

XLIV International Meeting on Fundamental Physics
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Lattice QCD

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Outline

1. Lattice QCD (why and what)
2. Precision Flavour Physics
3. Pushing the Frontiers

Motivation

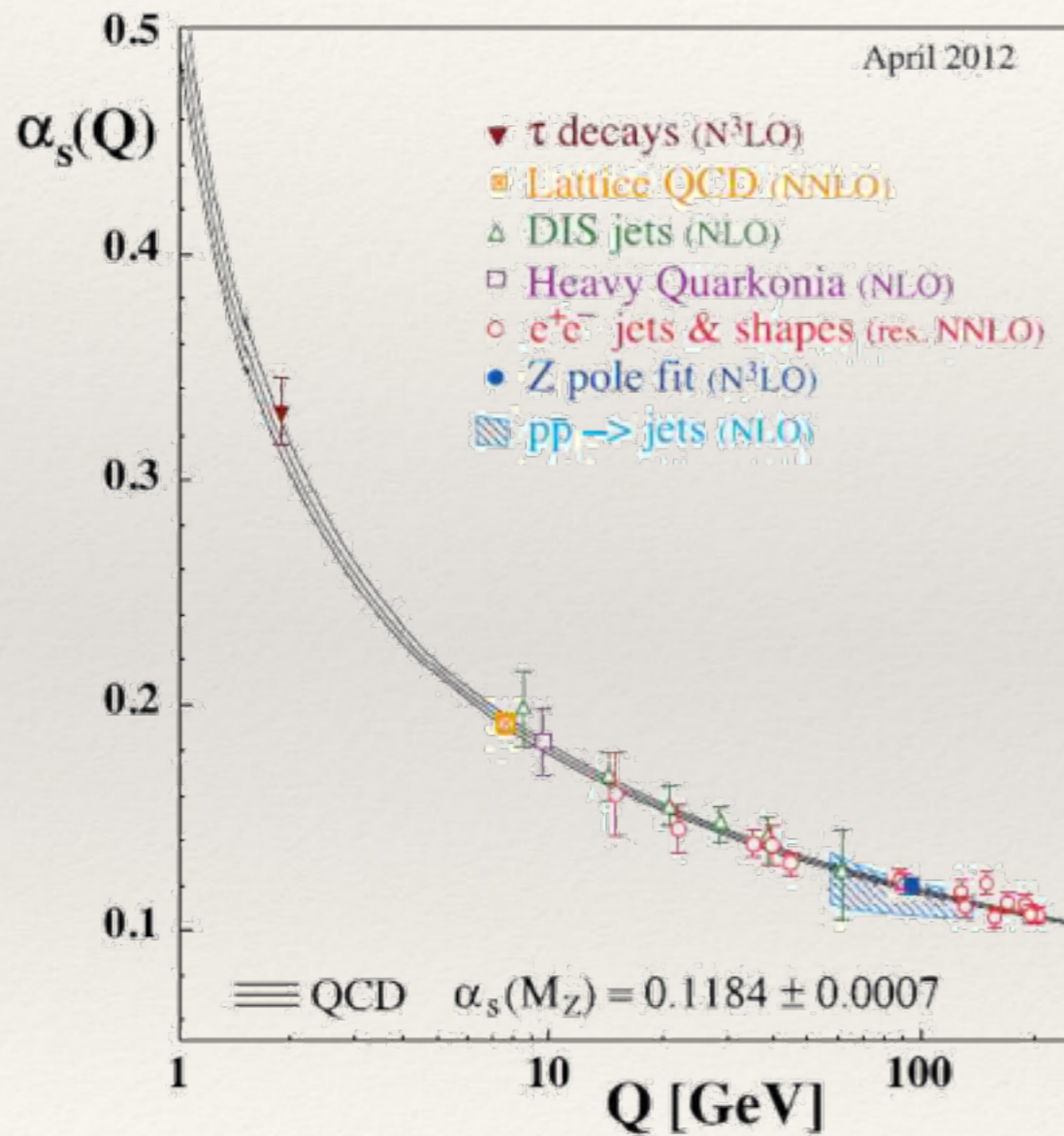
- Standard Model of elementary particle physics describes electromagnetic, weak and strong (QCD) interactions consistently in terms of a renormalisable quantum field theory
- but there is substantial phenomenological evidence that it can't be the whole story: dark matter, CP-violation, ... indicate that there must be sth. else
- despite decades of experimental and theoretical efforts we have not found a smoking gun

Motivation

- searches for new physics: direct vs. indirect search:
 - ‘bump in the spectrum’
 - SM provides correlation between processes
experiment + theory to over-constrain SM
- hadronic (QCD) uncertainties dominating error budget
- lattice QCD can in principle provide the relevant input and is becoming increasingly precise in its predictions

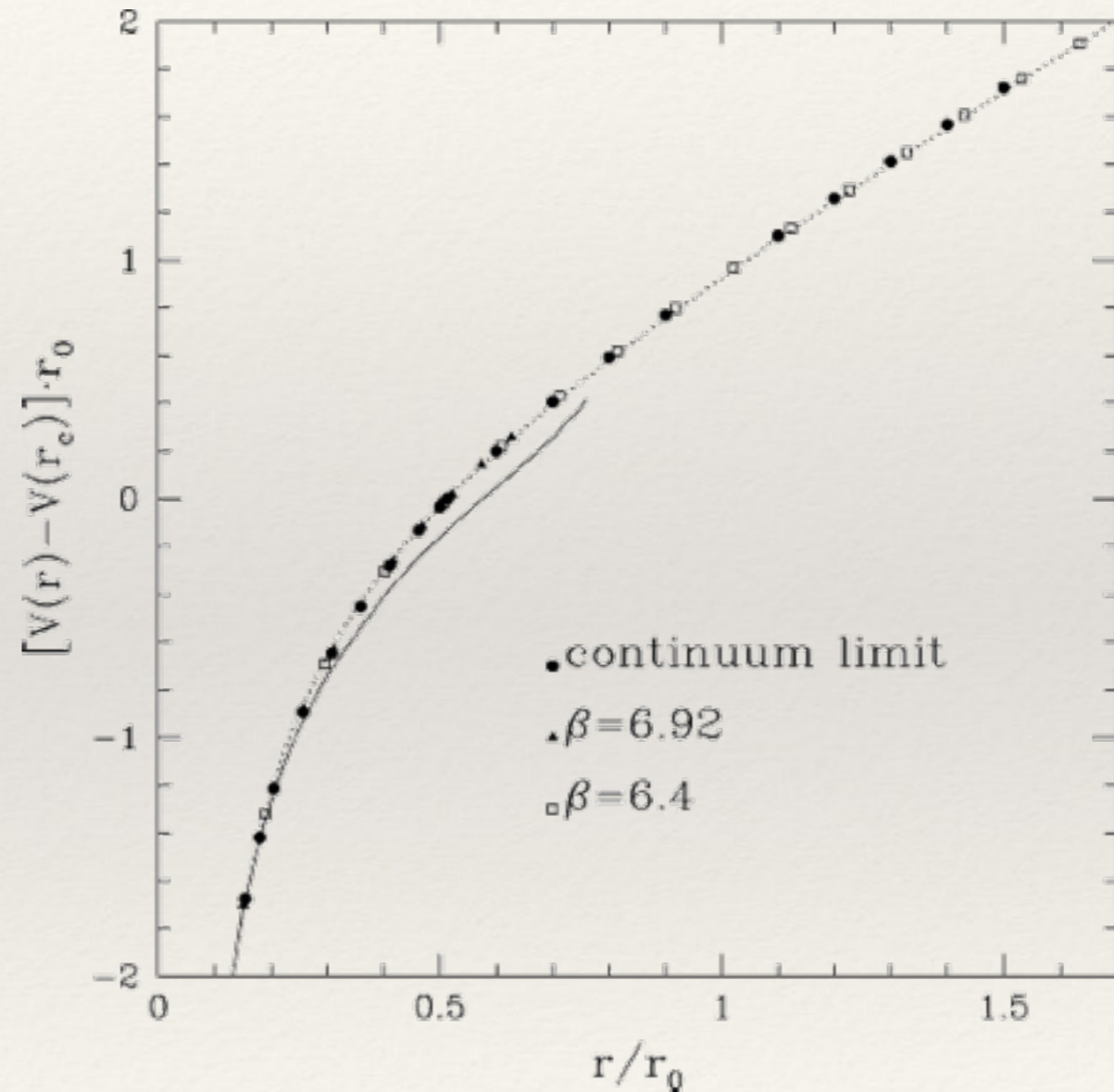
QCD

asymptotic freedom



PDG

confinement



Necco & Sommer NPB 622 (2002)

Lattice QCD

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + \sum_f \bar{\psi}_f (i\gamma^\mu D_\mu - m_f) \psi_f$$

Free parameters:

- gauge coupling $g \rightarrow \alpha_s = g^2/4\pi$
- quark masses $m_f = u, d, s, c, b, t$

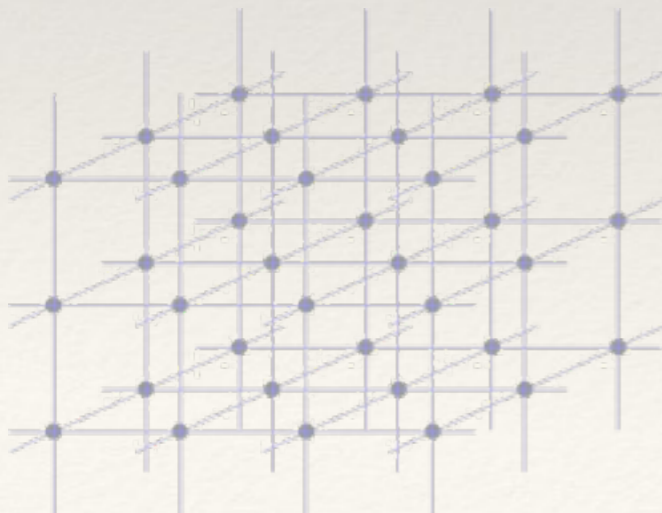
- Lagrangian of massless gluons and *almost massless quarks*
- what experiment sees are bound states, e.g. $m_\pi, m_P \gg m_{u,d}$
- underlying physics non-perturbative

Path integral quantisation:

$$\langle 0|O|0\rangle = \frac{1}{Z} \int \mathcal{D}[U, \psi, \bar{\psi}] O e^{-iS_{\text{lat}}[U, \psi, \bar{\psi}]}$$

$$\langle 0|O|0\rangle = \frac{1}{Z} \int \mathcal{D}[U, \psi, \bar{\psi}] O e^{-S_{\text{lat}}[U, \psi, \bar{\psi}]}$$

Euclidean space-time
Boltzmann factor



finite volume, space-time grid (IR and UV regulators)
 $\propto L^{-1} \quad \propto a^{-1}$

- well defined, finite dimensional Euclidean path integral
- from first principles

Lattice QCD

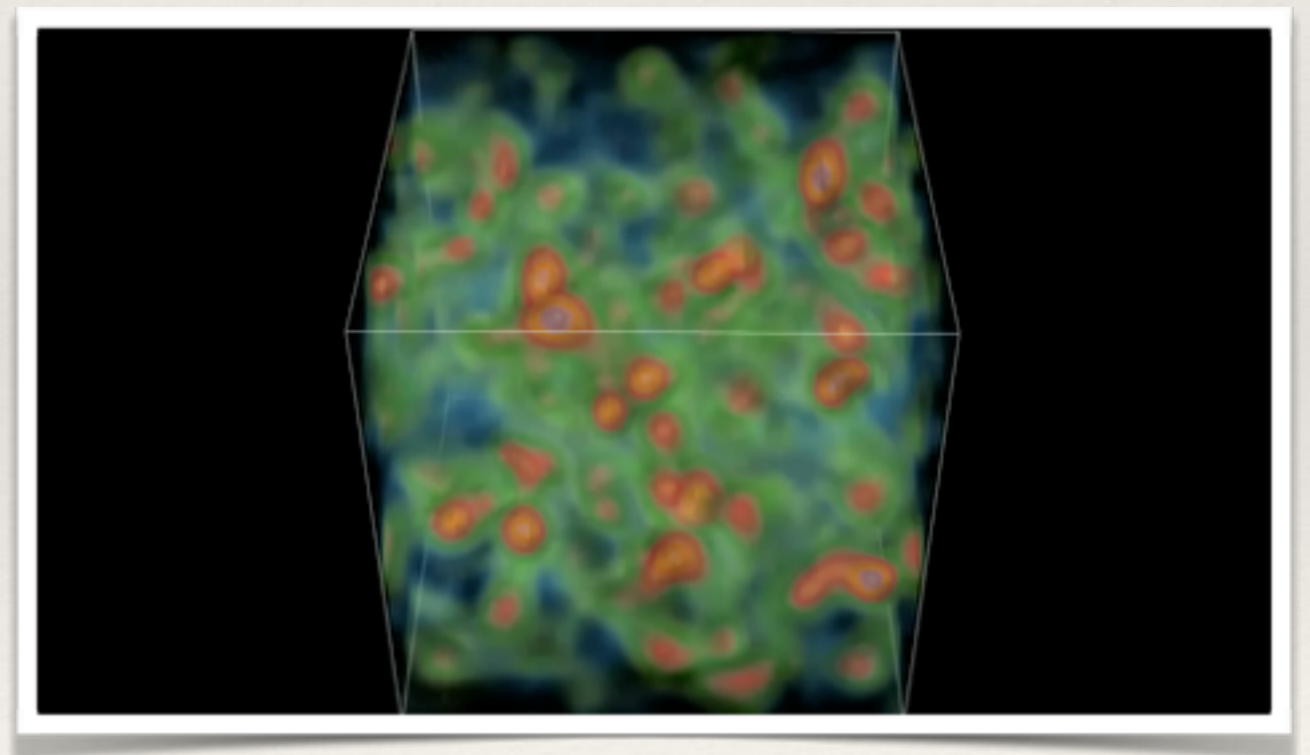
- gauge-invariant regularisation (Wilson 1974)
- naively: replace derivatives by finite differences, integrals by sums
- finite volume lattice path integral still over large number of degrees of freedom $>O(10^{10})$
- Evaluate discretised path integral by means of Markov Chain Monte Carlo on state-of-the-art HPC installations



State of the art of lattice QCD simulations

What we can do

- simulations of QCD with dynamical (sea) u, d, s, c quarks with masses as found in nature $\rightarrow N_f = 2, 2 + 1, 2 + 1 + 1$
- bottom only as valence quark
- cut-off $a^{-1} \leq 4\text{GeV}$
- volume $L \leq 6\text{fm}$

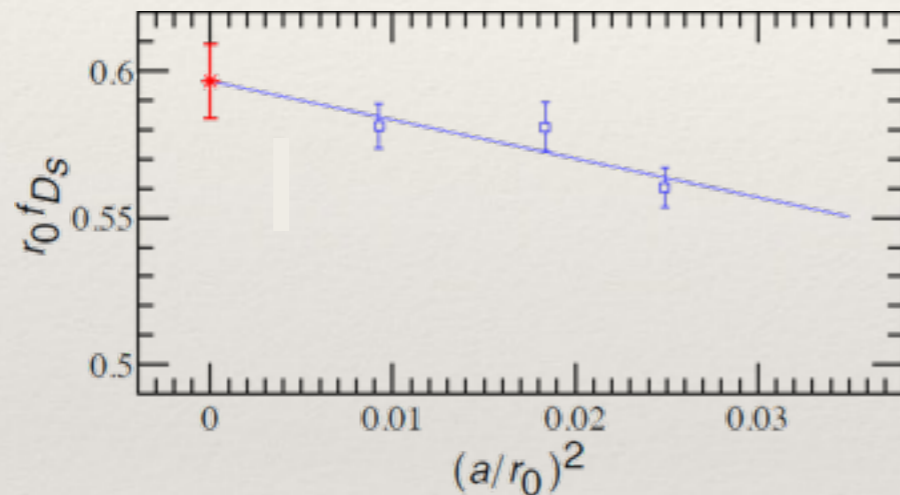


action density of RBC/UKQCD physical point DWF ensemble

lattice - systematics

In practice one needs to control a number of sources of systematic uncertainties, most notably:

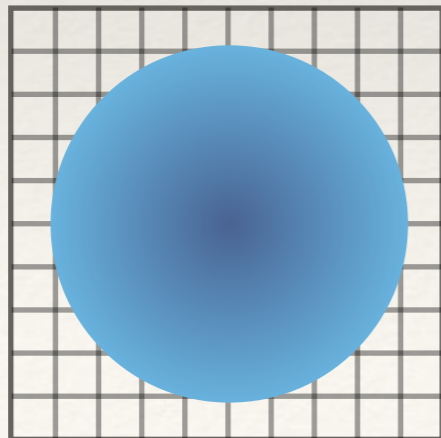
- **discr. errors** (lattice spacing a)



$$S_{\text{eff}} = \int d^4x \{ \mathcal{L}_0(x) + a\mathcal{L}_1(x) + a^2\mathcal{L}_2(x) + \dots \}$$

Symanzik 1982,1983

- **finite volume errors** (box size L)



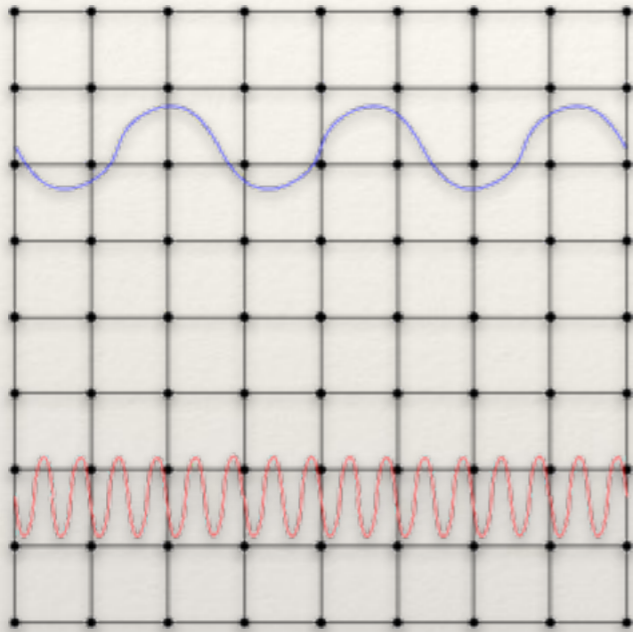
In QCD for simple ME $\propto e^{-m_\pi L} \propto O(1\%)$

more complicated for processes with several hadrons in initial or final state

Lüscher Commun.Math.Phys. 105 (1986) 153-188, Nucl. Phys. B354, 531 (1991)

quite crude, in practice more complicated

lattice size



need to keep

$$a^{-1} \ll \text{relevant scales} \ll L^{-1}$$

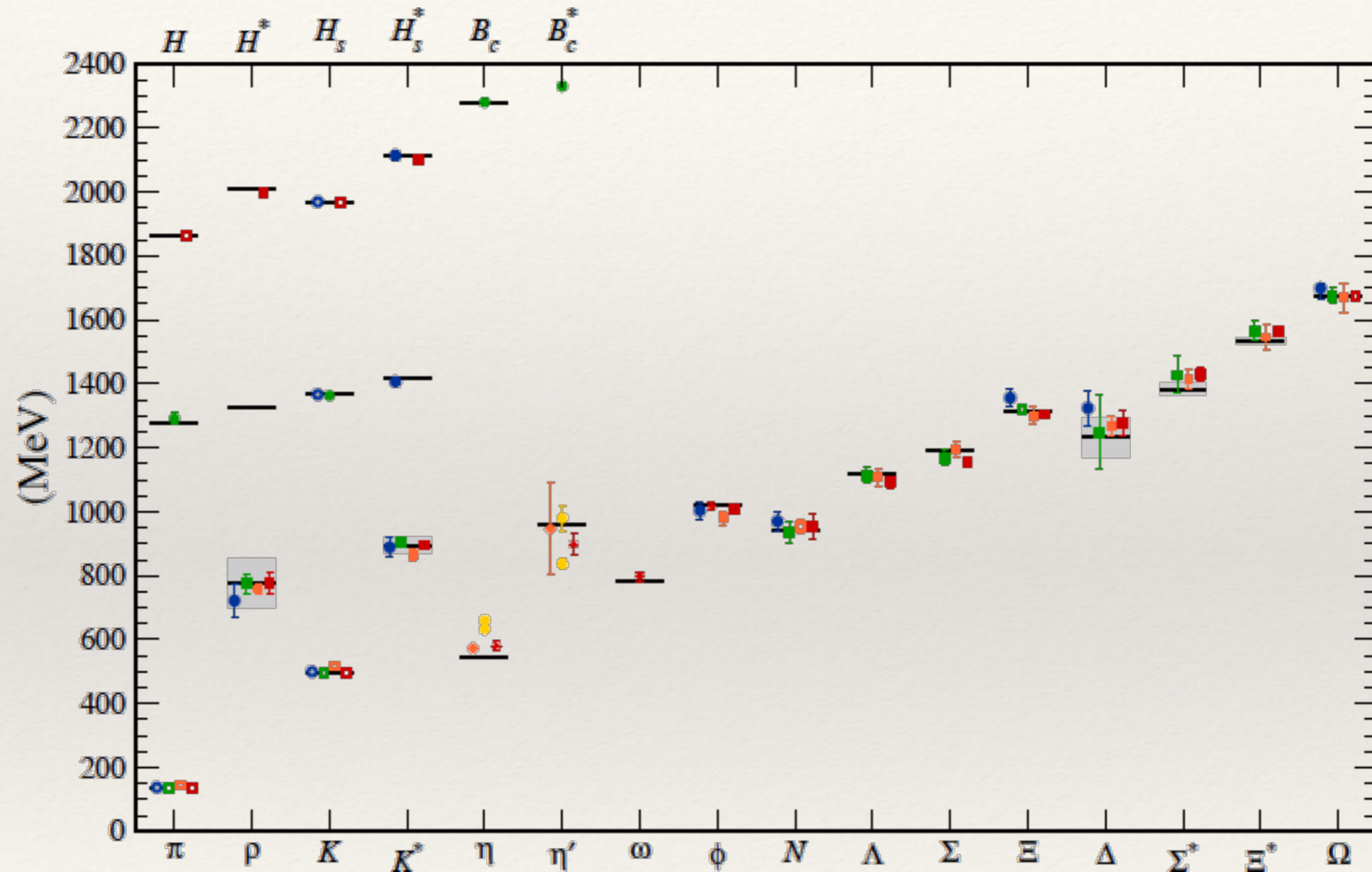
- for $m_\pi=140\text{MeV}$ the constraint for controlled finite volume effects of $m_\pi L \gtrsim 4$ suggests $L \approx 6\text{fm}$
- for charm quarks to be well resolved $am_c < 1$ e.g. a^{-1} larger than $\approx 2.5\text{GeV}$ needed
- lattices with $L/a \gtrsim 80$ needed

Fulfilling all the constraints is just starting to happen

(e.g. first $96^3 \times 192$ have been generated (MILC)) in the meantime most collaborations

- weaken the finite volume effects by simulating unphysical heavy pions
- extrapolate from coarser lattices relying on assumptions for functional form of cutoff effects

benchmark - the hadron spectrum



Kronfeld, Ann. Rev. of Nucl. Part. Sci 2012 62

Lattice pheno - what's possible

- **Standard:**

- meson ME with single incoming and / or outgoing pseudo-scalar states
 $\pi, K, D_{(s)}, B_{(s)} \rightarrow \text{QCD} - \text{vacuum}, \pi \rightarrow \pi, K \rightarrow \pi, D \rightarrow K, B \rightarrow \pi, \dots, B_K, (B_D), B_B$
- QCD parameters: quark masses, strong coupling constant
- meson / baryon spectroscopy of stable (in QCD) states

- **Challenging:**

- two initial / final hadronic states, one channel $\pi\pi \rightarrow \pi\pi, K\pi \rightarrow K\pi, K \rightarrow \pi\pi, \dots$
- elm. effects in spectra

- **Very challenging - new ideas needed/no clue:**

- multi-channel final states (hadronic D, B) (e.g. Hansen, Sharpe PRD86, 016007 (2012))
- long-distance contributions in e.g. rare Kaon decays, K, D -mixing (Bai et al. PRD113 2014)
- transition MEs with unstable in / out states (Briceño et al. arXiv:1406.5965)
- elm effects in matrix elements

Quark Flavour Physics

100 GeV

2 GeV

300 MeV

$\langle \text{had}_f | H_W | \text{had}_i \rangle$

$\langle 0 | H_W | \text{had}_i \rangle$

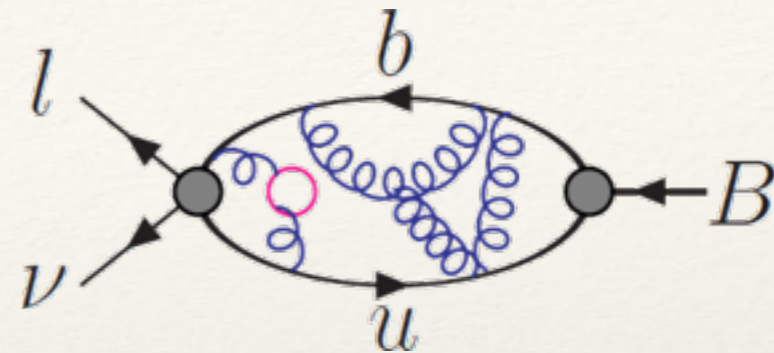
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

3x3 unitary matrix
4 unknown parameters

- quark mixing
- CP-violation (one complex phase)
- constraints on SM processes
- high energy reach
- inconsistencies -> failure of the SM?

Quark Flavour Physics

e.g tree level leptonic B decay:



Assumed factorisation: $\Gamma_{\text{exp.}} \stackrel{???}{=} V_{\text{CKM}}(\text{WEAK})(\text{EM})(\text{STRONG})$
currently Chiral EFT

$$\underbrace{\Gamma(B \rightarrow l\nu_l)}_{\text{theory}} = \underbrace{|V_{ub}|^2}_{\text{output}} \underbrace{\frac{m_B}{8\pi} G_F^2 m_l^2 \left(1 - \frac{m_l^2}{m_B^2}\right)^2}_{\text{theory prediction}} \underbrace{f_B^2}_{\text{theory prediction}}$$

Experimental measurement + theory prediction allows for extraction of CKM MEs

Flavour Physics

Determine CKM elements \longleftrightarrow (indirect) test of SM:

- over-determine elements of V_{CKM} and check consistency of CKM paradigm
- unitarity tests:
 - rows and columns are (in SM) complex unit vectors
 - rows (columns) are orthogonal to other rows (columns)violation of unitarity would indicate non-SM physics

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \stackrel{?}{=} 1$$

row-test

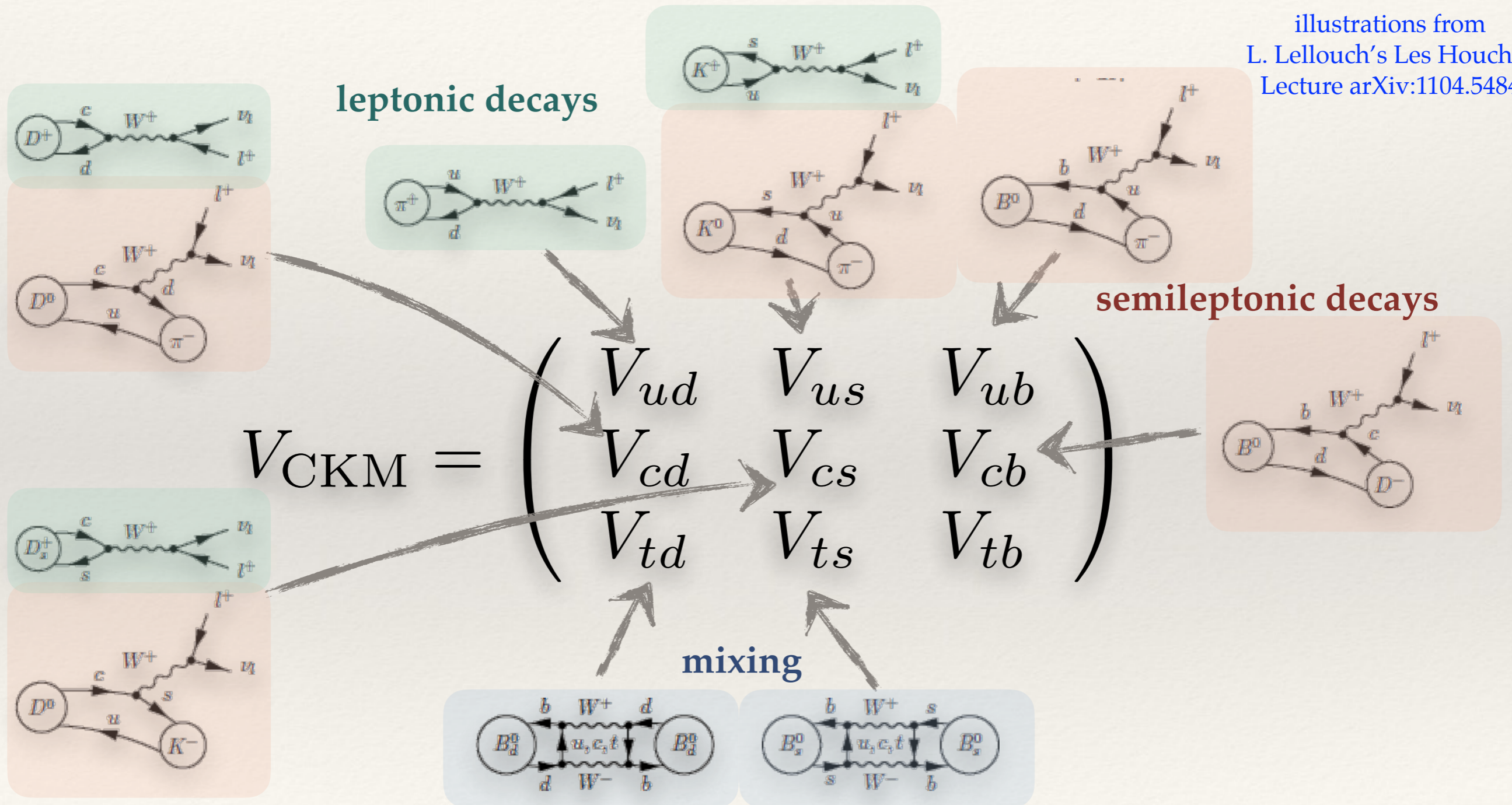
$$\sum_{U=u,c,t} V_{Ud}V_{Ub}^* = 0$$

triangle-test

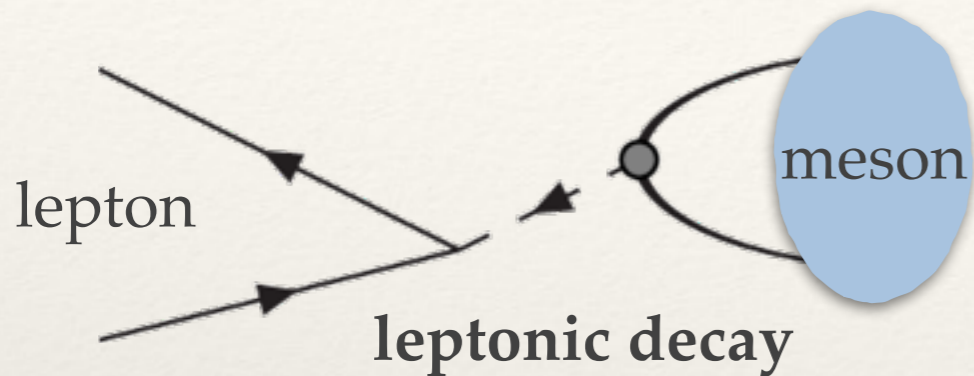
- in which channels is there still room for new physics
 - how much new physics would be compatible with measurements
 - what would be the properties of new physics

Lattice flavour physics and CKM

illustrations from
L. Lellouch's Les Houches
Lecture arXiv:1104.5484



“tree” kaon/pion decays

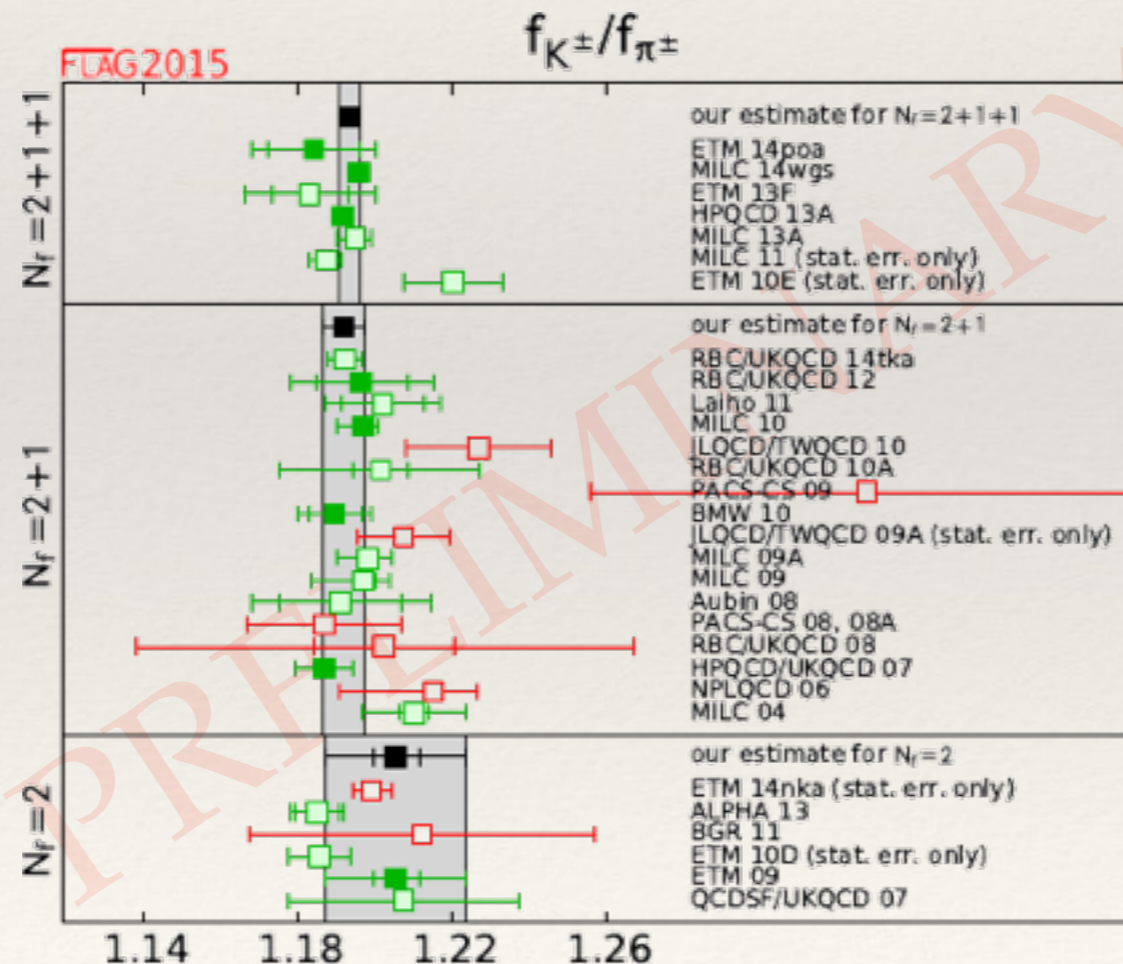


$$\Gamma(K \rightarrow \mu \bar{\nu}_\mu) = \frac{G_F^2}{8\pi} f_K^2 m_\mu^2 m_K \left(1 - \frac{m_\mu^2}{m_K^2}\right)^2 |V_{us}|^2$$

$$\langle 0 | \bar{s} / \bar{d} \gamma_\mu \gamma_5 u | K / \pi(p) \rangle = i f_{K/\pi} p_\mu$$

$$\frac{\Gamma(K \rightarrow \mu \bar{\nu}_\mu)}{\Gamma(\pi \rightarrow \mu \bar{\nu}_\mu)} = \frac{|V_{us}|^2}{|V_{ud}|^2} \left(\frac{f_K}{f_\pi}\right)^2 \frac{m_K (1 - m_\mu^2/m_K^2)^2}{m_\pi (1 - m_\mu^2/m_\pi^2)^2} \times 0.9930(35)$$

Marciano, Phys.Rev.Lett. 2004



1‰!!!

Standard calculations and results - FLAG

Flavour Lattice Averaging Group

“What’s currently the best lattice value for a particular quantity?”

FLAG-1 (Eur. Phys. J. C71 (2011) 1695)

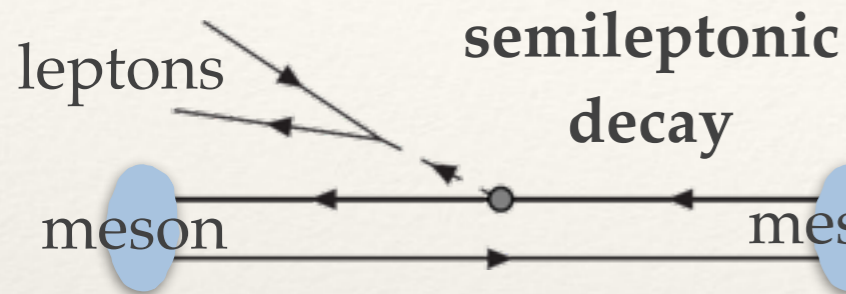
FLAG-2 (<http://itpwiki.unibe.ch/flag/>, Eur.Phys.J. C74 (2014) 2890)

- quantities:
 - FLAG-1: $m_{u,d}, m_s, f_K / f_\pi, f_+^{K\pi}(0), B_K, SU(2)$ and $SU(3)$ LECs
 - FLAG-2: FLAG-1 + $\alpha_s, f_{D(s)}, f_{B(s)}, B_{B(s)}, B, D$
- summary of results
 - evaluation according to FLAG quality criteria (colour coding)
 - averages of best values where possible
 - detailed summary of properties of individual simulations

FLAG-3 - working on it

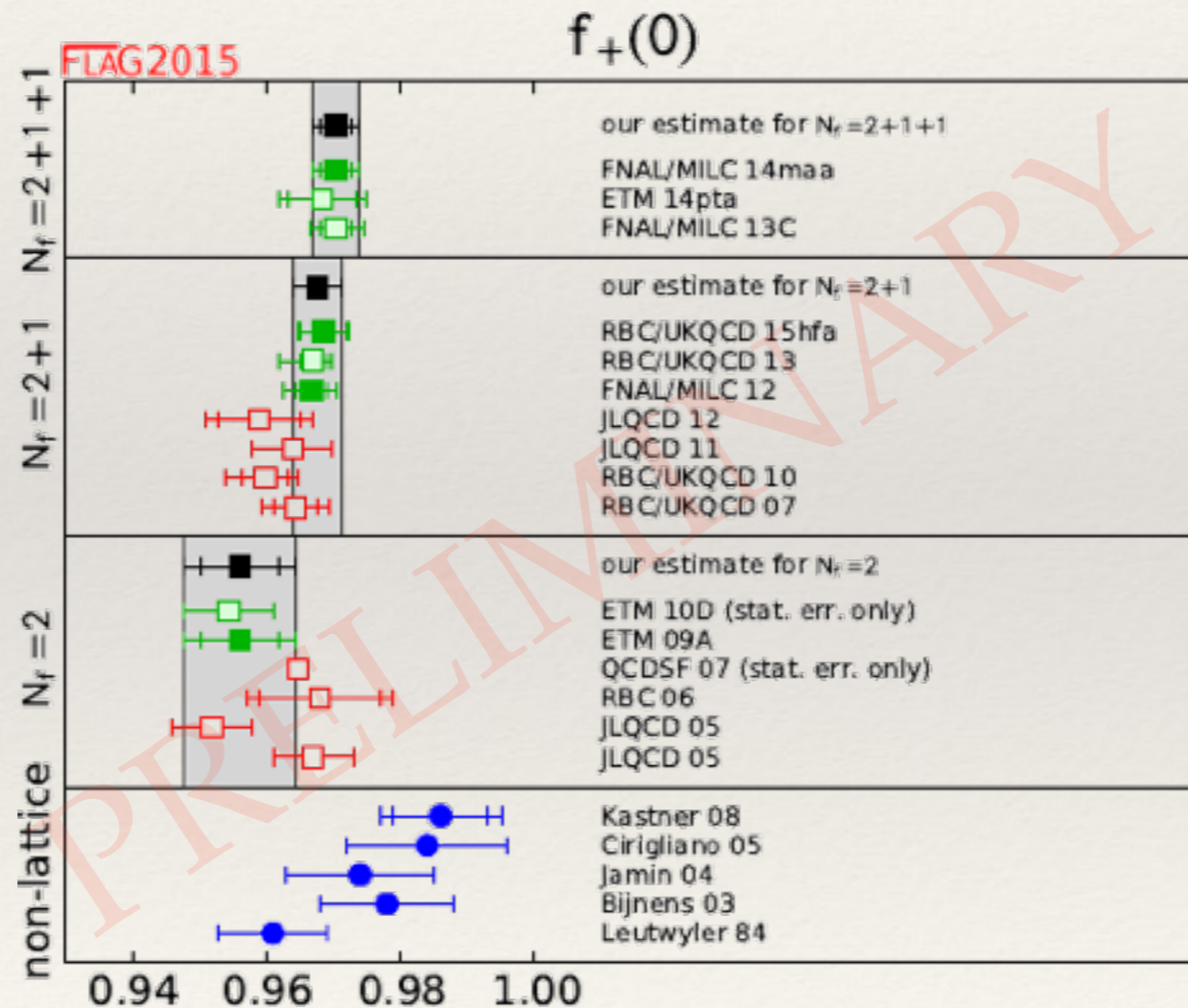
- FLAG-3: FLAG-2 + heavy quark masses

“tree” kaon/pion decays



$$\Gamma_{K \rightarrow \pi l \nu} = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^2} S_{EW} (1 + \Delta_{SU(2)} + \Delta_{EM})^2 I |f_+^{K\pi}(0)|^2 |V_{us}|^2$$

$$\langle \pi(p_\pi) | V_\mu(0) | K(p_K) \rangle = f_+^{K\pi}(q^2) (p_K + p_\pi)_\mu + f_-^{K\pi}(q^2) (p_K - p_\pi)_\mu$$



3‰!!!

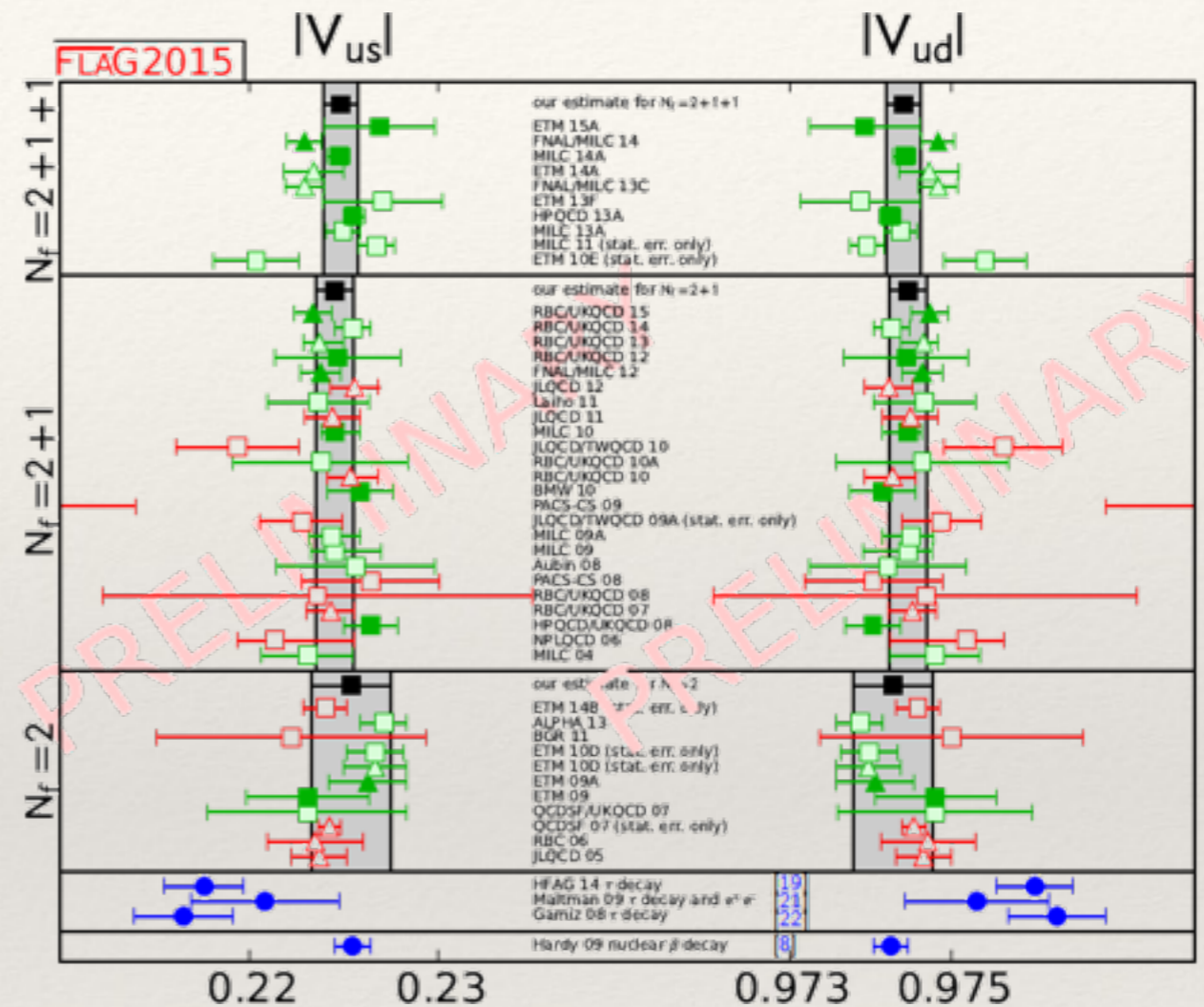
$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \approx |V_{ud}|^2 + |V_{us}|^2 \stackrel{?}{=} 1$$

Experimental results:

$$|V_{us}| f_+(0) = 0.2163(5)$$

$$\frac{f_K}{f_\pi} \frac{|V_{us}|}{|V_{ud}|} = 0.2758(5)$$

FLAVIA Kaon WG EPJ C 69, 399-424 (2010)
KTeV, Istra, KLOE



Numerical results from FLAG2, illustrations (preliminary) from FLAG3

First row unitarity:

- $f_+^{K\pi}(0)$ and $|V_{ud}|$ from experiment
- f_K/f_π and $|V_{ud}|$ from experiment
- $f_+^{K\pi}(0)$ and f_K/f_π from lattice

	$f_+(0), V_{ud} $	$f_K/f_\pi, V_{ud} $	combined
$N_f=2+1$	0.9993(5)	1.0000(6)	0.987(10)
$N_f=2$	1.0004(10)	0.9989(16)	1.029(35)

Eur.Phys.J. C74 (2014) 2890
[arXiv:1310.8555](https://arxiv.org/abs/1310.8555)

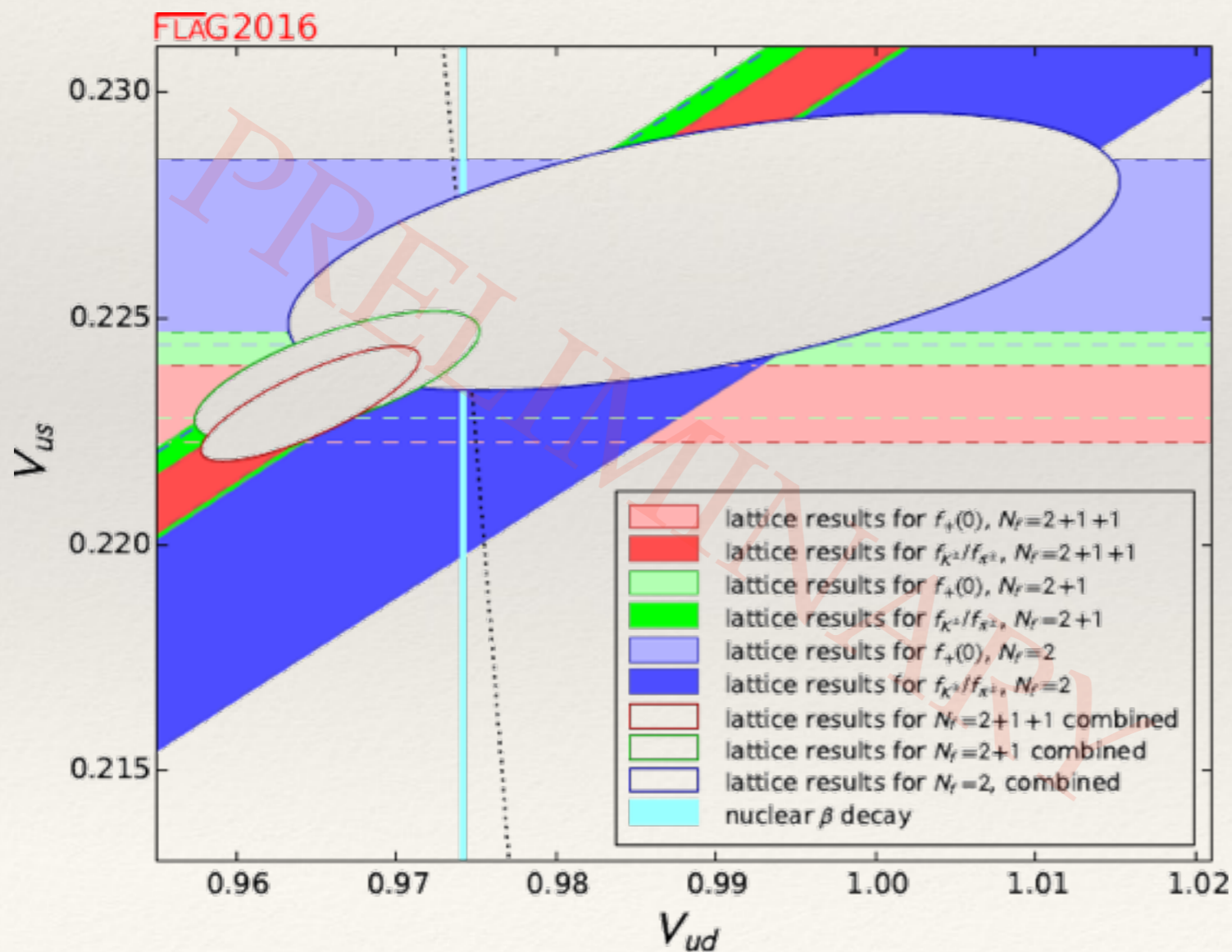
$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 \approx |V_{ud}|^2 + |V_{us}|^2 \stackrel{?}{=} 1$$

FLAG V_{us} Working Group (Boyle, Kaneko, Simula)

$$|V_{us}| f_+^{K^0 \pi^-}(0) = 0.2163(5)$$

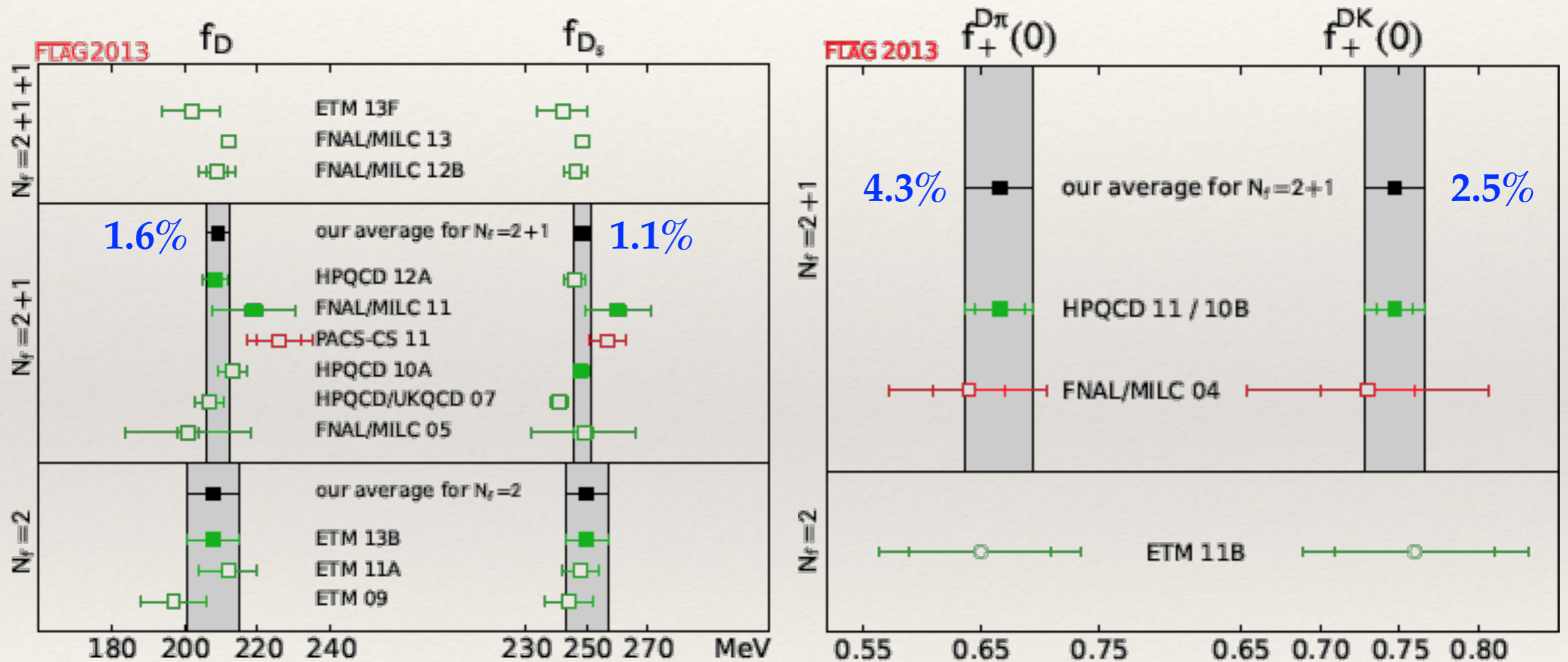
$$\frac{f_{K^+}}{f_{\pi^+}} \frac{|V_{us}|}{|V_{ud}|} = 0.2758(5)$$

FLAVIANet Kaon WG
EPJ C 69, 399-424 (2010)
[arXiv:1005.2323](https://arxiv.org/abs/1005.2323)

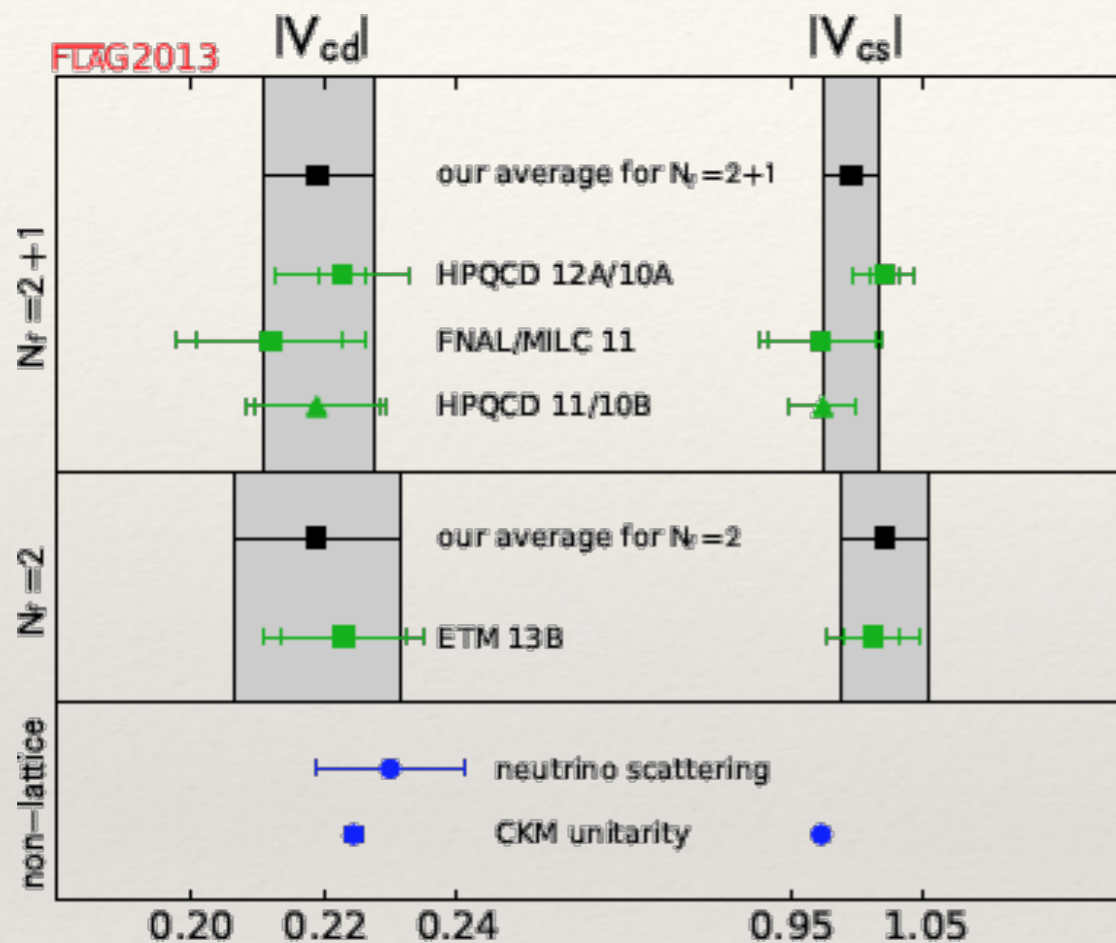


high precision test of
SM unitarity - no worrisome
tension at sub-percent-level
precision

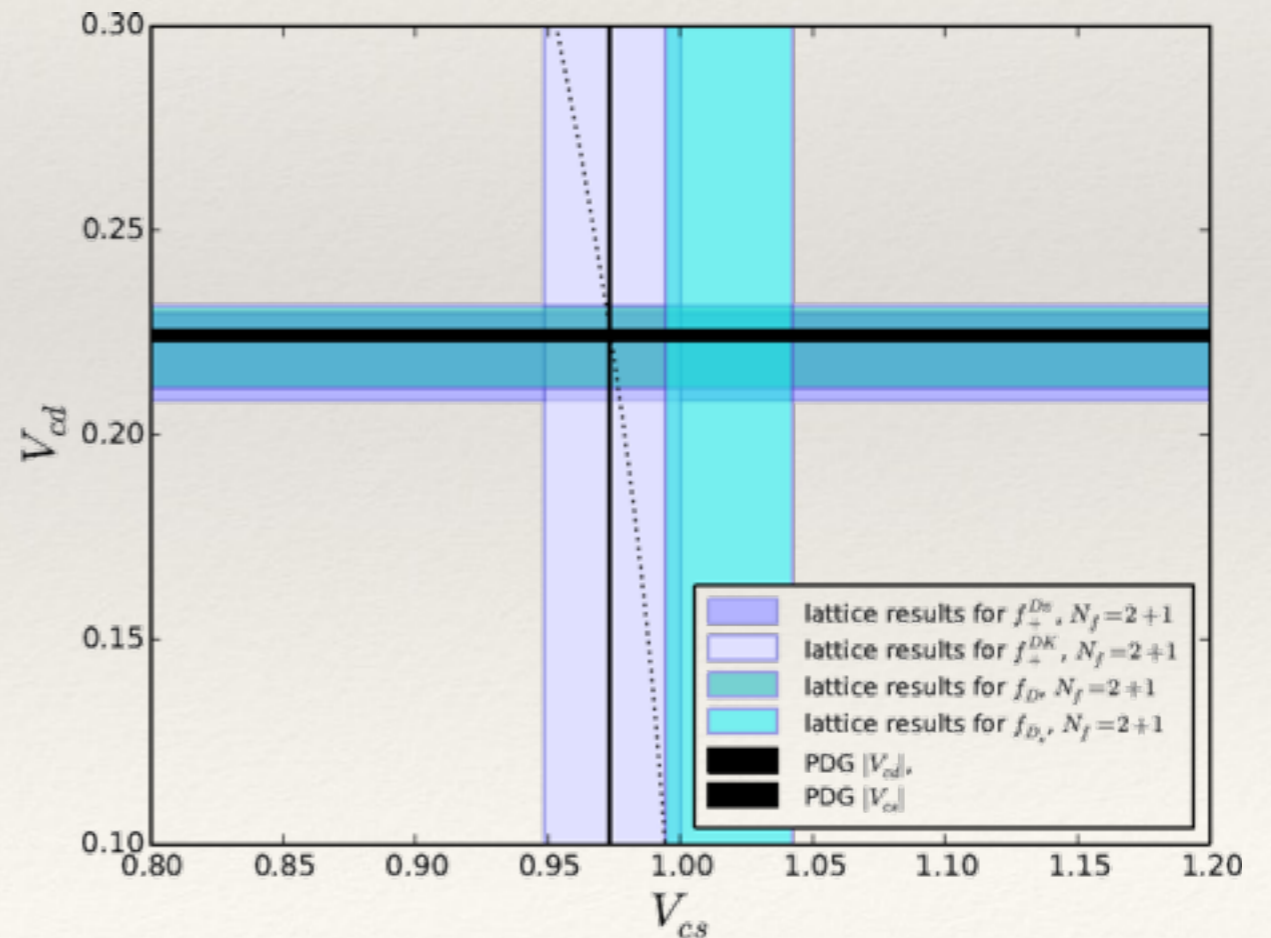
Leptonic $D_{(s)}$ meson decays



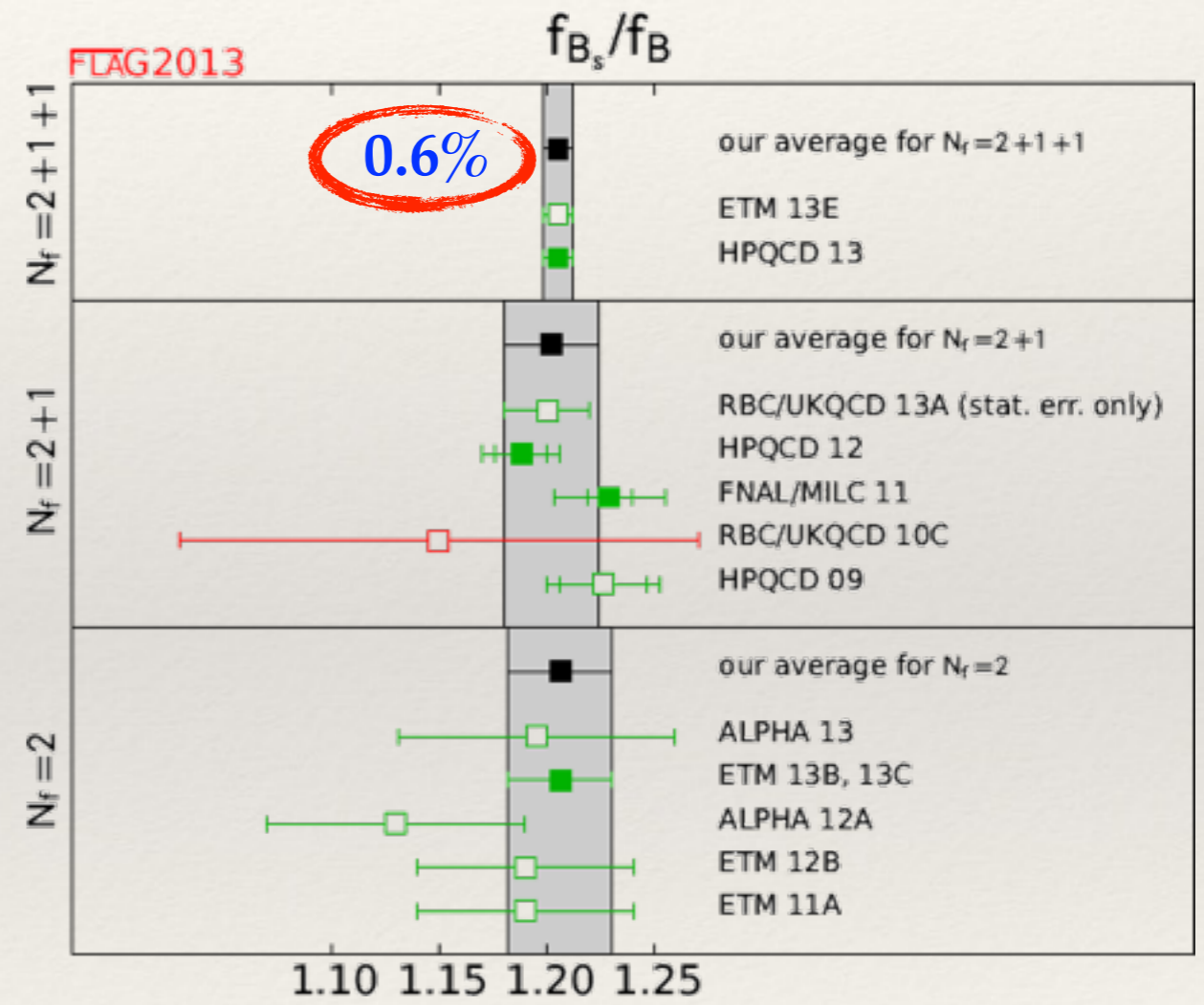
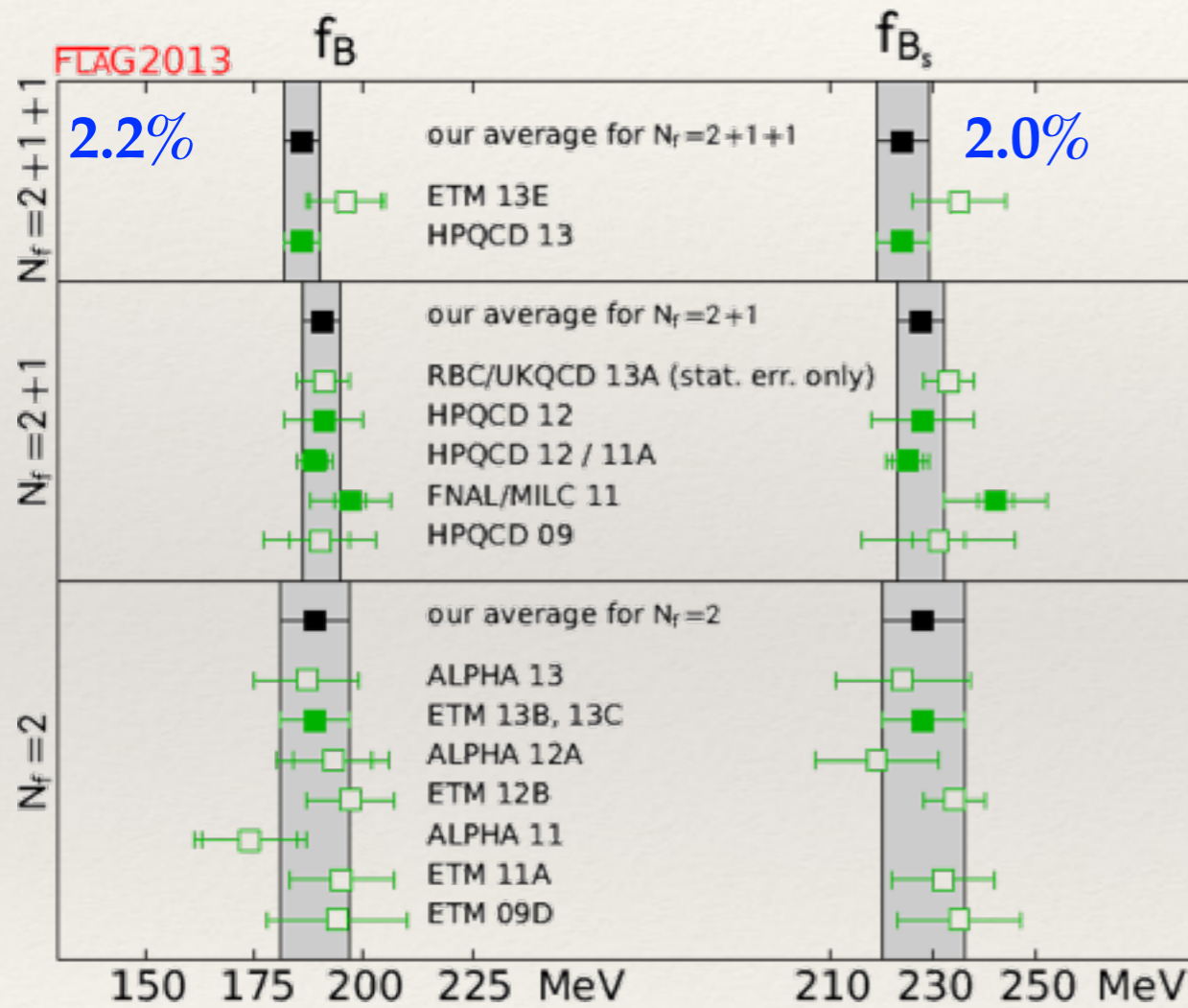
Results for $|V_{cd}|$ and $|V_{cs}|$



- $|V_{cs}|$ from leptonic decays is slightly larger than from semileptonic decays
- $|V_{cs}|$ from leptonic decays is at tension with CKM-unitarity by 1.9σ (\rightarrow HPQCD)



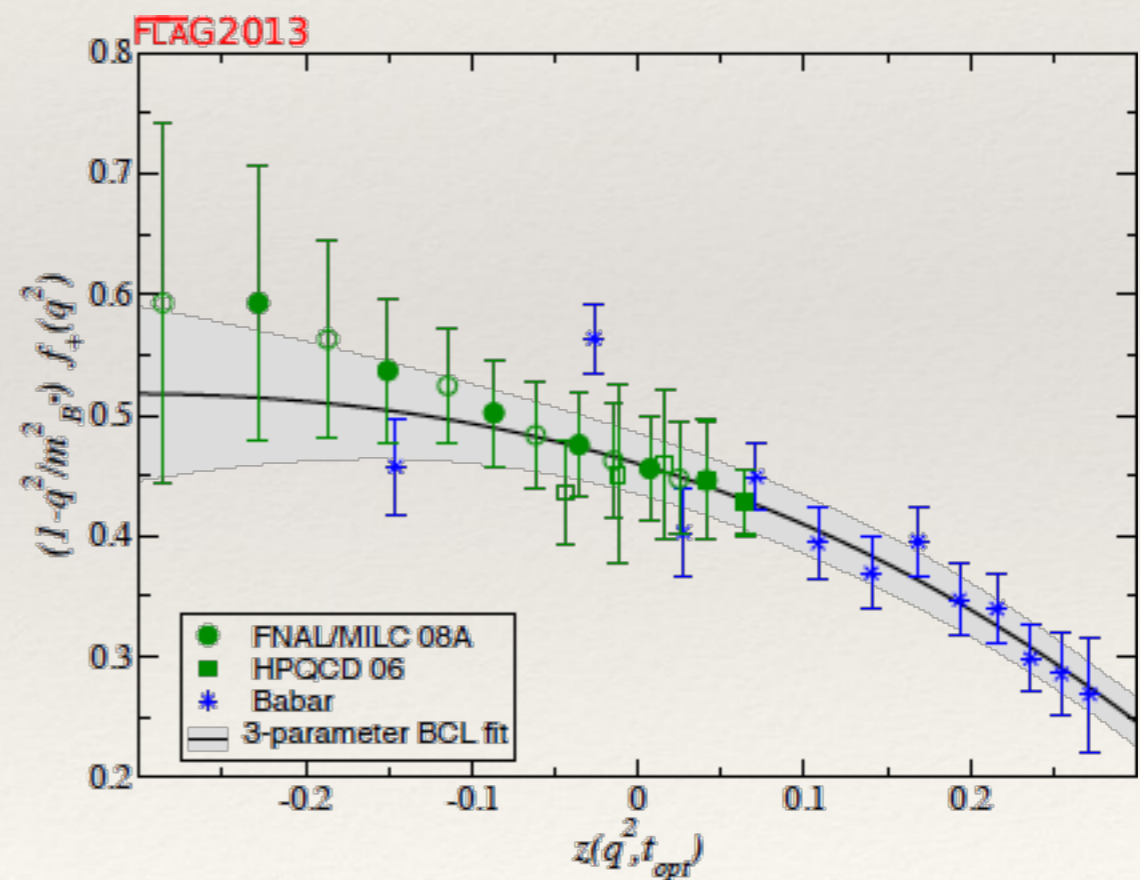
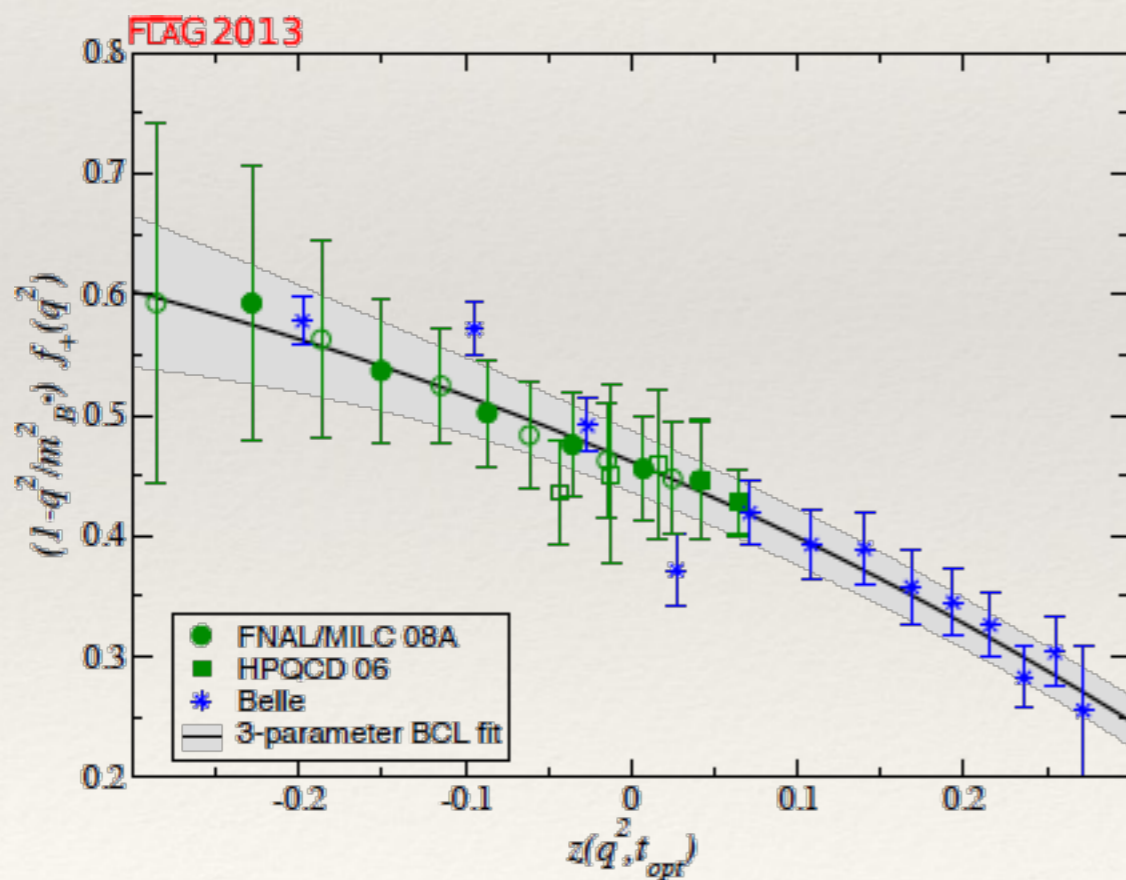
Leptonic beauty decays



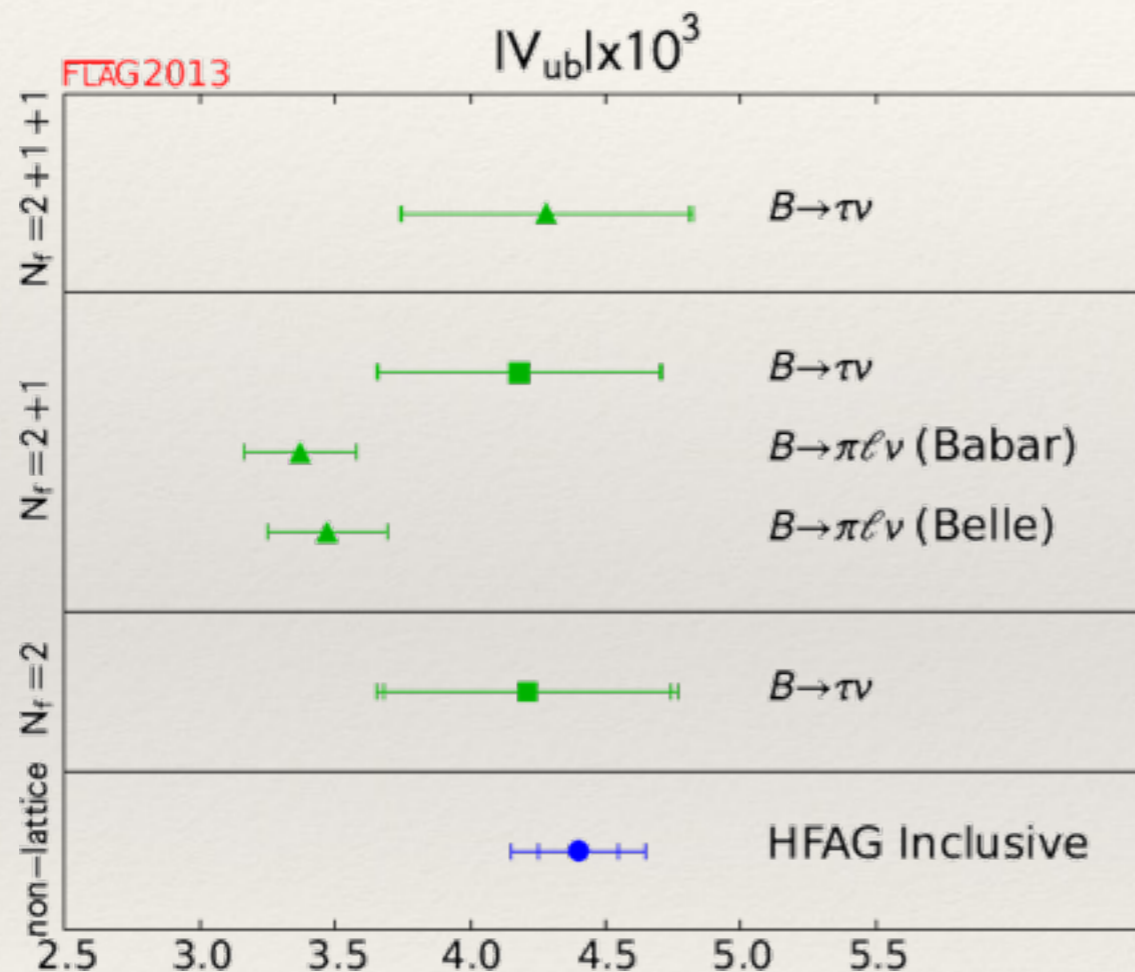
Semileptonic beauty decays

Kinematical reach limited in lattice QCD → extract value of V_{ub} from simultaneous analysis of exp. and lattice data

$$q^2 = (E_B - E_{\text{light}})^2 - (\vec{p}_B - \vec{p}_{\text{light}})^2 \quad \vec{p} = \frac{2\pi}{L} \vec{n}$$

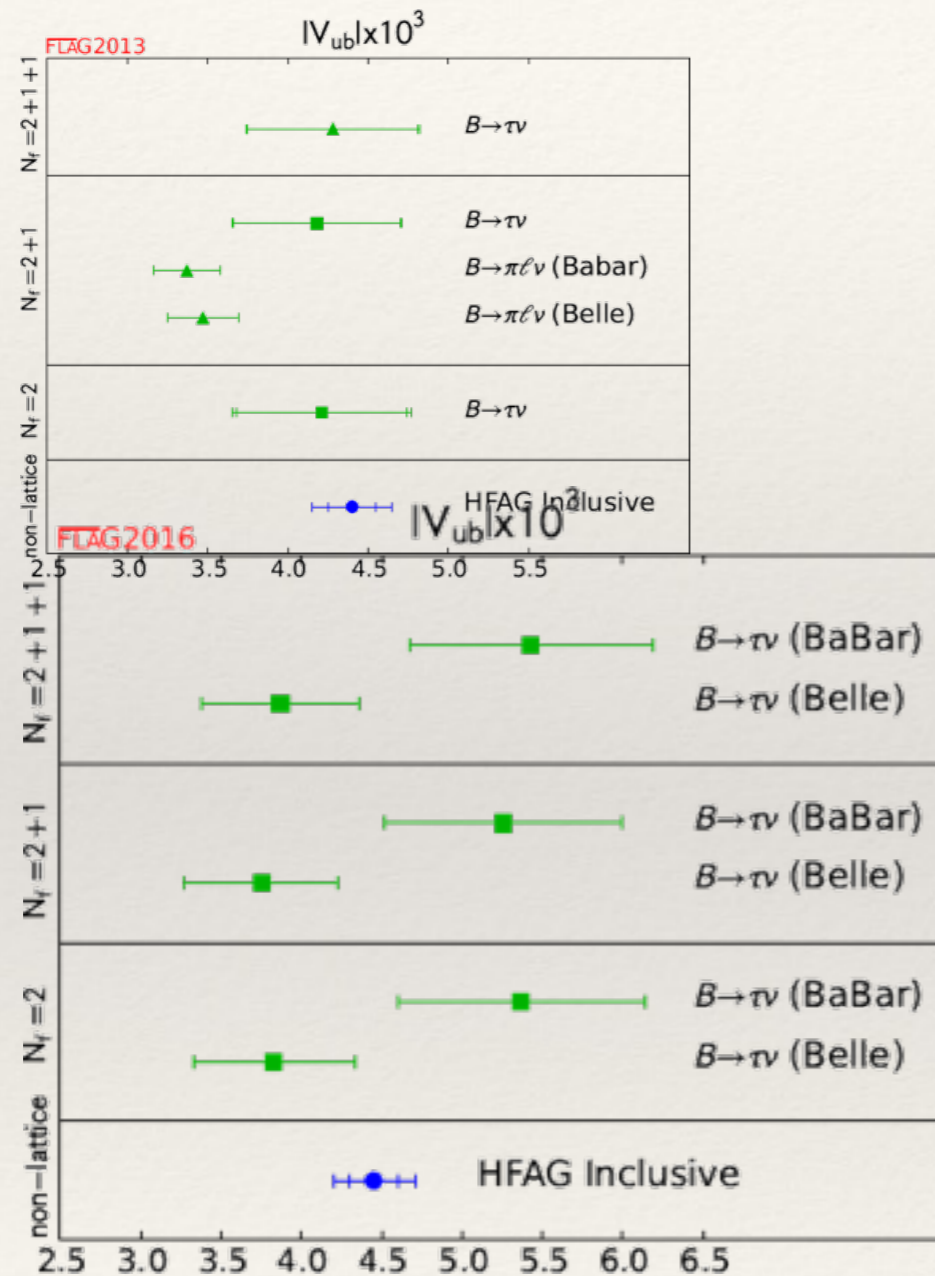


Results for $|V_{ub}|$



- tension between betw. incl. and excl. semilept. decays?
- lept. decay lies in between and agrees with both at 1.5σ

Results for $|V_{ub}|$



- slight tension in exp data
- looking forward to Belle II

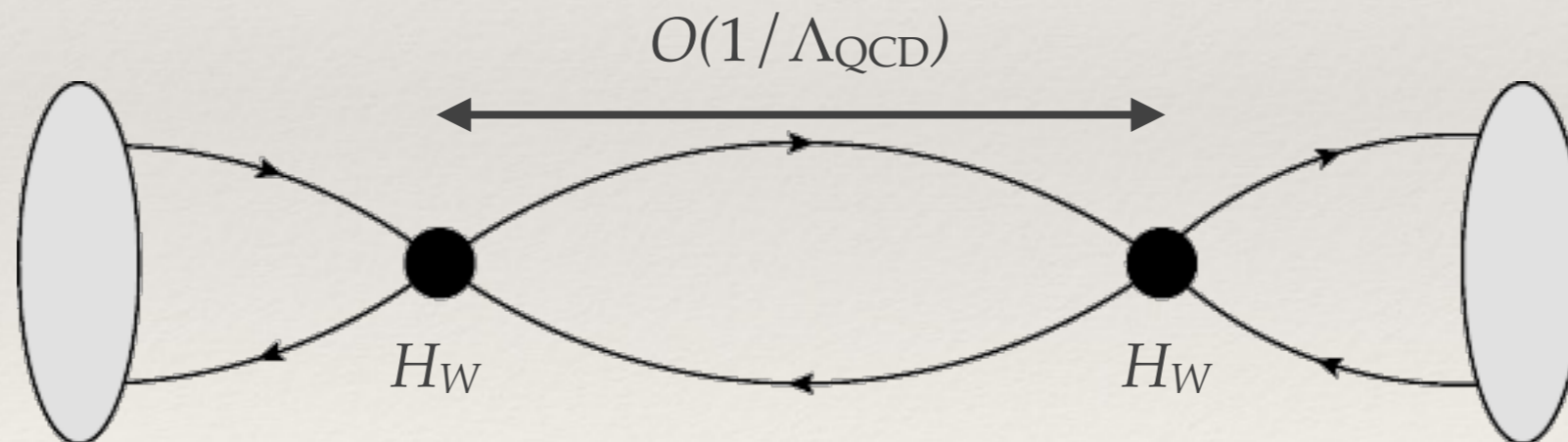
Beyond precision lattice QCD

Go beyond factorisation

$$\Gamma_{\text{exp.}} \stackrel{???}{=} V_{\text{CKM}}(\text{WEAK})(\text{EM})(\text{STRONG})$$

treat jointly in lattice QCD+QED

Go beyond short distance physics



Including QED in meson decay MEs

- Precision on MEs such that EM and strong isospin effects important
remember: so far mostly only QCD ($m_l=m_u=m_d, \alpha_{EM}=0$)
- we should go beyond EFT treatment (e.g. replace ChPT estimates)
- need to understand how this can be done conceptually
- already many results for spectroscopic quantities but not for matrix elements
- finite size effects with photons pose a substantial problem

Isospin corrections are important

- e.g. $K \rightarrow \pi l \nu$: $\Gamma_{K \rightarrow \pi l \nu} = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^2} S_{EW} (1 + \Delta_{SU(2)} + \Delta_{EM})^2 I |f_+^{K\pi}(0)|^2 |V_{us}|^2$

3%

Kastner & Neufeld

Eur. Phys. J. C 57 (2008) 541

- precision now such that corrections need to be improved:

Mode	$V_{us} f_+(0)$	% err	Approx contrib to % err			
			BR	τ	Δ	I
$K_{L e3}$	0.2163(6)	0.26	0.09	0.20	0.11	0.05
$K_{L \mu3}$	0.2166(6)	0.28	0.15	0.18	0.11	0.06
$K_{S e3}$	0.2155(13)	0.61	0.60	0.02	0.11	0.05
K_{e3}^\pm	0.2172(8)	0.36	0.27	0.06	0.23	0.05
$K_{\mu3}^\pm$	0.2170(11)	0.51	0.45	0.06	0.23	0.06

Moulson@CKM 2014

[arXiv:1411.5252](https://arxiv.org/abs/1411.5252)

QCD+QED

Action:

$$S[U, A, \bar{\psi}, \psi] = S_g[U; g] + S_\gamma[A] + \sum_f \bar{\psi}_f D[U, A; e, q_f, m_f] \psi_f$$
$$S_\gamma^{naive} = -\frac{a^4}{4} \sum_{\mu, \nu, x} (\partial_\mu A_{\nu, x} - \partial_\nu A_{\mu, x})^2$$

- MC simulation of discretised theory
- QCD has a mass gap \rightarrow finite volume effects $\propto e^{-m_\pi L}$ (for simple MEs)
- photon is massless and interacts over long range
 \rightarrow power-like finite volume effects $\propto 1/L, 1/L^2, \dots$ from exchange of photon around torus
- sufficiently large volumes currently not feasible, so use effective field theory to subtract finite volume effects

QCD+QED

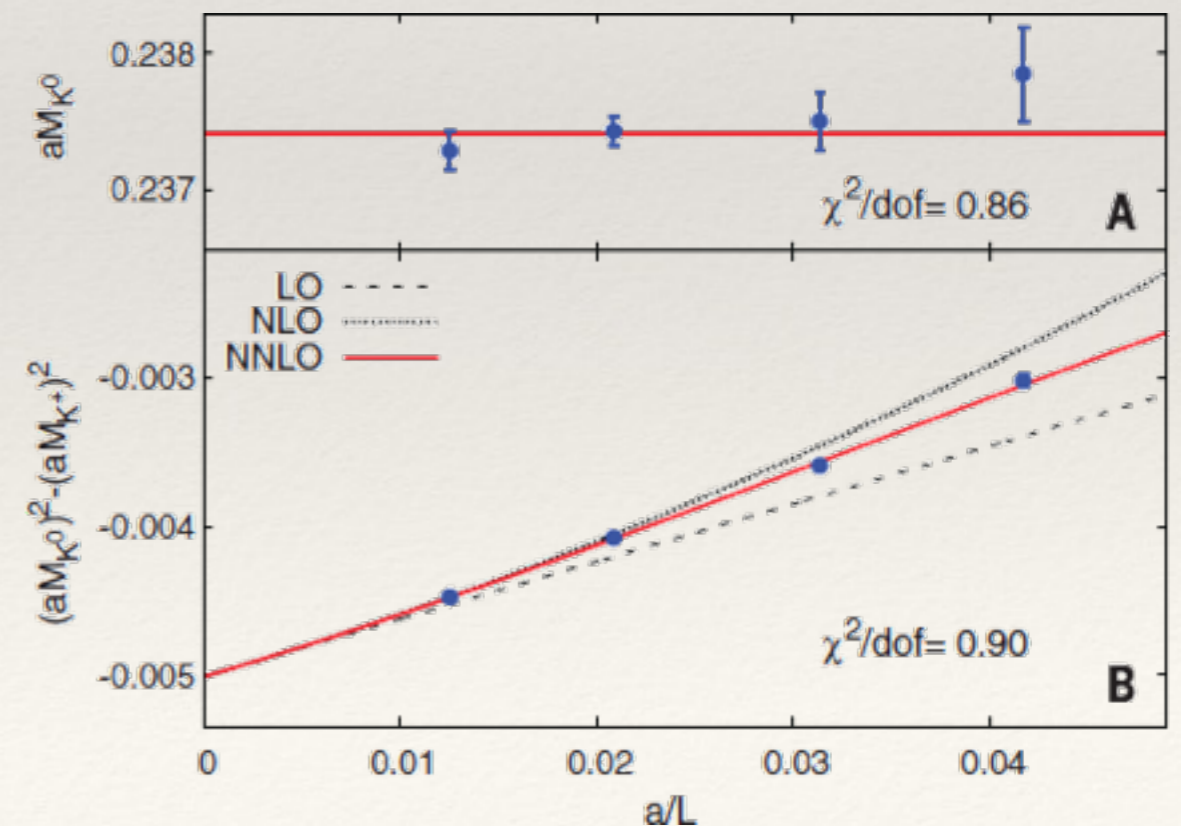
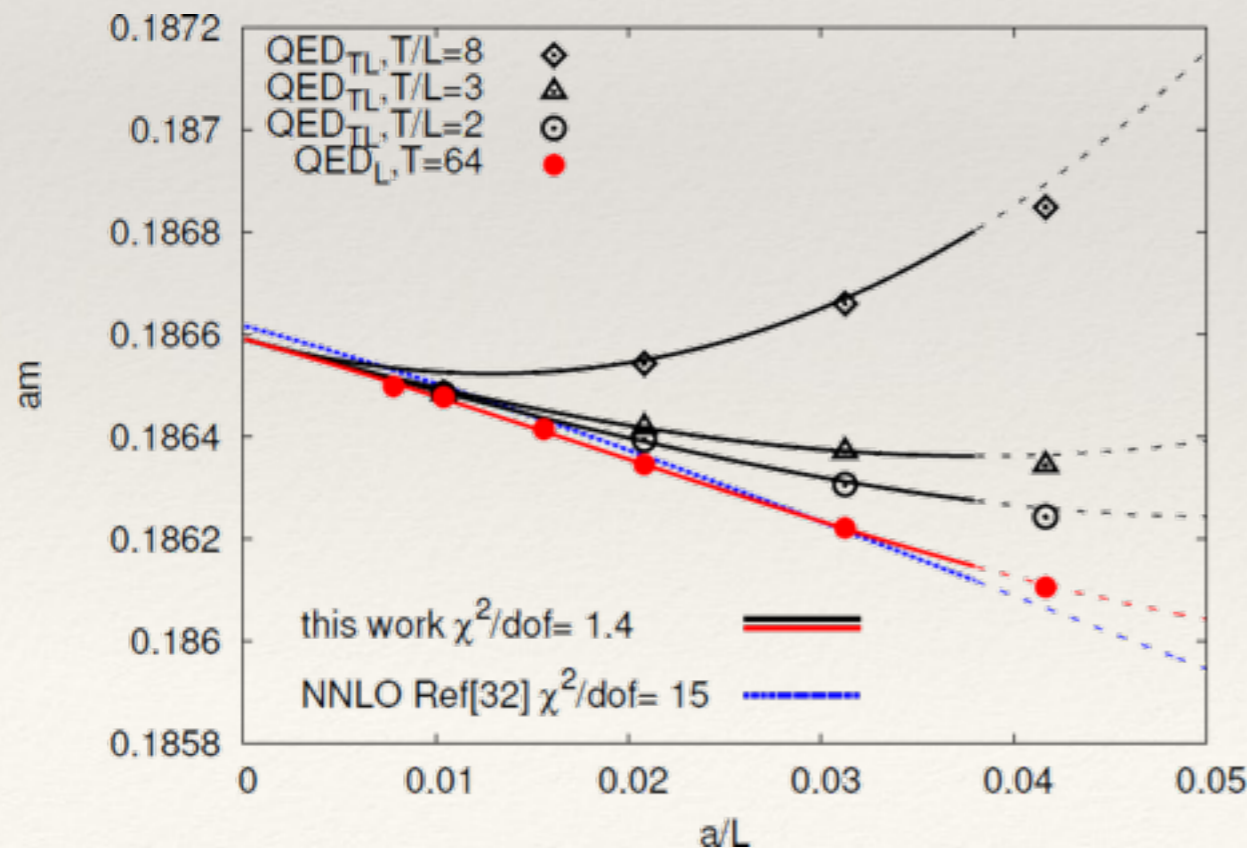
BMW Collaboration
 Science 347 (2015) 1452-1455
[arXiv:1406.4088](https://arxiv.org/abs/1406.4088)

Example: FV correction to mass of a spin-1/2 particle in QED

analytically compute the difference of the *finite volume* and *infinite volume* self energies Σ :

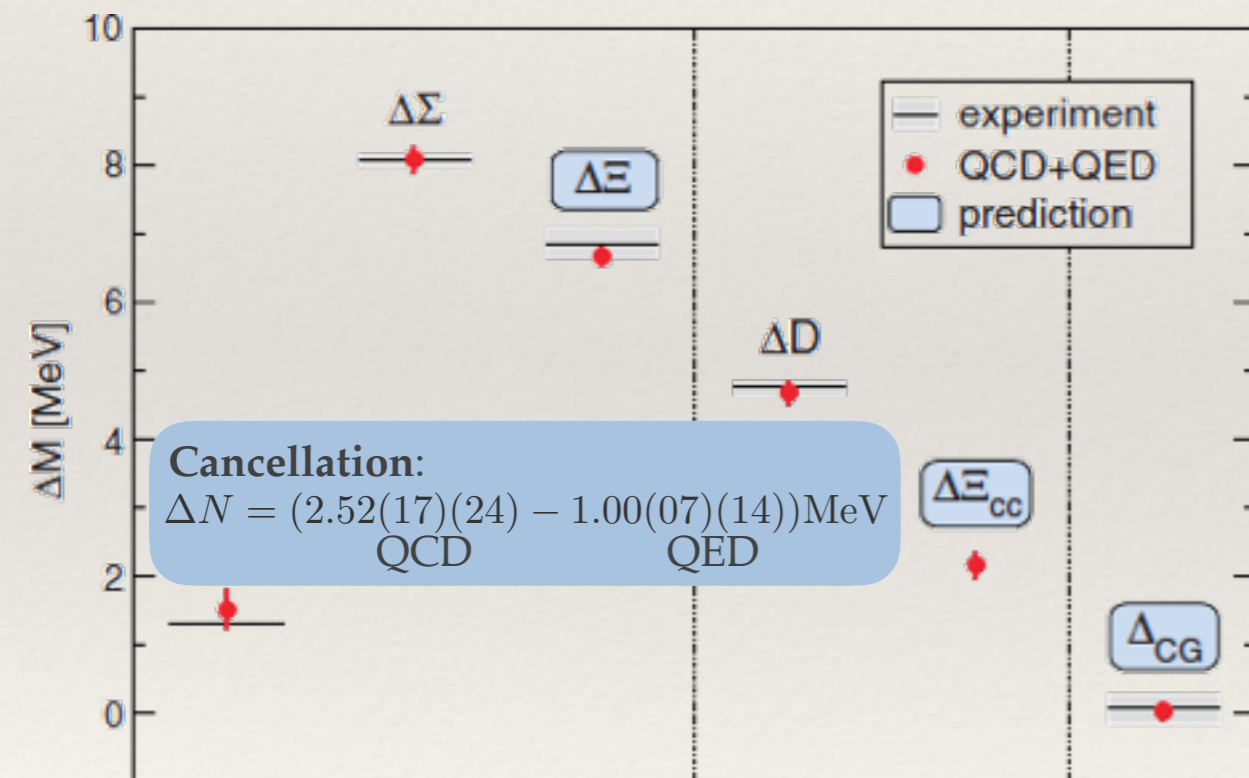
$$m^2(T, L) \stackrel{T, L \rightarrow \infty}{\propto} m^2 \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) - \frac{3\pi}{(mL)^3} \right] \right\}$$

leading behaviour universal in κ (structure- and spin-independent)



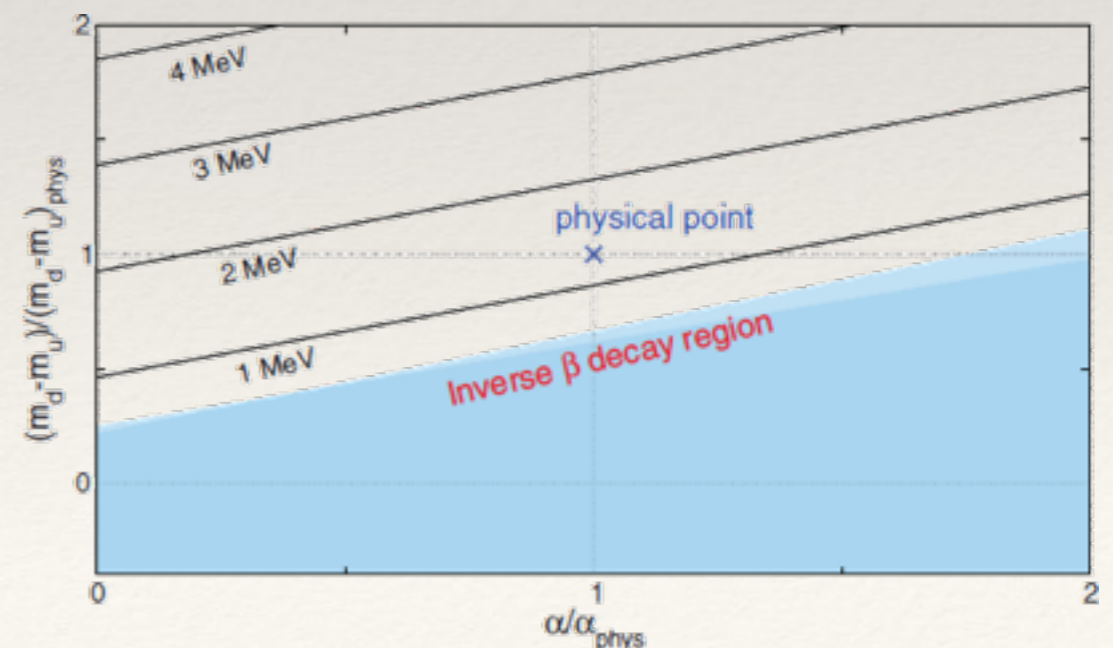
QCD+QED: baryon mass splitting

- relative neutron-proton mass difference found in nature 0.14%
- the value has significant implication for nature
 - smaller value \rightarrow inverse β -decay of H
 - much larger value \rightarrow faster β -decay for neutrons in BBN



BMW Collaboration
 Science 347 (2015) 1452-1455
[arXiv:1406.4088](https://arxiv.org/abs/1406.4088)

BMW carried out simulations of $N_f = 1+1+1+1$ QCD+QED simulations and determined the light baryon isospin splitting



Including QED in meson decay MEs

- leptonic decay at $O(\alpha^0)$:

$$\Gamma(\pi^+ \rightarrow l^+ \nu_l) = \frac{G_F^2 |V_{ud}|^2 f_\pi^2}{8\pi} m_\pi m_l^2 \left(1 - \frac{m_l^2}{m_\pi^2}\right)^2$$

- including elm. effects @ $O(\alpha)$:

$$\begin{aligned} \Gamma(\pi^+ \rightarrow l^+ \nu_l(\gamma)) &= \Gamma(\pi^+ \rightarrow l^+ \nu_l) + \Gamma(\pi^+ \rightarrow l^+ \nu_l \gamma) \\ &\equiv \Gamma_0 \quad + \quad \Gamma_1 \end{aligned}$$

IR div. cancel between terms on r.h.s.
between virtual and real photons
(Bloch Nordsieck)

Including QED in meson decay MEs

Carrasco et al. PRD 91 074506 (2015) [arXiv:1502.00257](https://arxiv.org/abs/1502.00257)

- cut on small photon momentum $< \Delta E \rightarrow \gamma$ sees point-like π
 $\Delta E \approx 20 \text{ MeV}$ experimentally accessible and π point like

$$\Gamma(\Delta E) = \lim_{V \rightarrow \infty} (\Gamma_0 - \Gamma_0^{\text{pt}}) + \lim_{V \rightarrow \infty} (\Gamma_0^{\text{pt}} + \Gamma_1^{\text{pt}}(\Delta E))$$

point approximation

$\Gamma(\pi^+ \rightarrow l^+ \nu_l)$
lattice and analytical
finite V

$\Gamma(\pi^+ \rightarrow l^+ \nu_l \gamma(\Delta E))$
analytically in $V \rightarrow \infty$

both terms separately IR finite, gauge invariant on its own

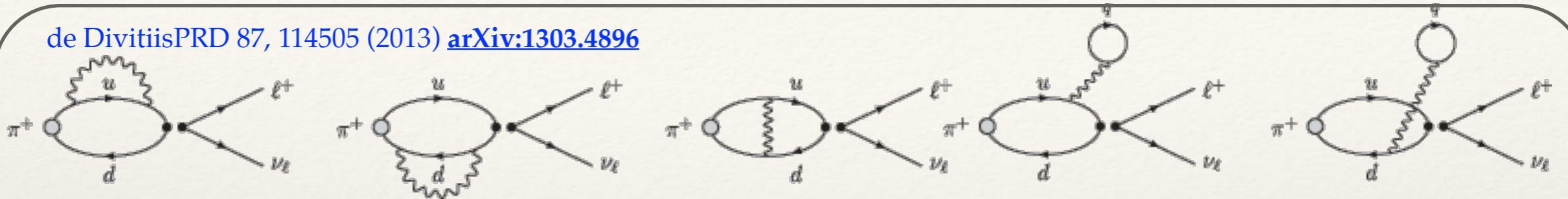
- analytical calculation for pt. approximation is done:

$$\mathcal{L}_{\pi-l-\nu_l} = i G_F f_\pi V_{ud}^* \{(\partial_\mu - ieA_\mu)\pi\} \left\{ \bar{\psi}_{\nu_l} \frac{1 + \gamma_5}{2} \gamma^\mu \psi_l \right\} + \text{H.C.}$$



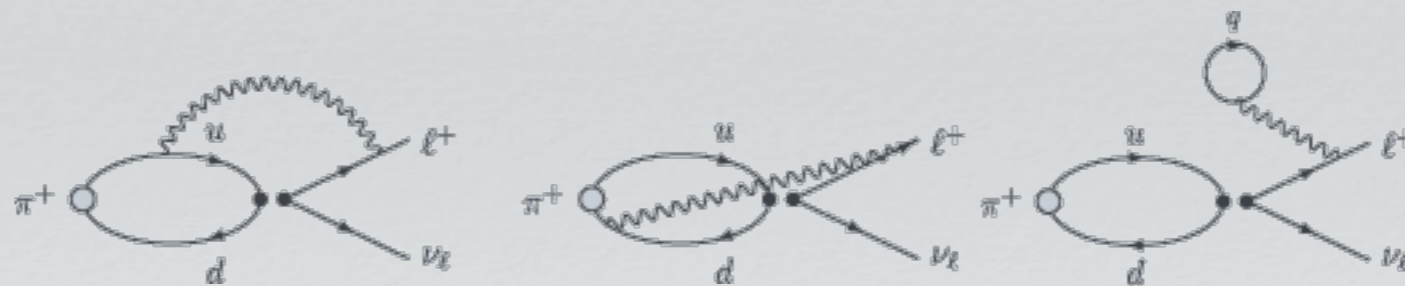
Including QED in meson decay MEs

de Divitiis PRD 87, 114505 (2013) [arXiv:1303.4896](https://arxiv.org/abs/1303.4896)



$$C_1(t) = -\frac{1}{2} \int d^3 \vec{x} d^4 x_1 d^4 x_2 \sum_{\vec{x}} \langle 0 | T \{ J_W^\nu(0) j_\mu(x_1) j_\mu(x_2) \phi^\dagger(\vec{x}, -t) \} | 0 \rangle \Delta(x_1, x_2)$$

$$C_0(t) + C_1(t) \simeq \frac{Z^\phi}{2m_\pi} e^{-m_\pi t} \mathcal{A} \simeq \frac{Z_0^\phi + \delta Z^\phi}{2(m_\pi^0 + \delta m_\pi)} e^{-m_\pi^0 t} (1 - \delta m_\pi t) (\mathcal{A}_0 + \delta \mathcal{A})$$

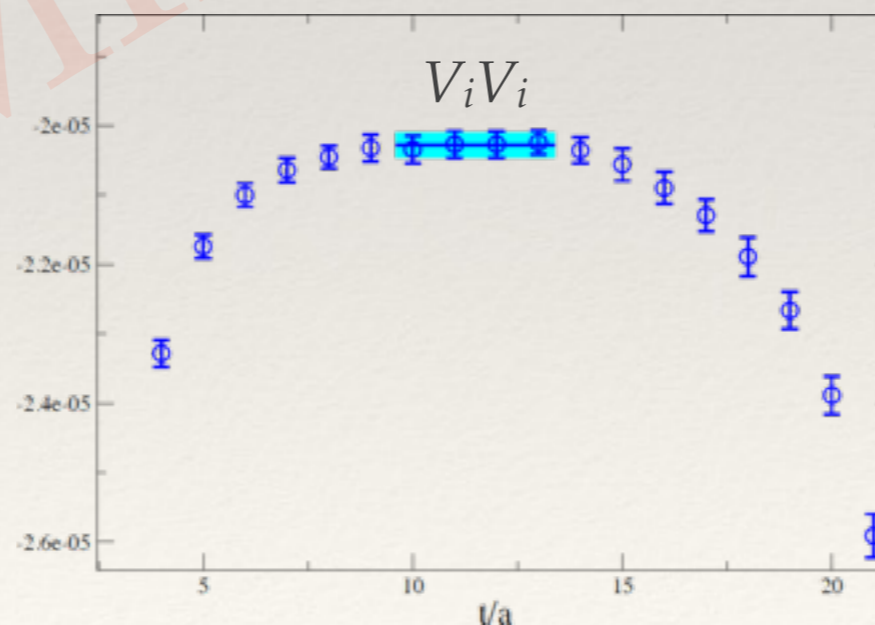
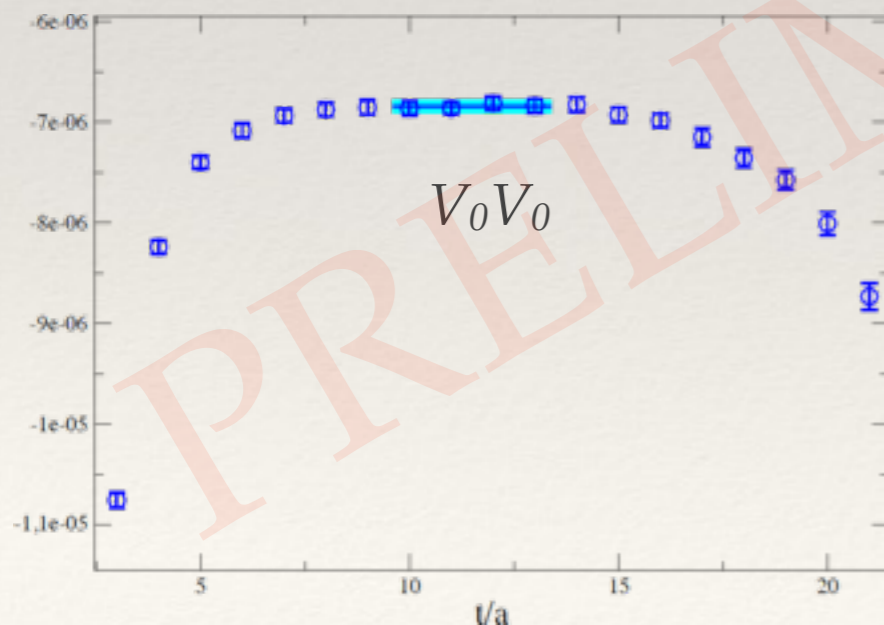
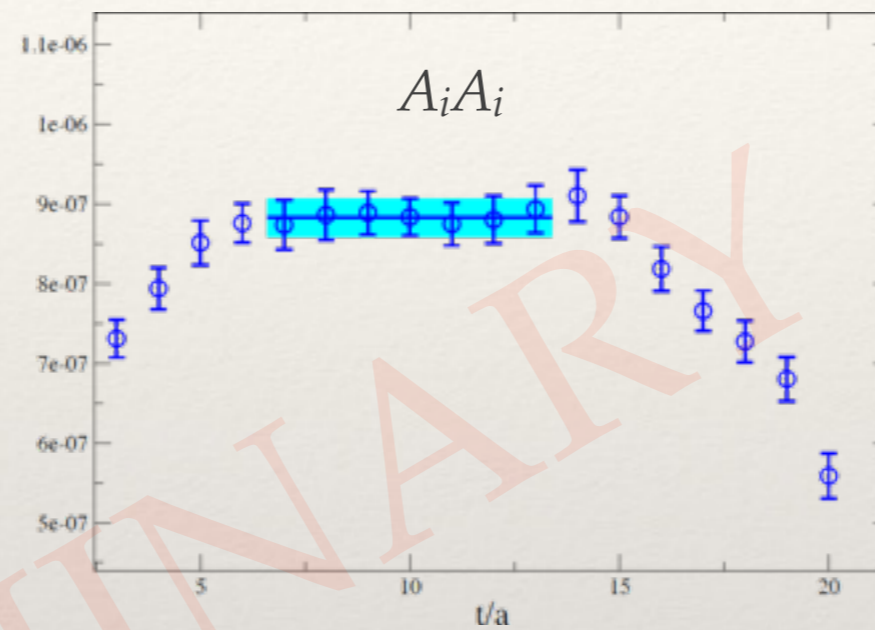
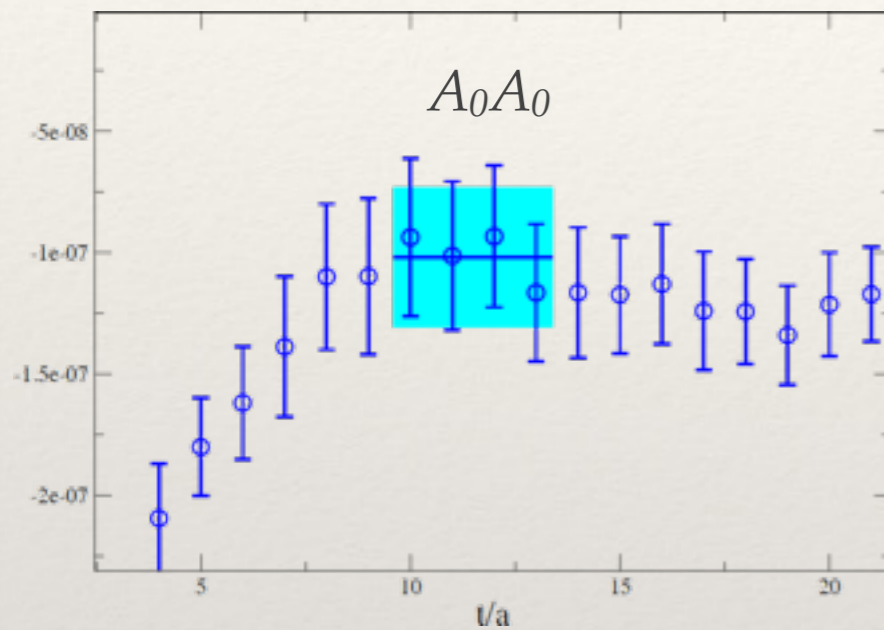


$$C_1(t)_{\alpha\beta} = -\int d^3 \vec{x} d^4 x_1 d^4 x_2 \langle 0 | T \{ J_W^\nu(0) j_\mu(x_1) \phi^\dagger(\vec{x}, -t) \} | 0 \rangle \\ \times \Delta(x_1, x_2) (\gamma_n u(1 - \gamma^5) S(0, x_2) \gamma_\mu)_{\alpha\beta} e^{E_l t_2 - i \vec{p}_l \cdot \vec{x}_2}$$

convergent by
energy / momentum conservation

Including QED in meson decay MEs

$24^3; m_\pi \approx 500 \text{ MeV}$



- first time ever conceptually clean attempt of calculation of leptonic decay at $O(\alpha)$
- disconnected pieces need to be included
- Γ_0 works, now needs to be combined wt. analytical results for $\Gamma_0^{\text{pt}}, \Gamma_1^{\text{pt}}(\Delta E)$
- $\sim 20\%$ stat. error would be sufficient for use in phenomenology

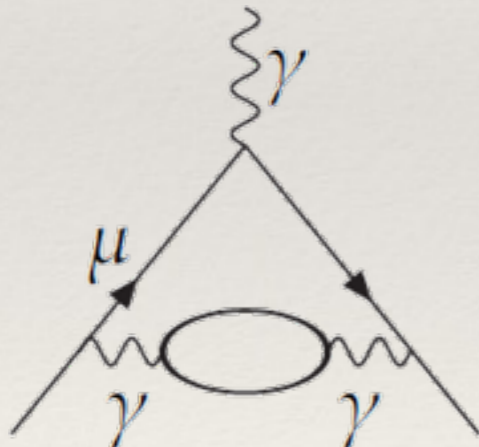
QCD+QED applications

- start with light flavour matrix elements $f_\pi, f_K, f_+(0), \dots$

$$\Gamma_{K \rightarrow \pi l \nu} = C_K^2 \frac{G_F^2 m_K^5}{192 \pi^2} S_{EW} (1 + \Delta_{SU(2)} + \Delta_{EM})^2 I |f_+^{K\pi}(0)|^2 |V_{us}|^2$$

$$\langle \pi(p_\pi) | V_\mu(0) | K(p_K) \rangle = f_+^{K\pi}(q^2) (p_K + p_\pi)_\mu + f_-^{K\pi}(q^2) (p_K - p_\pi)_\mu$$

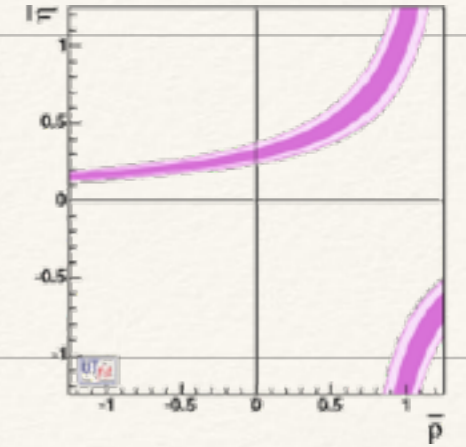
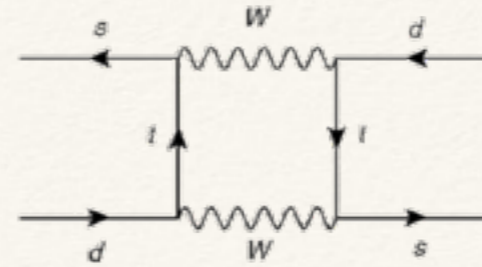
- lattice predictions of leading hadronic contribution to muon g-2



- lattice (isospin symmetric, $\alpha_{EM}=0$ is getting competitive with experimental determination ($e^+e^- \rightarrow hadrons$))
- next step would be inclusion of isospin breaking effects

- inclusion of QED effects will be one of the big challenges over the next years

Long distance effects in kaon physics - mixing

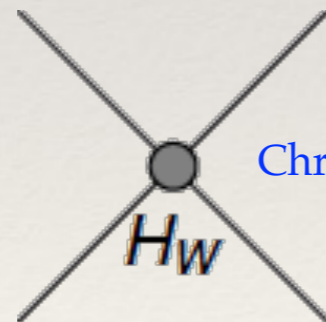


$$\epsilon_K = \frac{A(K_L \rightarrow (\pi\pi)_{I=0})}{A(K_S \rightarrow (\pi\pi)_{I=0})} = e^{i\Phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}\langle \bar{K}^0 | H_W^{\Delta S=2} | K^0 \rangle}{\Delta M_K} + \text{L.D. effects} \right)$$

Buras, Guadagnoli PRD 78 (2008)
Buras, Guadagnoli, Isidori, PLB 688 (2010)

Long Distance effects amount to $O(5\%)$, so certainly worth considering on the lattice

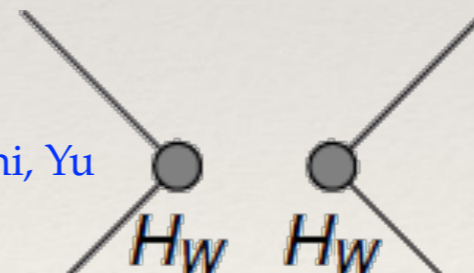
1st order Weak



single 4-quark OP,
length scale 10^{-18}m

Christ, Izubuchi, Sachrajda, Soni, Yu
arXiv:1212.5931

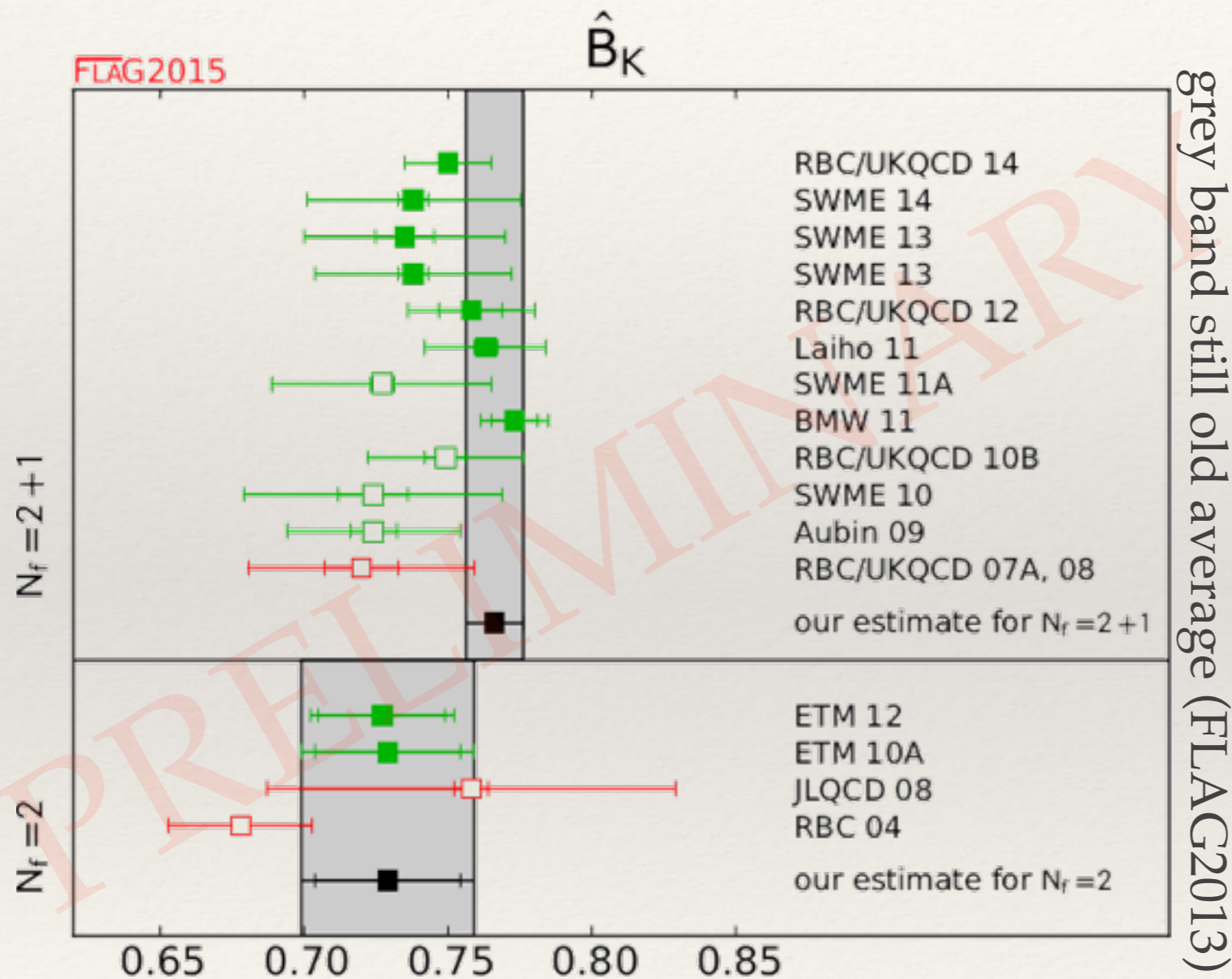
2nd order Weak



two 4-quark OP
length scale $1/\Lambda_{\text{QCD}}$

Neutral meson mixing - kaon - short distance

FLAG B_K Working Group Dimopoulos, Mawhinney, Wittig



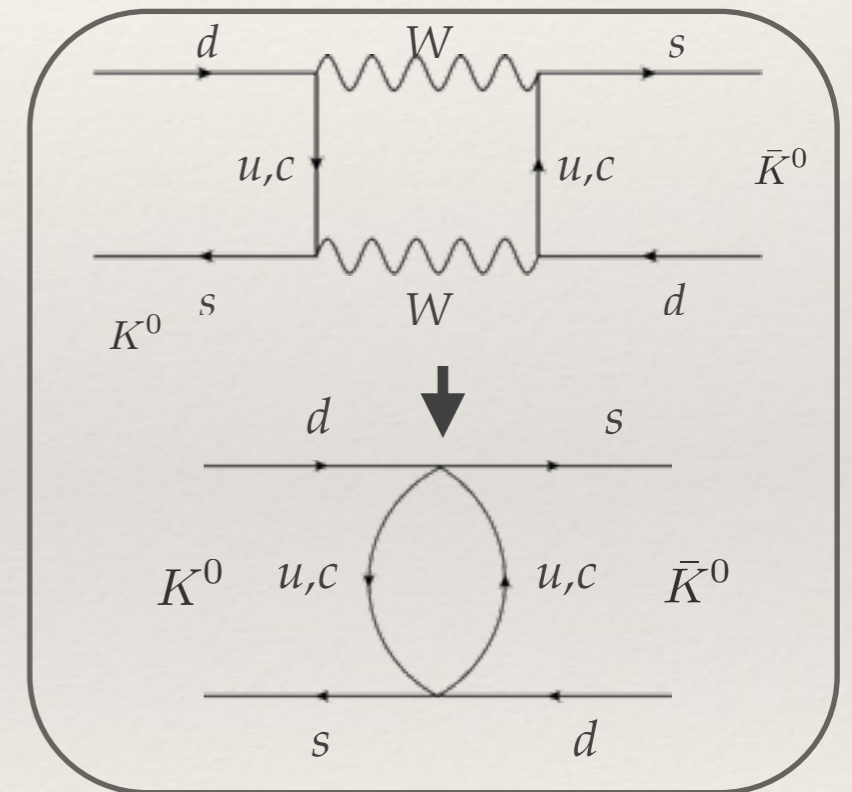
contrary to BSM bags here
excellent agreement at the
1.3%-level

Beyond short distance: e.g. ΔM_K

$$\Delta M_K = m_{K_S} - m_{K_L} = 2\text{Re}M_{00}$$

$$M_{00} = \mathcal{P} \sum_{\lambda} \frac{\langle \bar{K}^0 | H_W | \lambda \rangle \langle \lambda | H_W | K^0 \rangle}{m_K - E_{\lambda}}$$

- experimentally $\Delta M_K = 3.483(6) \times 10^{-12} \text{MeV}$ (PDG)
- suppressed by 14 orders of magnitude with respect to QCD \rightarrow poses strong BSM constraints (e.g. $(1/\Lambda)^2 \bar{s}d\bar{s}d$ BSM contribution) knowing ΔM_K at 10%-level $\rightarrow \Lambda \geq 10^4 \text{TeV}$
- SD about 70% of experimental value - rest LD?
- PT large contributions at $\mu \sim m_c$ where PT turns out to converge badly (NLO \rightarrow NNLO constitutes 36% correction) [Brod, Gorbahn PRL 108 121801 \(2012\) arXiv:1108.2036](#)



long distance effects – ΔM_K

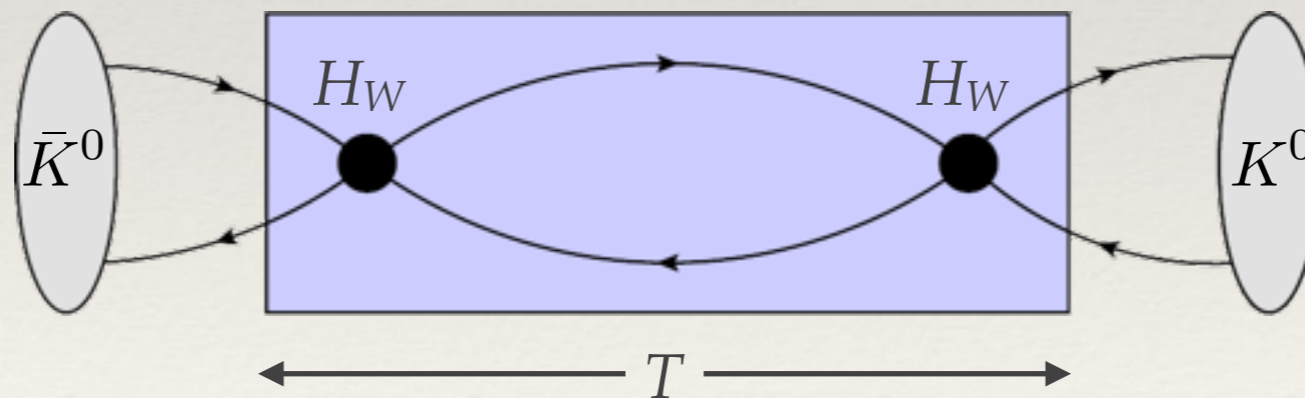
N. Christ et al. PRD 88 (2013) 014508 [arXiv:1212.5931](#)

Bai et al. PRL 113 (2014) 112003 [arXiv:1406.0916](#)

$$M_{\bar{0}0} = \mathcal{P} \sum_{\lambda} \frac{\langle \bar{K}^0 | H_W | \lambda \rangle \langle \lambda | H_W | K^0 \rangle}{m_K - E_{\lambda}}$$

$$\mathcal{A} = \langle 0 | T \left\{ K^0(t_f) \frac{1}{2} \int_{t_A}^{t_B} dt_2 \int_{t_A}^{t_B} dt_1 H_W(t_2) H_W(t_1) K^{0\dagger}(t_i) \right\} | 0 \rangle$$

Integrate operators (here H_W) over time interval where initial and final kaon dominate

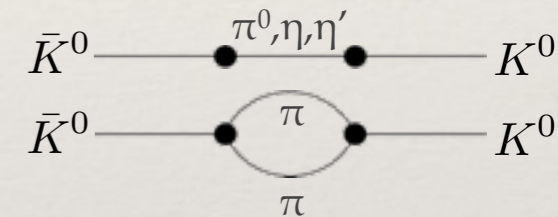


long distance effects – ΔM_K

N. Christ et al. PRD 88 (2013) 014508 [arXiv:1212.5931](#)
 Bai et al. PRL 113 (2014) 112003 [arXiv:1406.0916](#)

$$\mathcal{A} = N_K^2 e^{-M_K(t_f - t_i)} \sum_n \frac{\langle \bar{K}^0 | H_W | n \rangle \langle n | H_W | K^0 \rangle}{M_K - E_n} \left(-T - \frac{1}{M_K - E_n} + \frac{e^{(M_K - E_n)T}}{M_K - E_n} \right)$$

amplitude
irrelevant
exponential term
 Δm_K^{FV}
constant
needs to be subtracted



- multiple hadrons in intermediate states causing difficulties and need to be subtracted
- finite volume corrections from two-particle intermediate state can be sizeable

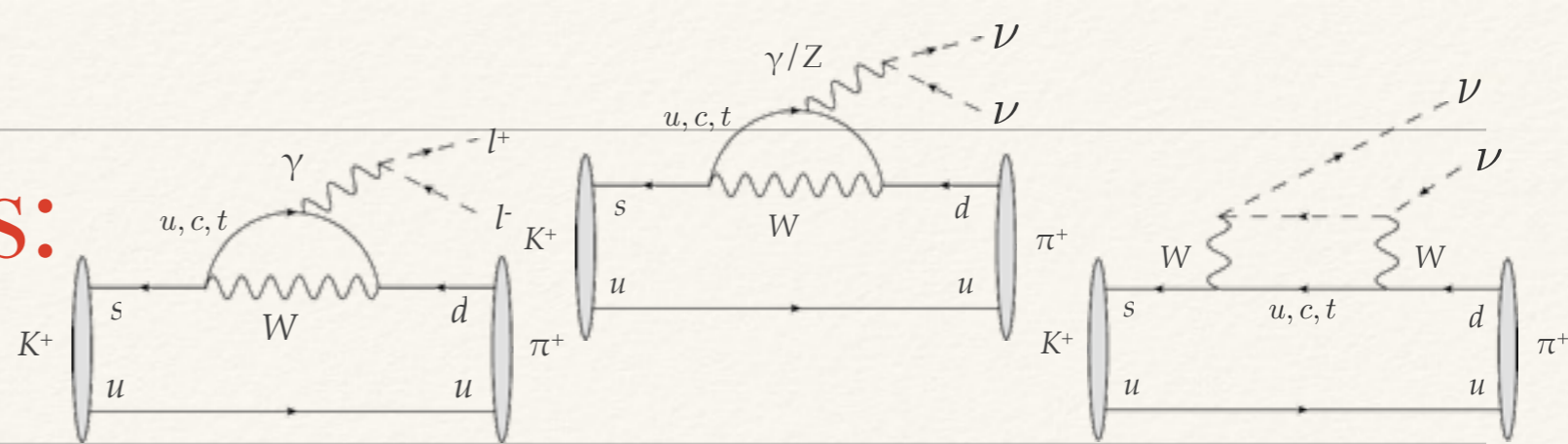
N. Christ et al. PRD91 (2015) 114510 [arXiv:1504.01170](#) also: Briceno, Hansen [arXiv:1502.04314](#)

extension of Lellouch-Lüscher correction to 2nd order weak MEs

$$\Delta^{\text{FV}}(\Delta M_K) = -\cot(\phi(M_K) + \delta_0(M_K)) \frac{d(\phi(E) + \delta_0(E))}{dE} \Big|_{E=M_K} |\langle \bar{K}^0 | H_W | \pi\pi, M_K \rangle^{\text{V}'}|^2$$

- what happens when the two H_W approach each other (GIM in action)?

long distance effects: Rare kaon decays



Two new experiments dedicated to rare kaon decays
NA62 (CERN) and KOTO (J-PARC) are running

- FCNC (W-W or γ/Z -exchange diagrams)
- deep probe into flavour mixing and SM/BSM due to suppression in the SM
- can determine V_{td} , V_{ts} and test SM

$$K_L \rightarrow \pi^0 \nu \bar{\nu}$$

- KOTO (J-PARC)
- direct CP violation
- exp. BR $\leq 2.6 \times 10^{-8}$
theory BR $3.0(3) \times 10^{-11}$
- GIM \rightarrow top dominated and charm suppressed, purely SD

$$K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

- NA62 (CERN)
- CP conserving
- exp. BR $1.73^{+1.15}_{-1.05} \times 10^{-10}$
theory BR $0.911(72) \times 10^{-10}$
- small LD contribution, candidate for lattice

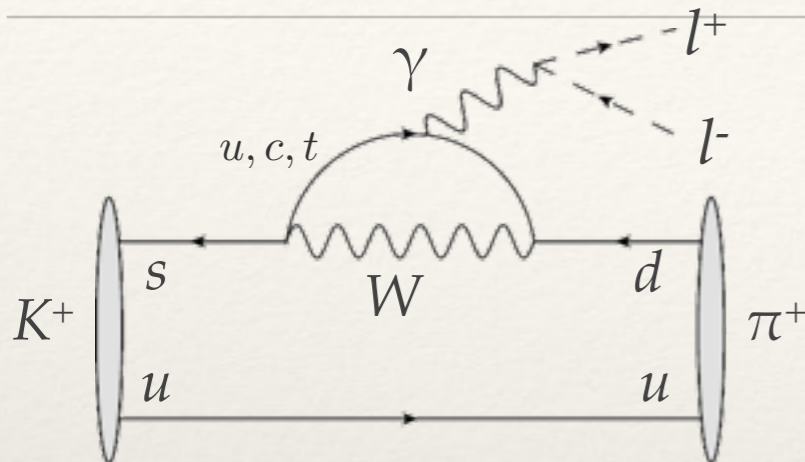
compute in lattice QCD

$$K^+ \rightarrow \pi^+ l^+ l^- \quad K_S \rightarrow \pi^0 l^+ l^-$$

- 1-photon exchange LD dominated
- indirect contribution to CP-violating rare K_L decay
- SM prediction mainly ChPT
- lattice can predict ME and LECs
- experimenters will be able to look at these channels as well

Rare kaon decays $K^+ \rightarrow \pi^+ l^+ l^-$

N. Christ et al. [arXiv:1507.03094](https://arxiv.org/abs/1507.03094)
[arXiv:1602.01374](https://arxiv.org/abs/1602.01374)



LD contribution given through $K \rightarrow \gamma^*$ contribution which is computed as

$$\mathcal{A}_\mu = (q^2) \int d^4x \langle \pi(p) | T [J_\mu(0) H_W(x)] | K(k) \rangle$$

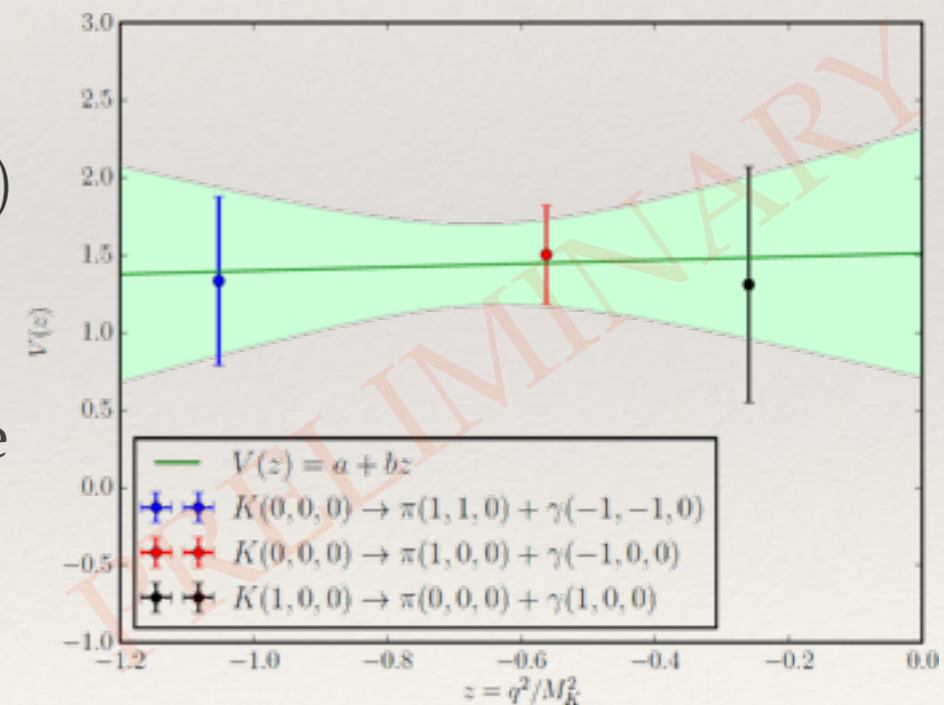
dominant operators: $Q_1^q = (\bar{s}_i \gamma_\mu^L d_i) (\bar{q}_j \gamma_\mu^L q_j)$, $Q_2^q = (\bar{s}_i \gamma_\mu^L q_i) (\bar{q}_j \gamma_\mu^L d_j)$

Decay amplitude in terms of elm. transition form factor

$$A_i = -\frac{G_F \alpha}{4\pi} V_i(z) (k+p)^\mu \bar{u}_l(p_-) \gamma_\mu \nu_l(p_+) \quad (i = +, S)$$

$$V_i(z) = a_i + b_i z + V_i^{\pi\pi}(z)$$

- the a_S and a_+ can be extracted from experiment or lattice
- a_S parameterises also the CP-violating contribution to the K_L decay



Summary/Conclusions

- considerable set of SM parameters, spectra and matrix elements now reliably and precisely predicted in full lattice QCD
- Flavour Lattice Averaging Group (FLAG)
- precision such that isospin breaking effects in matrix elements and spectra needs to be taken into account
- long distance effects:
 - neutral kaon system
 - rare kaon decays (new experimental facilities!!!)with loads of new questions and theoretical problems and potential impact on SM and BSM phenomenology

Summary/Conclusions

- this talk is by far not inclusive:
 - $K \rightarrow \pi\pi, g-2, \dots$
 - finite- T, μ
 - BSM
 - ...

34th International Symposium on
Lattice Field Theory

University of Southampton
24–30 July 2016



registration and abstract
submission open!

Looking forward to see you in Southampton!

<http://www.southampton.ac.uk/lattice2016/>

