

Heterotic Superpotentials and Moduli

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This talk is concerned with heterotic supergravity at $\mathcal{O}(\alpha')$, its four-dimensional effective supergravity and moduli.

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- String Compactifications.

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- String Compactifications.
- Gukov-Vafa-Witten superpotential and supersymmetry conditions.

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- First order deformations, holomorphic structures and moduli.

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- String Compactifications.
- Gukov-Vafa-Witten superpotential and supersymmetry conditions.
- First order deformations, holomorphic structures and moduli.
- Conclusions and outlook..

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String theory is ten-dimensional:

$$\mathcal{M}_{10} = \mathcal{M}_4 \times X_{\text{compact}} ,$$

where \mathcal{M}_4 is assumed Minkowski, and X is compact.

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First need to find massless spectrum, i.e. infinitesimal moduli!

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The Low energy theory of the heterotic string is a 10d $N = 1$ supergravity equipped with a $E_8 \times E_8$ gauge field A .

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Complications:

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Complications:

- torsional geometries not well understood, but some progress [Strominger 86, Becker et al 2003, Ivanov 2009, ..].

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- torsional geometries not well understood, but some progress [Strominger 86, Becker et al 2003, Ivanov 2009, ..].
- Complicated expressions to deal with, e.g. Bianchi Identity:

$$dH = -2i\partial\bar{\partial}\omega = \frac{\alpha'}{4}(\text{tr } F^2 - \text{tr } R^2) .$$

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Need a nicer description to deal with moduli [Anderson et al 10, Anderson et al 14, de la Ossa EES 14, Garcia-Fernandez et al 15].

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Four-dimensional heterotic theory has GVW-superpotential [Gukov et al 99, Becker et al 03, Cardoso et al 03, Lukas et al 05, ..]

$$W = \int_X (H + id\omega) \wedge \Omega ,$$

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where ω is the hermitian two-form (Kähler form), Ω is a complex top-form, $\Omega \in \Omega^{(3,0)}(X)$ encoding the complex structure,

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and where

$$\omega_{CS}^A = \text{tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right) .$$

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F-term conditions:

$$\delta W = W = 0 .$$

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- $\Rightarrow d\Omega = 0$ and so X is a *complex manifold*.

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- $\delta_1 \Omega = K\Omega + \chi^{(2,1)} \Rightarrow H = i(\partial - \bar{\partial})\omega$ [Strominger 86].

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Note: Also D-term conditions giving rise to (poly-)stability conditions on bundles [Anderson et al 09].

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Ignore D-terms and conformally balanced condition for this talk, and assume stable bundles.

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At the supersymmetric locus, the four-dimensional mass-matrix reads

$$V_{I\bar{J}} = e^{\mathcal{K}} \partial_I \partial_K W \partial_{\bar{J}} \partial_{\bar{L}} \bar{W} \mathcal{K}^{K\bar{L}} .$$

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$$\delta_2 \delta_1 W = 0 .$$

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Naive assumption:

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Naive assumption:

$$T\mathcal{M} = H^{(0,1)}(T^*X) \oplus H^{(0,1)}(TX) \oplus H^{(0,1)}(\text{End}(V))$$

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$$\begin{aligned} \delta_{12} W|_{\delta W=0} &= \int_X \frac{\alpha'}{2} \left(\text{tr} \delta_1 A \wedge \delta_2 (F \wedge \Omega) - \text{tr} \delta_1 \Theta \wedge \delta_2 (R \wedge \Omega) \right) \\ &+ \int_X d\tau_1 \wedge \delta_2 \Omega + \int_X \delta_2 (H + id\omega) \wedge \delta_1 \Omega \\ &+ \int_X (H + id\omega) \wedge \delta_2 \delta_1 \Omega . \end{aligned}$$

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It follows that

$$d\delta_2\Omega = 0 \Rightarrow \delta_2\Omega \in H^{(2,1)}(X) \Leftrightarrow \Delta_2 \in H^{(0,1)}(TX) ,$$

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$$d\delta_2\Omega = 0 \Rightarrow \delta_2\Omega \in H^{(2,1)}(X) \Leftrightarrow \Delta_2 \in H^{(0,1)}(TX),$$

Also get

$$\delta_2(F \wedge \Omega) = 0 \Leftrightarrow \Delta_2^a \wedge F_{a\bar{b}} dz^{\bar{b}} = \bar{\partial}\alpha_2,$$

where $\Delta_2 \in H^{(0,1)}(TX)$, $\alpha_2 \in \Omega^{(0,1)}(\text{End}(V))$.

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Similarly

$$\delta_2(R \wedge \Omega) = 0 \Leftrightarrow \Delta_2^a \wedge R_{a\bar{b}} dz^{\bar{b}} = \bar{\partial}\kappa_2.$$

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Note: Deformations $\delta_2\nabla = \kappa_2$ non-physical. Can be thought of as infinitesimal field redefinitions preserving Strominger system [de la Ossa EES 14].

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It follows that Δ_2 is in the kernel of [Anderson et al 10]

$$\mathcal{F} : H^{(0,1)}(TX) \rightarrow H^{(0,1)}(\text{End}(V))$$

$$\mathcal{R} : H^{(0,1)}(TX) \rightarrow H^{(0,1)}(\text{End}(TX)) .$$

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$$\mathcal{R} : H^{(0,1)}(TX) \rightarrow H^{(0,1)}(\mathbf{End}(TX)) .$$

Can equivalently be put in terms of holomorphic structure

$$\bar{\partial}_1 = \bar{\partial} + \mathcal{F} + \mathcal{R} ,$$

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It follows that Δ_2 is in the kernel of [Anderson et al 10]

$$\mathcal{F} : H^{(0,1)}(TX) \rightarrow H^{(0,1)}(\text{End}(V))$$

$$\mathcal{R} : H^{(0,1)}(TX) \rightarrow H^{(0,1)}(\text{End}(TX)) .$$

Can equivalently be put in terms of holomorphic structure

$$\bar{\partial}_1 = \bar{\partial} + \mathcal{F} + \mathcal{R} , \quad \text{Binachi Identities} \quad \Leftrightarrow \quad \bar{\partial}_1^2 = 0 .$$

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$\bar{\partial}_1$ defines an Atiyah algebroid [Atiyah 57]

$$0 \rightarrow \text{End}(V) \oplus \text{End}(TX) \rightarrow \mathcal{Q}_1 \rightarrow TX \rightarrow 0 .$$

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$$\begin{aligned} T\mathcal{M}_1 &= H^{(0,1)}(\mathcal{Q}_1) \\ &= H^{(0,1)}(\text{End}(V)) \oplus H^{(0,1)}(\text{End}(TX)) \oplus \ker(\mathcal{F} + \mathcal{R}) . \end{aligned}$$

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We also have the terms

$$\int_X \delta_2(H + id\omega) \wedge \delta_1\Omega + \int_X (H + id\omega) \wedge \delta_2\delta_1\Omega \in \delta_{12}W|_{\delta W=0}$$

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Algebra: \Rightarrow arrive at the following conditions

$$\bar{\partial}\tau_2^{(0,2)} = 0$$

$$\begin{aligned} 2\Delta_2^a \wedge i\partial_{[a}\omega_{b]}\bar{c}dz^{b\bar{c}} - \frac{\alpha'}{2}(\text{tr } \alpha_2 \wedge F - \text{tr } \kappa_2 \wedge R) \\ = \partial\tau_2^{(0,2)} + \bar{\partial}\tau_2^{(1,1)}. \end{aligned}$$

Technicality: Assume $H^{(0,1)}(X) = 0 \Rightarrow \partial\tau_2^{(0,2)}$ is $\bar{\partial}$ -exact.

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It follows that $x = (\Delta, \alpha, \kappa) \in H^{(0,1)}(\mathcal{Q}_1)$ is in the kernel of

$$\mathcal{H} : H^{(0,1)}(\mathcal{Q}_1) \rightarrow H^{(0,2)}(T^*X).$$

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The map \mathcal{H} defines the holomorphic double extension

$$0 \rightarrow T^*X \rightarrow \mathcal{Q}_2 \rightarrow \mathcal{Q}_1 \rightarrow 0 ,$$

with corresponding holomorphic structure

$$\bar{\partial}_2 = \bar{\partial}_1 + \mathcal{H} , \quad \text{Heterotic Bianchi Identity} \quad \Leftrightarrow \quad \bar{\partial}_2^2 = 0 .$$

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Note: \mathcal{Q}_2 as a holomorphic bundle is *self-dual*.

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Infinitesimal moduli [Anderson et al 14, de la Ossa EES 14]

$$T\mathcal{M}_2 = H^{(0,1)}(\mathcal{Q}_2) = H^{(0,1)}(T^*X) \oplus \ker(\mathcal{H}) .$$

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Infinitesimal moduli [Anderson et al 14, de la Ossa EES 14]

$$T\mathcal{M}_2 = H^{(0,1)}(\mathcal{Q}_2) = H^{(0,1)}(T^*X) \oplus \ker(\mathcal{H}) .$$

Get same kernel structure.

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Conclusions:

- Heterotic string is a nice playground for phenomenology, but the moduli problem is hard.

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Conclusions:

- Heterotic string is a nice playground for phenomenology, but the moduli problem is hard.
- From the heterotic superpotential, we derived the massless moduli space, and saw that it agrees with the 10d computation of [Anderson et al 14, de la Ossa EES 14] for the infinitesimal moduli space of solutions to the Strominger system.

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Conclusions:

- Heterotic string is a nice playground for phenomenology, but the moduli problem is hard.
- From the heterotic superpotential, we derived the massless moduli space, and saw that it agrees with the 10d computation of [Anderson et al 14, de la Ossa EES 14] for the infinitesimal moduli space of solutions to the Strominger system.
- We note that the heterotic anomaly condition may lead to lifting extra moduli, even in Calabi-Yau compactifications.

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Outlook, and work in progress:

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Outlook, and work in progress:

- So far mostly a mathematical investigation into the structure of $\overline{\mathcal{D}}_2$. Interesting to look for more phenomenological examples.

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Outlook, and work in progress:

- So far mostly a mathematical investigation into the structure of $\overline{\mathcal{D}}_2$. Interesting to look for more phenomenological examples.
- Further investigation into higher order deformations and obstructions corresponding to Yukawa couplings.

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- So far mostly a mathematical investigation into the structure of $\overline{\mathcal{D}}_2$. Interesting to look for more phenomenological examples.
- Further investigation into higher order deformations and obstructions corresponding to Yukawa couplings.
- Need Kähler potential to investigate the 4d theory outside of Minkowski vacua. Holomorphic structures usually come with natural Kähler metric (Weil-Peterson metric, etc). Clue for what Kähler metric is?

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$$dH + W_{NS5} = \frac{\alpha'}{4} (\text{tr } F^2 - \text{tr } R^2) .$$

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Outlook, and work in progress:

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$$dH + W_{NS5} = \frac{\alpha'}{4} (\text{tr } F^2 - \text{tr } R^2) .$$

⇒ Spoils holomorphic structure $\overline{\partial}_2$.

Thank you!

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Thank you for your attention!