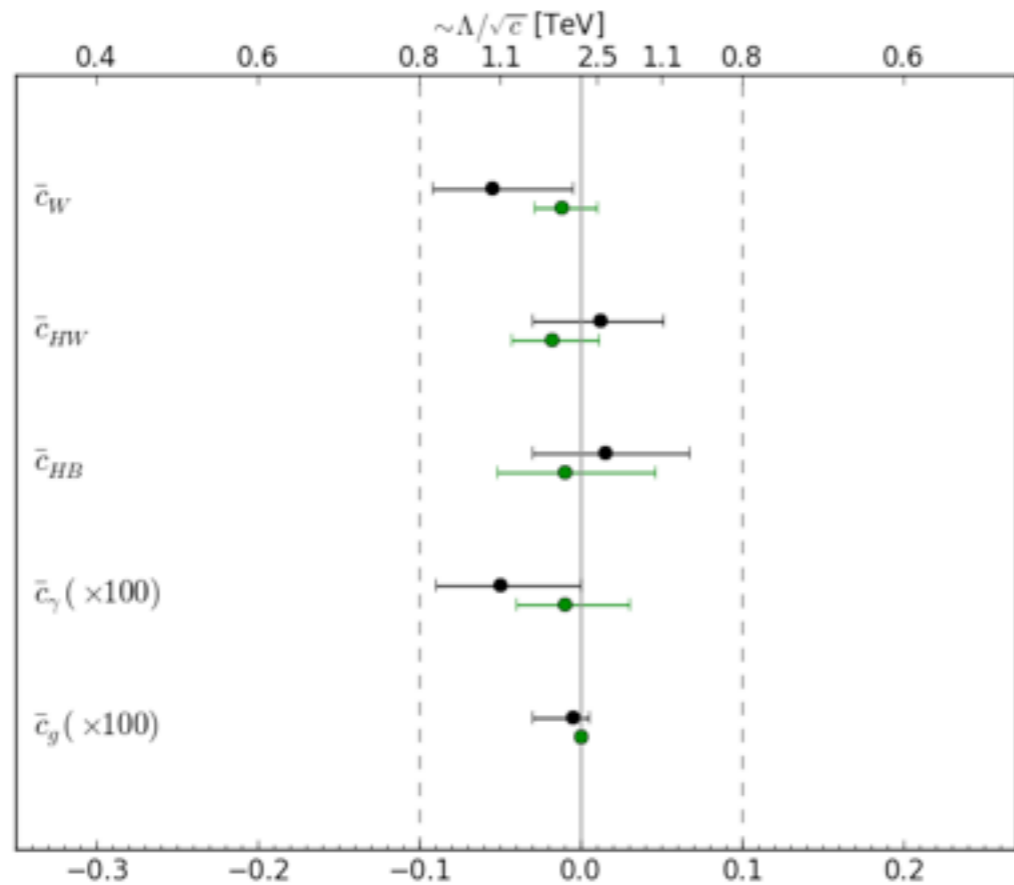
To what extent can we use this bin?

see also, **F. Riva, S. Dawson and M. McCullough talks,**

and Englert+Spannowsky. 1408.5147

Non-linear realization studies?

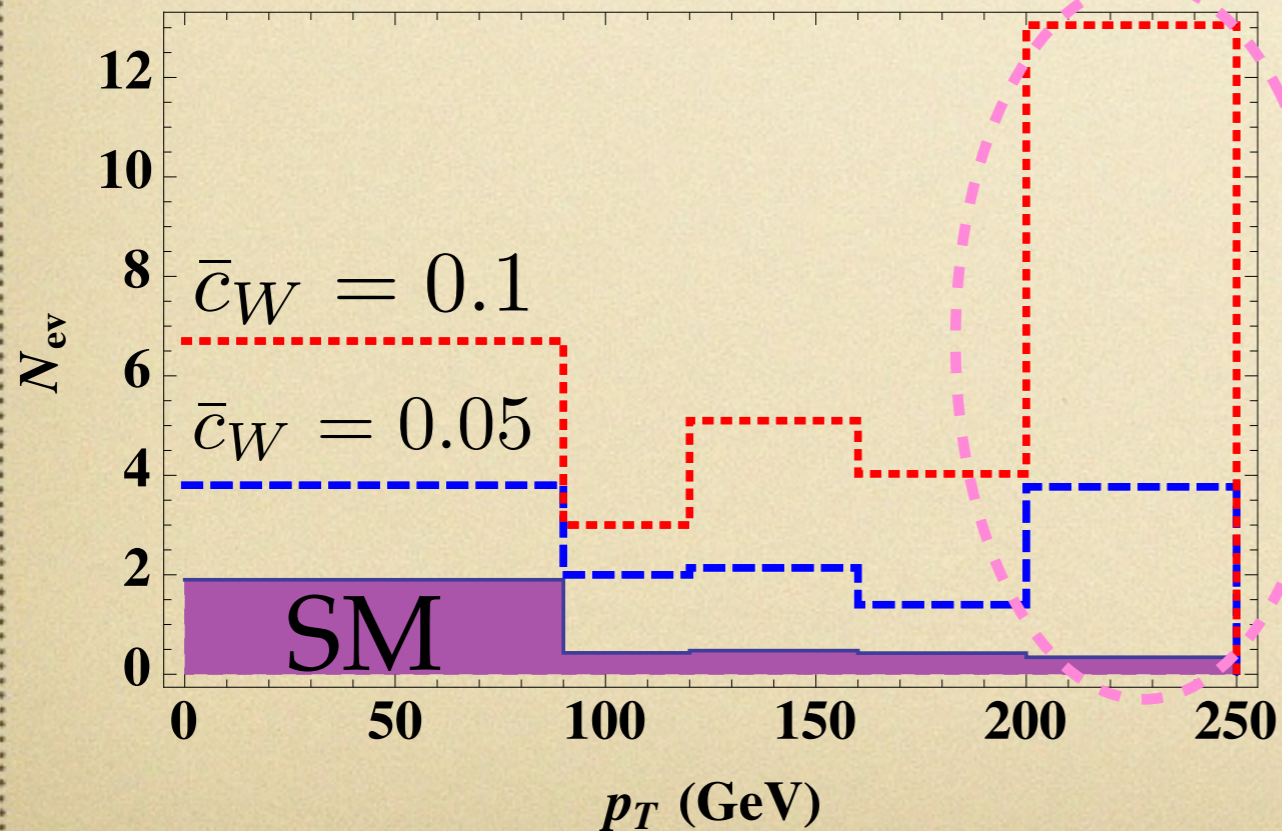




roughly speaking  
 $\sqrt{\hat{s}} \sim \mathcal{O}(800-1000) \text{ GeV}$

validity:  
 need to compare with  
 UV completions

LHC8 ATLAS VH





# Benchmarks for HEFTs

Masso and VS. 1211.1320

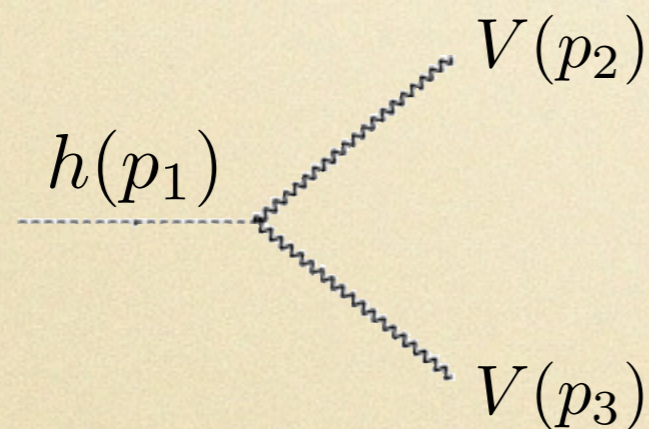
Gorbahn, No and VS. In preparation



# HEFT (linear realization) vs UV-completions

Kinematics most sensitive to operators with  
Lorentz structure *different* from SM

example of  $h$  to  $VV$



$$i\eta_{\mu\nu} \left( \underline{g_{hVV}^{(1)}} \left( \frac{\hat{s}}{2} - m_V^2 \right) + \underline{2g_{hVV}^{(2)}} m_V^2 \right)$$

$$- \underline{i g_{hVV}^{(1)}} p_3^\mu p_2^\nu$$

$$- \underline{i \tilde{g}_{hVV}} \epsilon^{\mu\nu\alpha\beta} p_{2,\alpha} p_{3,\beta}$$

looking for UV models  
generating  $c_W$ ,  $c_{HW}$ -types

$$g_{hww}^{(1)} = \frac{2g}{m_W} \bar{c}_{HW}$$

$$g_{hww}^{(2)} = \frac{g}{m_W} (\bar{c}_W + \bar{c}_{HW})$$



# HEFT (linear realization) vs UV-completions

UV models

Example 1.

tree-level operators

*radion/dilaton exchange*

Example 2.

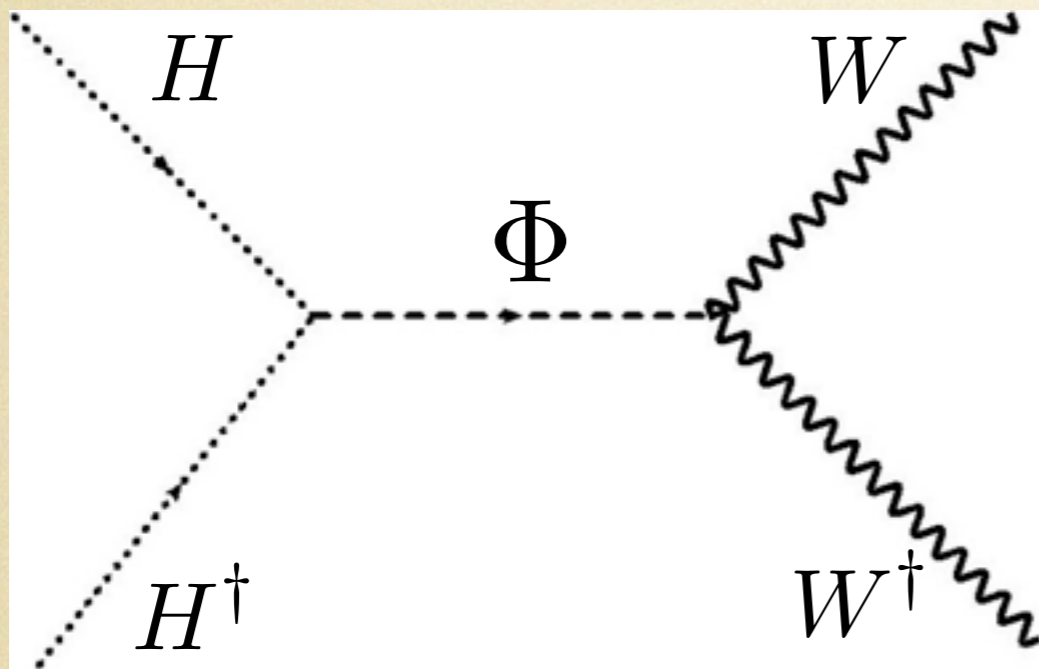
loop-induced operators

*2HDM and SUSY spartners*



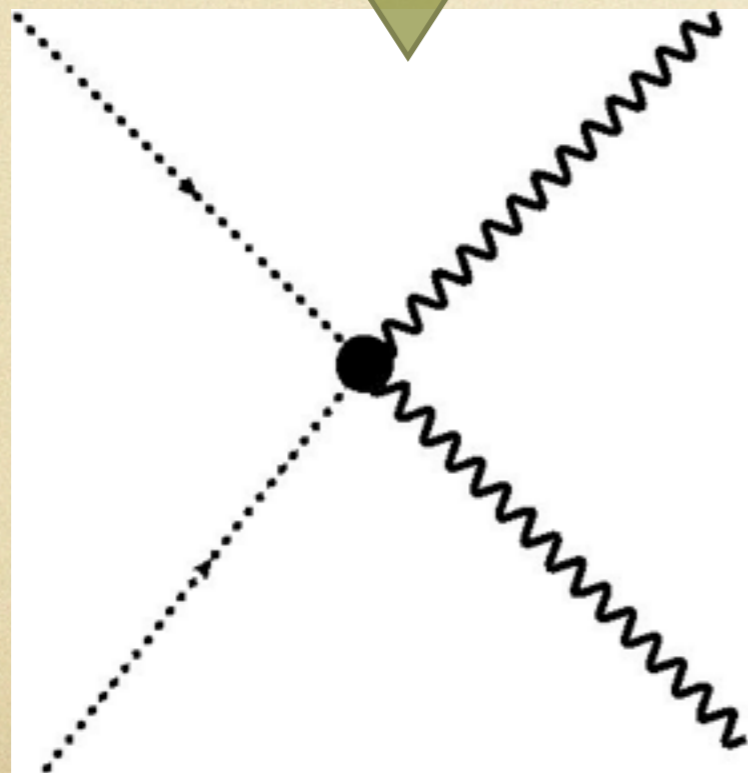
Example 1. Tree-level exchange

*radion/dilaton*



$$\frac{g_\Phi^2}{\hat{s} - M_\Phi^2} \simeq -\frac{g_\Phi^2}{M_\Phi^2} \left( 1 - \frac{\hat{s}}{M_\Phi^2} + \dots \right)$$

HEFT



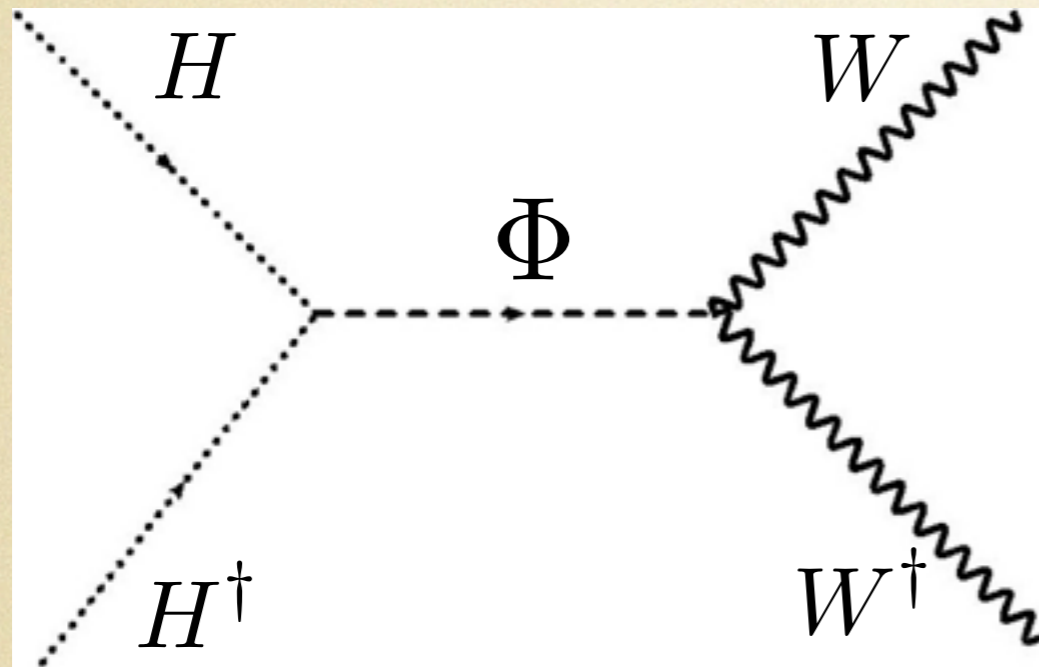
$$\hat{s} \lesssim M_\Phi^2$$

$$\bar{c}_W \simeq \left( \frac{m_H v}{\Lambda M_\Phi} \right)^2$$

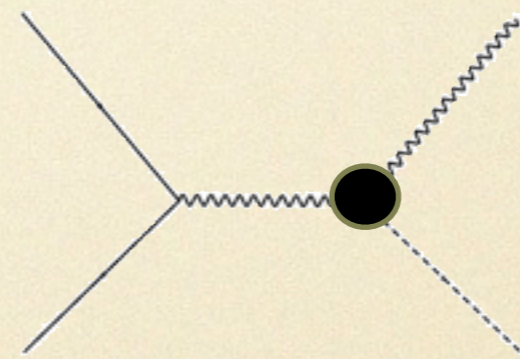


# Example 1. Tree-level exchange

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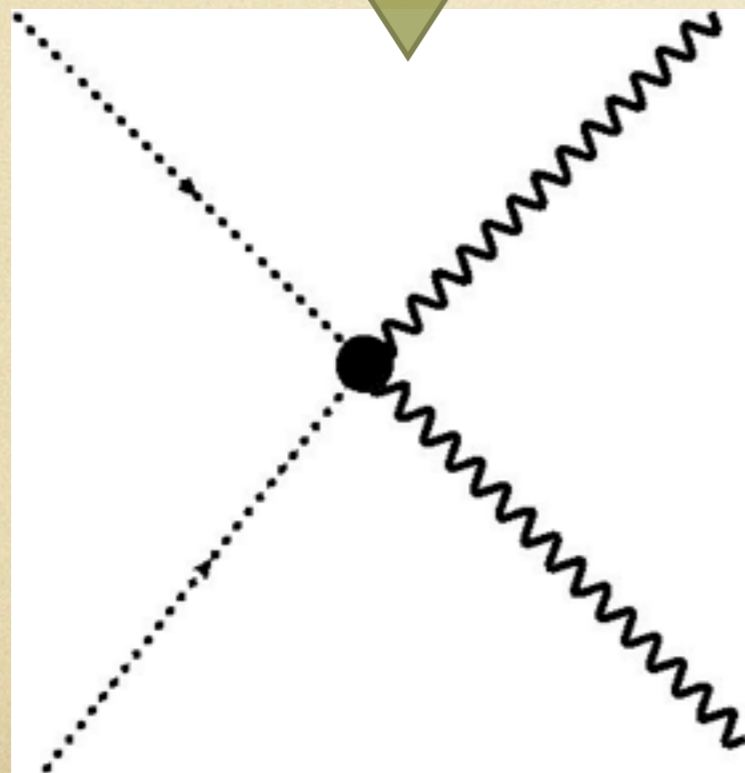


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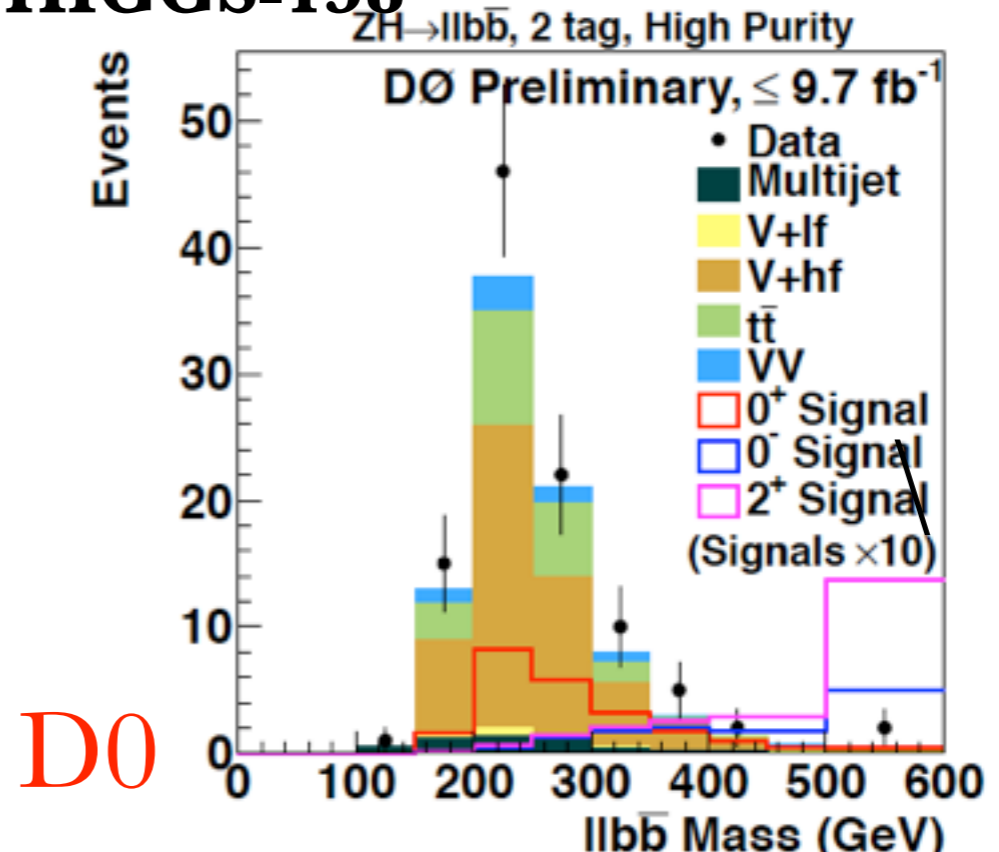


$m_V h$

HEFT



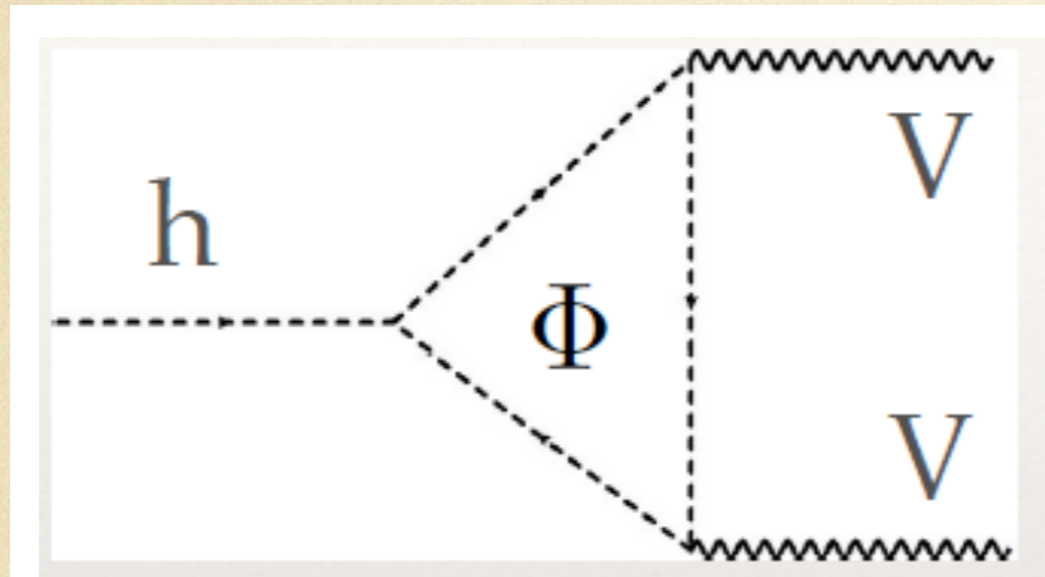
## HIGGS-138



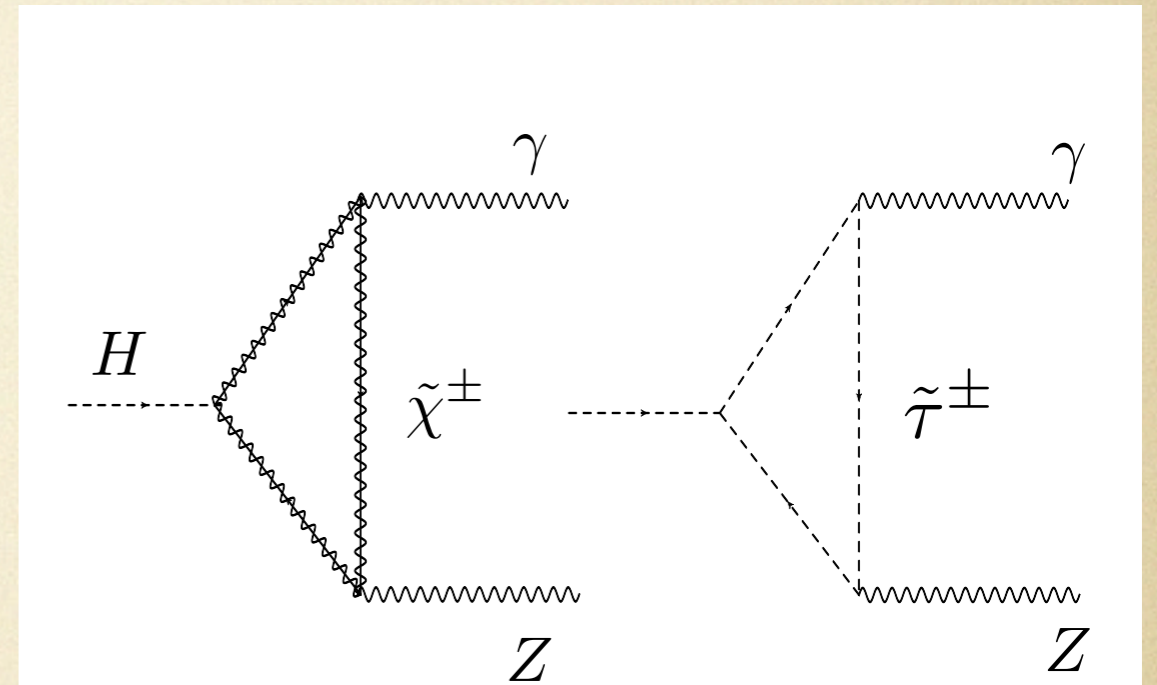
D0



## Example 2. Loop-induced



*2HDMs*



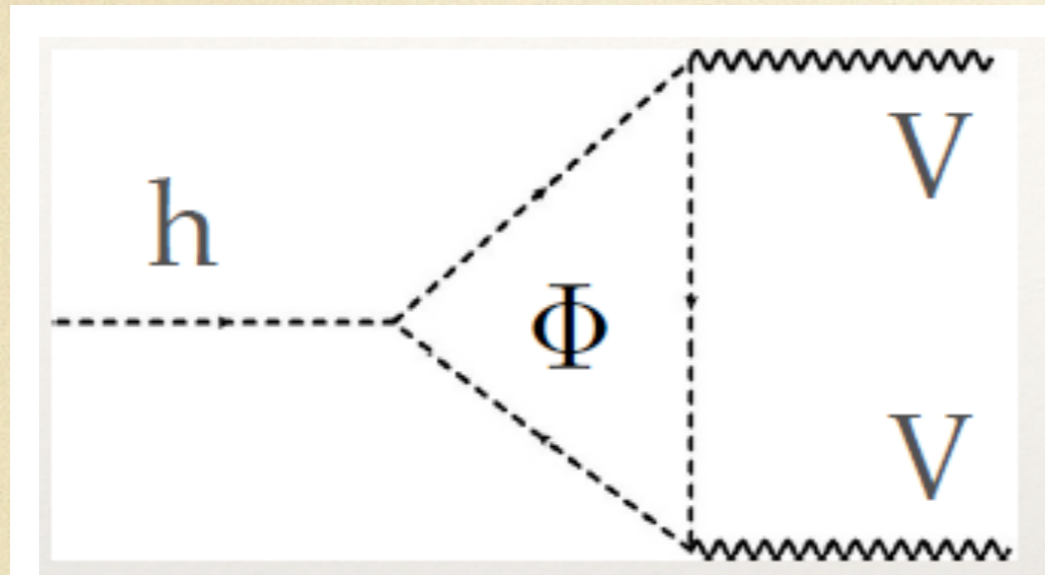
*SUSY spartners*

validity is now

$$\hat{s} \lesssim 4M_{\Phi}^2$$

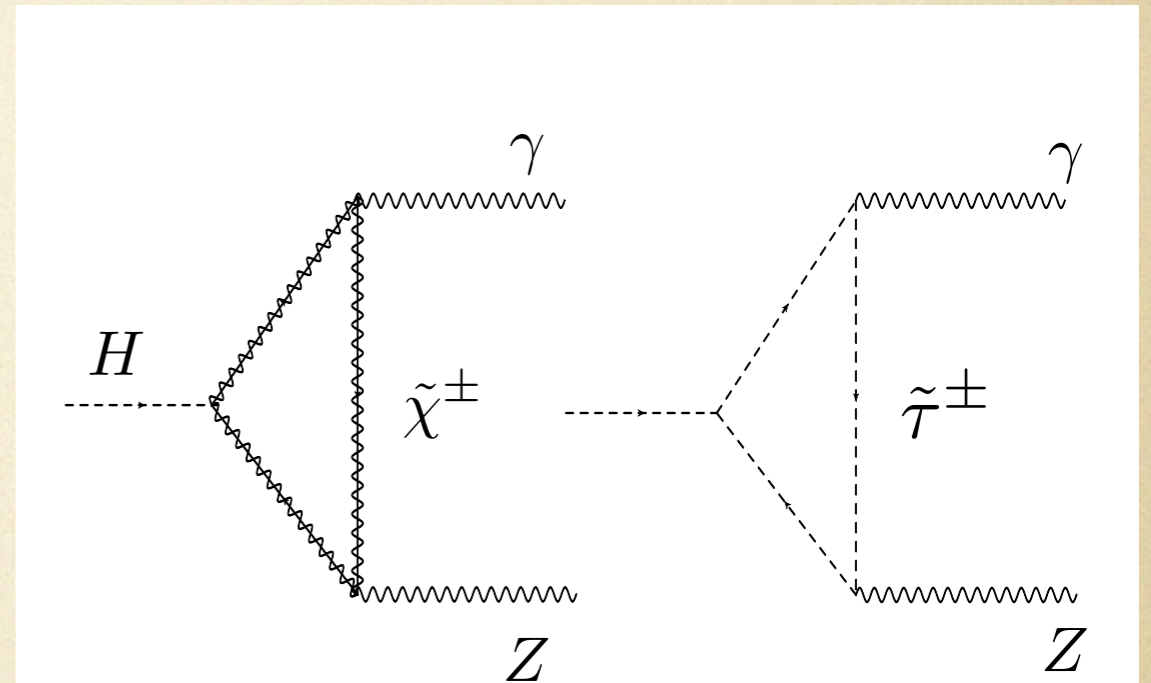


## Example 2. Loop-induced



*2HDMs*

Gorbahn, No and VS. In preparation



*SUSY spartners*

Masso and VS. 1211.1320

General predictions:

$$\bar{c}_W - \bar{c}_B = -(\bar{c}_{HW} - \bar{c}_{HB}) = 4\bar{c}_\gamma$$

$$\bar{c}_{HW} = -\bar{c}_W$$

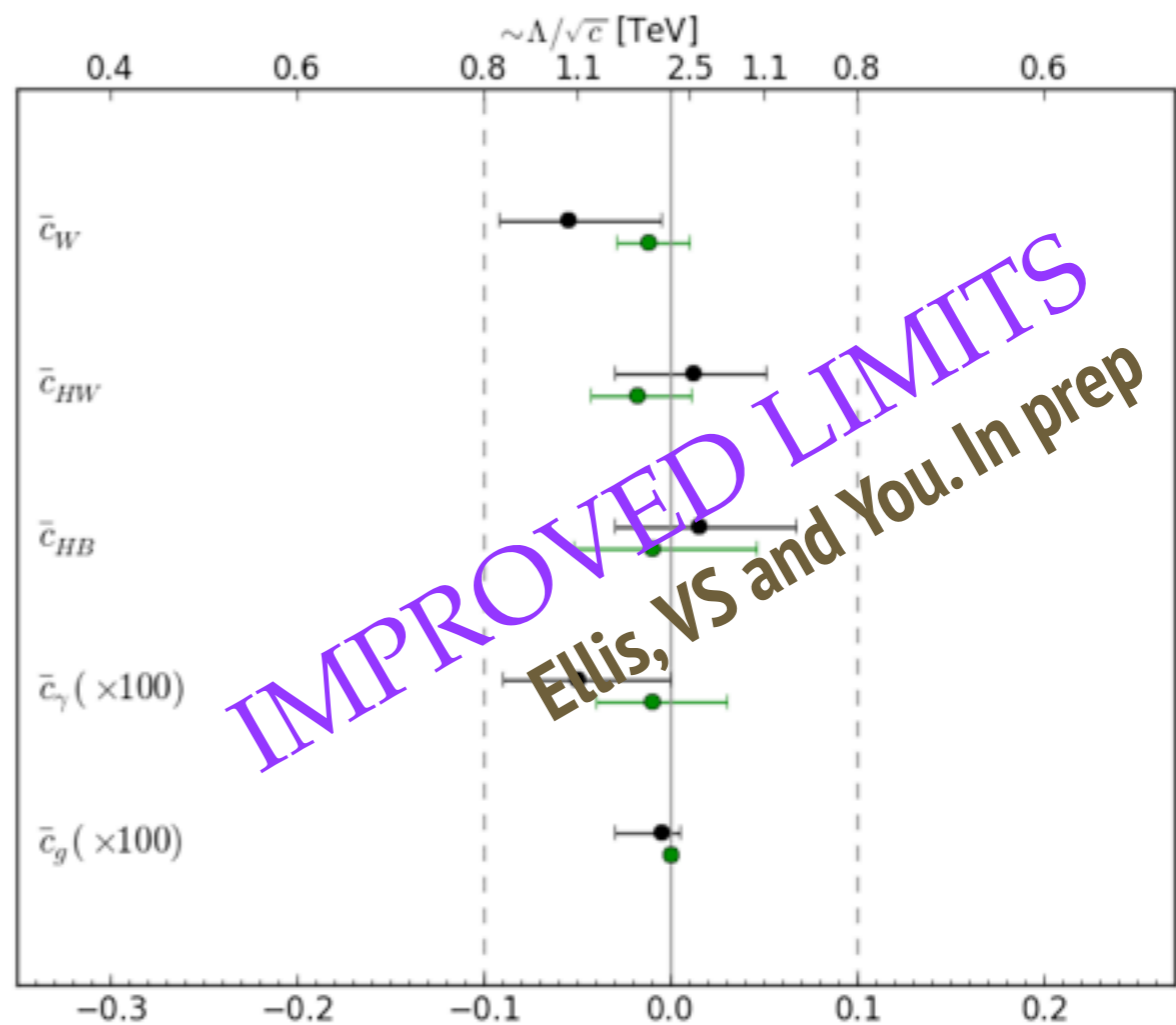
$$\bar{c}_{HB} = -\bar{c}_B$$



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$$\bar{c}_{HB} = -\bar{c}_B$$

## Matching to UV model

e.g. in the alignment limit

$$\bar{c}_{HW} = -\bar{c}_W = \frac{1}{6(16\pi^2)} (1 - x_0) \simeq 10^{-3} (1 - x_0)$$

$$\text{where } x_0 = \left( \frac{m_{H^0}}{m_{A^0}} \right)^2$$



# Conclusions

Absence of hints in direct searches  
EFT approach to Higgs physics

Higgs anomalous couplings:  
rates but also kinematic distributions

Complete global fit to Higgs physics  
enhanced using differential information

SM precision crucial: excess as **genuine** new physics

Exploring the validity of HEFT  
propose benchmarks

Benchmarks:  
correlations among coefficients, input for fit



# Framework for HDO studies

Feynrules HDOs involving Higgs and TGCs

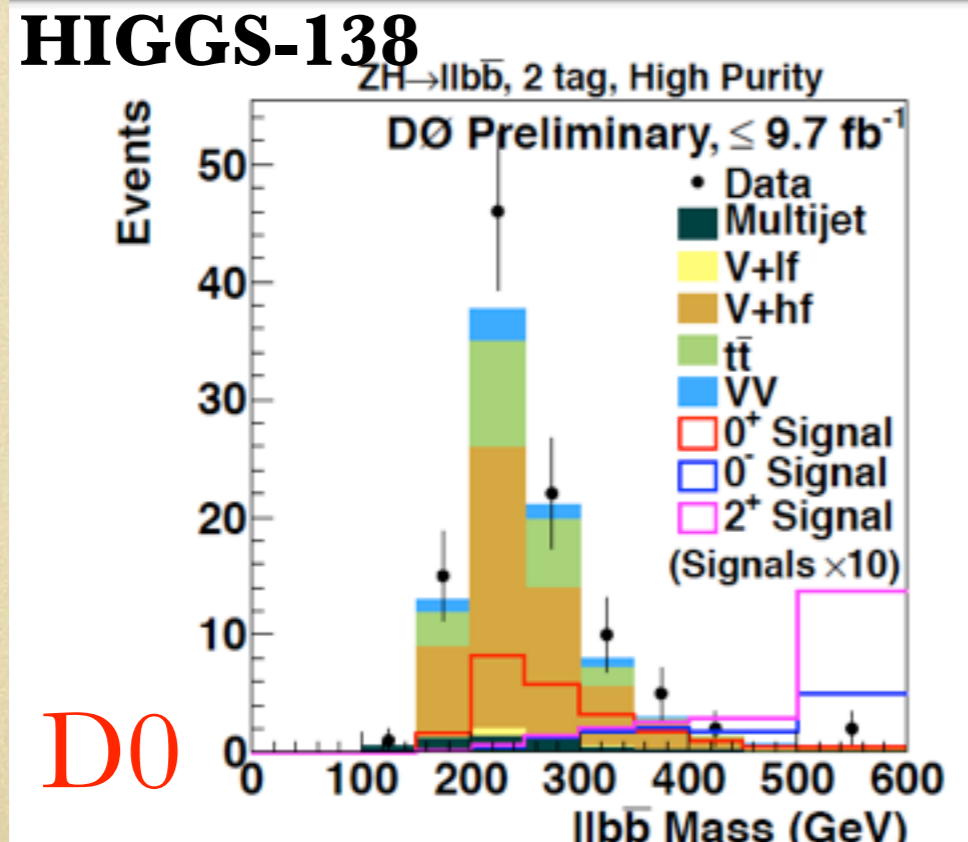
Alloul, Fuks, VS. 1310.5150

links to CalcHEP, LoopTools, Madgraph...

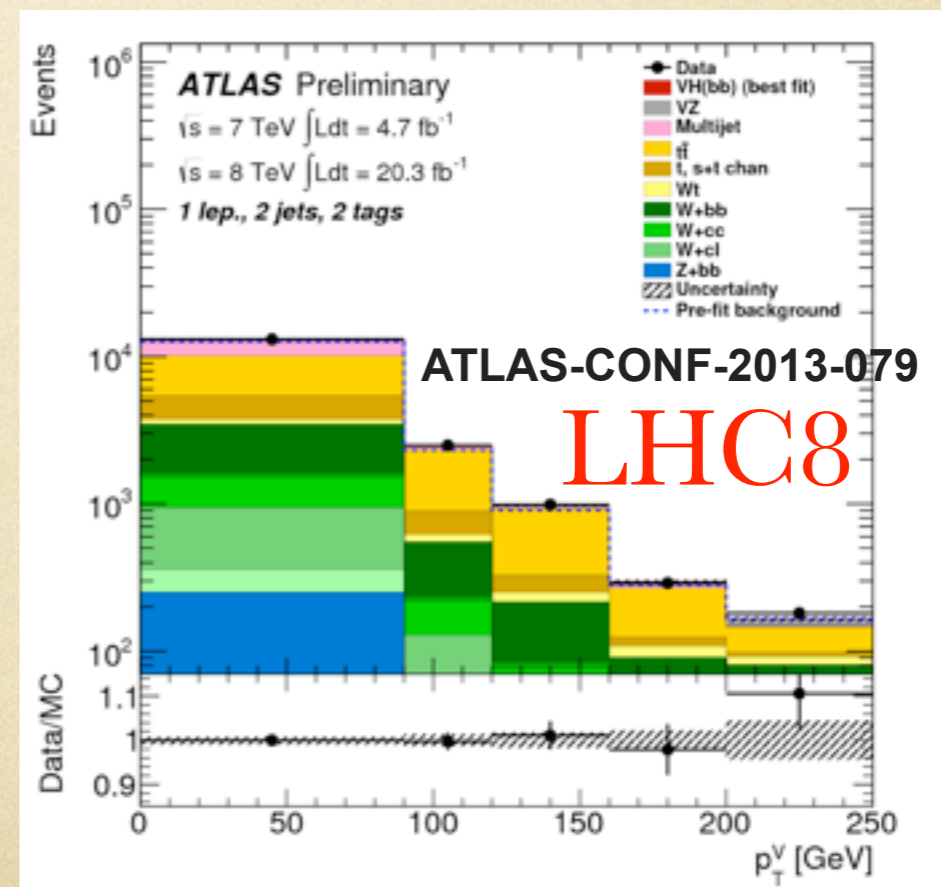
simulations: HDOs->Madgraph->Pythia... -> FastSim / FullSim

ex. Higgs in associated production

## HIGGS-138



$m_{Vh}$



$p_T^V$



# Framework for HDO studies

Feynrules HDOs involving Higgs and TGCs

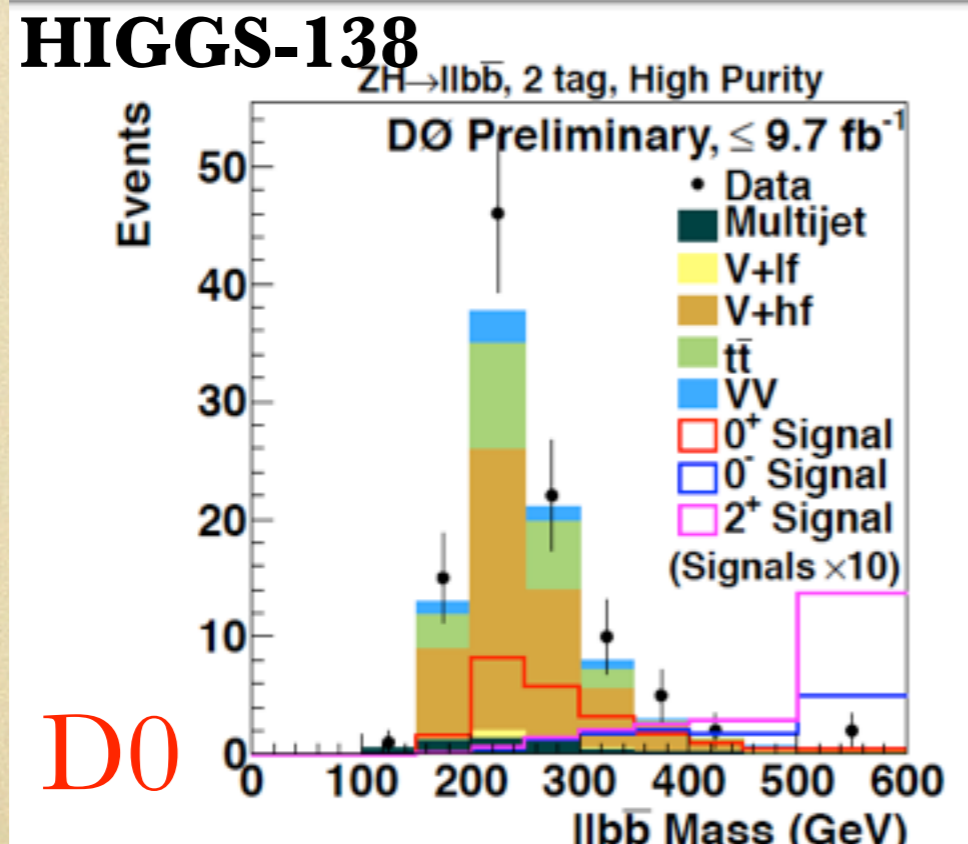
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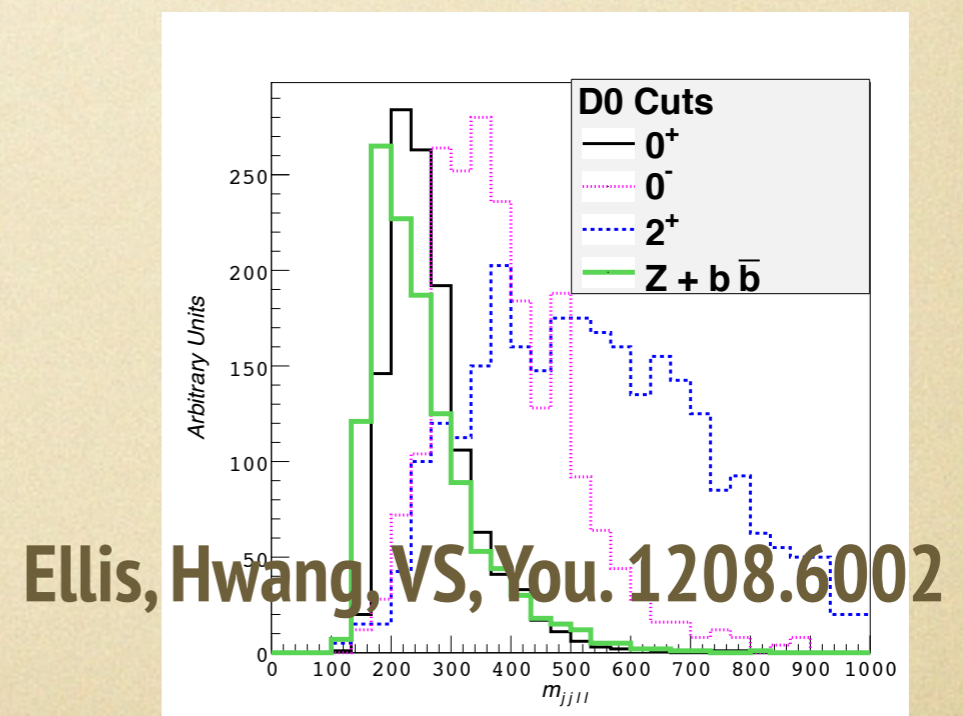
## HIGGS-138



D0

$m_{Vh}$

## Test other JCP



Ellis, Hwang, VS, You. 1208.6002



# Three CP-conserving operators affect TGCs

$$\begin{aligned} \mathcal{L}_{\text{TGC}}^{D=6} &= \frac{c_{WB} g_L g_Y}{m_W^2} B_{\mu\nu} W_{\mu\nu}^i H^\dagger \sigma^i H + \frac{i c_W g_L}{2m_W^2} \left( H^\dagger \sigma^i \overleftrightarrow{D}^\mu H \right) (D^\nu W_{\mu\nu})^i + \frac{c_{3W} g_L^3}{m_W^2} \epsilon^{ijk} W_{\mu\nu}^i W_{\nu\rho}^j W_{\rho\mu}^k \\ &+ \tilde{c}_{WB} \frac{g_L g_Y}{m_W^2} \tilde{B}_{\mu\nu} W_{\mu\nu}^i H^\dagger \sigma^i H + \frac{\tilde{c}_{3W} g_L^3}{m_W^2} \epsilon^{ijk} W_{\mu\nu}^i W_{\nu\rho}^j \tilde{W}_{\rho\mu}^k. \end{aligned} \quad (8)$$

$$\begin{aligned} \mathcal{L}_{\text{TGC}}^+ &= i(1 + \delta g_1^V) (W_{\mu\nu}^+ W_\mu^- - W_{\mu\nu}^- W_\mu^+) V_\nu + i(1 + \delta \kappa_V) V_{\mu\nu} W_\mu^+ W_\nu^- \\ &+ i \frac{\lambda_V}{m_W^2} W_{\mu\nu}^+ W_{\nu\rho}^- V_{\rho\mu} - g_5^V \epsilon_{\mu\nu\rho\sigma} (W_\mu^+ \partial_\rho W_\nu^- - \partial_\rho W_\mu^+ W_\nu^-) V_\sigma, \end{aligned}$$

$c_W$  and  $c_{WB}$   
affect Higgs physics  
and S-parameter,  
but more independent  
operators involved

dim-6 and TGCs

$$\begin{aligned} \delta \kappa_\gamma &= 4c_{WB}, \\ \delta \kappa_Z &= -4 \frac{g_Y^2}{g_L^2} c_{WB} - \frac{g_L^2 + g_Y^2}{2g_L^2}, \\ \delta g_Z &= -\frac{g_L^2 + g_Y^2}{2g_L^2} c_W, \\ \lambda_\gamma = \lambda_Z &= -6g_L^2 c_{3W}, \\ \tilde{\kappa}_\gamma &= 4\tilde{c}_{WB}, \\ \tilde{\kappa}_Z &= -4 \frac{g_Y^2}{g_L^2} \tilde{c}_{WB}, \\ \tilde{\lambda}_\gamma = \tilde{\lambda}_Z &= -6g_L^2 \tilde{c}_{3W}, \end{aligned}$$



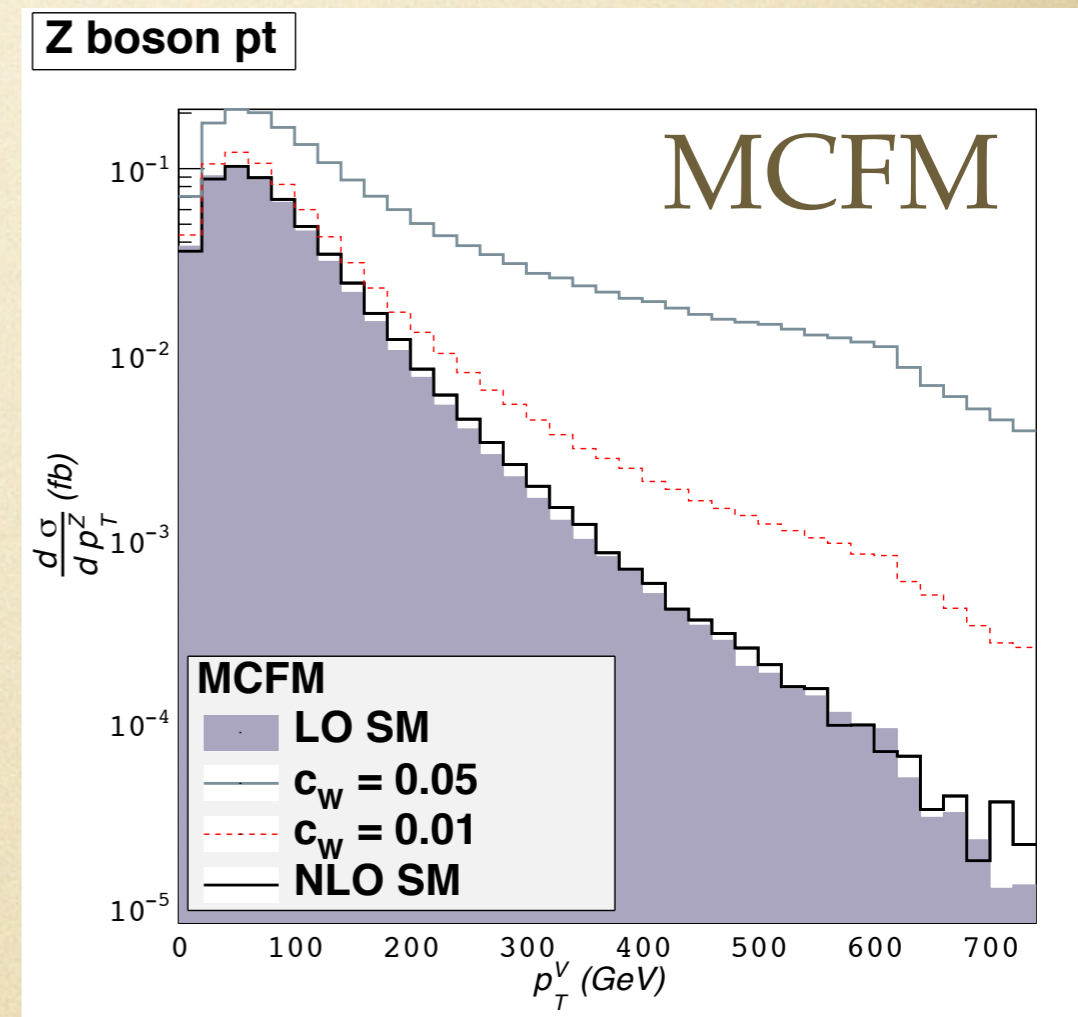
# Kinematics of associated production

comment 1:

pTV is more sensitive than mVH to QCD NLO  
but effect not yet at the level of operator values we can  
bound

comment 2:

Sensitivity to quadratic orders  
in  $c$ 's (dim-8) is less than current  
errors.



VS and Williams. In prep.



# Boring and necessary details

Bottom-up approach:  
operators w / SM particles and symmetries,  
plus the **newcomer**, the **Higgs**



# Boring and necessary details

Bottom-up approach:  
operators w / SM particles and symmetries,  
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## Realization of EWSB

Linear or non-linear



# Boring and necessary details

Bottom-up approach:  
operators w/ SM particles and symmetries,  
plus the **newcomer**, the **Higgs**



## Realization of EWSB

Linear or non-linear



**And the Higgs could be**

Weak doublet or singlet



Once this choice is made, expand...

$$\frac{1}{\Lambda^2}$$

Integrating out new physics

$$\frac{v^2}{f^2}$$

Non-linearity

$$U = e^{i\Pi(h)/f}$$

...order-by-order



For example, some operators  
Higgs-massive vector bosons

ex.

$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$

$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_B = (D_\mu \Phi)^\dagger (D_\nu \Phi) \widehat{B}^{\mu\nu}$$

$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}^{\mu\nu} \widehat{W}_{\mu\nu} \Phi$$

$$\mathcal{O}_{BB} = (\Phi^\dagger \Phi) \widehat{B}^{\mu\nu} \widehat{B}_{\mu\nu}$$



For example, some operators  
Higgs-massive vector bosons

ex.

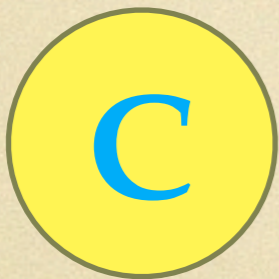
$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$

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$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}^{\mu\nu} \widehat{W}_{\mu\nu} \Phi$$

$$\mathcal{O}_{BB} = (\Phi^\dagger \Phi) \widehat{B}^{\mu\nu} \widehat{B}_{\mu\nu}$$



UV theory: tree-level or loop

may need a model bias

ex. SILH

$$\frac{2igc_{HW}}{m_W^2} (D^\mu \Phi^\dagger) \widehat{W}_{\mu\nu} (D^\nu \Phi)$$



redundancies trade off operators using EOM

**D** Choice of basis

And, finally

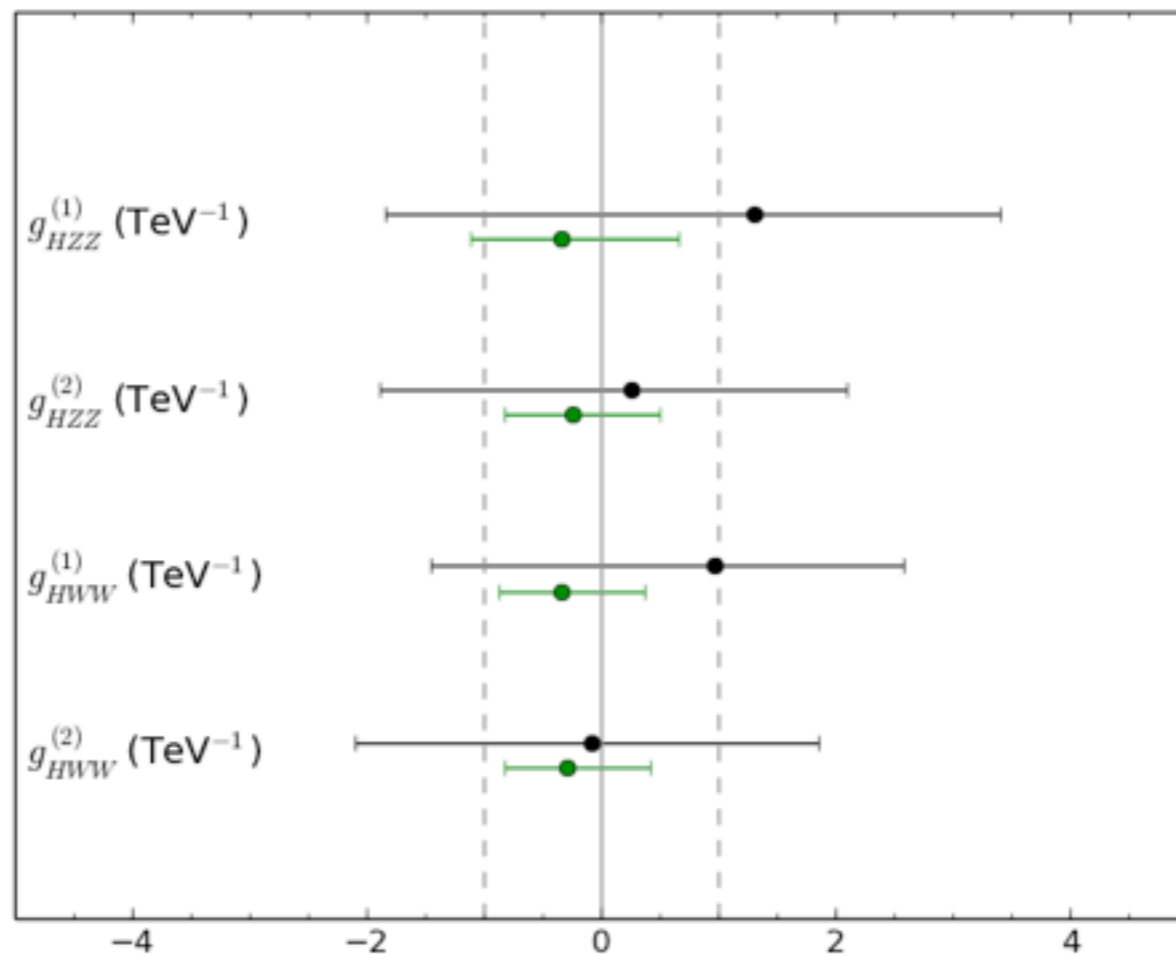
Observables as a function  
of HDOs coefficients



In summary

In terms of Higgs' anomalous couplings

$$\mathcal{L} \supset - \frac{1}{4} g_{HZZ}^{(1)} Z_{\mu\nu} Z^{\mu\nu} h - g_{HZZ}^{(2)} Z_\nu \partial_\mu Z^{\mu\nu} h \\ - \frac{1}{2} g_{HWW}^{(1)} W^{\mu\nu} W_{\mu\nu}^\dagger h - \left[ g_{HWW}^{(2)} W^\nu \partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.} \right],$$



black global fit  
green one-by-one fit







Global fit to **signal strengths**  
and **kinematic distributions**

Conclusions of the analysis

1. Breaking of blind directions requires information on associated production (AP)
2. Kinematic distributions in AP is as sensitive (or more) than total rates



# Global fit to signal strengths and kinematic distributions

## Conclusions of the analysis

1. Breaking of blind directions requires information on associated production (AP)

2. Kinematic distributions in AP is as sensitive (or more) than total rates