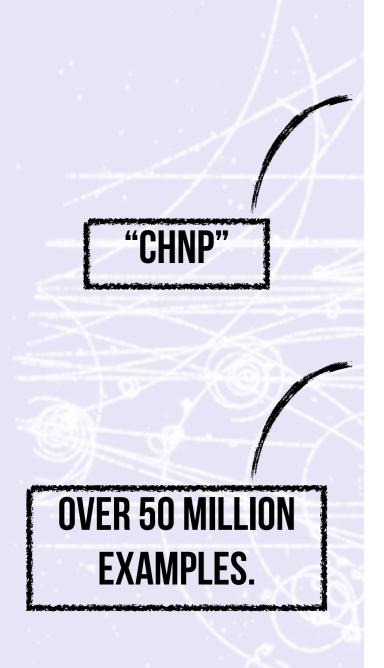
## Gauge Enhancement and Landscapes in G<sub>2</sub> Compactifications of M-theory

### JIM HALVERSON KAVLI INSTITUTE FOR THEORETICAL PHYSICS

based on works with Dave Morrison (1412.4123 and 1507.xxxxx appear.)

### A REALLY GREAT G2 MATH PAPER



#### G<sub>2</sub>-MANIFOLDS AND ASSOCIATIVE SUBMANIFOLDS VIA SEMI-FANO 3-FOLDS

ALESSIO CORTI, MARK HASKINS, JOHANNES NORDSTRÖM, AND TOMMASO PACINI

Asstract. We provide a significant extension of the twisted connected sum construction of  $G_2$ —manifolds, is Riemannian 7-manifolds with holonomy group  $G_2$ , first developed by Kovalev; along the way we address some foundational questions at the heart of the twisted connected sum construction. Our extension allows us to prove many new results about compact  $G_2$ —manifolds and leads to new perspectives for future research in the area. Some of the main contributions of the paper are:

- We correct, clarify and extend several aspects of the K3 "matching problem" that occurs as a key step in the twisted connected sum construction.
- (ii) We show that the large class of asymptotically cylindrical Calabi-Yau 3-folds built from semi-Fano 3-folds (a subclass of weak Fano 3-folds) can be used as components in the twisted connected sum construction.
- (iii) We construct many new topological types of compact G<sub>2</sub>-manifolds by applying the twisted connected sum to asymptotically Calabi-Yau 3-folds of semi-Fano type studied in [18].
- (iv) We obtain much more precise topological information about twisted connected sum G<sub>2</sub>-manifolds; one application is the determination for the first time of the diffeomorphism type of many compact G<sub>2</sub>-manifolds.
- (v) We describe "geometric transitions" between G<sub>2</sub>-metrics on different 7-manifolds mimicking "flopping" behaviour among semi-Fano 3-folds and "conifold transitions" between Fano and semi-Fano 3-folds.
- (vi) We construct many G2-manifolds that contain rigid compact associative 3-folds.
- (vii) We prove that many smooth 2-connected 7-manifolds can be realised as twisted connected sums in numerous ways; by varying the building blocks matched we can vary the number of rigid associative 3-folds constructed therein.

The latter result leads to speculation that the moduli space of G<sub>2</sub>-metrics on a given 7manifold may consist of many different connected components, and opens up many further questions for future study. For instance, the higher-dimensional gauge theory invariants proposed by Donaldson may provide ways to detect G<sub>2</sub>-metrics on a given 7-manifold that are not deformation equivalent.

#### 1. Introduction

In this paper we construct a large number of new compact  $G_2$ -manifolds, that is Riemannian 7-manifolds (M, g) whose holonomy group is the compact exceptional Lie group  $G_2$ , using the so-called twisted connected sum construction; since any  $G_2$ -manifold is Ricci-flat this yields many Ricci-flat metrics on compact 7-manifolds. As an alternative to Joyce's original pioneering construction of compact  $G_2$ -manifolds via "orbifold resolutions" [41,42], Kovalev (based on a suggestion of Donaldson) developed the twisted connected sum construction [44] as a way to obtain a compact  $G_2$ -manifold by combining a pair of (exponentially) asymptotically cylindrical (ACyl) Calabi-Yau 3-folds. Loosely speaking, this method seeks to construct  $G_2$ -manifolds

Key words and phrases. Differential geometry, Einstein and Ricci-flat manifolds, special and exceptional holonomy, noncompact Calabi-Yau manifolds, compact G<sub>2</sub>-manifolds, Fano and weak Fano varieties, lattice polarised K3 surfaces, calibrated submanifolds, associative submanifolds, differential topology. 1207.4470

KOVALEV'S TWISTED CONNECTED SUMS (TCS)

CONSTRUCTS
SOME ASSOCIATIVE
SUBMANIFOLDS.

I cannot recommend this paper enough. Pedagogical and readable!

**MATH LITERATURE** 

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   [Aganagic, Vafa], [Gukov, Sparks], [Beasley, Witten]

and many many others! (please let me know if I acc. didn't include you, I'm still learning this literature!)

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see [Acharya, Gukov] review and references therein!

## THIS TALK

neither strings nor pheno, unfortunately!

Still a lot to learn about TCS G2 before thinking pheno.

Instead, present and discuss physics of TCS G<sub>2</sub> and develop a language / set of techniques for **global** models, specifically as it regards gauge enhancement.

#### • G<sub>2</sub> Review

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#### Gauge Enhancement

Need singularities. Can't use standard CY trick. What to do? Example: Coulomb branches, monopoles, instantons, etc.

# G2 REVIEW

both math and physics

My apologies that this has to be so fast. This is a lightning review.

Two slides on math.

One slide on physics, since you've heard it.

## G2 Manifolds

- $G_2$  is an exceptional Lie group with rk=2 and dim=14.
- A  $G_2$  manifold X is a 7-manifold with hol( $G_2$ ). No CAG. No Yau's thm.
- Characterized by a  $G_2$  structure that determines a metric  $g_{\Phi}$  and  $G_2$  three-form  $\Phi$ . These give a  $G_2$  action on every tangent space.
- Note well: only one type of metric modulus, from  $\Phi$ .
- The following are equivalent:  $Hol(g_{\Phi}) \subseteq G_2$

$$\Phi$$
 TORSION FREE  $abla\Phi=0, ext{ and }$   $d\Phi=d\star\Phi=0.$ 

• Holonomy is exactly  $G_2$  iff above and X has finite fundamental group.

## G2 Calibrations

 G<sub>2</sub> manifolds can have calibrated submanifolds, associative 3-mflds, coassociative 4-mflds.

[HARVEY, LAWSON]

- They are volume minimizing reps of their hom. classes.
- Can compute their volumes, even though we don't know the metric.

$$Vol(M_3) = \int_{M_3} \Phi \qquad Vol(M_4) = \int_{M_4} \star \Phi$$

 CHNP: first construction theorem for associatives in compact G<sub>2</sub> manifolds.

(rough: rigid one to each rigid holomorphic curve in build. block non-rigid for each appropriate sLag in building block).

## 11D SUGRA + M-branes on G2

- Key point: gives d=4 N=1 theory.
- KK reduction: reduce metric mod  $\Phi$  as well as  $C_3$  and  $C_6$ .

Summary:  $b_2(X)$  abelian vector multiplets  $b_3(X)$  massless uncharged chirals

### [PAPADOPOULOS, TOWNSEND]

- M2 (M5) on two (five) -cycles give charged particles (monopoles).
- Also from wrapped M-branes: instantons, domain walls, strings, etc.
- Key: at most abelian gauge symmetry!
   so we need singularities for non-abelian theories (NAGS)
   or massless charged matter (MCM)
   Smoothing is the Higgs mechanism. How do we reliably un-Higgs?
- Local work: NAGS, vector matter, chiral matter from codim 4, 6, 7.

## **A LANDSCAPE**

from twisted connected sum G<sub>2</sub> manifolds.

Classical, on smooth  $G_2$  at large volume.

There's been huge math progress recently, and there are now over 50 million smooth, compact  $G_2$  manifolds.

What does M-theory on them yield?
What can we learn about this landscape?
What is still to be done?

## Twisted Connected Sums

#### HOW DO WE USE THE TCS CONSTRUCTION TO GET A G2 MANIFOLD?

[KOVALEV]

- Idea: "appropriately" glue two "appropriate" building blocks  $M_{+-}$  X inherits  $G_2$  structure from blocks. Show torsion free.
- $M_{\pm}=V_{\pm} imes S_{\pm}^{1}$  where  $V_{\pm}$  is an ACyl Calabi-Yau threefold.
- ullet i.e.  $V_+$  asymptotes to "Calabi-Yau cylinder"  $\mathbb{C}^* imes S$  , S a K3.

$$z = e^{t+i\theta}$$
 on  $\mathbb{C}^*$ 

• This gives a natural  $G_2$  structure at each asymptotic end:

$$\Phi = d\varphi \wedge dt \wedge d\theta + d\varphi \wedge \omega_S + d\theta \wedge \operatorname{Re}(\Omega_S) + dt \wedge \operatorname{Im}(\Omega_S)$$

- Change CS of S to get K3  $\Sigma$  such that  $\omega_{\Sigma} = \operatorname{Re}(\Omega_{S}), \operatorname{Re}(\Omega_{\Sigma}) = \omega_{S}$   $\operatorname{Im}(\Omega_{\Sigma}) = -\operatorname{Im}(\Omega_{S})$
- Key: then  $(\varphi, t, \theta, S)$  to  $(\theta, -t, \varphi, \Sigma)$  leaves  $\Phi$  invariant!

## Kovalev's Theorem

#### (LOVALEV) (DUH!)

- This is the key observation for the gluing. It means we can glue  $G_2$  structures on  $M_{+}$  to get  $G_2$  structure on X.
- K3 diffeomorphism r satisfying:

$$r^*(\omega_{S_-}) = \operatorname{Re}(\Omega_{S_+})$$
 $r^*(\operatorname{Re}(\Omega_{S_-})) = \omega_{S_+}$ 
 $r^*(\operatorname{Im}(\Omega_{S_-})) = -\operatorname{Im}(\Omega_{S_+})$ 

• Gluing map:

$$F: \qquad M_{+} \cong S_{-}^{1} \times \mathbb{R}^{+} \times S_{+}^{1} \times S_{+} \qquad \longrightarrow \qquad S_{+}^{1} \times \mathbb{R}^{+} \times S_{-}^{1} \times S_{-} \cong M_{-},$$
$$(\theta_{-}, t, \theta_{+}, x) \qquad \longmapsto \qquad (\theta_{+}, T + 1 - t, \theta_{-}, r(x))$$

Kovalev's Theorem: (rough: see paper for more)

Given such  $M_{\pm}$  and r, F, can glue to get TCS seven-manifold X. X has a natural  $G_2$  form related to  $G_2$  forms of  $M_{\pm}$  and it has a torsion free def within its cohomology class. With that metric, X is  $G_2$  mfld.

• People (e.g. CHNP) have gotten really good at constructing ( $M_{\pm}$ , r, F) to the tune of **50** million examples. Their progress:  $V = Z \setminus S$ , Z weak Fano.

## TCS Topology ICHNPI

- Let's be specific about the topology since it matters for physics.
- Typically get  $V_{\pm}$  from other alg. threefolds  $Z_{\pm}$  as  $V_{\pm} = Z_{\pm} \setminus S_{\pm}$
- There are natural restriction maps with kernel K<sub>±</sub> and image N<sub>±</sub>.  $\rho_{\pm}: H^2(V_{\pm}, \mathbb{Z}) \longrightarrow H^2(S_{\pm}, \mathbb{Z})$
- Then for the second and third cohomology we have:

$$H^{2}(X,\mathbb{Z}) = (N_{+} \cap N_{-}) \oplus K_{+} \oplus K_{-}$$
$$H^{3}(X,\mathbb{Z}) \supset H^{3}(Z_{+},\mathbb{Z}) \oplus H^{3}(Z_{-},\mathbb{Z}) \oplus K_{+} \oplus K_{-}$$

- Note: some two-forms and three-forms "come together."
- Thm: if C is rig hol. curve in V then def CxS<sup>1</sup> is rig associative in X.

# Broad Assessment of TCS G<sub>2</sub> Landscape [J.H., MORRISON]

• If M on TCS G2 X is vacuum of broken NAGT, nearly all Higgs branches.

(semi-Fano building blocks of CHNP large numbers have K=0, together with simplest gluing gives  $b_2(X) = 0$ ).

MATH ARTIFACT OF GLUING DIFFICULTY.

- Fluxes: never have to turn them on in  $G_2$ . [BEASLEY, WITTEN]
- Instantons: due to CHNP rigid associated theorem, for the first time we have M2-instantons with no deformation modes. Must be concerned about Wilson line modulini!
- Some G<sub>2</sub> transitions can be understood in terms of top trans in V<sub>±</sub>.

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- Particle charges: intersections in X via intersections in 6-mfld.

	$Q_1$	$Q_2$	$Q_3$
$\Psi_{12}^k$	-2	-1	-1
$\Psi^k_{13}$	-1	-2	-1
$\Psi^k_{14}$	-1	-1	-2
$\Psi_{23}^k$	1	-1	0
$\Psi^k_{24}$	1	0	-1
$\Psi^k_{34}$	0	1	-1

4 (vector-pairs of chiral multiplets) of each type.

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#### HARVEY, MOORE

• One V has 24 rigid holomorphic curves, giving 24 assoc.

$$W \supset 4(A_1e^{-\Phi_1} + A_2e^{-\Phi_2} + A_3e^{-\Phi_3} + A_4e^{\Phi_1-\Phi_2} + A_5e^{\Phi_1-\Phi_3} + A_6e^{\Phi_2-\Phi_3})$$

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# Circle of Conifolds Transition [J.H., MORRISON]

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   Collapse S<sup>2</sup> by collapsing associative. Circle of conifolds
- Deform. Since S<sup>3</sup> is sLag in V, it's associative in X. (use other CHNP thm)
- But we know what we're Higgsing! Field theory predicts def top.

## GAUGE ENHANCEMENT

how do you get it, globally!?

Families of metrics in non-compact examples are great, but we don't have that luxury in compact examples and want a language / approach appropriate for compact examples.

Related: in the last example we ran into an important problem but didn't really discuss it. Let's do that now.

### The Problem

#### WHY DOES GAUGE ENHANCEMENT WORK FOR CY (F-TH)? WHAT IS G2 OBSTRUCTION?

- CY singularities are well understood. (defining equations, CAG, etc)
- More specifically: know how certain singularities relate to families of two-cycles going to zero volume via variation in Kahler moduli.

$$vol(C) = \int_C J \mapsto 0$$

Dimensionality of family M determines spacetime quantum numbers.

#### [WITTEN] [ASPINWALL, KATZ, MORRISON]

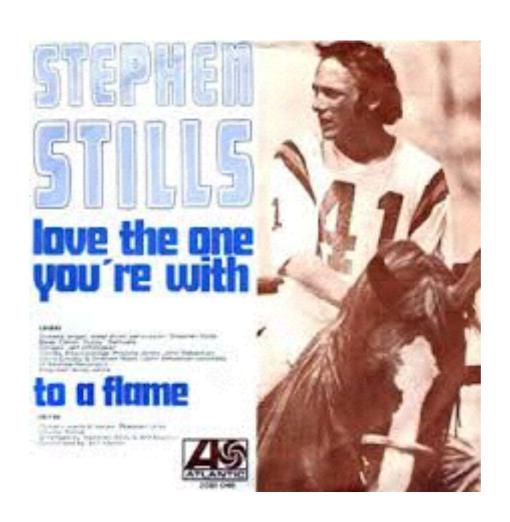
- So varying J and studying dim(M), we can identify limits that give massless charged matter and / or massless charged W-bosons.
- Problem: G<sub>2</sub> has no calibrated two-cycles! What to do!?

# A Proposal [J.H., N WHEN ROCK AND ROLL FITS YOUR TALK, YOU'VE GOT TO USE IT

#### [J.H., MORRISON]



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**HOLOMORPHIC CURVES** 

ASSOC AND COASSOC SUBMANIFOLDS

## A Proposal

#### [J.H., MORRISON]

#### USE THE CALIBRATIONS YOU HAVE! 3- AND 4-MFLDS.

- Recall: associative 3-mflds and coassociative 4-mflds are calibrated, so we can control their volumes as a function of moduli.
- Idea: define cones of assoc. and coassoc. analogous to Kahler cone.
   Get singularities by collapsing associatives or coassociatives.
- Math trick) assoc or coassoc collapse via collapsing two-cycle in them.
   particle masses to zero.
- Physics option) Use other signatures of symmetry breaking. (e.g. defects)

$$M_w = g v$$
  $M_m = v / g$   $T_{ANO} = 2\pi v^2$ 

Some defects arise from calibrated cycles. (e.g. strings, inst, dom walls)
 Some not from not-calibrated cycles (e.g. monopoles, but they're still useful).

## WHAT DO WE WANT?

Ideally, a natural 3 or 4-cycle associated with symmetry

• Consider any non-trivial class  $[\sigma] \in H^2(X,\mathbb{R})$  then:

$$[\sigma] \cup [\sigma] \cup [\Phi] < 0$$

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- NO! That's the limit of infinite gauge coupling since.

$$Vol(D_{\Sigma}) = \int_{D_{\Sigma}} \Phi = -\int \sigma \wedge \sigma \wedge \Phi = \int \sigma \wedge *\sigma \sim \frac{1}{g^2}$$

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$$Vol(D_{\Sigma}) = \int_{D_{\Sigma}} \Phi = -\int \sigma \wedge \sigma \wedge \Phi = \int \sigma \wedge *\sigma \sim \frac{1}{g^2}$$

• Upshot: finite g, and if  $D_{\Sigma}$  exists, a place to wrap gauge instantons!

DEDUCE ADDITIONAL FEATURES OF X FROM PHYSICS, THEN SEE WHAT TO CALIBRATE.

• Break G to  $H=U(1)^{rk(G)}$ , for G a simple Lie group.

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- Furthermore: though for v non-zero **ρ** not a modulus the zero size instanton still solves EOM. (e.g. use "constrained instantons" of Affleck)

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$$[\Sigma] \in H_2(X) \quad [\tilde{\Sigma}] \in H_5(X)$$

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• More is required by breaking from G. Need W-bosons, monopoles, instantons from wrapped branes. Requires submanifolds of these classes, call them  $\tilde{\Sigma}$   $\Sigma$   $D_{\Sigma}$ 

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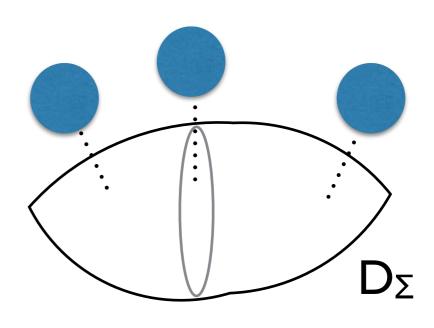
but by the Joyce lemma we also have  $[D_{\Sigma}] \equiv -[\tilde{\Sigma}] \cap [\tilde{\Sigma}] \in H_3(X, \mathbb{Z})$ 

- Instanton behavior follows from  $vol(D_{\Sigma}) \sim \frac{1}{a^2}$

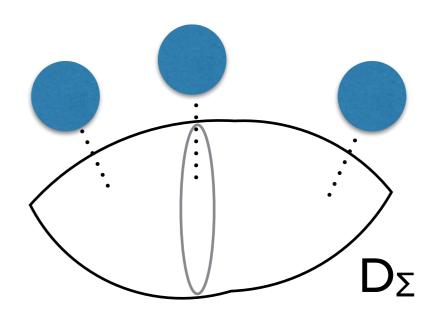
- Geometrically, we have three volumes, of  $\Sigma$   $\tilde{\Sigma}$   $D_{\Sigma}$  and two physical parameters, g and v (f moduli) Overconstrained system?
- More specifically:  $M_W \propto g_{YM} |v| \propto vol(\Sigma)$  and  $M_M \propto \frac{|v|}{g_{YM}} \propto vol(\tilde{\Sigma})$
- ullet The gauge coupling is computed by the volume of  $D_{\Sigma}$  so

$$M_M \propto \frac{|v|}{g_{YM}} \propto \frac{M_W}{g_{YM}^2} \propto vol(\tilde{\Sigma}) \propto vol(\Sigma)vol(D_{\Sigma})$$

ullet Volume relation suggests  $ilde{\Sigma}$  fibered over  $D_{\Sigma}$  by curves of class  $\Sigma$ 



 Upshot: by rather general global G<sub>2</sub> and Coulomb branch considerations, we've landed on a fibration structure expected from the standard picture of S<sup>2</sup> fibered over a three-mfld.

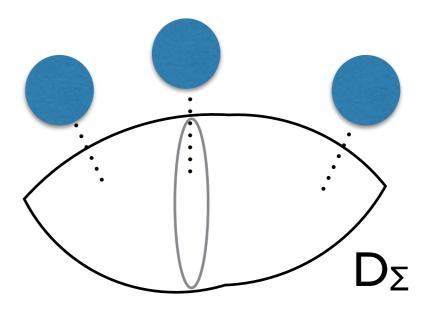


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- What to calibrate to zero, though!?

No adjoint chiral in the story yet.  $D_{\Sigma}$  can have topology. If  $b_1(D_{\Sigma})$  non-zero, we have a two-sphere fibration over each one-cycle.

Assoc. rep? Calibrate that to zero ...

One-cycle in singular limit associated to adj chiral.

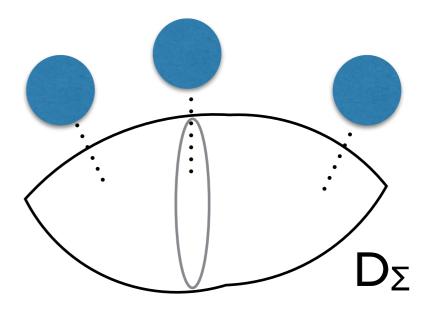


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• Joyce example: SU(2) to U(1) gives  $S^2$  fib over  $T^3$ . 3 adj chiral mult.

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   Math trick (two-spheres collapse)
   Diff. physics (topological defects)
- Coulomb branch: rather general physics arguments and a simple G<sub>2</sub> fact lead to standard fibration picture and three-cycles to collapse.

# Thanks so much to the organizers for a truly great conference!