

④ SUGRA in SSP II. Higher N and higher D 4.1-④  
(con-shell)

4.1. WARM-UP: SYM IN FLAT SSP.

SPACE TIME GAUGE FIELD  $\Leftrightarrow$  1-FORM

$D=3,4,\dots,10$

$$F_2 = \frac{1}{2} dx^\mu \wedge dx^\nu F_{\mu\nu} \equiv \frac{1}{2} e^a{}_\mu e^b{}_\nu F_{ab} = dA - A \wedge A$$

$$\delta F_2 = D\delta A, \quad A = dx^\mu A_\mu = e^a{}_\mu A_a$$

IN FLAT SPACE  $e^a{}_\mu(x) = dx^\mu \delta_\mu^a$ .

• FERMIONIC FIELD IN Adj of gauge group,  $\chi^\alpha$

$$D\chi^\alpha = d\chi^\alpha + [A, \chi^\alpha]$$

$$DD\chi^\alpha = [F_2, \chi^\alpha] = dD\chi^\alpha - A \wedge D\chi^\alpha - D\chi^\alpha \wedge A$$

be careful with diff. form! (or work with components)

$$[A, \chi^\alpha] = D(A\chi^\alpha - \chi^\alpha A) = A \wedge d\chi^\alpha + d\chi^\alpha \wedge A + [DA, \chi^\alpha]$$

Spacetime action in 10D

$$S = \int d^{10}x \left( \frac{1}{4g^2} t_2 F_{\mu\nu} F^{\mu\nu} + i t_2 \chi^\alpha \overleftrightarrow{D}_\mu \chi^\alpha \right)$$

$\mu=0,1,\dots,9; \alpha=1,\dots,16$

Lower D actions may be obtained by dimensional reduction.

What we need to know about 10D KG coeffs (generalized Pauli matrices)

$$\begin{aligned} \sigma_{\alpha\beta}^a &= \sigma_{\beta\alpha}^a, & \tilde{\sigma}^{b\alpha\beta} &= \tilde{\sigma}^{b\beta\alpha} & [16 \times 16] \\ (\tilde{\sigma}^{ab})_\alpha{}^\beta &= \gamma^{ab} \delta_\alpha{}^\beta + (\sigma^{ab})_\alpha{}^\beta & & & \\ \sigma_{a(\alpha\beta} \sigma_{\gamma)}^\alpha &\equiv 0; & \alpha, \beta, \gamma &= 1, \dots, 16 \end{aligned}$$

No  $C^{\alpha\beta}$ ;  $C_{\alpha\beta}$ ! (in MW repr.)

$$\delta S = \int d^{10}x \left( \left( -\frac{1}{g^2} D_\mu F^{\mu\nu} - i \sigma_{\alpha\beta}^\nu \chi^\alpha \chi^\beta \right) \delta A_\nu + 2i \delta \chi^\alpha \overleftrightarrow{D}_\mu \chi^\alpha + \partial_\mu (\dots) \right)$$

SUSY (ex: to obtain  $\delta_\epsilon S = 0$ )

$$\begin{aligned} \delta A_\mu &= \epsilon \sigma_\mu \chi \\ \delta \chi^\alpha &= -\frac{i}{4g^2} (\epsilon \sigma^{\mu\nu})^\alpha{}_\beta F_{\mu\nu} \end{aligned}$$

Notice the  $\sigma_{\alpha\beta}^\nu \chi^\alpha \chi^\beta \delta_\epsilon A_\nu = 0$  due to  $\sigma_{a(\alpha\beta} \sigma_{\gamma)}^\alpha = 0$

SSM Lagrangian 100.  $t_2 X^T \epsilon_1 [SA_1, X^T] = -t_2 (X^T \epsilon_1^T X^T + X^T X^T \epsilon_1^T - t_2^{-1} X^T \epsilon_1^T SA_1$

$$S = \int_{t_0}^{t_1} dt [F_m F_m^T + i \epsilon_1^T X^T \epsilon_1^T D_t X^T] + [A_1, X^T]$$

$$SS = \int_{t_0}^{t_1} dt \left( \frac{1}{2} t_2 F_m^T D_t S A_m + z_1 t_2 X^T \epsilon_1^T D_t X^T + i \epsilon_1^T X^T \epsilon_1^T [SA_1, X^T] \right) - \frac{1}{2} t_2 D_t F_m^T S A_m$$

$$SA_1 = \epsilon_1^T X^T$$

$$S X^T = -\frac{i}{2} \epsilon_1^T \epsilon_1^T X^T$$

$$S \frac{1}{2} t_2 F_m^T F_m^T = -\frac{1}{2} t_2 F_m^T F_m^T \epsilon_1^T D_t X^T$$

$$S : t_2 X^T \epsilon_1^T D_t X^T = 2 \frac{1}{2} t_2 \epsilon_1^T \epsilon_1^T D_t X^T \cdot F_m = \frac{1}{2} t_2 \epsilon_1^T D_t X^T F_m + \frac{1}{2} t_2 \epsilon_1^T \epsilon_1^T D_t X^T F_m$$

$$t_2 \{X^T, X^T\} \epsilon_1^T \epsilon_1^T = \epsilon_1^T t_2 \{X^T, X^T\} \epsilon_1^T \epsilon_1^T$$

SSP  $A_1 = \Delta Z^T A_m(Z) = E^T A_a + E^T A_b$

100  $E_a = \Delta X^a - i \theta \delta^a \theta$ ;  $E^a = \Delta \theta^a$   $a = 1, \dots, 16$

$$F_{ab} = E_a^T E_b^T F_{ab} + \frac{1}{2} E_a^T A_a E_b^T F_{ab} = -(F)^T$$

$$0 = D F = -i E_a^T E_a^T E^T \epsilon_1^T \epsilon_1^T F_{ab} + E_a^T E_a^T \epsilon_1^T D_a F_{ab} - i E_a^T E_a^T \epsilon_1^T \epsilon_1^T F_{ab} + \frac{1}{2} E_a^T E_a^T E^T (D_a F_{ab} - 2 D_b F_{ab}) + \frac{1}{2} E_a^T E_a^T \epsilon_1^T \epsilon_1^T D_b F_{ab}$$

$$D_a W^T = \frac{1}{2} \epsilon_1^T \epsilon_1^T F_{ab}$$

$$D_a F_{ab} = -2 \epsilon_1^T \epsilon_1^T D_b W^T$$

$$\epsilon_1^T \epsilon_1^T W^T = i \epsilon_1^T \epsilon_1^T D_a W^T = \frac{1}{2} \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T D_a F_{ab} = -i \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T D_b W^T$$

$$i \epsilon_1^T \epsilon_1^T D_a W^T = -\frac{i}{2} \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T D_b W^T = + \frac{1}{2} \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T D_b W^T$$

$$D X^T = \Delta X^T + [A_1, X^T]$$

$$D D X^T = [F, X^T] = \frac{d D X^T}{dt} - [F A_1, D X^T] + [A_1, X^T]$$

$$0 = \epsilon_1^T \epsilon_1^T D_a W^T = \epsilon_1^T \epsilon_1^T D_a D_b W^T + \epsilon_1^T \epsilon_1^T [F_{ab}, W^T] = \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T D_b W^T + \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T [F_{ab}, W^T]$$

$$M = D F = 0$$

$$I = D M = \frac{1}{2} E_a^T E_a^T E^T \epsilon_1^T \epsilon_1^T F_{ab} + \frac{1}{2} D_a (F_{ab}) + \frac{1}{2} E_a^T E_a^T \epsilon_1^T \epsilon_1^T D_b F_{ab}$$

$$0 = 3 \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T = \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T + 2 \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T$$

$$3 \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T = \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T \Rightarrow 3 \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T = 10 \epsilon_1^T \epsilon_1^T \epsilon_1^T \epsilon_1^T$$



10D SYM IN SSP  $\sum_0^{(10|16)}$  ( $Z^M = (x^\mu, \theta^\alpha)$   $\mu=0,1,\dots,9$   $\alpha=1,\dots,16$ )

$$A = E^\alpha A_\alpha + E^\alpha A_\alpha = dZ^M A_M(Z)$$

$$E^\alpha = dx^\alpha - i d\theta^\alpha \delta_{\alpha\beta}^a \theta^\beta, \quad E^a = d\theta^a$$

SSP CONSTRAINTS / STRUCTURE EQS OF  $\sum_0^{(10|16)}$   $T^a = dE^a = -i E^\alpha \wedge E^\beta \delta_{\alpha\beta}^a$   
 and  $R^{ab} = 0$  so we gauge fix  $\omega^{ab} = 0$   $T^\alpha = dE^\alpha = 0$

WHAT CONSTRAINTS WE HAVE TO IMPOSE ON  $F = dA$ ?

SUGGESTION:

we expect to have  $\delta_\epsilon A = i_\epsilon F = E^\alpha E^\beta \delta_{\alpha\beta}^a \chi = i_\epsilon (E^\alpha \wedge E^\beta \delta_{\alpha\beta}^a \chi)$

(THIS IS JUST SUGGESTION: WE ARE WITH RIGID SUSY, WHERE  $i_\epsilon E^a \neq 0$  BUT  $\delta_\epsilon E^a = 0$ )

SO WE IMPOSE

$F_{\alpha\beta} \delta_{\alpha\beta}^a = 0$  essential  
 $F_{\alpha\beta} \delta_{\alpha\beta}^a = 0$  CONVENTIONAL  
 $F_{\alpha\beta} = 0$   $\{D_\alpha, D_\beta\} = 2i \delta_{\alpha\beta}^a D_a$   
 DETERMINES  $A_a$  IN TERMS OF  $A_\alpha$ , ALSO  $F_{\alpha\beta} = \delta_{\alpha\beta}^a W^a$  WORKS!  
 ACTUALLY, AS WE WILL SEE IN A MOMENT,

$$F = dA - A \wedge A = E^\beta \wedge E^\alpha F_{\alpha\beta} + \frac{1}{2} E^\beta \wedge E^\alpha F_{ab} = -(F)^\dagger$$

$$0 = DF = -i E^\alpha \wedge E^\beta \delta_{\alpha\beta}^a F_{ab} + E^\beta \wedge E^\alpha (D_\alpha F_{\beta b} - i \delta_{\alpha\beta}^a F_{ab}) + \frac{1}{2} E^\beta \wedge E^\alpha \wedge E^\gamma (D_\alpha F_{\beta\gamma} - 2 D_{[\alpha} F_{\beta\gamma]}) + \frac{1}{2} E^\alpha \wedge E^\beta \wedge E^\gamma D_\alpha F_{\beta\gamma}$$

dim  $3/2$ :  $\sim E^\alpha \wedge E^\beta \wedge E^\gamma$   
 $(M_{\alpha\beta\gamma} = 0)$

THE GENERAL SOLUTION  $\delta_{\alpha\beta}^a F_{ab} = 0$

$F_{\alpha\beta} = \delta_{\alpha\beta}^a W^a$

REMEMBER:  $\delta_{\alpha\beta}^a \delta_{\gamma\delta}^b \delta = 0$  in  $D=10!$  (3,4,6)

dim 2:  $\sim E^\alpha \wedge E^\beta$

$$D_\alpha F_{\beta b} - i \delta_{\alpha\beta}^a F_{ab} = 0$$

$$\delta_{\alpha\beta}^a D_{[\alpha} W^{\beta]} = D_{[\alpha} W^{\beta]} \delta_{\alpha\beta}^a$$

To solve this eqs we should appreciate that

$\delta_{\alpha\beta}^a, \delta_{\alpha\beta}^{a_1 a_2 a_3 a_4}$  - symm.  
 $\delta_{\alpha\beta}^{abc}$  - asymm.

decompose

$$D_\beta W^\beta = a_0 \delta_\beta^\beta + a_{ab} \delta_{ab}^\beta + a_{a_1 \dots a_4} \delta_{a_1 \dots a_4}^\beta$$

and find that eq. implies  $a_0 = 0, a_{a_1 \dots a_4} = 0$

$$a_{bd} (\delta^{bd} \delta_{\alpha\beta}^a) = +2 a_{ab} \delta_{\alpha\beta}^a$$

$D_\alpha W^\beta = \frac{i}{2} \delta_{\alpha\beta}^a F_{ab}$

ACTUALLY THESE ARE ALL THE INDEPENDENT RELATIONS FOR MAIN SUPERFIELDS

AS CAN BE PROVED BY STUDYING IDENTITIES FOR IDENTITIES

$DM_3 = 0$

where  $M_3 \equiv DF$  ARE L.H.S.-S OR R.H.S.-S



EX: TO PROVE THE DEPENDENCE OF HIGHER DIM. IS IS  
 $M_{abc} = 0$  and  $M_{aba} = 0$

PROOF: WRITE "IDENTITY FOR IDENTITY"  $0 = DM_3$  setting  $M_{pqr} = 0, M_{apq} = 0$ :

$$0 = \frac{1}{2} E^a \wedge E^b \wedge E^c E^d (-i \delta_{ap}^c M_{abc} + \frac{1}{2} D_a M_{pab}) + \frac{1}{2} E^b \wedge E^c \wedge E^d E^a \delta_{ap}^c M_{abc} + \dots$$

- If we set  $M_{abd} = 0$ , this gives  $\delta_{ap}^c M_{abc} = 0 \Rightarrow M_{abc} = 0$   
Thus dim 3 id. is dependent provided dim 5/2 id. is.
- For  $M_{aba}$  we have  $\delta_{ap}^c M_{cjab} = 0$

contracting with  $\frac{3}{2} \delta_{a1}^{b2}$  we find  $M_{cab} = -(\delta_{ca}^b)^p M_{pcb} = -M_{cab} + \delta_{apc} \delta^{cpq} M_{qcb}$   
 $3M_{cab} = \delta_{apc} \delta^{cpq} M_{qcb} \Rightarrow M_{cab} = 0$

(AND CONVENIENT) BUT IN OUR SIMPLE CASE IT IS NOT DIFFICULT TO STUDY ALL  $M_{aba}$   $M_{abc}$

dim 5/2:  $D_a F_{ab} = -2 \delta_{[a\alpha\beta} D_{b]} W^\beta$  dim 3:  $D_a F_{bc} = 0$

THESE CONSTRAINTS DESCRIBE 10D SYM MODEL ON THE MASS SHELL (ON-SHELL)

THE FERMIONIC EQS CAN BE EXTRACTED FROM

$$D_\beta W^\alpha = \frac{1}{2} \delta_{\beta}^{ab} F_{ab}$$

Take  $D_\gamma$  derivative, symmetrize  $(\beta\gamma)$  use dim 5/2 (BI)

$$i \delta_{\beta\gamma}^a D_a W^\alpha = \frac{i}{2} \delta_{\beta\gamma}^{bc} D_b F_{bc}$$

THE RESULTING EQ IS SOLVED BY (SIMPLEST WAY: TO CONTRACT WITH  $\delta_a^\alpha$ )

$$i \delta_{\beta\gamma}^a D_a W^\alpha = 0$$

$$-2 \delta_{[\beta\gamma]c} D_c W^\alpha$$

THE BOSONIC EQS CAN BE OBTAINED FROM  $D_\alpha$  OR ONE FERMIONIC EQS.

$$0 = \delta_{ap}^q D_\gamma D_a W^\alpha = \delta_{ap}^q D_a D_\gamma W^\alpha + \delta_{ap}^q \{F_{\gamma a}, W^\alpha\}$$

$\delta_{ap}^q D_a D_\gamma W^\alpha \leftarrow \frac{1}{2} F_{bc} \delta_{ap}^q \delta_{\gamma}^a$       $\delta_{ap}^q \{F_{\gamma a}, W^\alpha\} \leftarrow \delta_{ap\gamma} \{W^\alpha, W^\alpha\}$

$\frac{1}{2} \delta_{ap}^q \delta_{\gamma}^a D_b F_{bc} - i \delta_{\beta\gamma}^b D_b F_{ca} \leftarrow \delta_{ap\gamma} \delta_{abd} = 0 \leftarrow -2 \delta_{\beta\gamma}^a \delta_{abd} \{W^\alpha, W^\alpha\}$

$$\Rightarrow D_a F_{ab} = 2i \delta_{bap} \{W^\alpha, W^\alpha\}$$



SYM IN FLAT D=4 N=1 SSP

BY ANALOGY WITH 10D SYM, THE CONSTRAINTS ARE

$F_{\alpha\beta} = 0$  ,  $F_{\dot{\alpha}\dot{\beta}} = 0$  ,  $F_{\alpha\dot{\beta}} = 0$   
 $\{D_\alpha, D_\beta\} = 0$  ESSENTIAL ,  $\{D_{\dot{\alpha}}, D_{\dot{\beta}}\} = 0$  CONVENTIONAL  
 $\{D_\alpha, D_{\dot{\beta}}\} = 2i\delta_{\alpha\dot{\beta}} \partial_a$   
 ALLOW FOR THE EXISTENCE OF CHIRAL SUPERFIELDS WITH VALUES IN A NONTRIVIAL REPRESENTATION OF THE GAUGE GROUP  
 GRASSMANN ANALYTICITY

BIS: for  $F = E^b{}_\lambda E^\alpha F_{ab} + E^b{}_\lambda \bar{E}^{\dot{\alpha}} F_{\dot{a}b} + \frac{1}{2} E^b{}_\lambda E^{\dot{\alpha}} F_{ab}$

$$\begin{aligned}
 0 = D F &= 2i E^\alpha{}_\lambda E^{\dot{\beta}} E^{\dot{\gamma}} F_{(\alpha\dot{\beta})\dot{\gamma}} + 2i E^\alpha{}_\lambda \bar{E}^{\dot{\alpha}} E^{\dot{\beta}} F_{(\alpha\dot{\beta})\dot{\gamma}} + \\
 &+ E^b{}_\lambda E^{\dot{\alpha}} E^{\dot{\beta}} \partial_a F_{\dot{a}b} + E^b{}_\lambda \bar{E}^{\dot{\alpha}} E^{\dot{\beta}} \bar{\partial}_{\dot{a}} F_{\dot{a}b} + E^b{}_\lambda E^{\dot{\alpha}} \bar{E}^{\dot{\beta}} (\partial_{\dot{a}} F_{\dot{a}b} + \bar{\partial}_{\dot{a}} F_{\dot{a}b} - 2i\delta_{\dot{a}b} F_{\dot{a}b}) \\
 &+ \frac{1}{2} E^\alpha{}_\lambda E^{\dot{\beta}} E^{\dot{\gamma}} (\partial_a F_{bc} - 2\partial_{[b} F_{a]c}) + \frac{1}{2} E^\alpha{}_\lambda E^{\dot{\beta}} \bar{E}^{\dot{\gamma}} (\bar{\partial}_{\dot{a}} F_{bc} - 2\bar{\partial}_{[\dot{a}} F_{b]c}) + \frac{1}{3!} E^\alpha{}_\lambda E^{\dot{\beta}} E^{\dot{\gamma}} \partial_a F_{bc}
 \end{aligned}$$

dim  $3/2$ :  $F_{(\alpha\dot{\beta})\dot{\gamma}} = 0 \Rightarrow F_{ab} = i\delta_{ba} W^\alpha$ , and c.c.  $F_{\dot{a}b} = -i\delta_{b\dot{a}} W^\alpha$

$$F = i E^b{}_\lambda E^\alpha (\delta_b^\alpha W)_\lambda - i E^b{}_\lambda \bar{E}^{\dot{\alpha}} (W \delta_b^{\dot{\alpha}})_\lambda + \frac{1}{2} E^b{}_\lambda E^{\dot{\alpha}} F_{ab} = -F^\dagger$$

dim 2:  $\forall E^\alpha{}_\lambda E^{\dot{\beta}} E^{\dot{\gamma}} \partial_a W^\alpha = 0$  c.c.  $\bar{\partial}_{\dot{a}} W^\alpha = 0$

$$\begin{aligned}
 \forall E^b{}_\lambda E^{\dot{\alpha}} \bar{E}^{\dot{\beta}} \partial_a W^\alpha & \quad 2i\delta_{\dot{a}b} F_{ab} = i\delta_{\dot{a}b} \partial_a W^\alpha - i\partial_a W^\beta \delta_{b\dot{\alpha}} \\
 \Leftrightarrow & \quad F_{\alpha\dot{\beta}\dot{\gamma}} = \epsilon_{\dot{\beta}\dot{\gamma}} \partial_a W^\alpha - \epsilon_{\dot{\beta}\dot{\gamma}} \bar{\partial}_{\dot{a}} W^\alpha \\
 \Leftrightarrow & \quad \partial^\alpha W_\alpha = \bar{\partial}^{\dot{\alpha}} W_{\dot{\alpha}}
 \end{aligned}$$

$$F_{ab} = \frac{1}{4} \partial_a W^\beta (\delta_{ab})_\beta^\alpha + \frac{1}{4} (\delta_{ab})_\beta^\alpha \bar{\partial}_{\dot{a}} W^\alpha$$

Notice:  $\partial^\alpha W_\alpha + \bar{\partial}^{\dot{\alpha}} W_{\dot{\alpha}} = 2\mathbb{D}$  generically!  $\neq 0$   
 ITS LEADING ( $\theta=0$ ) COMPONENT THIS IS THE AUXILIARY FIELD OF SYM

$$\partial^\alpha W_\alpha = \mathbb{D} \quad , \quad \bar{\partial}^{\dot{\alpha}} W_{\dot{\alpha}} = \mathbb{D}$$

SUPERFIELD EQS OF MOTION AS  $\mathbb{D} = 0$

Then  $\partial^\alpha W_\alpha = 0 \Rightarrow \partial^{\alpha\dot{\alpha}} W_{\dot{\alpha}} = 0$  (COV. FERMIONIC EQ.)



Georgie field theory in SSP of unit 5G

$$F = dA = A \wedge A = \frac{1}{2} d^2 x^\mu d^2 x^\nu F_{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta} = F^I (T_I)^a_{bc} \dots$$

$$= \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} A_{\mu} A_{\nu} A_{\alpha} A_{\beta} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} A_{\mu} A_{\nu} A_{\alpha} A_{\beta} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} A_{\mu} A_{\nu} A_{\alpha} A_{\beta} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} A_{\mu} A_{\nu} A_{\alpha} A_{\beta}$$

$$F^T = -F$$

$$F_{\mu\nu} = 0 = F_{\mu\nu}, \quad F_{\mu\nu} = 0$$

$$F = E^i A^j F_{ij} + E^i A^j \partial^k F_{ik} + \frac{1}{2} \epsilon^{ijkl} F_{ij} F_{kl}$$

$$0 = D F = -T^i A^j F_{ij} - T^i A^j \partial^k F_{ik} + \epsilon^{\mu\nu\alpha\beta} F_{\mu\nu} F_{\alpha\beta} + \epsilon^{\mu\nu\alpha\beta} F_{\mu\nu} F_{\alpha\beta} + \epsilon^{\mu\nu\alpha\beta} F_{\mu\nu} F_{\alpha\beta}$$

$$\sim E^i A^j \partial^k F_{ik}$$

$$F_{\alpha\beta} = i \epsilon_{\alpha\beta} W^i$$

$$\sim E^i A^j \partial^k F_{ik}$$

$$0 = i \epsilon_{\alpha\beta} D_{\mu} W^{\mu}$$

$$D_{\mu} W^{\mu} = 0$$

$$\sim E^i A^j \partial^k F_{ik}$$

$$0 = -i D_{\mu} W^{\mu} \epsilon_{\mu\nu\alpha\beta} + i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} - 2 i \epsilon_{\mu\nu\alpha\beta} F_{\mu\nu}$$

$$D^{\mu} W_{\mu} - D^{\nu} W_{\nu} = 0$$

$$F_{\alpha\beta} = \frac{1}{4} D_{\mu} W^{\mu} \epsilon_{\alpha\beta} + \frac{1}{4} \epsilon_{\alpha\beta} D_{\mu} W^{\mu}$$

$$\sim E^i A^j \partial^k F_{ik}$$

$$0 = \frac{1}{2} D_{\mu} F_{\mu} + i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} + \frac{1}{2} R \epsilon_{\mu\nu\alpha\beta} W^{\mu}$$

$$D_{\mu} F_{\mu} = -2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} - R \epsilon_{\mu\nu\alpha\beta} W^{\mu}$$

$$D_{\mu} F_{\mu} = 2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} + R \epsilon_{\mu\nu\alpha\beta} W^{\mu}$$

$$\Rightarrow D_{\mu} F_{\mu} = 2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} + 2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} = -2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu} + 2 i \epsilon_{\mu\nu\alpha\beta} D_{\mu} W^{\nu}$$

U.1 - (21)  
S. 17. - (15)  
S. 16