

The End & Other Tall Tales

@ A. Padilla, Nottingham

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Sequestering:

- 1) What it is not: soln of CC problem (yet!)
- 2) What it is: a way around Weinberg no-go
- 3) Mechanism
 - a) Cancelling vacuum energy loops
 - b) Keeping nonzero mass gap in matter QFT
 - c) Symmetries & protection mechanism
- 4) Features & predictions
- 5) Doomsday cosmology
- 6) Questions & future directions

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Cosmological Constant problem

You hear:

- 1) "Why is it not M_{pl}^4 ?"
- 2) "Why is it not 0?"
- 3) "Why is it $10^{-120} M_{pl}^4$?"

THIS IS ALL NONSENSE !!!

We do not yet have a fully consistent UV complete algorithm for computing vacuum energy; maybe we do not even have a unique vacuum

The best we CAN do: compute in QFT coupled to gravity + see what happens

SO DO IT, BUT DO IT "RIGHT": IE USE THE RULES OF QFT TO THE LETTER

(AND, SPIRIT...!)

Cosmological constant - aka vacuum energy diverges:

$$\Lambda_{\text{vac}} = \bigcirc + \bigcirc \text{ with diagonal line} + \bigcirc \text{ with two diagonal lines} + \dots \quad \text{matter loops}$$
$$+ \bigcirc \text{ with jagged line} + \bigcirc \text{ with jagged line and diagonal line} + \dots \quad \text{gravity + matter loops}$$

Must be renormalized! Finite part NOT calculable but MUST be measured!

But: if determined at N-loops level and fitted to observations (like Higgs mass!) it needs to be completely refitted - from scratch - at N+1 loops

JUST LIKE THE HIGGS MASS ...

MEASURED CC IS BADLY RADIATIVELY
UNSTABLE!

This is the ONLY meaningful formulation
of CC problem in QFT of matter + gravity

The problem is very serious due to the
famous Weinberg no-go theorem! (1989)

(Alternative philosophy: forget naturalness &
use anthropic arguments - but I will ignore
this here - we all choose our own poison)

Weinberg no-go: prohibits dynamical adjustment
in EFT + gravity @ a nonzero mass gap

Adjustment + radiative stability \Leftrightarrow Conformal symmetry

$$\Lambda_{\text{eff}} = \left(\Lambda_{\text{vac}} + \Lambda_{\text{classical}} \right) e^{4\phi}$$

either of the two options is bad:

1) $\Lambda_{\text{vac}} + \Lambda_{\text{classical}} \ll (\text{cutoff})^4$: TUNING

2) $\phi \rightarrow -\infty$: "ADJUSTMENT" ?

But then, since $m_{\text{phys}} = m e^{\phi}$,

$$m_{\text{phys}} \rightarrow 0!$$

NOT OUR WORLD!

SEQUESTERING

$$S = \int d^4x \sqrt{g} \left\{ \frac{M_{pl}^2}{2} R - \Lambda - \lambda^4 \mathcal{L}(g^{mn}, \phi) \right\} + \sigma \left(\frac{\Lambda}{\lambda^4 M^4} \right)$$

Postulate: Λ , λ dynamical GLOBAL variables

\therefore LAGRANGE MULTIPLIERS JUST LIKE
IN THE ISOPERIMETRIC PROBLEM

$\int d^4x \sqrt{g} \Lambda$: Legendre transformation: trades $\int d^4x \sqrt{g}$
for new INDEPENDENT VARIABLE Λ

Variational eqs fix Λ to $\frac{1}{4} \langle T_m^{m \text{const}} \rangle_{vac}$

because λ is just the engineering scale

$$\therefore \frac{1}{4} \langle T_m^{m \text{const}} \rangle_{vac} = \Lambda_{vac} = 0 + \textcircled{0} + \textcircled{0} + \dots$$

NO INTERNAL GRAVITON LINES !!!

- ∴ We work @ QFT of matter @ (semi)classical gravity: gravity is a spectator, a probe only
- ∴ OK: this is the original version of the problem recognized by Pauli & by Zeldovich
- ∴ The Weinberg NO GO precludes a dynamical solution even in this restricted setup

NOTE: can define $\bar{g}_{\mu\nu} = \lambda^2 g_{\mu\nu}$ to rewrite the QFT of matter as

$$S_{\text{matter}} = - \int d^4x \sqrt{\bar{g}} \mathcal{L}_{\text{matter}}(\bar{g}^{\mu\nu}, \phi)$$

Since the theory is Poincaré-inv & diff-inv - as long as we insist that the UV regulator is also included in $\mathcal{L}_{\text{matter}}$ - the form of the matter action is PRESERVED by all loop corrections

So:

$$\Lambda_{\text{vac}} = \bigcirc + \bigcirc + \bigcirc + \dots = \lambda^4 \left(\begin{array}{l} \text{dimensionful} \\ \text{factor} \end{array} \right)$$

$$\text{Thus: } \Lambda_{\text{vac}} = \langle 0 | \mathcal{L}_{\text{matter}} | 0 \rangle = \frac{1}{4} \langle 0 | T^{\mu}_{\mu} | 0 \rangle$$

Variational Eqs

$$M_{Pl}^2 G^{\mu}_{\nu} = T^{\mu}_{\nu} - \Lambda \delta^{\mu}_{\nu}$$

$$\frac{\sigma'}{\lambda^4 \mu^4} = \int d^4x \sqrt{g} \quad \Lambda \frac{\sigma'}{\lambda^4 \mu^4} = \frac{1}{4} \int d^4x T^{\mu}_{\mu}$$

In the vacuum, $T^{\mu}_{\mu} \rightarrow \langle 0 | T^{\mu}_{\mu} | 0 \rangle$ regularized; & elimination of σ' gives us

$$\Lambda = \frac{1}{4} \frac{\int d^4x \sqrt{g} \langle 0 | T^{\mu}_{\mu} | 0 \rangle}{\int d^4x \sqrt{g}} = \frac{1}{4} \langle T^{\mu}_{\mu} \rangle$$

$\therefore \Lambda$ is SPACETIME average of the T^{μ}_{μ} - integrated over all times, past & future! (GLOBAL CONSTRAINT!)

$$M_{Pl}^2 G^{\mu}_{\nu} = T^{\mu}_{\nu} - \frac{1}{4} \langle T^{\alpha}_{\alpha} \rangle \delta^{\mu}_{\nu}$$

Define now $T^{\mu}_{\nu} = -\Lambda_{vac} \delta^{\mu}_{\nu} + \underbrace{T^{\mu}_{\nu}}_{\text{NONCONST PART}}$

Since $\langle \Lambda_{vac} \rangle \equiv \Lambda_{vac} \equiv -\frac{1}{4} T^{\alpha}_{\alpha} (\text{const})$:

$$M_{pl}^2 G^{\mu\nu} = T^{\mu\nu} - \frac{1}{4} \langle T^{\alpha}_{\alpha} \rangle \delta^{\mu\nu}$$

Λ_{vac} completely cancelled from the source of the curvature irrespective of the loop order

IN THIS THEORY VACUUM ENERGY INVISIBLE TO $G^{\mu\nu}$
 \therefore Trick: global variable Λ is tied to Λ_{vac} and always exactly cancels it; the residual finite part is $-\frac{1}{4} \langle T^{\alpha}_{\alpha} \rangle$ - historic average of energy density of all NONCONSTANT sources (γ, B, DM, \dots)

Note: canonically normalizing all matter fields shows that physical dimensional parameters are

$$\frac{M_{phys}}{M_{pl}} = \lambda \frac{M_{bare}}{M_{pl}}$$

So $m_{phys} \neq 0 \iff \lambda \neq 0$; but by

$$\lambda^4 \mu^4 = \frac{\sigma'}{\int d^4x \sqrt{g}} \implies \int d^4x \sqrt{g} < \infty \quad \text{FINITE! UNIVERSE!}$$

Recapitulate:

- 1) Vacuum energy completely cancelled!
- 2) Residual "c.c." determined by the historic average of all nonconstant matter

$$\Lambda_{\text{eff}} = \frac{1}{4} \langle T_{\alpha}^{\alpha} \rangle$$

This is NONLOCAL IN TIME! —but: NO PATHOLOGIES in (semi)classical gravity

Recall QFT: if Q is divergent, it CANNOT be predicted: one regulates it (Λ_{vac}), picks the counterterm (Λ), cancels the divergent part out MEASURES the finite remainder!

Λ : codimension-0 parameter — need cod. 0 detector to measure

THE WHOLE UNIVERSE IS THE ONLY DETECTOR

3) Need $\int d^4x \sqrt{g} < \infty$ to keep nonren QFT mass gap

UNIVERSE LARGE, BUT FINITE IN SPACE-TIME!

- predictions:
- 1) will collapse
 - 2) finite spatial sections (CMB?)
 - 3) cannot accelerate forever!

Symmetries & Naturalness

Whenever \exists magical cancellation look for symmetries

There are 2 approx symmetries

1) approx scaling symms

$$\lambda \rightarrow \Omega \lambda$$
$$g_{\mu\nu} \rightarrow g_{\mu\nu} / \Omega^2$$
$$\Lambda \rightarrow \Lambda \Omega^4$$

2) approx shift symmetry

$$\Lambda \rightarrow \Lambda + \alpha \lambda^4$$
$$\mathcal{L} \rightarrow \mathcal{L} - \alpha$$

Weakly broken by unit: $\delta S \sim \alpha \lambda^4 \int d^4x \sqrt{g} \sim \alpha \left(\frac{m_{\text{pl}}}{M_{\text{pl}}}\right)^4$

- \therefore Scaling ensures cancellation happens at all scales
- \therefore Shift performs the cancellation
- \therefore Since symmetries are approx, $\Lambda_{\text{eff}} \neq 0 \ll M_{\text{pl}}^4$
- \therefore It is naturally small since in the conformal limit $m_{\text{pl}} \rightarrow 0$ symmetries enhanced!

PROTECTION MECHANISM!

Further properties & consequences

- 1) Consistent @ inflation - can make a big old universe @ slowly rolling ϕ
- 2) Phase transition contributions to Λ_{eff} automatically small! (Reason: PT takes only a fraction of universes lifetime and its correction to Λ_{eff} is weighted by $\int d^4x \sqrt{g}$!)
- 3) $\Lambda_{\text{eff}} = \frac{1}{4} \langle T^\alpha_\alpha \rangle \lesssim M_{\text{pl}}^2 H_*^2$ where H_* is the Hubble at the onset of collapse: $H_* \lesssim H_{\text{now}}$

$$\Lambda_{\text{eff}} \leq 10^{-12} \text{ eV}^4$$

as long as we live in a big old universe (we do!)

- 4) Offers a new perspective on the cosmic coincidence "problem": can we explain cosmic acceleration NOW?

Yes!

Cosmology:

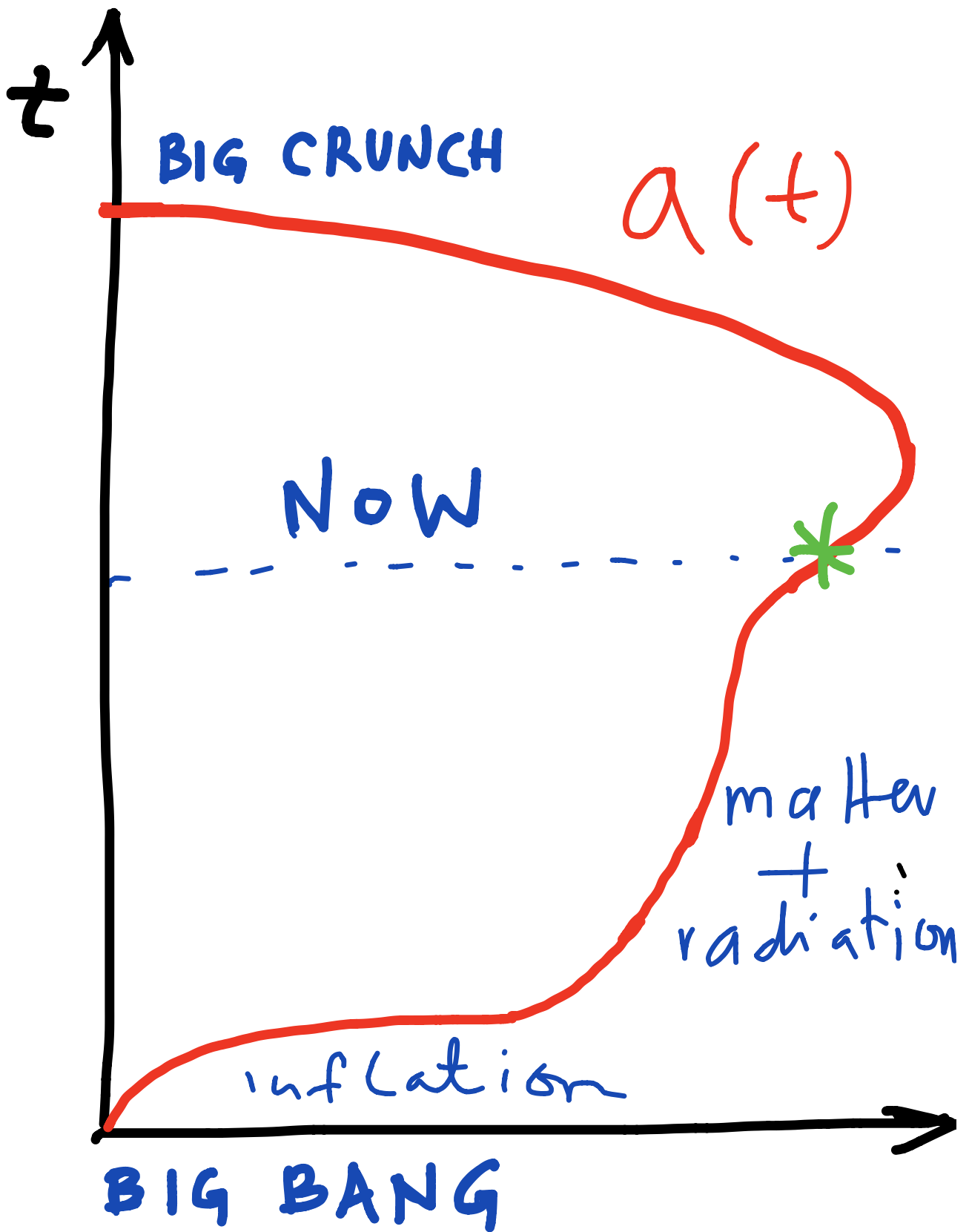
- 1) Need a FINITE universe s.t. $\int \sqrt{g} < \infty$
to ensure mass gap $M_{\text{phys}} \neq 0$
- 2) This universe needs to get large
and old - since it is large and old
- 3) This can be done @ inflation - NONETERNAL
- 4) then some time in the future collapse
can occur

Getting the collapse to happen at the
right time is tuned, to be sure

But: the choice of the time of cosmic doom is

- 1) Radiatively stable
- 2) Predicts epoch of acceleration before collapse

This can explain "why now" problem!



To condemn the universe one needs $\rho < 0$

$$3M_p^2 \left(H^2 + \frac{1}{a^2} \right) = \rho_{\text{usual}} + \rho_{\text{trigger}}$$

Having $\rho_{\text{trigger}} < 0$ and dominating can flip $H > 0 \rightarrow H < 0$: eg an unstable V

A generic potential is radiatively unstable which can reintroduce loop sensitivity

Linde '87: LINEAR POTENTIAL: $V = m_*^3 \varphi$

- 1) form & slope are radiatively stable
- 2) $\varphi \rightarrow \varphi + c$ & $\Lambda \rightarrow \Lambda - m_*^3 c$ shift symmetry \rightarrow loop corrections DO NOT move the initial value of φ

The point is that field eqs are

$$M_{pl}^2 G^{\mu}_{\nu} = T^{\mu}_{\nu} - \frac{1}{4} \langle T \rangle \delta^{\mu}_{\nu}$$

so the shift of $\varphi \equiv$ shift of $\Lambda \rightarrow$ it cancels out!!!

$t \nearrow \varphi \nearrow V \searrow$: when $V < 0$
 sufficiently, collapse occurs $\rightarrow H = 0$

If m_{ϕ} is big this will happen too soon

So we need small m_{ϕ}

Hence for most of the expansion

- 1) φ & $V(\varphi)$ are negligible
- 2) φ is in slow roll \rightarrow changes little

Near the onset of collapse - by definition - φ dominates; before this time $\Delta\varphi$ is

$$\frac{\varphi_t - \varphi_{in}}{M_{pl}} = \int_{t_{in}}^{t_t} dt \frac{\dot{\varphi}}{M_{pl}} \simeq -O(1) \frac{m_{\text{slope}}^3 M_{pl}}{M_{pl}^2 H_t^2}$$

$$\therefore \Delta\varphi < M_{pl}$$

since $m_{\text{slope}}^3 = M_{pl} H_0^2$ @ $H_0 < H_t$ and $H_0 \simeq$ Hubble now

Since $\left(\frac{\dot{\varphi}^2}{m_{\text{slope}}^2 \varphi}\right)_+ = O(1) \frac{M_{\text{slope}}^6}{M_{\text{PL}} H_+^4} < 1$
 when this happens the universe is dominated
 by a slowly rolling scalar field: **DE!**

This goes on until the end of slow roll
 $\varphi \leq M_{\text{PL}}$, which happens at

$$t_{\text{ase}} \sim \frac{1}{H_{\text{ase}}} \approx \frac{1}{H_+} + O(1) \sqrt{\frac{M_{\text{PL}}}{M_{\text{slope}}^3}} \approx \frac{1}{H_0}$$

$$\therefore \# \text{ of e-folds } \mathcal{N} \approx \frac{H_+}{H_0} !$$

More importantly:

$$\frac{\varphi_{\text{axl}} - \varphi_{\text{in}}}{M_{\text{PL}}} = -O(1) \frac{M_{\text{slope}}^3 M_{\text{PL}}}{M_{\text{PL}}^2 H_{\text{axl}}^2} \approx -O(1) \frac{H_0^2}{H_{\text{axl}}^2}$$

$$\therefore \varphi_{\text{axl}} - \varphi_{\text{in}} > M_{\text{PL}} !$$

In GR φ_{in} is arbitrary, and so when it is small for late stage of acceleration will not last long enough

Initial condition tuning for quintessence models

Not so here! At the onset of collapse

$$\langle \rho_\varphi \rangle = \langle \dot{\varphi}^2 \rangle - 4m_{\text{slow}}^3 \langle \varphi \rangle \simeq -O(1) m_{\text{slow}}^3 \varphi_{\text{age}}$$

AND

$$\langle T \rangle = \langle 3p - \rho \rangle = -O(1) M_{\text{pl}}^2 H_{\text{age}}^2$$

$$\therefore \frac{\varphi_{\text{age}}}{M_{\text{pl}}} \simeq -O(1) \frac{M_{\text{pl}} H_{\text{age}}^2}{m_{\text{slow}}^3} \simeq -O(1) \frac{H_{\text{age}}^2}{H_0^2} \lesssim -1$$

Therefore

$$\varphi_{in} \simeq \left[O(1) \frac{H_0^2}{H_{\text{age}}^2} + O(1) \frac{H_{\text{age}}^2}{H_0^2} \right] M_{\text{pl}} \gtrsim M_{\text{pl}}!$$

Consistency requires $\rho_{in} \gtrsim M_{pl} \rightarrow$ tuning
the age of the universe automatically yields
a prediction that prior to collapse there
is a stage of accelerated expansion for
an e-fold or few

But ... there is more ...

field eqs require that $\langle R \rangle = 0$

$$FRW: \int dt a^3 \left(\dot{H} + 2H^2 + \frac{\kappa}{a^2} \right) = 0$$

Integrate by parts:

$$a^3 H \Big|_{t_{bang}}^{t_{crunch}} = \int dt a^3 \left(H^2 - \frac{\kappa}{a^2} \right)$$

In big old universes LHS is small, $\ll 0$

So there must be a cancellation on the
RHS as well! $\rightarrow \kappa = +1!$

So: we have 2 predictions for cosmology:

1) $w \neq -1$ DE: Linear potential

$$w = -1 + O(1) \frac{H_0^4}{H_t^4}$$

2) $k=+1 \rightarrow \Omega_k \text{ now} < 0 \rightarrow \Omega_k \Big|_{\text{age}} \simeq -1$

But there is only ONE dimensional parameter!

$$w = -1 + O(1) [\Omega_k(\text{now})]^2$$

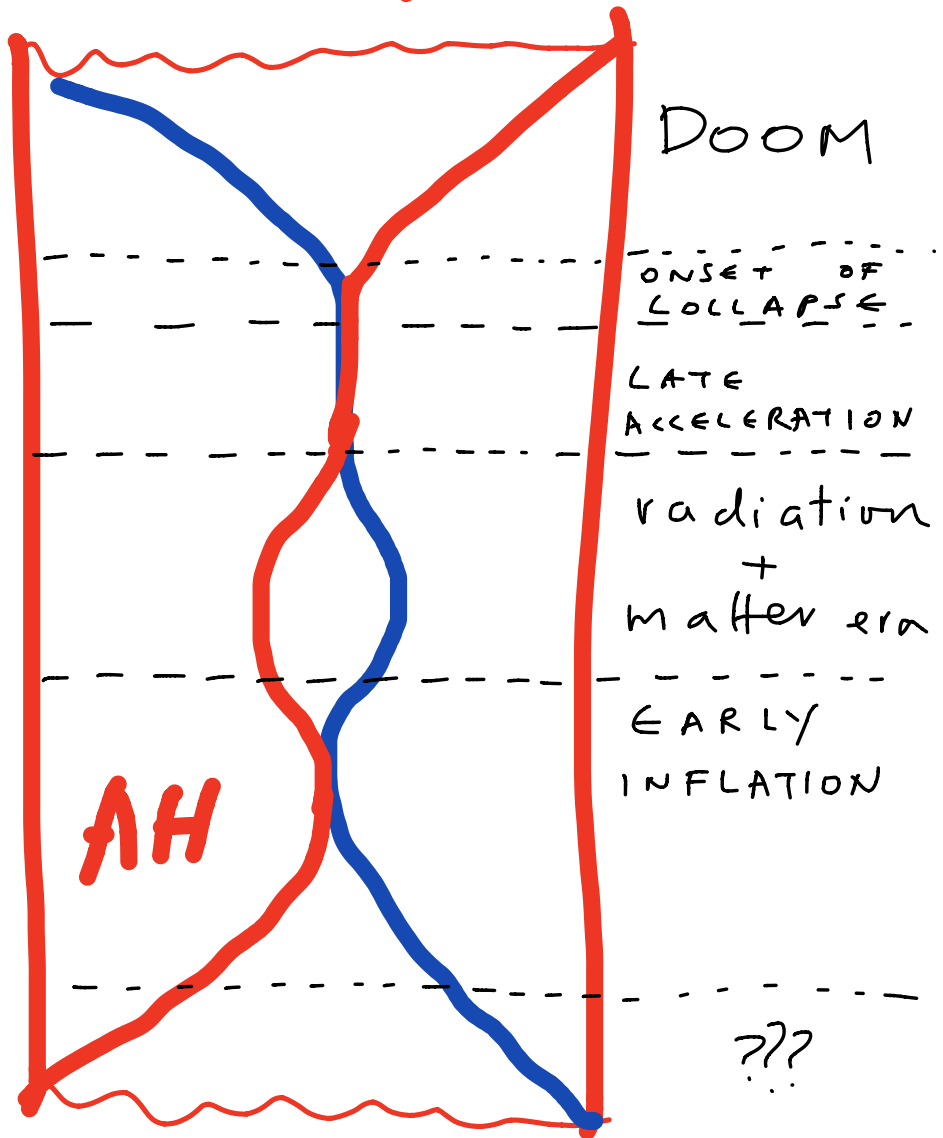
This is a strong prediction of the setup
could be tested observationally

Note: this was sloppy; a more precise estimate needed

Work in progress, @ A. Moss, A. Padilla, J. Pearson

The portrait of the world

I_+



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SUMMARY

∴ FOUND A WAY AROUND WEINBERG NO GO

- 1) Cancels all loop corrections to Λ
- 2) Maintains $M_{\text{pl}} \neq 0$
- 3) Radiatively stable - aka technically natural
- 4) Maximally minimal modification of GR+QFT
∴ Poincaré & diff. invariant

∴ No new local DOFs

∴ Cosmology differs: $w \neq -1$, $\Omega_K < 0$, universe finite, will collapse

- 5) phase transitions tamed
- 6) Inflation (finite!) OK!
- 7) Cosmic coincidence → why accelerating now!

- a) Microscopic origin & UV completion?
- b) Protection from gravity loops?
- c) Uniqueness? Or not? Further predictions?