

A holographic model of Weyl semi-metal

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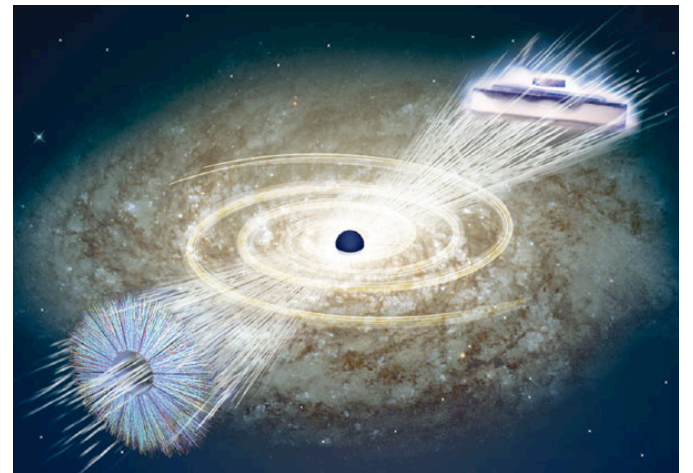
Based on

Karl Landsteiner, YL, **1505.04772**

Karl Landsteiner, YL, Ya-Wen Sun, **1511.05505**

AdS/CMT

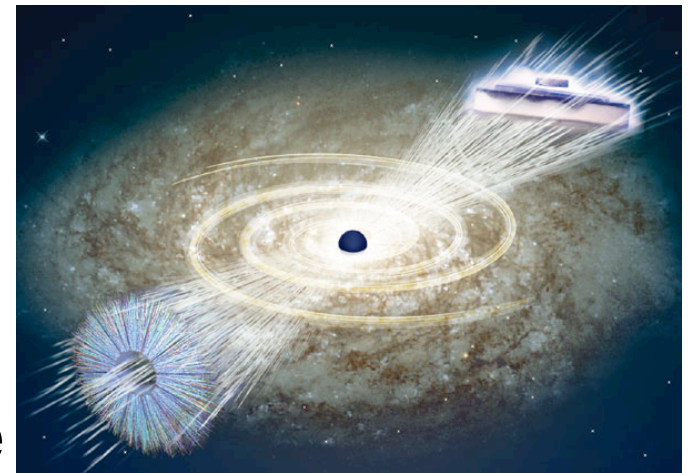
- Many phenomena observed in condensed matter systems can be studied by general relativity.



- Superconducting/superfluid phase transition
- Fermi surface, non Fermi liquid
- Lattice/impurity effect ...

[Hartnoll, Horowitz, Herzog, H.Liu, McGreevy, Schalm, Zaanen, Tong, et al.]

AdS/CMT

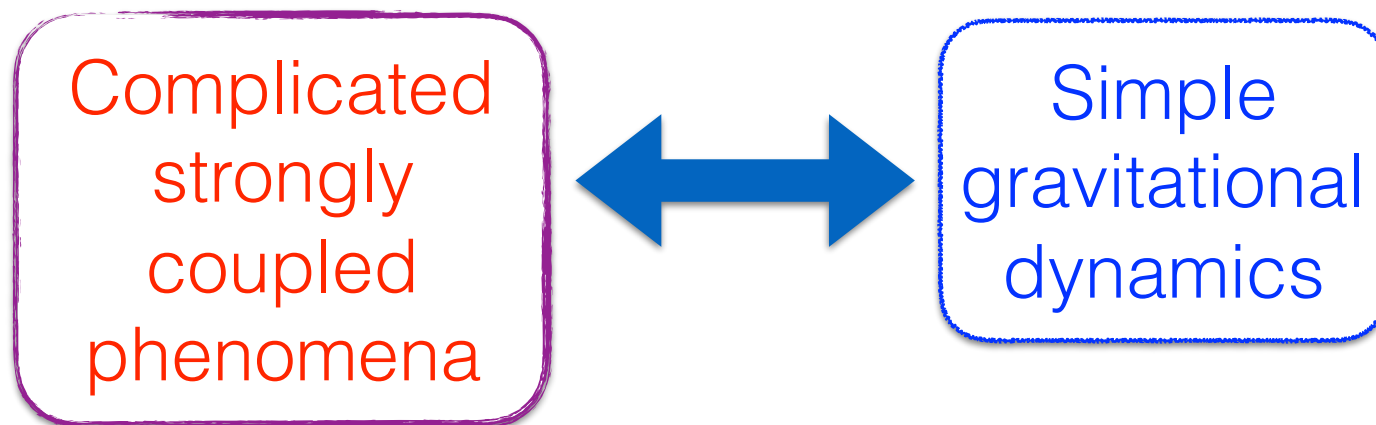


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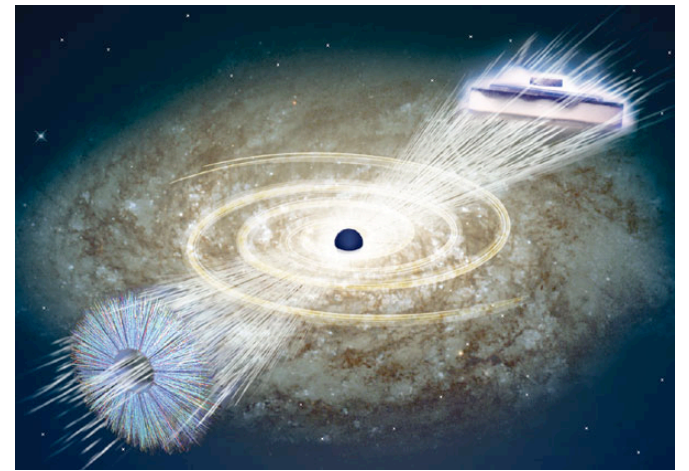
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[Hartnoll, Horowitz, Herzog, H.Liu, McGreevy, Schalm, Zaanen, Tong, et al.]

- It is a consequence of **Gauge/gravity duality**



AdS/CMT



- Many phenomena observed in condensed matter systems can be studied by general relativity.

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[Hartnoll, Horowitz, Herzog, H.Liu, McGreevy, Schalm, Zaanen, Tong, et al.]

* AdS/other examples of condensed matter system?

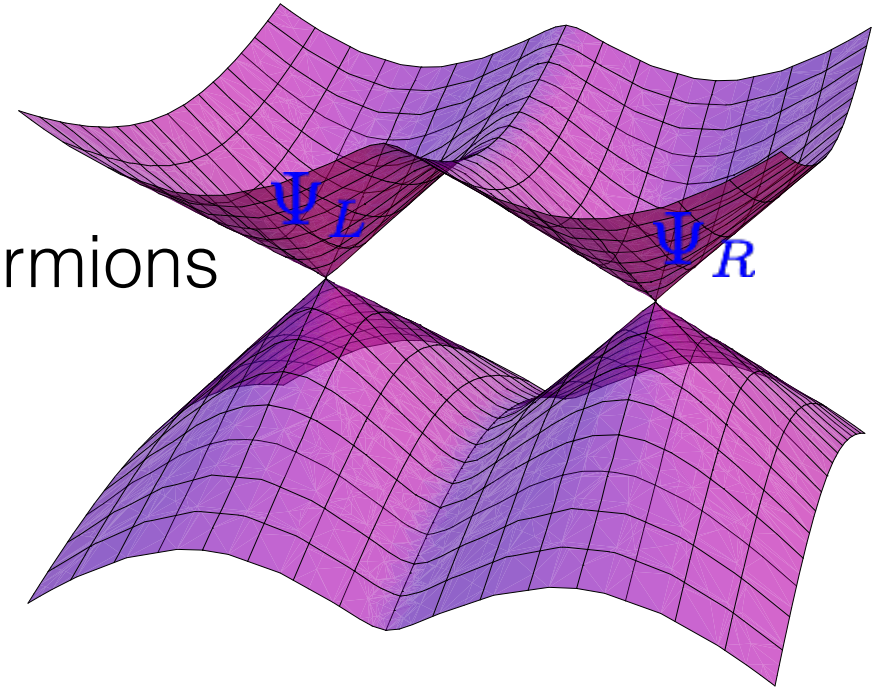
- **Weyl semi-metal (a topological state)**

Outline

- Weyl semi-metal (WSM): QFT model
- Holographic model of WSM
- Summary and open questions

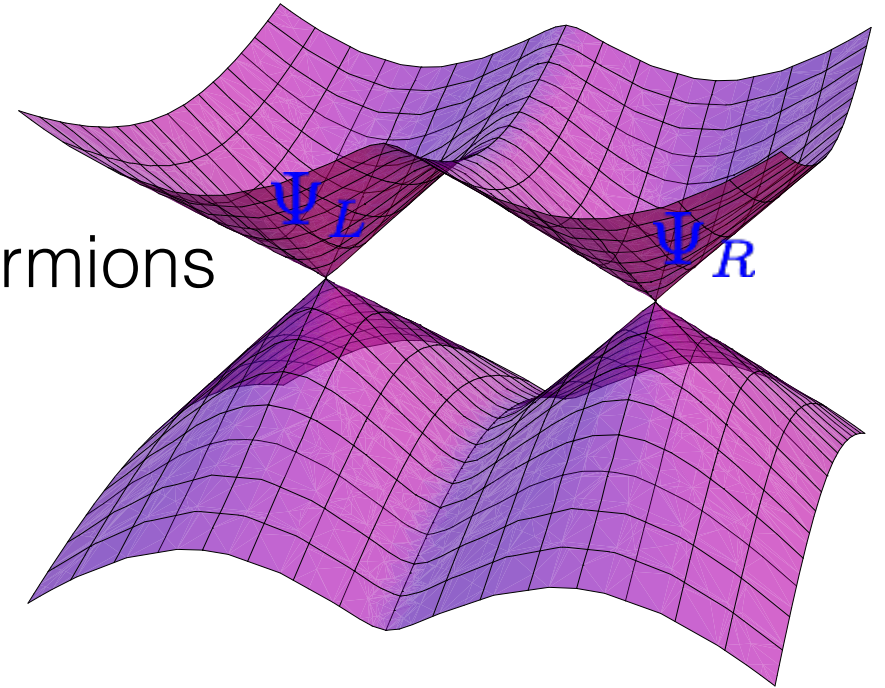
Weyl semi-metal

- Low energy excitation: Weyl fermions
- conduction and valence band (linear) touch



Weyl semi-metal

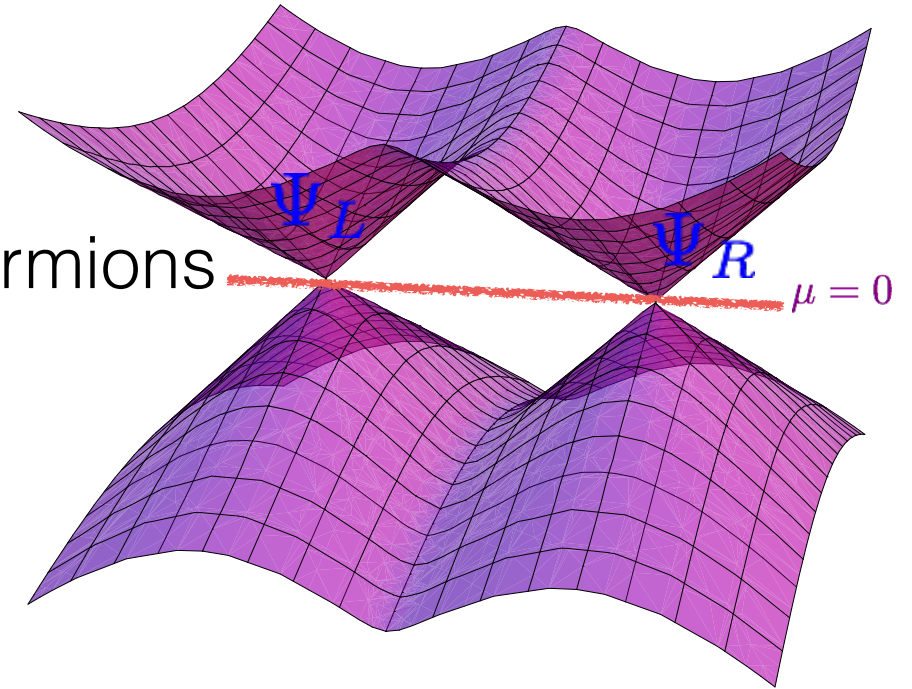
- Low energy excitation: Weyl fermions
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- Weyl points:
 - ◆ Close to Weyl points, Weyl equation
 - ◆ Weyl points appear in +/- pairs with opposite chiralities [Nielsen, Ninomiya]
 - ◆ Weyl points are topologically protected in momentum space
- Experimentally realisation: TaAs (2015)...

Weyl semi-metal

- Low energy excitation: Weyl fermions
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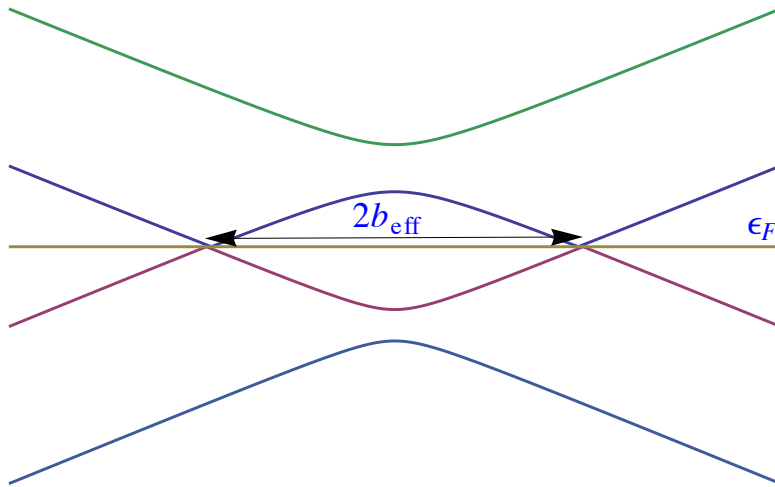
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QFT of WSM

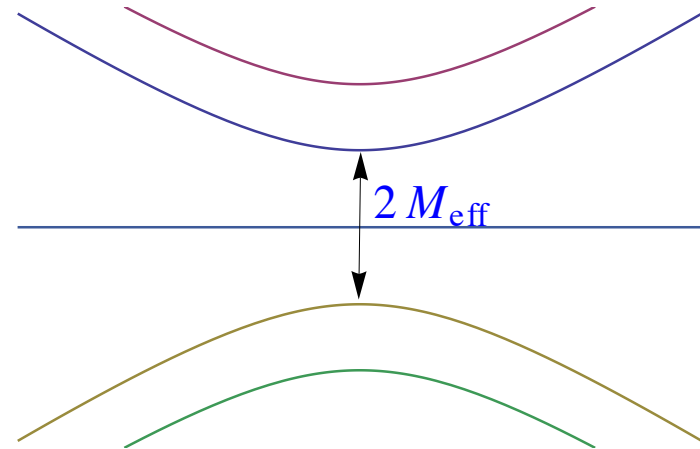
$$\mathcal{L} = \bar{\Psi} (i\gamma^\mu \partial_\mu + M - \gamma_5 \gamma_z b) \Psi .$$

spectrum

$$M < b : \quad b_{\text{eff}} = \sqrt{b^2 - M^2}$$



$$M > b : \quad M_{\text{eff}} = \sqrt{M^2 - b^2}$$



QFT of WSM

- Topological phase

$$M < b: \quad b_{\text{eff}} = \sqrt{b^2 - M^2} \quad \mathcal{L}_{\text{eff}} = \bar{\psi} (i\gamma^\mu \partial_\mu - \gamma_5 \gamma_z b_{\text{eff}}) \psi$$

constant axial gauge field
axial gauge transformation
axial anomaly
electric current

$$A_z^5 = b_{\text{eff}}$$

$$\theta_5 = b_{\text{eff}} z$$

$$W = \int d^4x \theta_5 F \wedge F$$

$$J^\mu = \frac{\delta W}{\delta A_\mu}$$

- Anomalous Hall Effect (AHE) [Haldane, 1987]

$$\mathbf{J} = \frac{e^2}{2\pi^2} \mathbf{b}_{\text{eff}} \times \mathbf{E}$$

QFT of WSM

$$M > b: \quad M_{\text{eff}} = \sqrt{M^2 - b^2} \quad \mathcal{L}_{\text{eff}} = \bar{\psi} (i\gamma^\mu \partial_\mu + M_{\text{eff}}) \psi$$

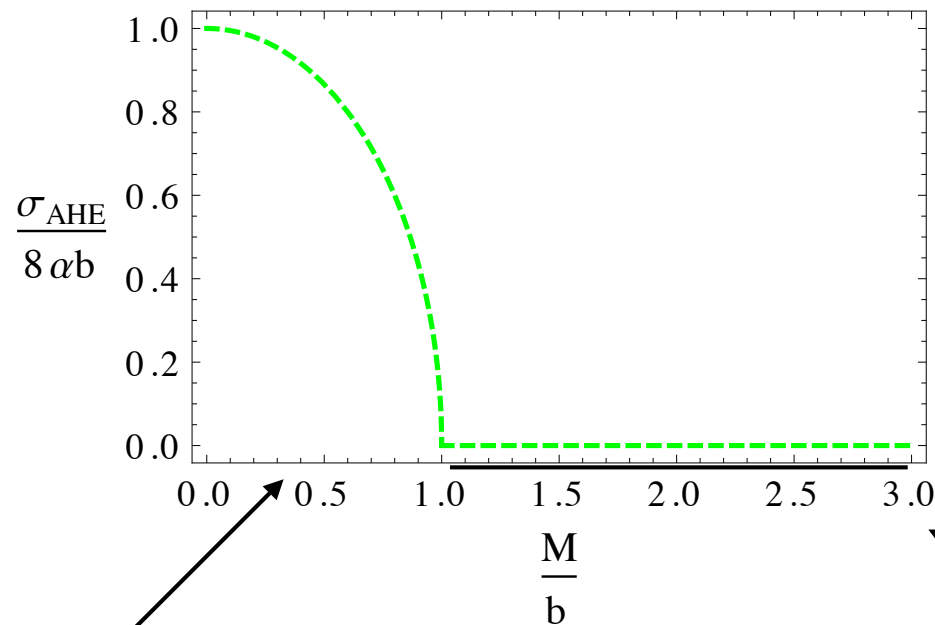
- gapped phase with vanishing AHE.
- More generally, more Dirac cones can be present
topologically trivial semi-metal

$$\mathcal{L} = \bar{\Psi} (i\gamma^\mu \partial_\mu + M - \gamma_5 \gamma_z b) \Psi + \sum_{j=1}^N \bar{\psi} (i\gamma^\mu \partial_\mu) \psi$$

QFT of WSM

$$\mathcal{L} = \bar{\Psi} (i\gamma^\mu \partial_\mu + M - \gamma_5 \gamma_z b) \Psi .$$

- Topological phase can be characterised by a **nonzero** anomalous Hall conductivity.



topological phase

Trivial phase

Motivation for Holographic WSM

- How does WSM work in strongly coupled case?

without quasiparticle,
no notion of Berry phase (Weyl points)

A holography model for WSM can teach us qualitative lessons!

AdS/CFT: dictionary

[Maldacena; Witten;
Gubser, Klebanov,
Polyakov]

Bulk AdS

4+1D gravity
weakly coupled
metric
gauge field
local symmetry
scalar field

Boundary field theory

3+1D field theory
strongly coupled
stress tensor
current
global symmetry
scalar operator

$$e^{iS_{\text{bulk}}}|_{\phi \rightarrow J} = \langle \exp(i \int d^4x \sqrt{-g_0} J \mathcal{O}) \rangle$$

Model WSM from gravity

- Weakly coupled theory $\mathcal{L} = \bar{\Psi} (i\gamma^\mu \partial_\mu + M - \gamma_5 \gamma_z b) \Psi$.

$$\partial_\mu J^\mu = 0$$

$$\partial_\mu J_5^\mu = \frac{1}{16\pi^2} 3\mathcal{F} \wedge \mathcal{F} + 2M\bar{\Psi}\gamma_5\Psi$$

- **Ingredients** for AdS gravity model:
 - Gravity with negative cosmological constant
 - One gauge field dual to electric current
 - One gauge field dual to axial current
 - One scalar field charged under axial gauge symmetry

Holographic WSM

- Holographic model

$$\begin{aligned}\mathcal{L} = & \frac{1}{2\kappa^2} \left(R + \frac{12}{L^2} \right) - \frac{1}{4} \mathcal{F}^2 - \frac{1}{4} F_5^2 \\ & + \frac{\alpha}{3} A_5 \wedge (F_5 \wedge F_5 + 3\mathcal{F} \wedge \mathcal{F}) + \\ & + |(\partial_\mu - iqA_\mu^5)\Phi|^2 - V(\Phi)\end{aligned}$$

- CS structure = anomaly form
- the dimension of the dual scalar operator is chosen to be 3, mass deformation

Holographic WSM

- Currents:

$$J^\mu = \frac{\delta S}{\delta A_\mu(r = \infty)} = \lim_{r \rightarrow \infty} \sqrt{-g} \left(\mathcal{F}^{\mu r} + 4\alpha \epsilon^{r\mu\nu\rho\alpha} A_\nu^5 \mathcal{F}_{\rho\alpha} \right)$$

$$J_5^\mu = \frac{\delta S}{\delta A_\mu^5(r = \infty)} = \lim_{r \rightarrow \infty} \sqrt{-g} \left(F_5^{\mu r} + 4\alpha \epsilon^{r\mu\nu\rho\alpha} A_\nu^5 F_{\rho\alpha}^5 \right)$$

- Wald identity

$$\partial_\mu J^\mu = 0$$

$$\partial_\mu J_5^\mu = \left(\frac{\alpha}{3} \left[F_5 \wedge F_5 + 3\mathcal{F} \wedge \mathcal{F} \right] - iq\sqrt{-g} \left[\Phi(D_r\Phi)^* - \Phi^*(D_r\Phi) \right] \right) \Big|_{r \rightarrow \infty}$$

Holographic WSM

- **Ansatz (T=0)**

$$ds^2 = u(-dt^2 + dx^2 + dy^2) + \frac{dr^2}{u} + h dz^2,$$

$$\Phi = \phi,$$

$$A^5 = A_z^5 dz.$$

- **Near UV**

Metric: $ds^2|_{r \rightarrow \infty} = \frac{dr^2}{r^2} + r^2(-dt^2 + d\vec{x}^2)$

vector gauge field: $A_\mu = 0$

axial gauge field: $A_\mu^5|_{r \rightarrow \infty} = b\delta_\mu^z$

scalar field: $r\Phi_\mu|_{r \rightarrow \infty} = M$

Holographic WSM

At zero temperature: 3 distinct classes of solutions
(leading order solution @ IR)

- **M/b < 0.744** (Topological phase) $A_z^5(0) = b_{\text{eff}}, \Phi(0) = 0$

- **M/b = 0.744** (Critical point) $A_z^5 = r^\beta, \Phi(0) = \phi_0$

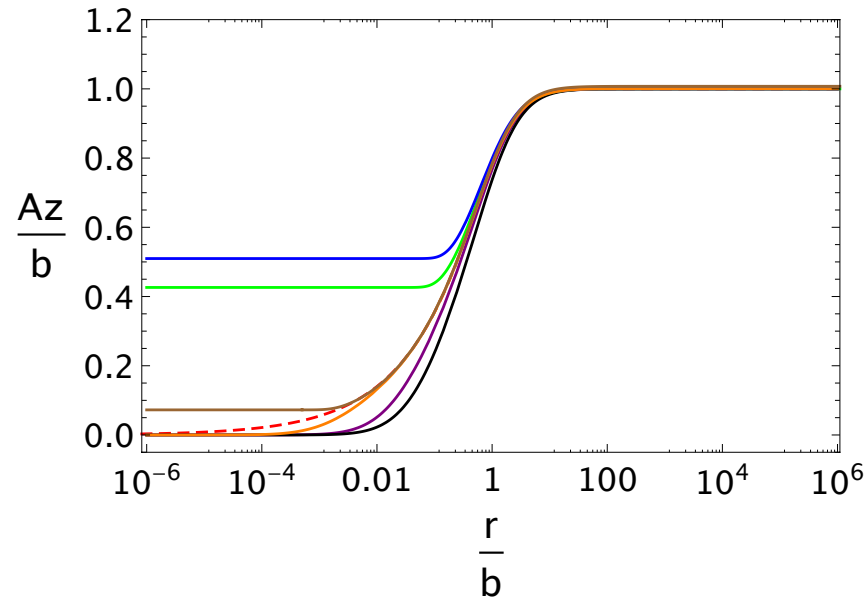
$$ds^2 = u_0 r^2 (-dt^2 + dx^2 + dy^2) + \frac{dr^2}{u_0 r^2} + h_1 r^{2\beta} dz^2,$$

- **M/b > 0.744** (Trivial phase) $A_z^5(0) = 0, \Phi(0) = \phi_{\text{min}}$

$$\frac{dV}{d\phi}(\phi_{\text{min}}) = 0$$

Holographic WSM

Bulk profile of
axial gauge field



$$M/b = 0.695$$

$$M/b = 0.712$$

$$M/b = 0.743$$

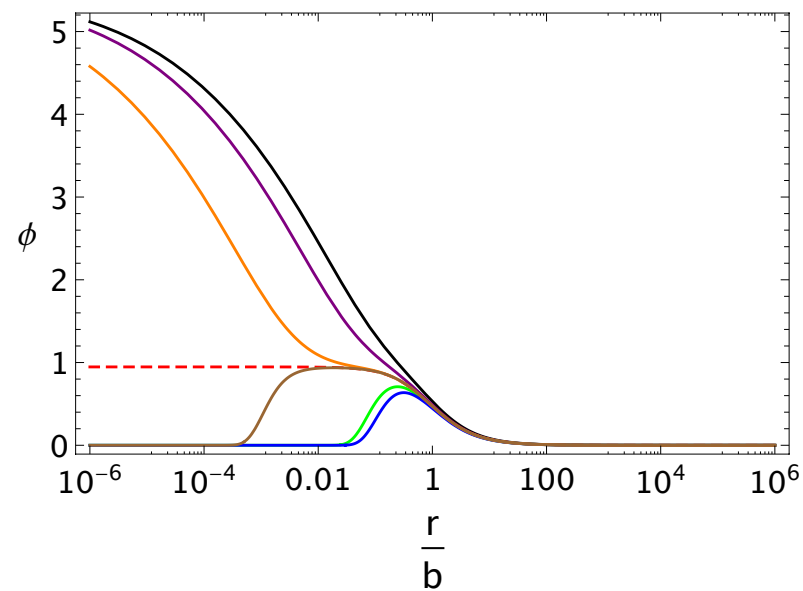
$$M/b = 0.744$$

$$M/b = 0.745$$

$$M/b = 0.778$$

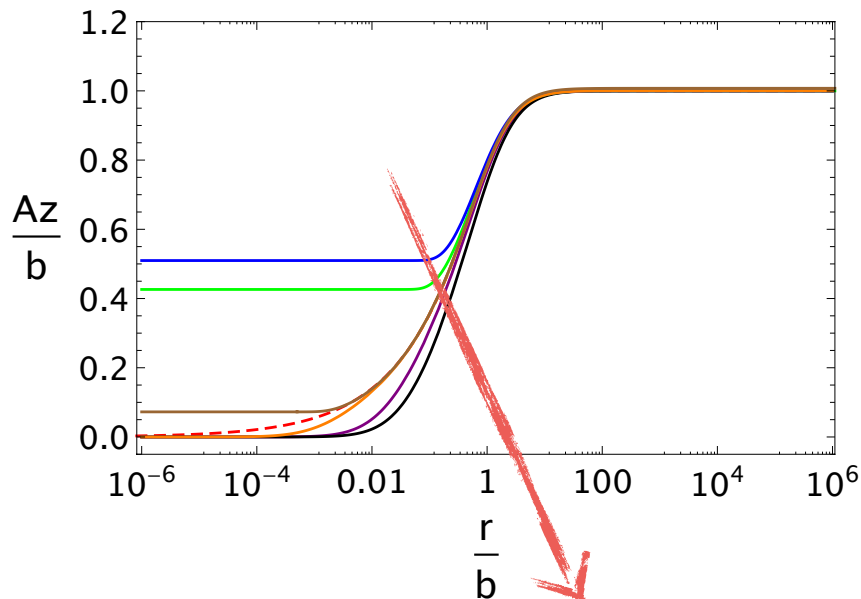
$$M/b = 0.856$$

Bulk profile of
scalar field



Holographic WSM

Bulk profile of axial gauge field



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$M/b = 0.712$

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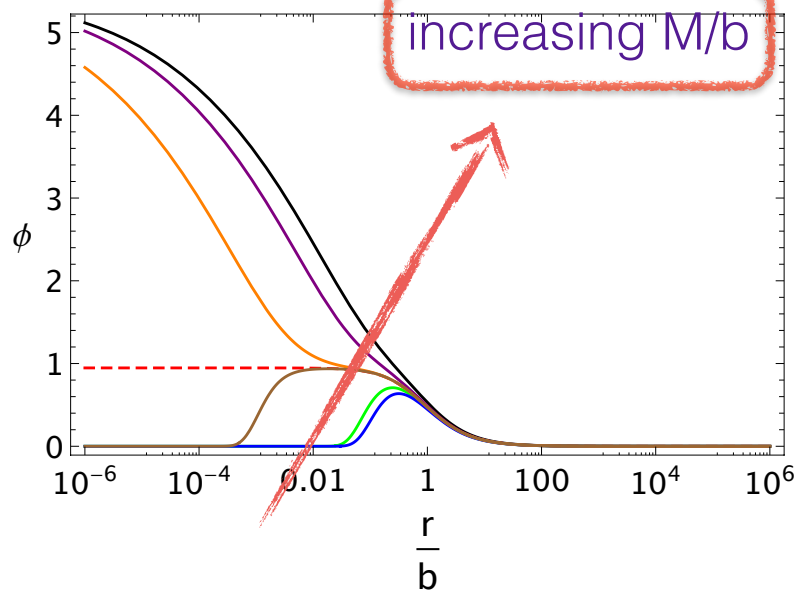
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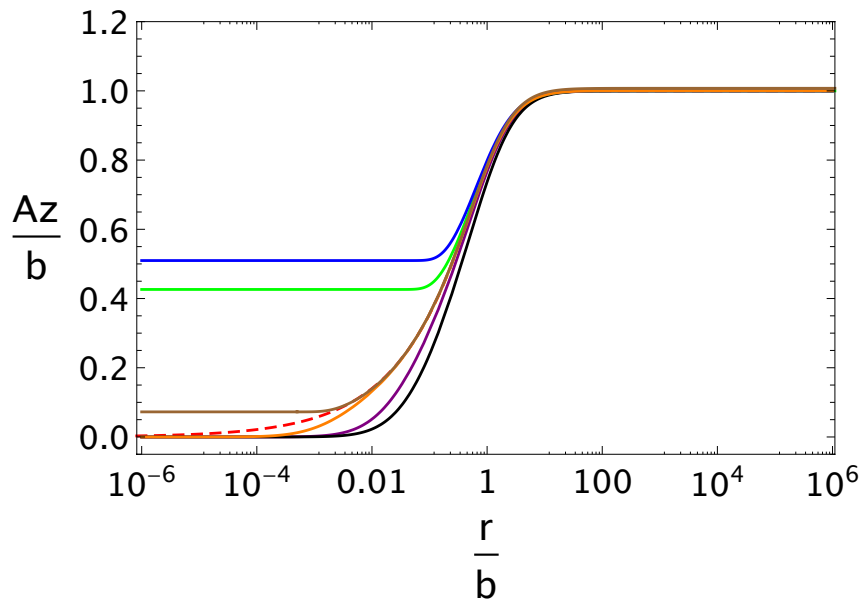
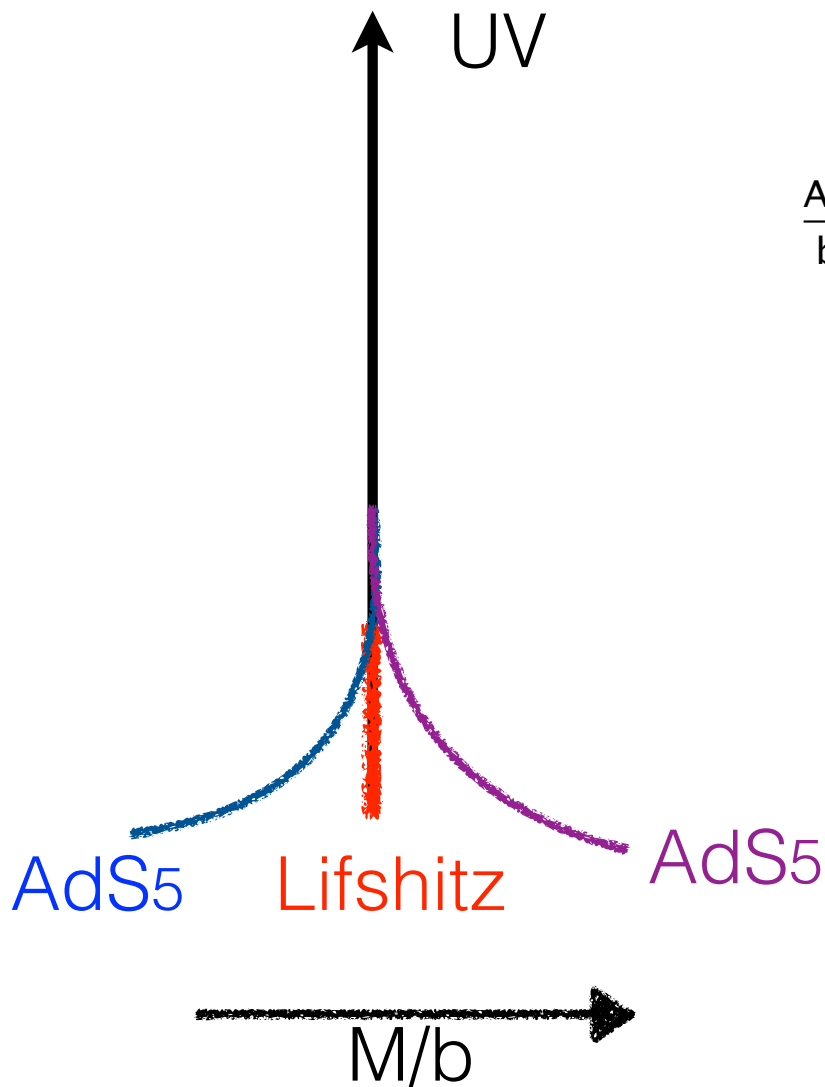
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Bulk profile of scalar field



increasing M/b

Holographic WSM



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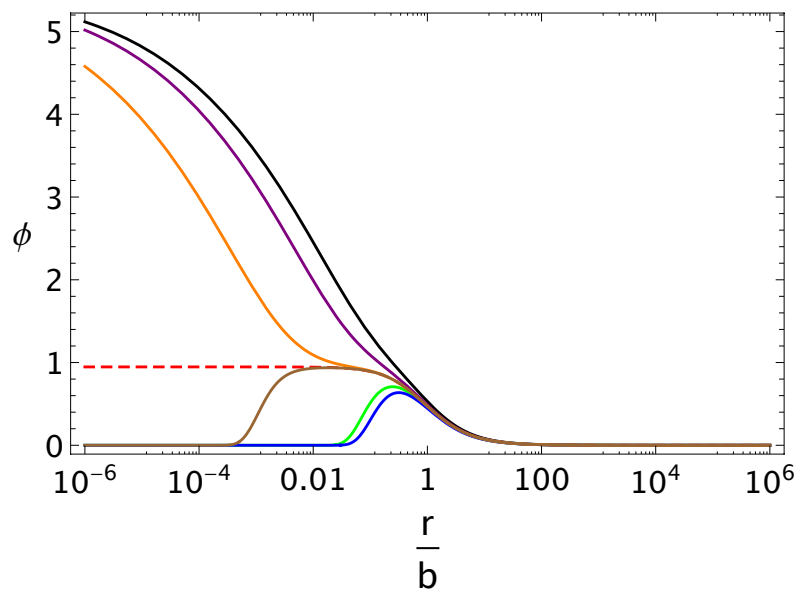
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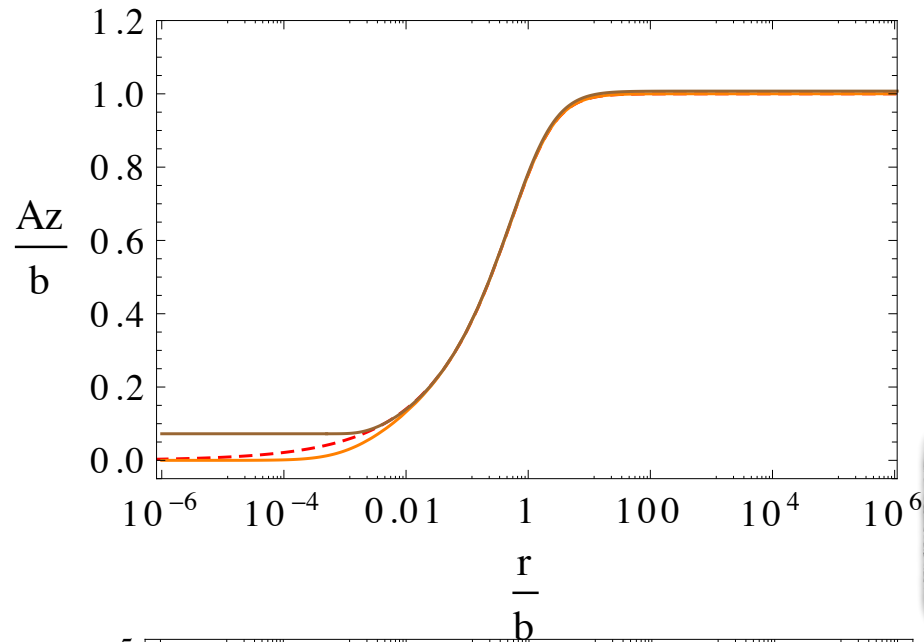
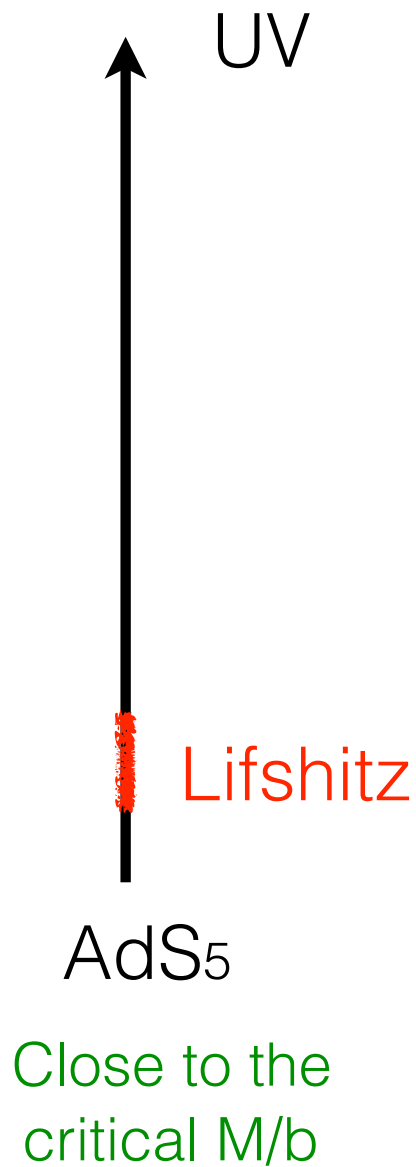
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Holographic WSM



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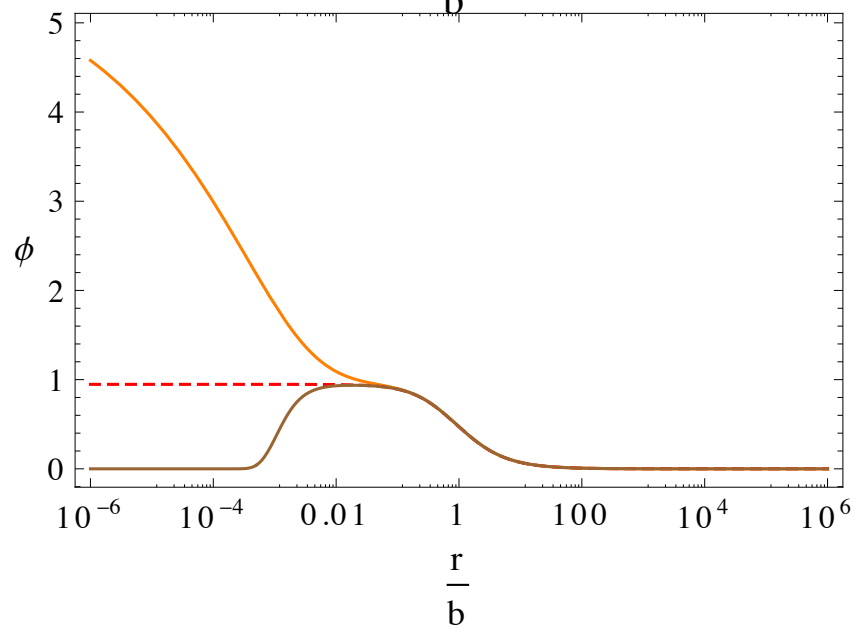
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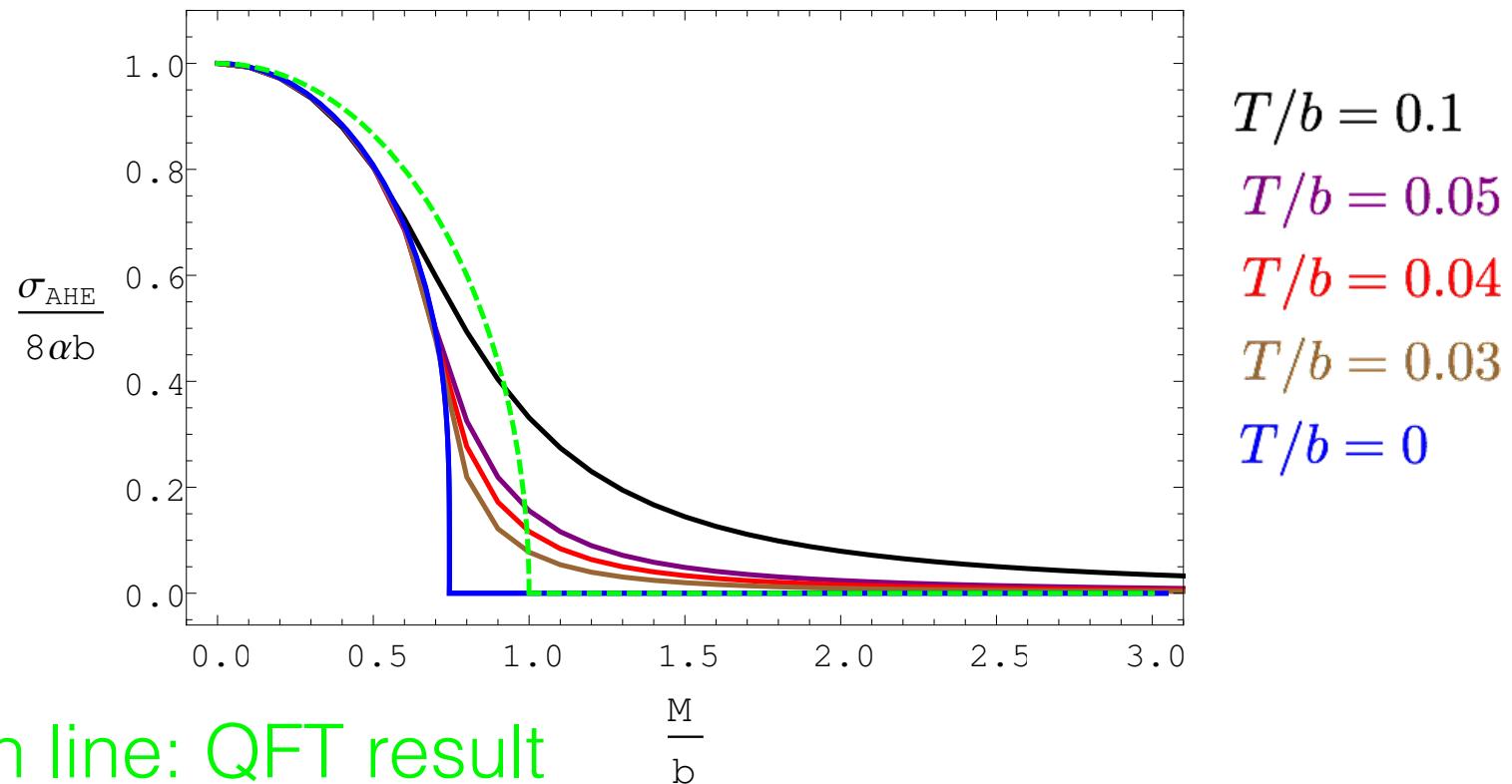
$M/b = 0.856$



Holographic WSM

- Order parameter of topological state of matter: AHE

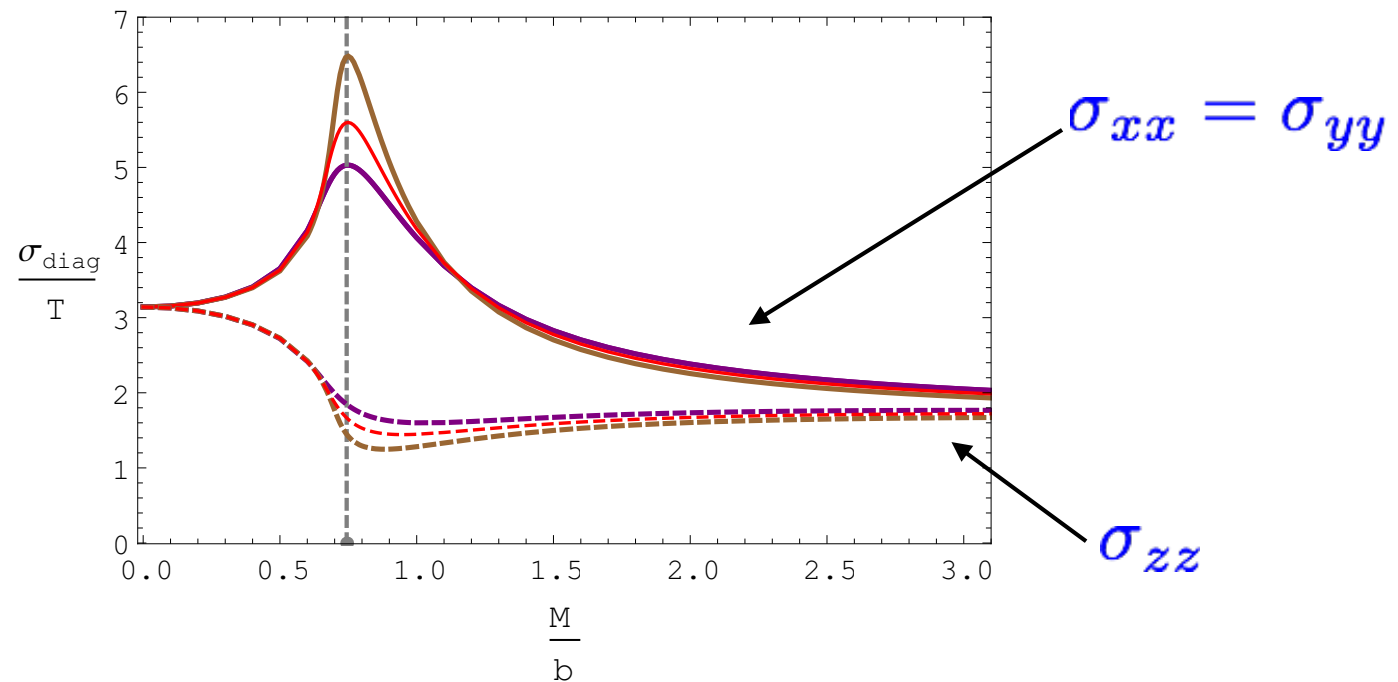
$$\sigma_{\text{AHE}} = 8\alpha A_z^5(0)$$



Holographic WSM

- Diagonal conductivities at $T=0$: $\sigma_{xx} = \sigma_{yy} = \sigma_{zz} = 0$
- Diagonal conductivities at $T>0$:

$T/b = 0.03$
 $T/b = 0.04$
 $T/b = 0.05$



Summary

- A holographic model for Weyl semi-metal using gauge/gravity duality.
- Order parameter = AHE
- Varying M/b , a quantum phase transition between topological non trivial state and topological trivial state.
- Diagonal conductivities (peak/dip behaviour)

Open questions

- Holographic realisation of phase transition between topological non-trivial state and insulator?
- Holographic surface state (Fermi arc)?
- Effective field theory description for the topological phase transition from holographic WSM?
- ...

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Thank You!

Free energy

