

#### Franco D. Albareti

*PhD student under the supervision of* 

**Prof. Antonio L. Maroto and Prof. Francisco Prada** 



Instituto de Física Teórica UAM/CSIC 19<sup>th</sup> November 2015



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Acknowledgements



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# **Outline**

#### Introduction

#### • Modes

#### • Results

#### • Observational effects

#### • Conclusions









IV Post. Meet. Theor. Phys.

# $\Box \phi + V'(\phi) = 0$

 $oldsymbol{V'}(oldsymbol{\hat{\phi}})\,+\,oldsymbol{V''}(oldsymbol{\hat{\phi}})\,oldsymbol{\delta}\phi\,+\,rac{1}{2}oldsymbol{V'''}(oldsymbol{\hat{\phi}})\,oldsymbol{\delta}\phi^2\,+\,...$ 

# $\Box \phi + V'(\phi) = 0$

 $oldsymbol{V'}(oldsymbol{\hat{\phi}})\,+\,oldsymbol{V''}(oldsymbol{\hat{\phi}})\,oldsymbol{\delta}\phi\,+\,rac{1}{2}oldsymbol{V'''}(oldsymbol{\hat{\phi}})\,oldsymbol{\delta}\phi^2\,+\,...$ 

Take  $\langle 0|...|0 \rangle$ 

$$egin{aligned} & igcap \hat{\phi} + m{V'}_{ ext{eff}} \left( \hat{\phi} 
ight) = m{0} \ & m{V'}(\hat{\phi}) + m{V''}(\hat{\phi}) raket{\delta \phi} + rac{1}{2} m{V'''}(\hat{\phi}) raket{\delta \phi^2} + ... \end{aligned}$$

Take 
$$\langle 0|...|0
angle$$

$$\begin{array}{l} \textbf{Introduction} \\ \Box \, \hat{\phi} + V'_{eff} \left( \hat{\phi} \right) = 0 \\ V'(\hat{\phi}) + V''(\hat{\phi}) \langle \delta \phi \rangle + \frac{1}{2} V'''(\hat{\phi}) \langle \delta \phi^2 \rangle + \dots \\ O \quad \text{Quantum fluctuations} \\ (1-loop) \\ \text{Take} \quad \langle 0 | \dots | 0 \rangle \end{array}$$

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# Introduction $\Box \, \hat{\phi} + V'_{\text{eff}} \left( \hat{\phi} \right) \, = \, \mathbf{0}$ $V'(\hat{\phi}) + rac{1}{2}V'''(\hat{\phi}) \langle \delta \phi^2 angle + ...$



1-loop effective potential for  $\hat{\phi}$ 

$${V}_{ ext{eff}}^{\circ}\left(\hat{\phi}
ight) = V(\hat{\phi}) + V^{\circ}(\hat{\phi})$$

$$V^{\circ}(\hat{\phi})\,=\,rac{1}{2}\int_{0}^{m^{2}(\phi)}\mathrm{d}m^{2}\,\langle0|\delta\phi^{2}|0
angle$$

1-loop effective potential for  $\hat{\phi}$ 

$${V}_{ ext{eff}}^{\circ}\left(\hat{\phi}
ight)\,=\,V(\hat{\phi})\,+\,V^{\circ}(\hat{\phi})$$

$$V^{\circ}(\hat{\phi}) \,=\, rac{1}{2} \int_{0}^{m^2(\hat{\phi})} \mathrm{d}m^2 \, \langle 0 | \delta \phi^2 | 0 
angle$$
 Quadratic operator

#### **Flat space-time**

$$m{V}^{ extsf{o}}_{ extsf{eff}} = m{V} + rac{\hbar}{2} \int rac{ extsf{d}^3 m{k}}{(2\pi)^3} \sqrt{m{k}^2 + m^2}$$
  
Vaccum energy of  $rac{1}{2} \hbar m{\omega}$ 

**Flat space-time** 

$$m{V}^{ extsf{o}}_{ extsf{eff}} = m{V} + rac{\hbar}{2} \int rac{ extsf{d}^3 m{k}}{(2\pi)^3} \sqrt{m{k}^2 + m^2}$$
  
Vaccum energy of  $rac{1}{2} \hbar m{\omega}$ 

$$oldsymbol{V}_{ extbf{eff}}^{ extbf{o}} = oldsymbol{V} + rac{\hbar}{64\pi^2} oldsymbol{m}^4 \log\left(rac{oldsymbol{m}^2}{oldsymbol{\mu}^2}
ight)$$

Coleman & Weinberg '73 Including W, Z and top contributions



Flat FRW  
$$\boldsymbol{V}_{eff}^{\circ} = \boldsymbol{V} + \frac{\hbar}{2} \int \frac{\mathrm{d}^3(\boldsymbol{k}/\boldsymbol{a})}{(2\pi)^3} \sqrt{(\boldsymbol{k}/\boldsymbol{a})^2 + \boldsymbol{m}^2}$$

Comoving scales get Komoving scales det Komovi

$$m{V}_{ ext{eff}}^{\circ} = m{V} + rac{\hbar}{64\pi^2} \, m{m}^4 \log\left(rac{m{m}^2}{m{\mu}^2}
ight)$$
  
Physical renormalization scale

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# Non-homogeneus effect?

#### Perturbations

 $\mathrm{d}s^2\,=\,a^2(\eta)\left\{\left[1+2\Phi(\eta,x)
ight]\mathrm{d}\eta^2-\left[1-2\Psi(\eta,x)
ight]\mathrm{d}x^2
ight\}$ 

Static weak gravitational field (Solar System...)

Cosmological effects due to the LSS



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# **Outline**



#### • Modes

#### • Results

#### • Observational effects

#### • Conclusions

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## Introduction

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Modes1-loop effective potential for 
$$\hat{\phi}$$
 $V_{ ext{eff}}^{\circ}(\hat{\phi}) = V(\hat{\phi}) + V^{\circ}(\hat{\phi})$  $V^{\circ}(\hat{\phi}) = rac{1}{2} \int_{0}^{m^{2}(\hat{\phi})} \mathrm{d}m^{2} \langle 0|\delta\phi^{2}|0
angle$ Quadratic operator



# this is all about classical gravity

# with quantum fields

# QFT in curved spacetimes

(Birrel & Davies '82)

1) Quantize the fluctuations canonically

$$\delta \phi = \int rac{\mathrm{d} k^3}{(2\pi)^3} \left( a_k \, \delta \phi_k \, + \, a_k^\dagger \, \delta \phi_k^* 
ight)$$

#### 2) Find the modes

• Solve KG to first order (in metric perturbations...)

• Using a WKB ansatz

$$\delta \phi_k \!=\! f_k \, e^{i\, heta_k}$$

$$\omega^2 \gg \mathcal{H}^2,\, 
abla^2(\Phi,\Psi)$$

1) Quantize the fluctuations canonically

$$\delta \phi = \int rac{\mathrm{d} k^3}{(2\pi)^3} \left( a_k \, \delta \phi_k \, + \, a_k^\dagger \, \delta \phi_k^st 
ight)$$

#### 2) Find the modes

- Solve KG to first order (in metric perturbations...)
- Using a WKB ansatz
- Boundary conditions Match the perturbed modes to the unperturbed ones at  $\eta = 0$ (adiabatic vacuum)

1) Quantize the fluctuations canonically

$$\delta \phi = \int rac{\mathrm{d} k^3}{(2\pi)^3} \left( a_k \, \delta \phi_k \, + \, a_k^\dagger \, \delta \phi_k^st 
ight)$$

2) Find the modes

$$\delta \phi_k \!=\! f_k \, e^{i\, heta_k}$$

3) Compute

 $\langle \mathbf{0} | oldsymbol{\delta} oldsymbol{\phi}^2 | \mathbf{0} 
angle$ 

4) Regularize & Renormalize
# Modes

#### 4) Regularize & Renormalize

- Fourier space
- Expand in powers of p
- Dimensional regularization for the integration over quantum modes k

#### Renormalization?

 $\langle \mathbf{0} | oldsymbol{\delta} oldsymbol{\phi}^2 | \mathbf{0} 
angle$ 

- The same UV behaviour than in flat space-time
- Contributions from  $oldsymbol{\Phi},\, oldsymbol{\Psi}$  to the  $\,V_{\,
  m eff}^{\circ}\,$  are finite.

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$$m{V}_{ ext{eff}}^{ ext{o}} = m{V} + rac{\hbar}{64\pi^2} m{m}^4 \log\left(rac{m{m}^2}{m{\mu}^2}
ight) + rac{\hbar}{16\pi^2} m{m}^4 \left(m{H}_{m{\Phi}} + m{H}_{m{\Psi}}
ight)$$

$$m{V}_{ ext{eff}}^{\circ} = m{V} + rac{\hbar}{64\pi^2} m{m}^4 \log\left(rac{m{m}^2}{m{\mu}^2}
ight) + rac{\hbar}{16\pi^2} m{m}^4 \left(m{H}_{m{\Phi}} + m{H}_{m{\Psi}}
ight)$$
Homogeneous
contribution

Divergent

 Does not depend on geometry to this order

$$V_{eff}^{o} = V + \frac{\hbar}{64\pi^{2}} m^{4} \log\left(\frac{m^{2}}{\mu^{2}}\right) + \frac{\hbar}{16\pi^{2}} m^{4} (H_{\Phi} + H_{\Psi})$$
Homogeneous  
contribution
Non-homogeneous  
contribution
Oivergent
Divergent
Does not depend on  
geometry to this order
Non-local

Why not effective action?

$$\Gamma_{oldsymbol{Q}} = \int \mathrm{d}^4 x \left( - oldsymbol{V}_{ ext{eff}} 
ight) + ... \left( \partial \phi 
ight)^2 + ...$$

$$\omega^2 \gg \mathcal{H}^2, \, 
abla^2(\Phi,\Psi)$$

Higgs boson ~ 125 GeV mass Solar System ~ 10<sup>-25</sup> GeV

Cosmology ~ 10<sup>-39</sup> GeV



Non-homogeneous effect, the field value which minimizes the potential is different for each point of spacetime.

#### Perturbations

 $\mathrm{d}s^2\,=\,a^2(\eta)\left\{\left[1+2\Phi(\eta,x)
ight]\mathrm{d}\eta^2-\left[1-2\Psi(\eta,x)
ight]\mathrm{d}x^2
ight\}$ 

Static weak gravitational field (Solar System...)

Cosmological effects due to the LSS

#### **Perturbations**

 $\mathrm{d} s^2\,=\,a^2(\eta)\left\{\left[1+2\Phi(\eta,x)
ight]\mathrm{d} \eta^2-\left[1-2\Psi(\eta,x)
ight]\mathrm{d} x^2
ight\}$ 

gravitational field (Solar System.)

Static weak

Cosmological effects due to the LSS

Solar System (momentum space)

$$H_{\Phi+\Psi} = \left(rac{\sin(p\,t)}{p\,t} - \cos(p\,t)
ight) \left(rac{\Phi(p)+\Psi(p)}{2}
ight)$$

$$H_{\Phi-\Psi}\,=\,\left(rac{\sin(p\,t)}{p\,t}-1
ight)\left(rac{\Phi(p)-\Psi(p)}{2}
ight)$$

Solar System (momentum space)

$$H_{\Phi+\Psi} = \left(rac{\sin(p\,t)}{p\,t} - \cos(p\,t)
ight) \left(rac{\Phi(p)+\Psi(p)}{2}
ight)$$

$$H_{\Phi-\Psi} = \left(rac{\sin(p\,t)}{p\,t} - 1
ight) \left(rac{\Phi(p) - \Psi(p)}{2}
ight)$$

$$H_{\Phi+\Psi}, H_{\Phi-\Psi} \xrightarrow{p \to 0} 0$$

There is no shift in the spacetime mean value of Higgs VEV

ResultsNewtonian  
potentialSolar System (momentum space)
$$H_{\Phi+\Psi} = \left(\frac{\sin(pt)}{pt} - \cos(pt)\right) \left(\frac{\Phi(p) + \Psi(p)}{2}\right)$$
 $H_{\Phi-\Psi} = \left(\frac{\sin(pt)}{pt} - 1\right) \left(\frac{\Phi(p) - \Psi(p)}{2}\right)$  $H_{\Phi+\Psi}, H_{\Phi-\Psi} \xrightarrow{p \to 0} 0$ There is no shift in the spacetime mean  
value of Higgs VEV

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Solar System (real space, momentum space)

• Newtonian potential  $\Phi_N(r) = -rac{1}{r} \, G \, M \, igstarrow \Phi_N(p) = -rac{4\pi}{p^2} \, G \, M$ 

Solar System (real space)

$$egin{aligned} H^N_{\Phi+\Psi} &= \left(rac{r}{t}
ight) \, \Phi_N \ & imes \, heta(t^2-r^2) \ &H^N_{\Phi-\Psi} &= rac{1}{2} \left(rac{r}{t}-1
ight) \, \Phi_N \left(1-\gamma
ight) \end{aligned}$$

Solar System (real space)

$$egin{aligned} H^N_{\Phi+\Psi} &= \left(rac{r}{t}
ight) \, \Phi_N \ H^N_{\Phi-\Psi} &= rac{1}{2} \left(rac{r}{t}-1
ight) \, \Phi_N \left(1-\gamma
ight) \ egin{aligned} imes \, extsf{ ilde heta} &= rac{1}{2} \left(rac{r}{t} extsf{ ilde heta} 
ight) \, \Phi_N \left(1-\gamma
ight) \ egin{aligned} imes \, imes \,$$

Solar System (real space)

$$\begin{split} H^{N}_{\Phi+\Psi} &= \left(\frac{r}{t}\right) \Phi_{N} \\ H^{N}_{\Phi-\Psi} &= \frac{1}{2} \left(\frac{r}{t} - 1\right) \Phi_{N} \left(1 - \gamma\right) \\ \end{split} \\ \\ \mathsf{Remarks} \\ \end{split}$$

•  $\frac{r}{t}$   $\Rightarrow$  Boundary effects from the bc's of the modes

$$egin{aligned} H^N_{\Phi+\Psi} &= \left( egin{aligned} r & \Phi_N \ H^N_{\Phi-\Psi} &= rac{1}{2} \left( egin{aligned} r & -1 \ t & -1 \ \end{array} 
ight) \Phi_N \left( 1 - \gamma 
ight) & & & \mathcal{H}^N_{\Phi-\Psi} \ \mathcal{F}_{ausality} \end{aligned}$$
Remarks

•  $rac{r}{t}$  ightarrow Boundary effects from the bc's of the modes  $t 
ightarrow \infty$  Franco D. Albareti

$$egin{aligned} H^N_{\Phi+\Psi}&=0\ H^N_{\Phi-\Psi}&=-rac{1}{2}\Phi_N\left(1-\gamma
ight)\ extscale{Remarks} & extscale{Eddington}\ extscale{Farmeter} \end{aligned}$$

•  $rac{r}{t}$  ightarrow Boundary effects from the bc's of the modes  $t 
ightarrow \infty$  Franco D. Albareti

$$egin{aligned} H^N_{\Phi+\Psi}&=0\ H^N_{\Phi-\Psi}&=-rac{1}{2}\Phi_N\left(1-\gamma
ight)\ H^N_{\Phi-\Psi}&=rac{$$

Solar System (real space, momentum space)

• Newtonian potential  $\Phi_N(r) = -rac{1}{r} \, G \, M \, igstarrow \Phi_N(p) = -rac{4\pi}{p^2} \, G \, M$ 

Solar System (real space, momentum space)

• Newtonian potential 
$$\Phi_N(r) = -rac{1}{r} \, G \, M \, igstarrow p_N(p) = -rac{4\pi}{p^2} \, G \, M$$

• General potential

$$egin{aligned} \Phi(r) &= -rac{1}{r}\sum_{l=0}^{\infty}\sum_{m=l}^{l}rac{Q_{lm}}{r^l}\sqrt{rac{4\pi}{2l+1}}Y_{lm}( heta,\phi) \ &igoplus \ &igoplus \ \Phi(p) &= -rac{4\pi}{p^2}\sum_{l=0}^{\infty}i^l\sum_{m=l}^{l}rac{Q_{lm}}{(2p)^l}\sqrt{rac{4\pi}{2l+1}}Y_{lm}( heta_p,\phi_p) \end{aligned}$$

Newtonian results

$$egin{aligned} H^N_{\Phi+\Psi} &= 0 \ H^N_{\Phi-\Psi} &= -rac{1}{2} \Phi_N \left(1-\gamma
ight) \ Faddington \ Farameter \end{aligned}$$

General results

$$egin{aligned} H_{\Phi+\Psi}^{oldsymbol{X}}&=0\ H_{\Phi-\Psi}^{oldsymbol{X}}&=-rac{1}{2}\Phi_{oldsymbol{X}}\left(1-\gamma
ight)\ ext{Eddington}\ ext{Eddington}\ ext{parameter} \end{aligned}$$

# Outline

#### Introduction





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#### Introduction





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$$\Delta_{
m ~Higgs}= rac{3\lambda}{8\pi^2}\Delta\Phi\left(1-\gamma
ight)$$

$$rac{\Delta \mu}{\mu} = -\Delta_{
m Higgs}$$

$${
m Proton-to-electron}\ {
m mass\ ratio}\ {\Delta\mu\over\mu} < 10^{-16}\ {
m Atomic\ clocks}\ {
m on\ Earth}\ {
m Huntemann,\ et\ al.}$$

2014



$$\Delta^{i}_{
m \, Higgs} pprox rac{3\lambda}{8\pi^{2}} \Delta \Phi \left(1-\gamma
ight)$$

- Higgs self-interactions
- Vector bosons
- Top quark

$$\Delta^{i}_{
m Higgs} pprox rac{3\lambda}{8\pi^{2}} \Delta\Phi\left(1-\gamma
ight) imes n_{
m eff} imes \ v \ v \ {
m Bosons,} \ {
m Fermions}$$

- Higgs self-interactions
- Vector bosons
- Top quark

$$\Delta^{i}_{\mathrm{Higgs}} \approx rac{3\lambda}{8\pi^{2}} \Delta \Phi \left(1-\gamma\right) \times n_{\mathrm{eff}} \times \left(rac{g_{i}}{\lambda}
ight) \times \left(rac{m_{i}}{m_{\mathrm{Higgs}}}
ight)^{4}$$
  
Bosons,  
Fermions Mass/coupling factors

- Higgs self-interactions
- Vector bosons
- Top quark

Work in progress...

# Outline Outline Introduction Modes Results

# Observational effects

#### • Conclusions



#### Conclusions
## Conclusions

Metric perturbations contribute to the finite part of the Higgs 1-loop effective potential

• This leads to a space-time dependent Higgs VEV

which translates into variations on the masses of all the elementary particles.

Within the Solar System, constraints on the Eddington
parameter can be obtained from measurements of the proton-to-electron mass ratio.

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