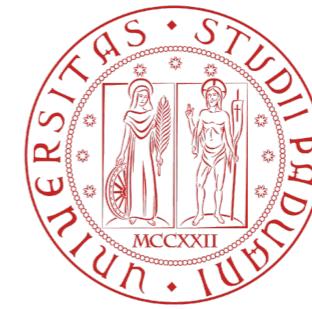




DIPARTIMENTO  
DI FISICA  
E ASTRONOMIA  
Galileo Galilei



# Effective field theory of 3d $\mathcal{N} = 2$ CFT's from holography

Luca Martucci

University of Padova

- |           |            |  |
|-----------|------------|--|
| based on: | 1803.xxxx  | with Stefano Cremonesi & Stefano Lanza |
|           | 18xx.xxxx  | with Stefano Cremonesi                 |
|           | 1603.04470 | with Alberto Zaffaroni                 |

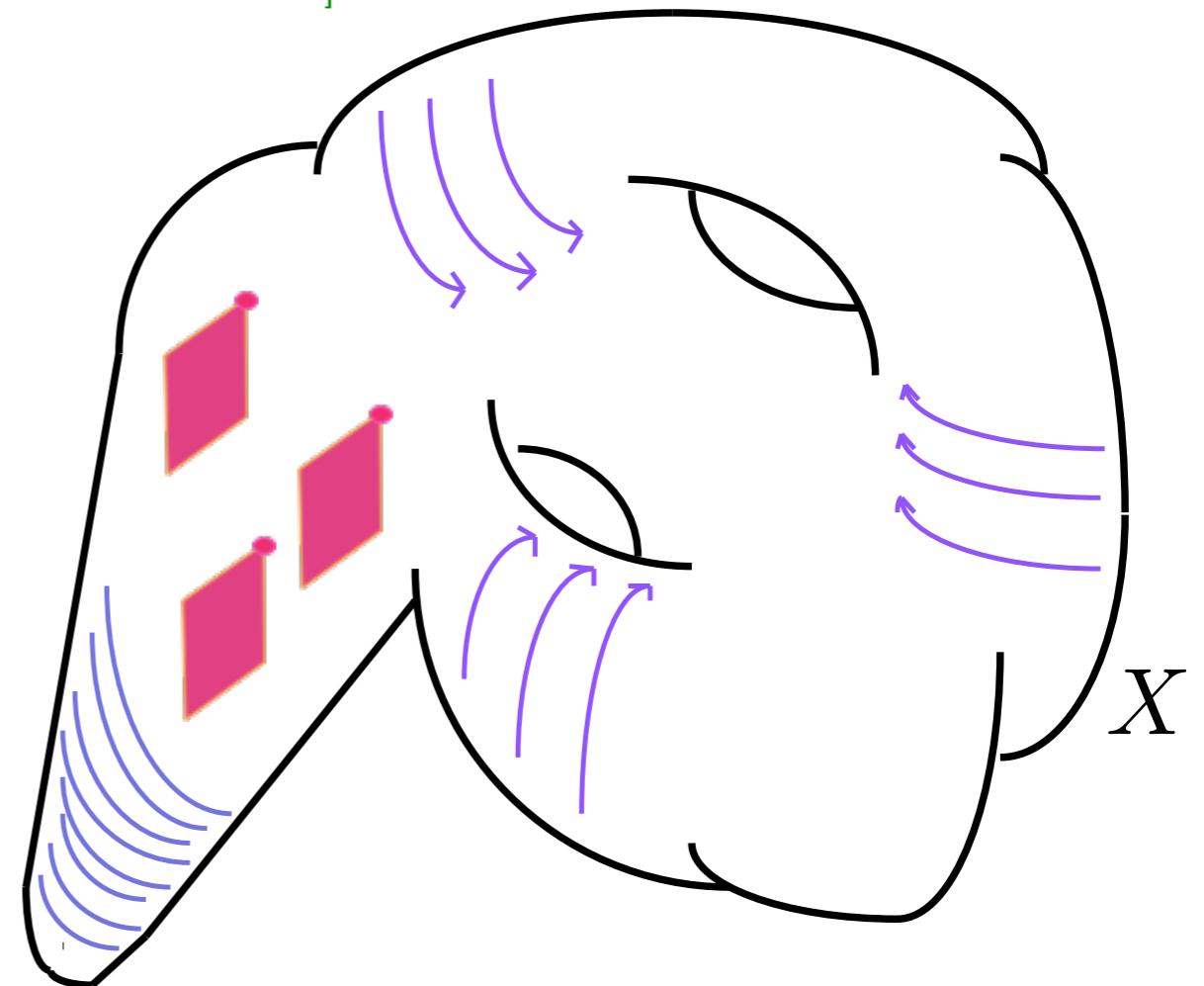
A motivation:  
EFTs for warped compactifications?

- M-theory compactifications to d=3 are generically **warped**

$$ds_{11}^2 = H^{-\frac{2}{3}} ds_{\mathbb{R}^{1,2}}^2 + H^{\frac{1}{3}} ds_X^2 \quad [\text{Becker-Becker '96}]$$

warping generated by:

- \* mobile M2-branes
- \*  $G_4$ -flux



- M/F-theory duality  $\rightarrow$  IIB warped compactifications to d=4

[Grana-Polchinski, Gubser, Giddings-Kachru-Polchinski '00]

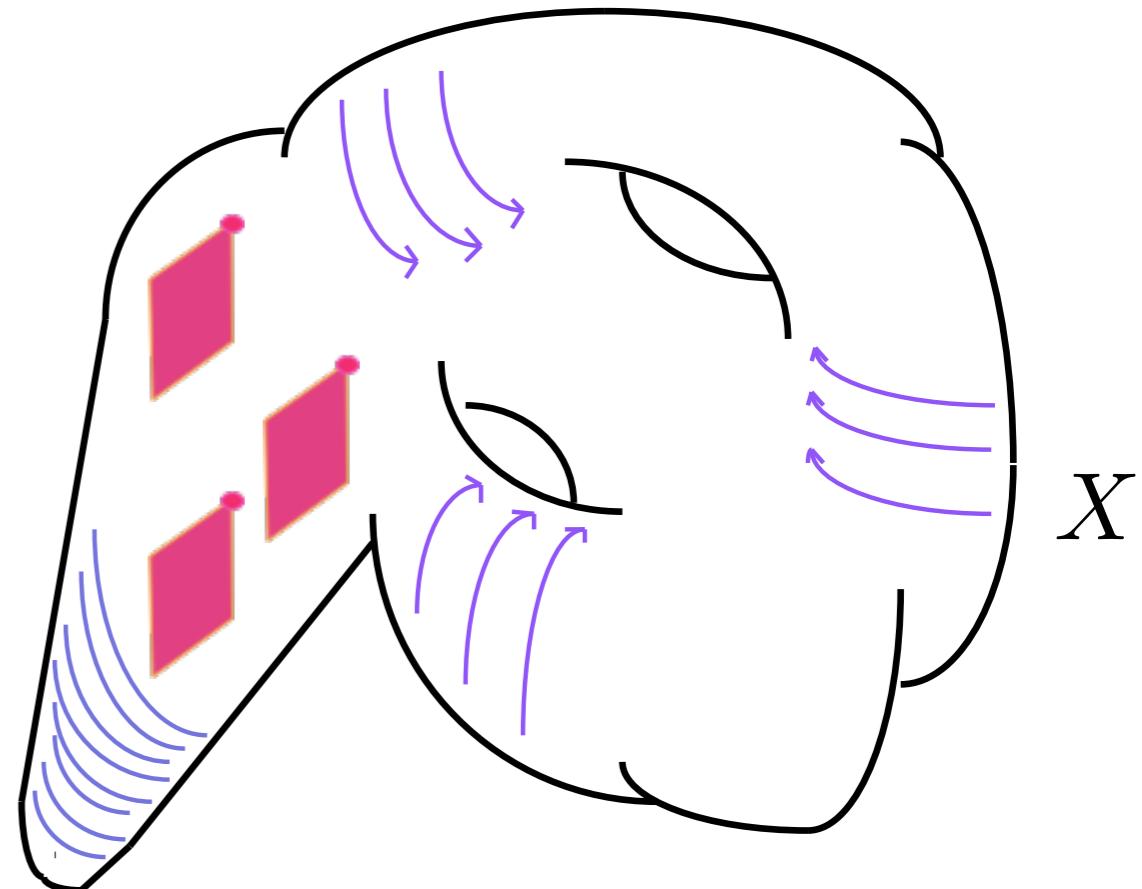
• Tree-level moduli space  $\mathcal{M}$ :

- \* M2-brane positions
- \* Kähler + axionic moduli

• Effective  $K(\Phi, \bar{\Phi})$  explicitly depends on:

- \* internal flux  $G_4$  [LM `14-`16]
- \* detailed form of the CY metric

**'NON-TOPOLOGICAL' !**



E.g: Kähler moduli space metric includes:

[Frey-Roberts `13] [Cownden-Frey-David Marsh-Underwood `16] [LM `16]

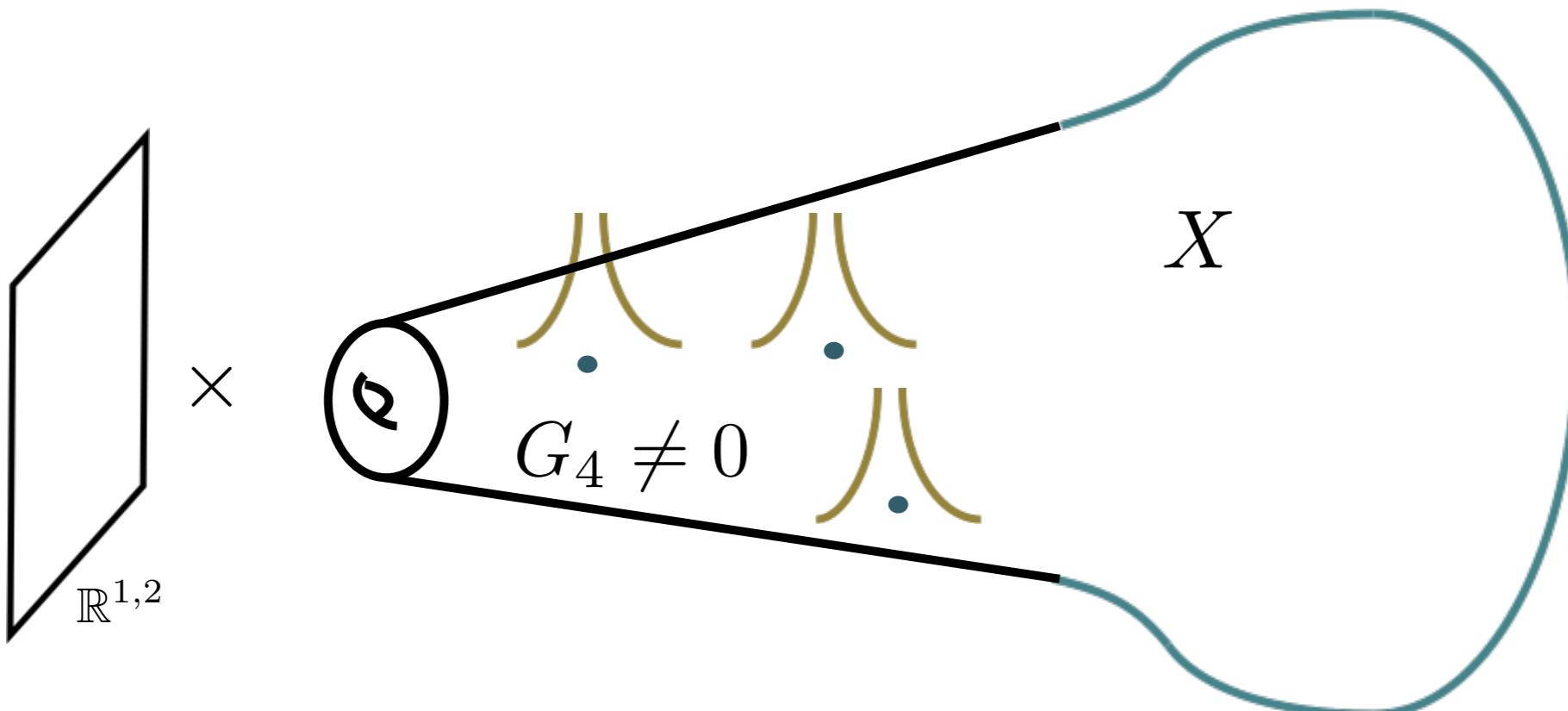
$$\int_X H\omega_a \wedge *_X \omega_b + \int_X d\Lambda_a \wedge *_X d\Lambda_b \xrightarrow{\Delta_X^{-1}} *_X (\omega_a \wedge G_4)$$

• Further corrections from M-theory higher derivative terms

[Grimm-Pugh-Weissenbacher `14-`15]  
 [Grimm-Mayer-Weissenbacher `17]

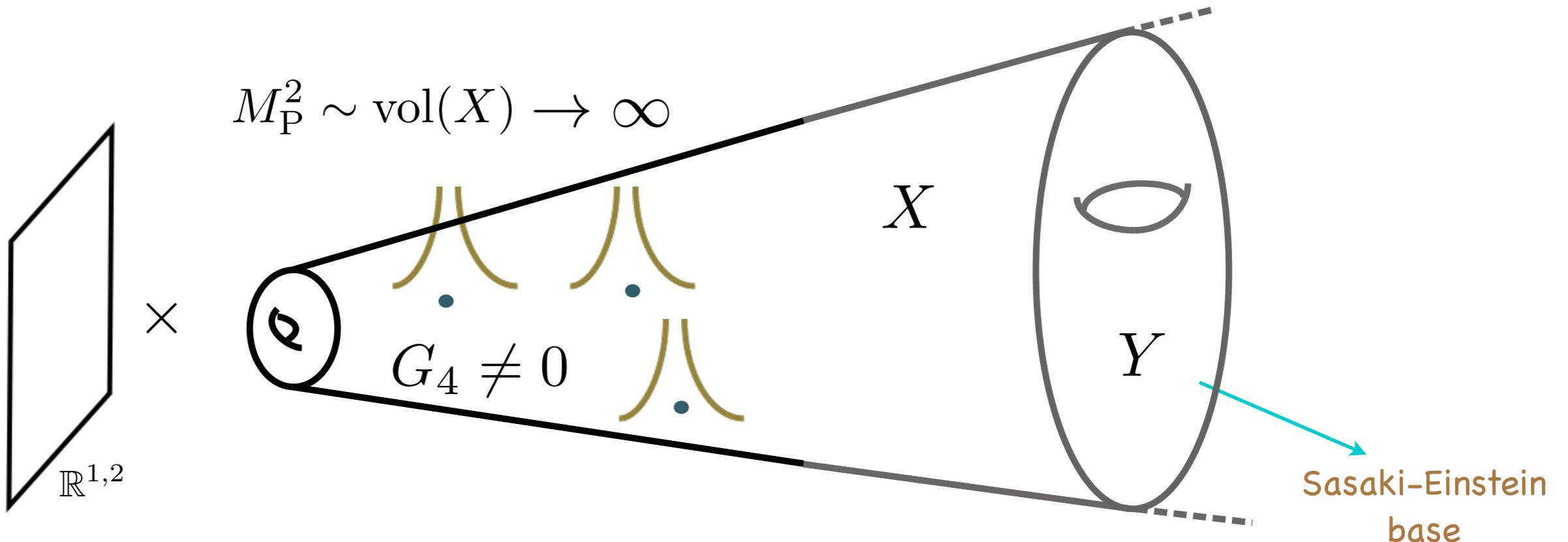
# Computable EFTs?

- I will consider local models ...



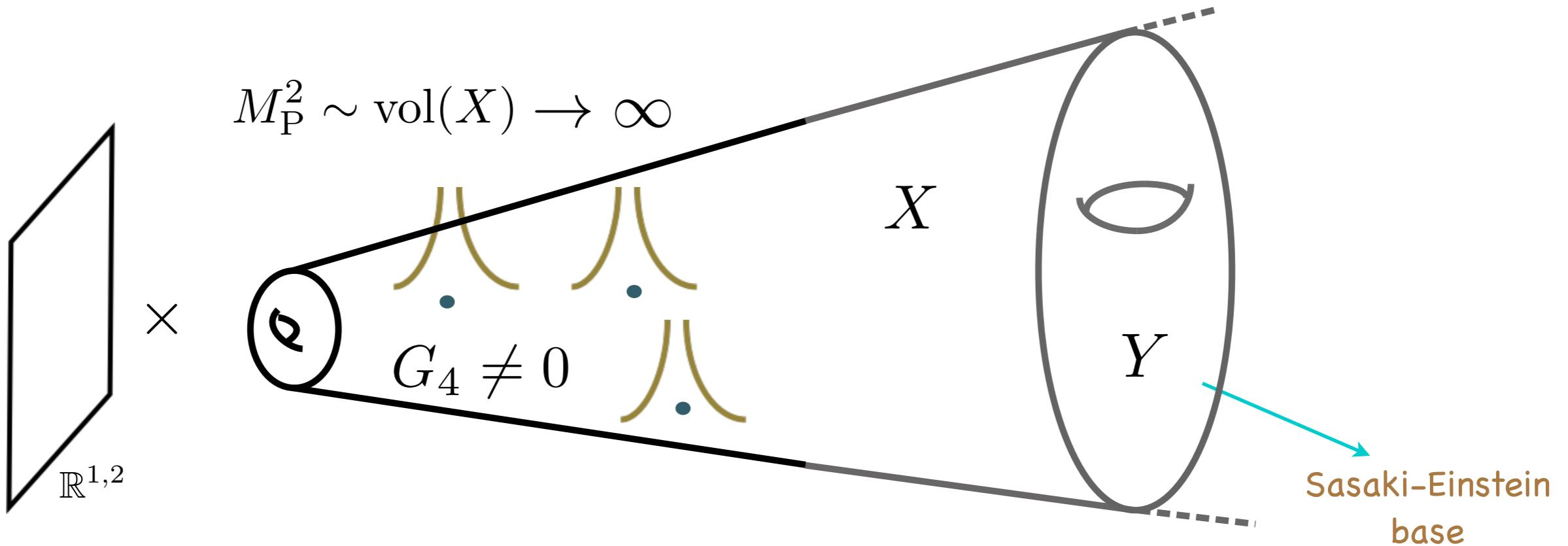
# Computable EFTs?

- I will consider local models ...



# Computable EFTs?

- I will consider local models ...



- ... in the near-horizon limit:

$$H \sim a + \frac{L^6}{r^6} \xrightarrow{a \rightarrow 0} H \sim \frac{L^6}{r^6}$$

asymptotically  
AdS<sub>4</sub> × Y

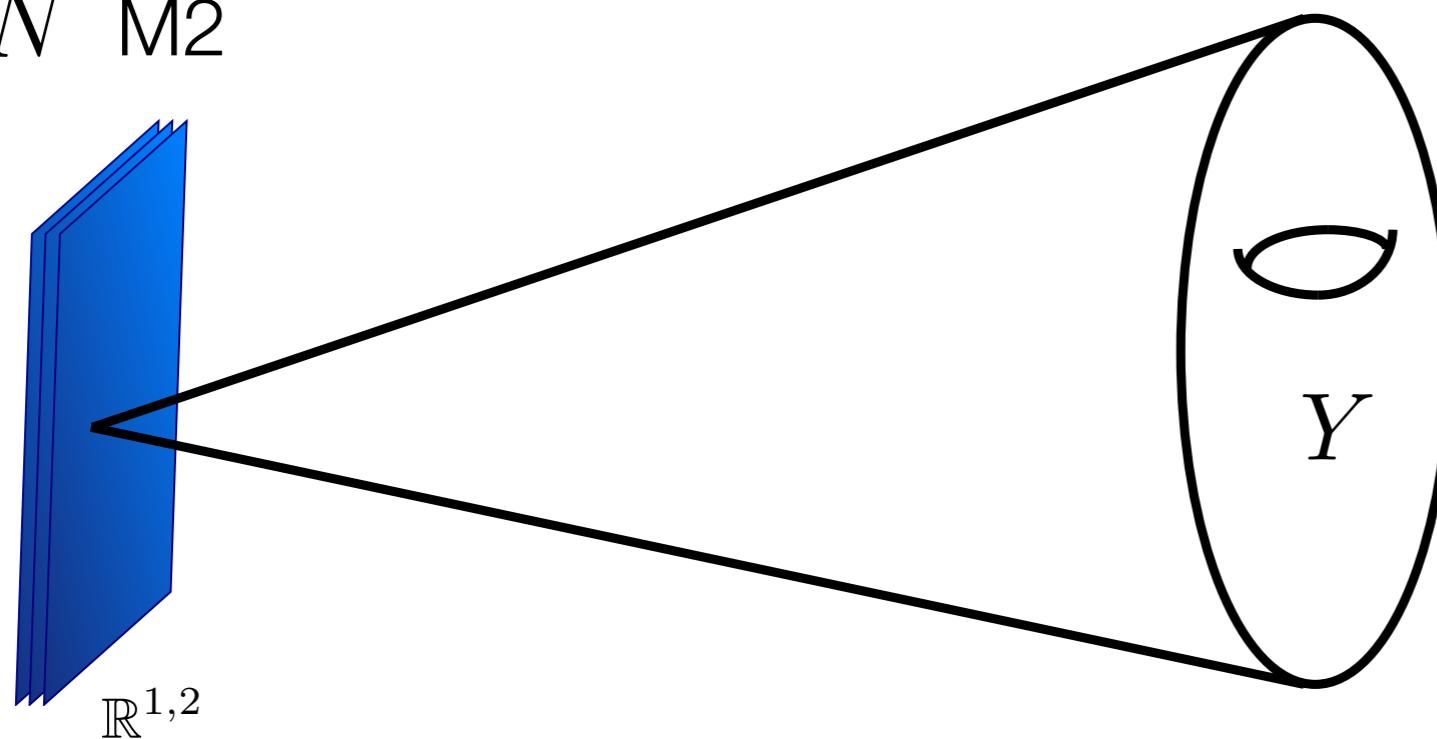


(Rigid) holographic EFT for dual  $\mathcal{N} = 2$  CFT<sub>3</sub>

# Large-N EFTs from supergravity

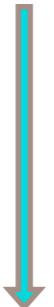
# Dual $\mathcal{N} = 2$ CFTs

$N$  M2



generalizations of  
 $\mathcal{N} = 6$  ABJM theory  
[Aharony-Bergman-Jafferis-Maldacena '08]

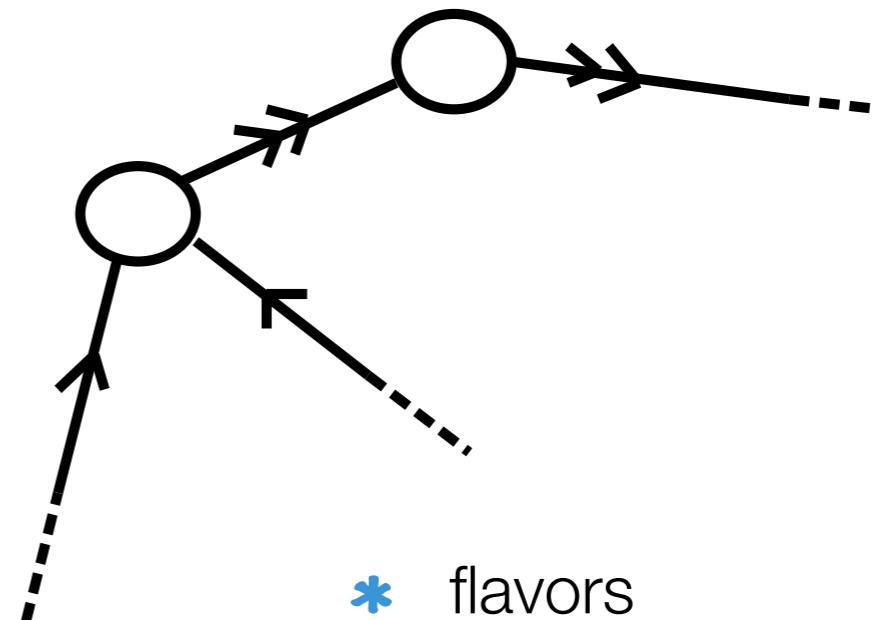
- 3d  $\mathcal{N} = 2$  quiver models



RG-flow

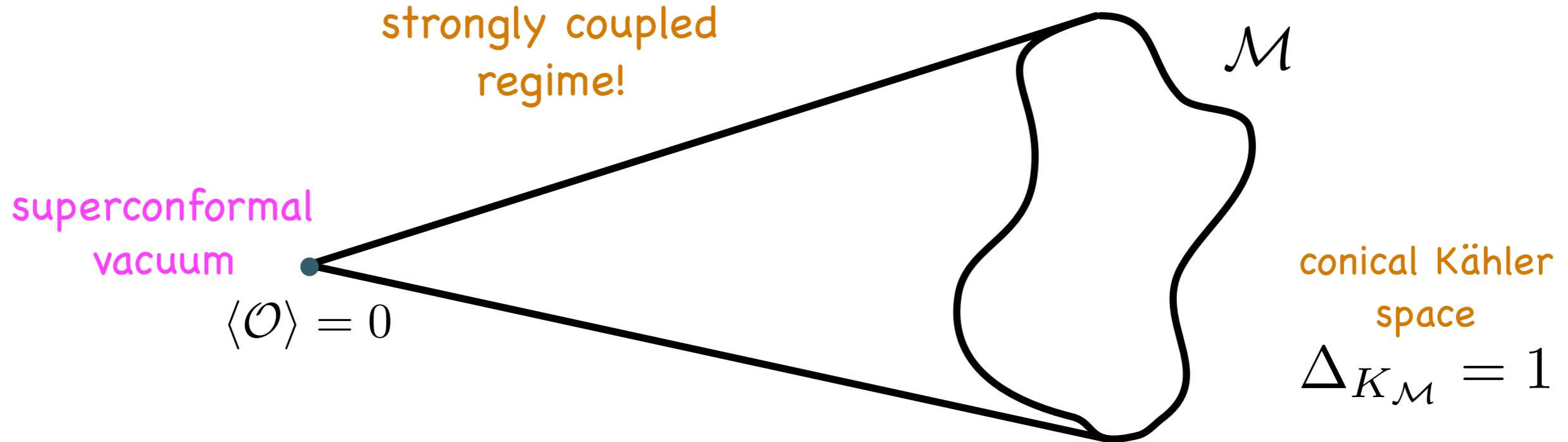
3d  $\mathcal{N} = 2$  SCFTs

strongly coupled!

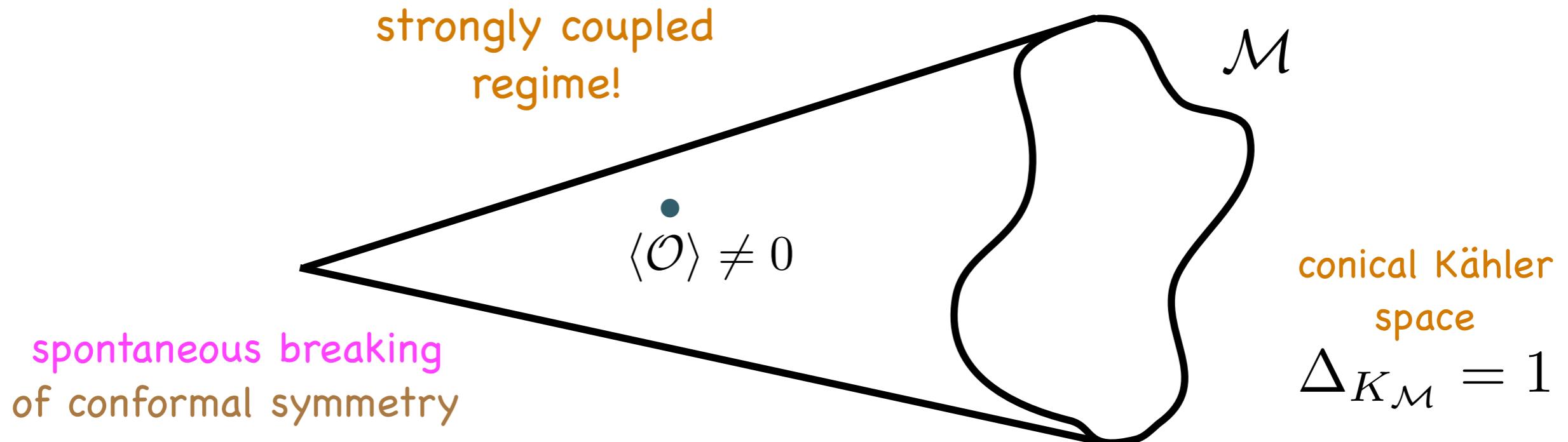


- \* flavors
- \* unequal ranks  $\sim N$
- \* monopole operators
- \* non-perturbative effects

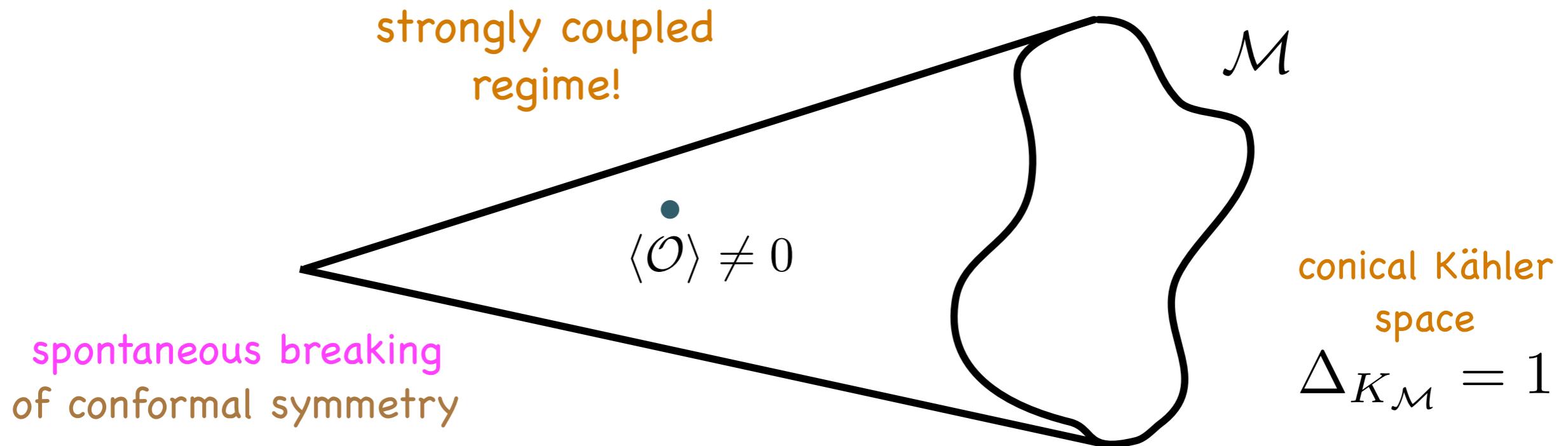
# $\mathcal{N} = 2$ SCFT moduli space



# $\mathcal{N} = 2$ SCFT moduli space



# $\mathcal{N} = 2$ SCFT moduli space



- Complex/algebraic structure well studied

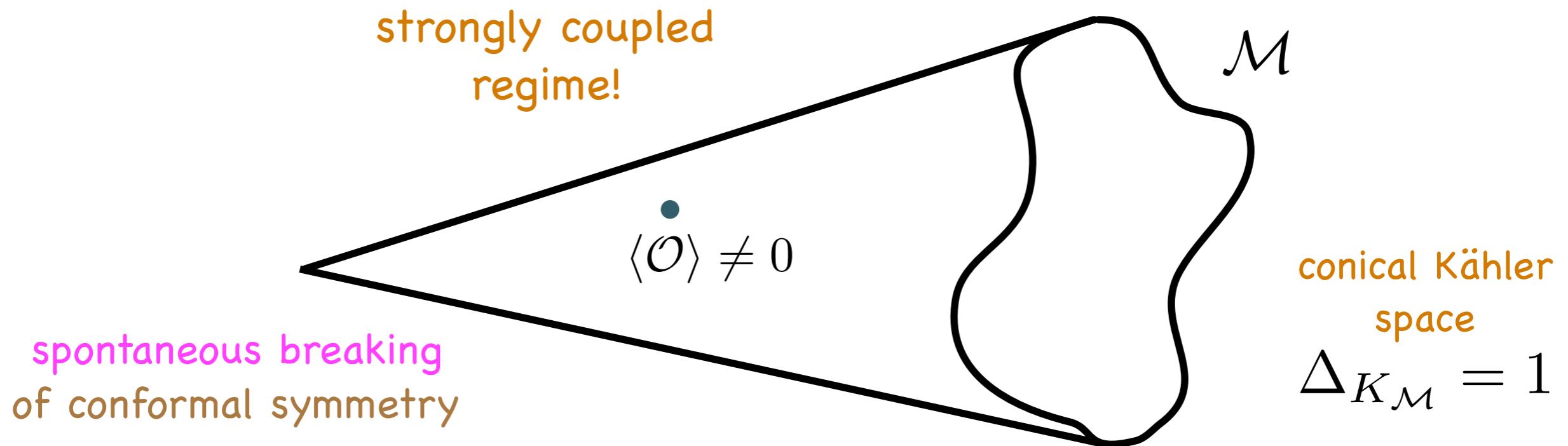
- \* semi-classical analysis

[De Boer-Hori-Oz '97] [Aharony-Hanany-Intriligator-Seiberg-Strassler '97] [Tong '00] ...

[Martelli-Sparks '08] [Hanany-Zaffaroni '08] [Benini-Closset-Cremonesi '13] ... [Intriligator-Seiberg '13] ...

- \* chiral rings and Hilbert series [Cremonesi, Hanany, Makkareya, Zaffaroni ... ]

# $\mathcal{N} = 2$ SCFT moduli space



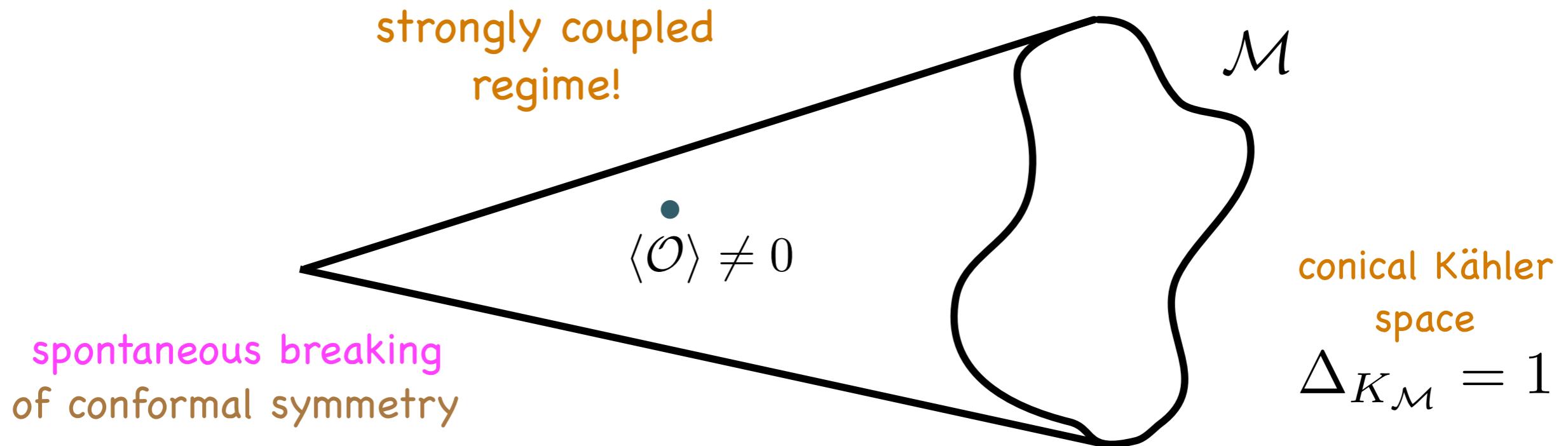
Low energy EFT ?

[De Boer-Hori-Oz '97] [Aharony-Hanany-Intriligator-Seiberg-Strassler '97]  
... [Intriligator-Seiberg '13] ...

$$\mathcal{L}_{\text{eff}} = \int d^4\theta K_{\mathcal{M}}(\Phi, \bar{\Phi}) = -g_{i\bar{j}}(\phi, \bar{\phi}) \partial^\mu \phi^i \partial_\mu \bar{\phi}^{\bar{j}} + \dots$$

non-protected quantity

# $\mathcal{N} = 2$ SCFT moduli space



- Low energy EFT ?

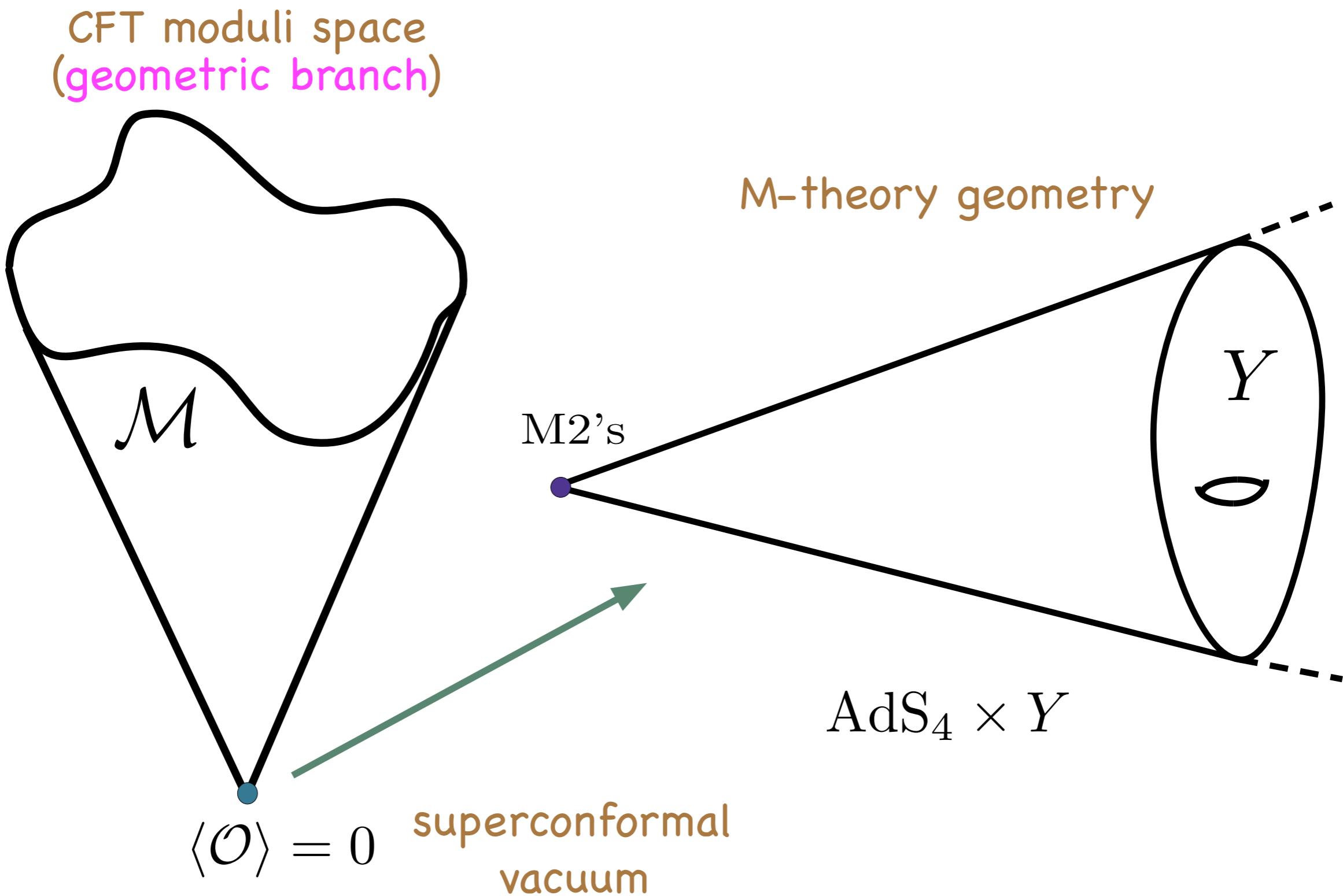
[De Boer-Hori-Oz '97] [Aharony-Hanany-Intriligator-Seiberg-Strassler '97]  
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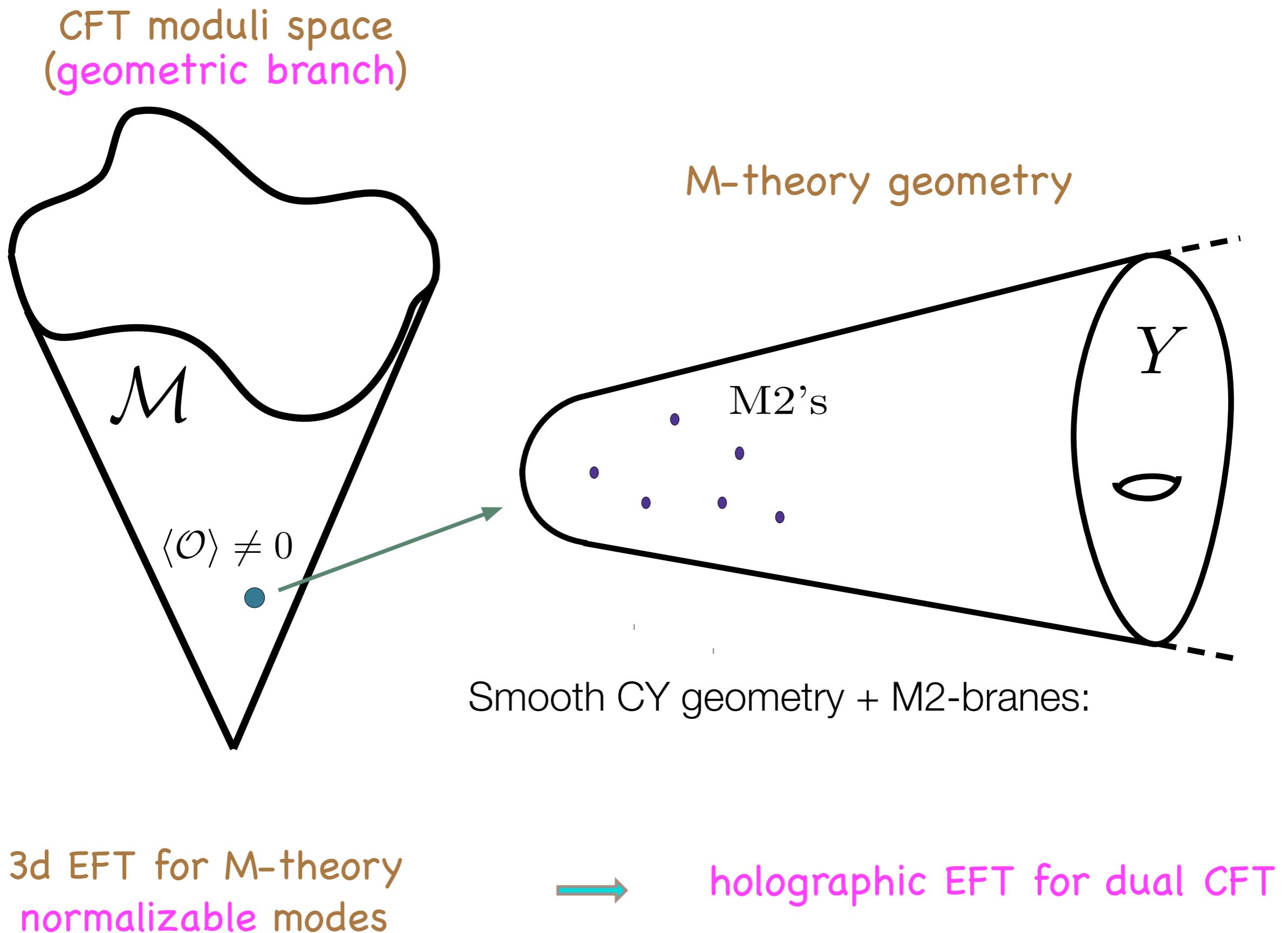
non-protected quantity

- At  $N \gg 1 \longrightarrow$  holographic EFT from supergravity!

# Computation of HEFT



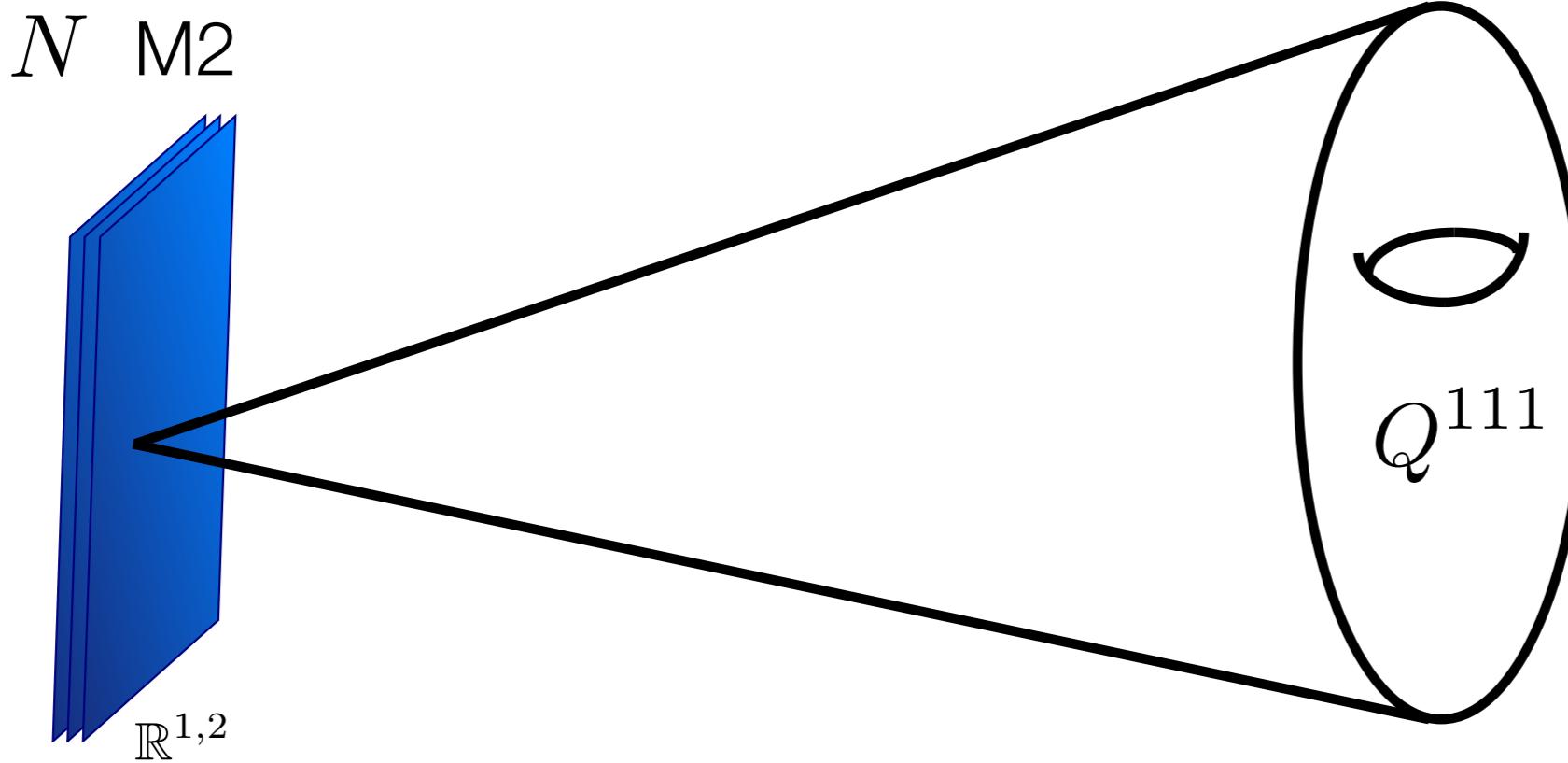
# Computation of HEFT



Example 2:

the  $Q^{111}$  model

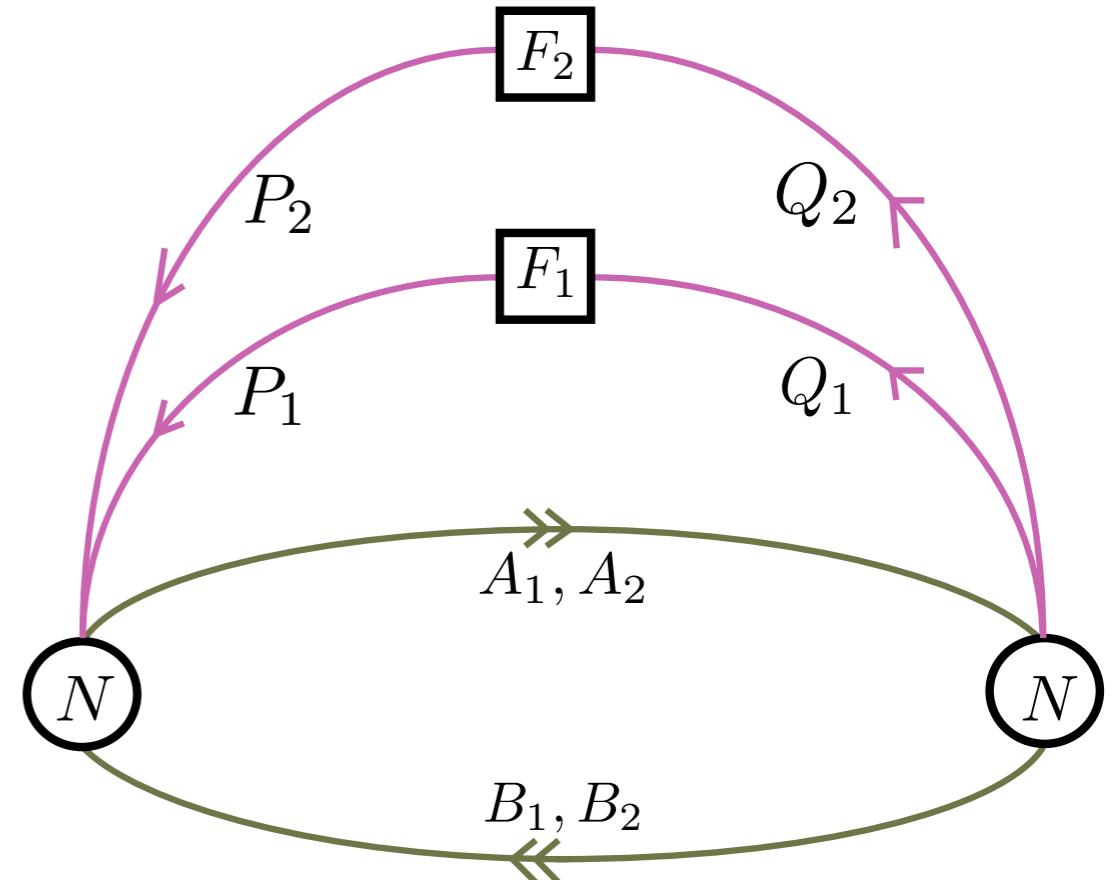
• M-theory origin:



• UV quiver gauge theory:

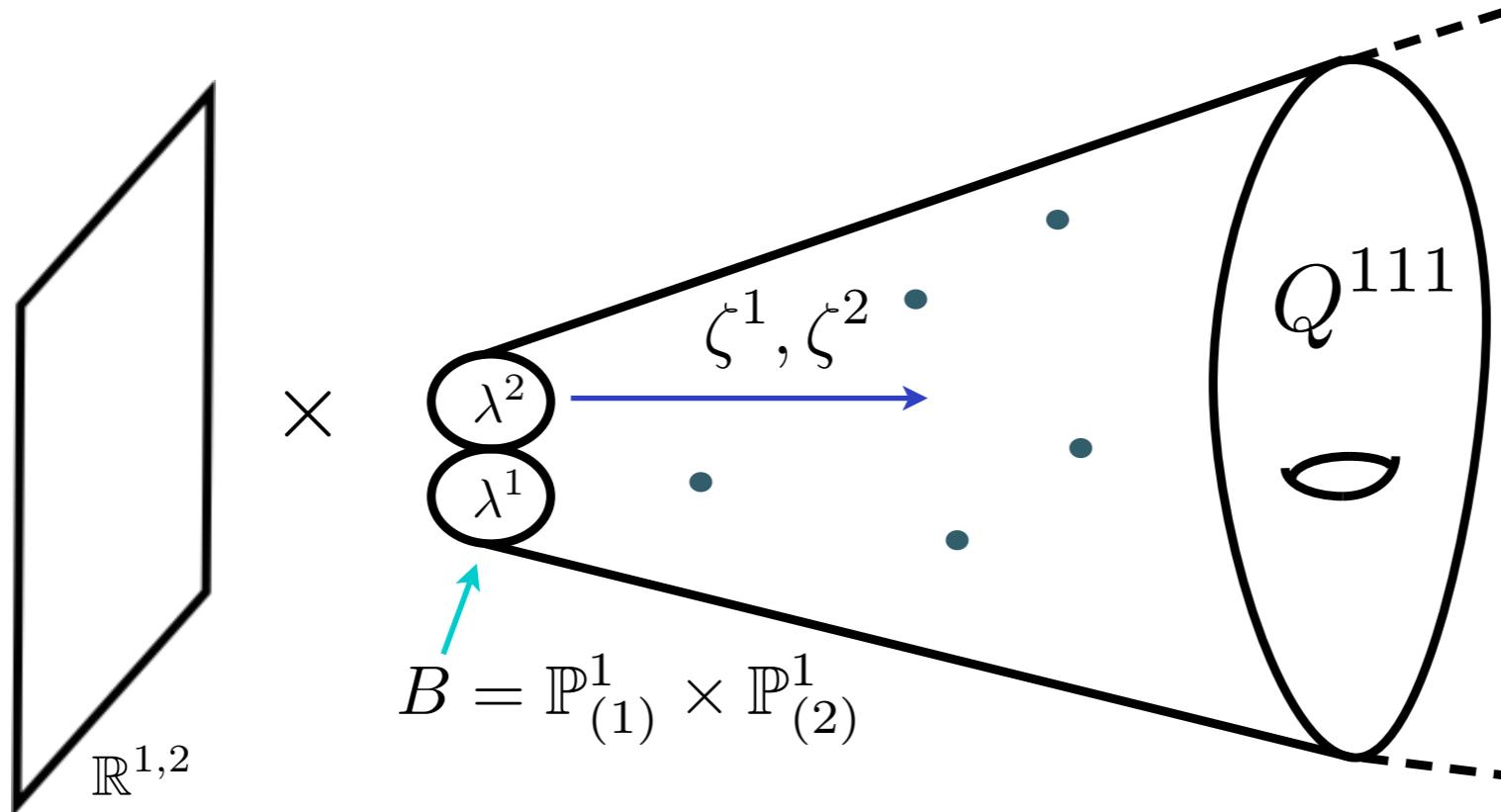
$$* \quad G = \frac{U(N) \times U(N)}{U(1)_B} \times U(1)_{F_1 - F_2}$$

$$* \quad W = \text{Tr}(A_1 B_1 A_2 B_2 - A_1 B_2 A_2 B_1) \\ + P_1 A_1 Q_1 + P_2 A_2 Q_2$$



[Benini-Closset-Cremonesi '09]  
[Jafferis '09]

# HEFT of the $Q^{111}$ model



$$X = \mathcal{L}_B \oplus \mathcal{L}_B$$

$\downarrow$

$$\mathcal{O}_B(-1, -1)$$

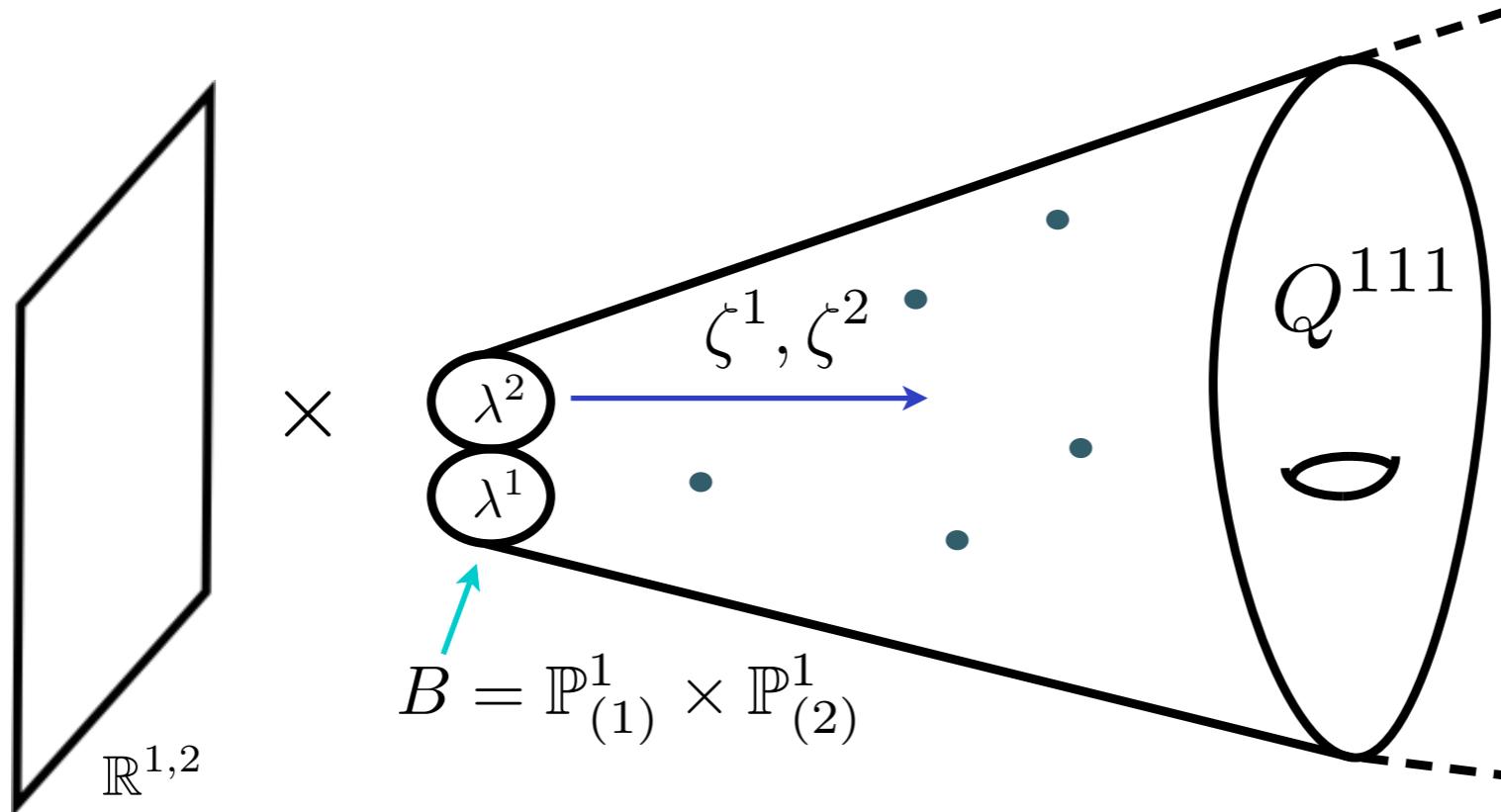
cf. [Cvetič-Gibbons-Lu-Pope '01]  
 [Benishti-Rodríguez Gómez-Sparks '10]

$$z^i = (\lambda^1, \lambda^2, \zeta^1, \zeta^2)$$

$$v_a = \text{vol}(\mathbb{P}^1_a)$$

Moduli	chiral fields	
M2-brane positons	$z_I^i = (\zeta_I^1, \zeta_I^2, \lambda_I^1, \lambda_I^2)$	$I = 1, \dots, N$
$v_1, v_2 + C_6$ -axions	$\rho_1, \rho_2$	

# HEFT of the $Q^{111}$ model



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$$z^i = (\lambda^1, \lambda^2, \zeta^1, \zeta^2)$$

$$v_a = \text{vol}(\mathbb{P}^1_a)$$

- HEFT more easily described in terms of **dual vector multiplets**  $V_a$ :

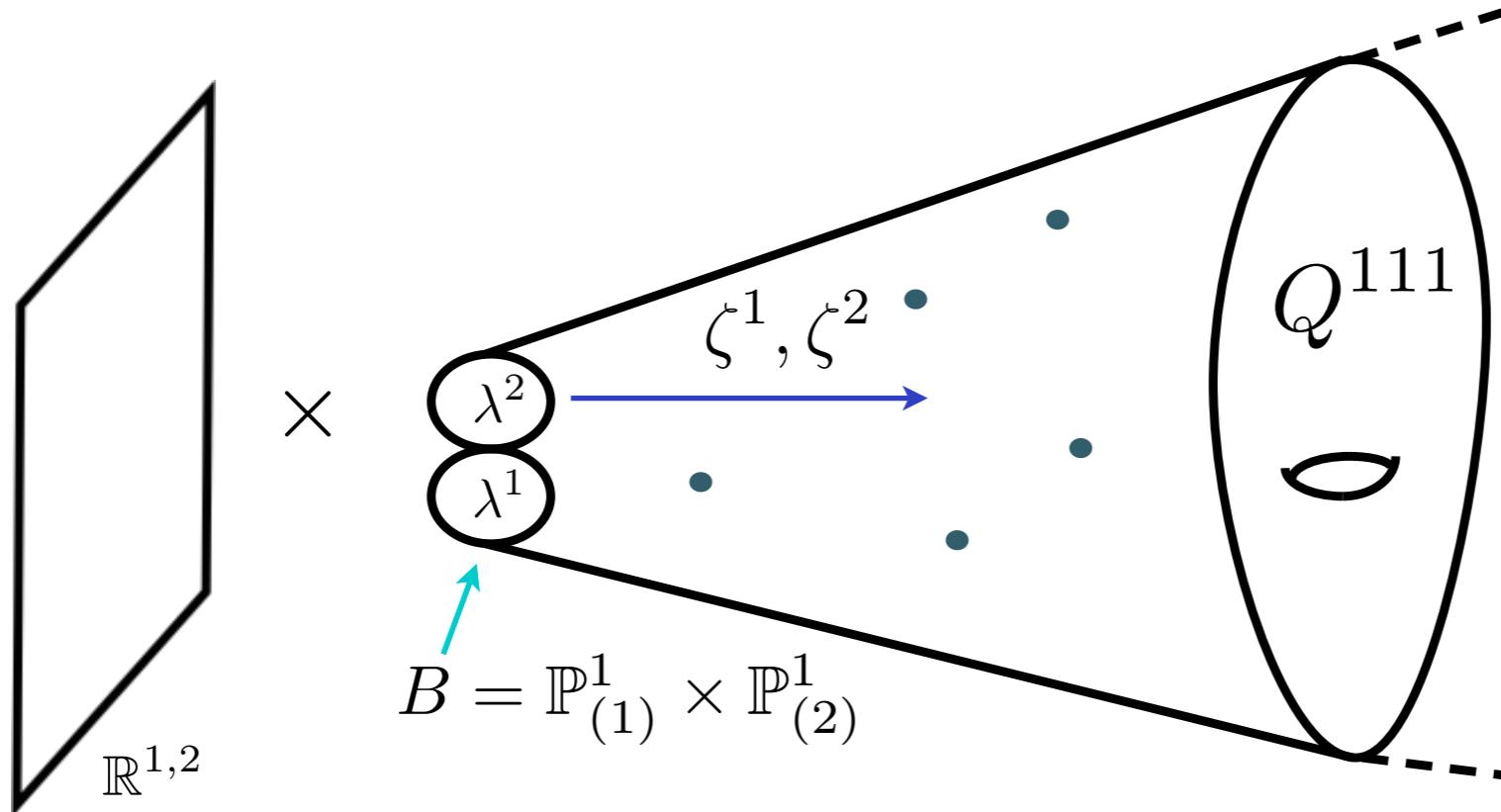
\*  $\mathcal{L}_{\text{HEFT}} = \int d^4\theta \mathcal{F}(z, \bar{z}, \Sigma)$  with

$$\Sigma_a = D\bar{D}V_a = v_a + \dots$$

$$\mathcal{F}(z, \bar{z}, \Sigma) = \sum_{I=1}^N k_X(z_I, \bar{z}_I; \Sigma)$$

$$J_X = i\partial\bar{\partial}k_X$$

# HEFT of the $Q^{111}$ model



$$X = \mathcal{L}_B \oplus \mathcal{L}_B$$

$\downarrow$

$$\mathcal{O}_B(-1, -1)$$

cf. [Cvetič-Gibbons-Lu-Pope '01]  
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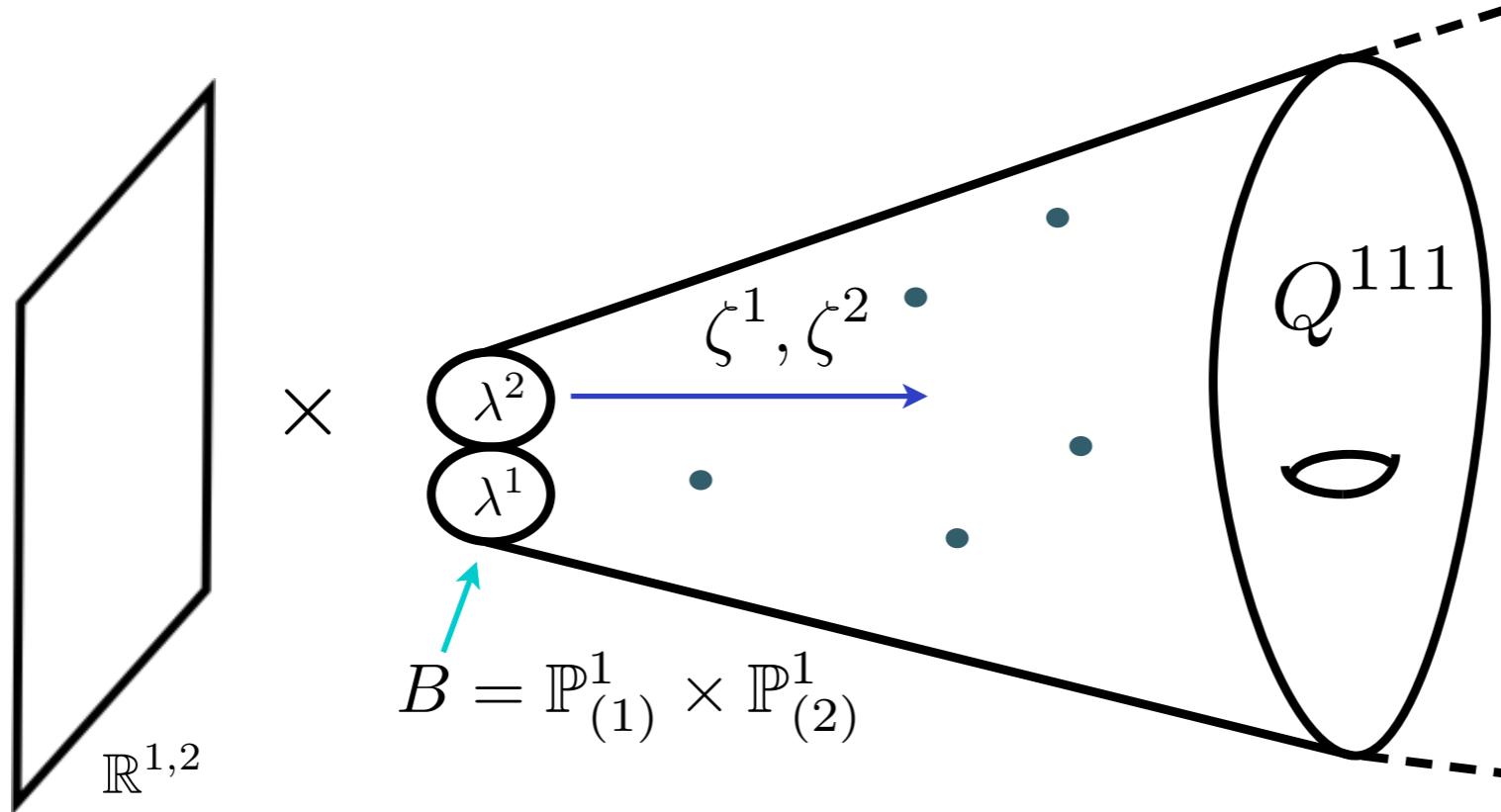
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$$\mathcal{F}(z, \bar{z}, \Sigma) = \sum_{I=1}^N k_X(z_I, \bar{z}_I; \Sigma)$$

- Chiral description:**  
[Hitchin-Karhede-Lindstrom-Rocek '88]  
[de Boer-Hori-Oz '97]

$K_{\mathcal{M}} = \mathcal{F} - \text{Re} \rho^a \Sigma_a$  with  $\text{Re} \rho^a = \frac{\partial \mathcal{F}}{\partial \Sigma_a}$

# HEFT of the $Q^{111}$ model



$$X = \mathcal{L}_B \oplus \mathcal{L}_B$$

$\downarrow$

$$\mathcal{O}_B(-1, -1)$$

cf. [Cvetič-Gibbons-Lu-Pope '01]  
[Benishti-Rodríguez Gómez-Sparks '10]

$$z^i = (\lambda^1, \lambda^2, \zeta^1, \zeta^2)$$

$$v_a = \text{vol}(\mathbb{P}_a^1)$$



Kähler potential:

$$K_{\mathcal{M}} = \sum_{I=1}^N k_0(z_I, \bar{z}_I; v_1, v_2)$$

with  $k_0(z, \bar{z}; v_1, v_2) = s + \frac{A_-^2 - v_1 v_2}{A_+ - A_-} \log \left( 1 + \frac{s}{A_-} \right) - \frac{A_+^2 - v_1 v_2}{A_+ - A_-} \log \left( 1 + \frac{s}{A_+} \right)$

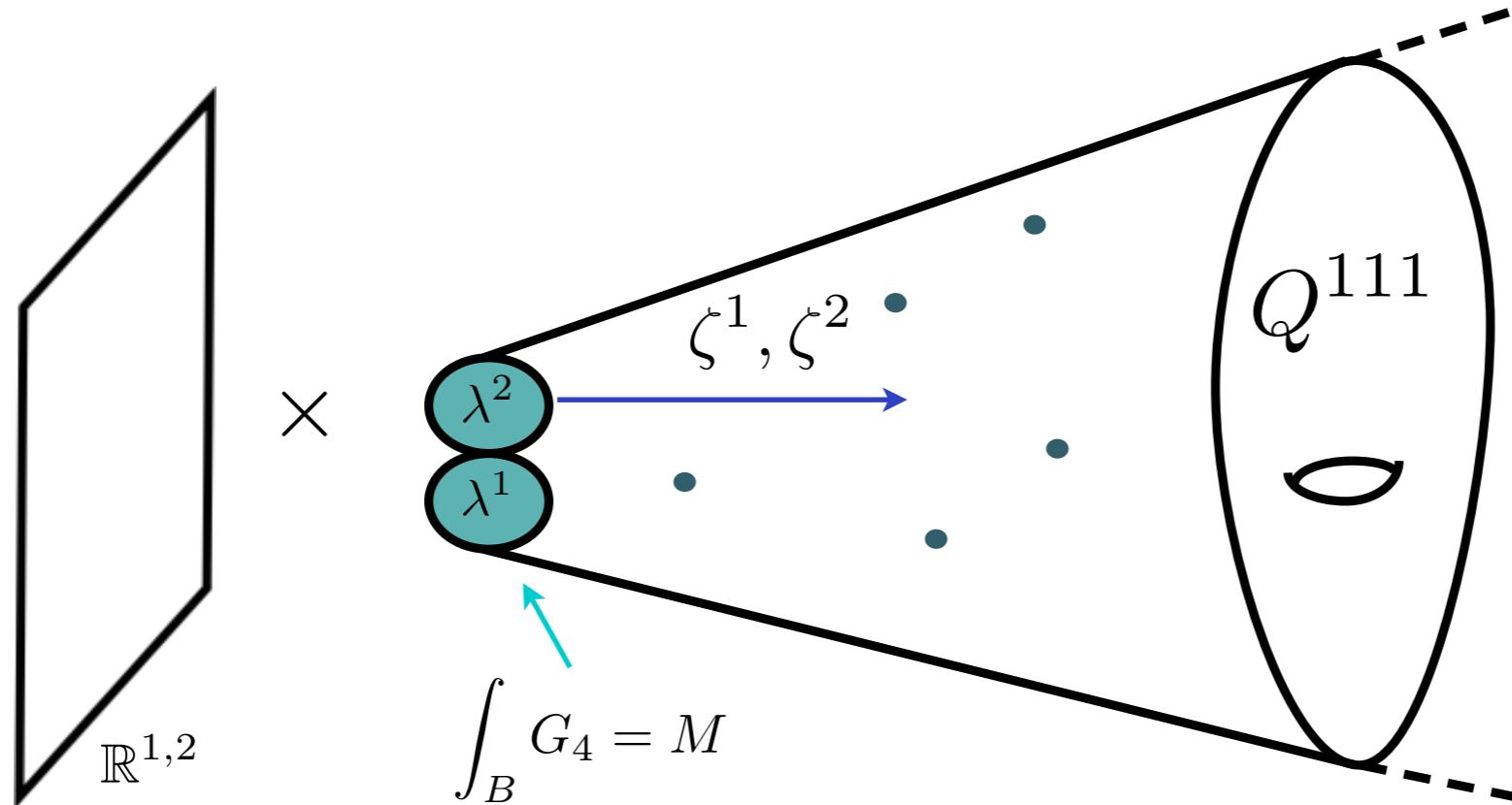
\*  $\text{Re } \rho^1 = \sum_{I=1}^N \left[ \frac{v_2 - A_-}{A_+ - A_-} \log \left( 1 + \frac{s_I}{A_-} \right) - \frac{v_2 - A_+}{A_+ - A_-} \log \left( 1 + \frac{s_I}{A_+} \right) + k_{\mathbb{P}^1}(\lambda_I, \bar{\lambda}_I) \right]$

\*  $A_{\pm} = \frac{1}{3} \left( 2v_1 + 2v_2 \pm \sqrt{4v_1^2 - 10v_1 v_2 + 4v_2^2} \right)$

\*  $s = (|\zeta_1|^2 + |\zeta_2|^2) e^{K_B(\lambda, \bar{\lambda})}$

Holographic (large-N)  
EFT for dual CFT

# HEFT of the $Q^{111}$ model



$$X = \mathcal{L}_B \oplus \mathcal{L}_B$$

$\downarrow$

$$\mathcal{O}_B(-1, -1)$$

cf. [Cvetič-Gibbons-Lu-Pope '01]  
[Benishti-Rodríguez Gómez-Sparks '10]

$$z^i = (\lambda^1, \lambda^2, \zeta^1, \zeta^2)$$

$$v_a = \text{vol}(\mathbb{P}_a^1)$$

- One can turn on an (explicit) **supersymmetric**  $G_4$ -flux

$$* K_{\mathcal{M}}^{\text{flux}} = \sum_{I=1}^N k_0(z_I, \bar{z}_I; v) + \frac{1}{2} \int_X k_0(z, \bar{z}; v) (G_4 \wedge G_4)(z, \bar{z}; v)$$

CFT interpretation under study

cf. [Gukov-Vafa-Witten '99]  
[Cvetič-Gibbons-Lu-Pope '01]  
[Herzog-Klebanov '01]  
[Martelli-Sparks '09]

# Superconformal symmetry

- Definite scaling dimensions:

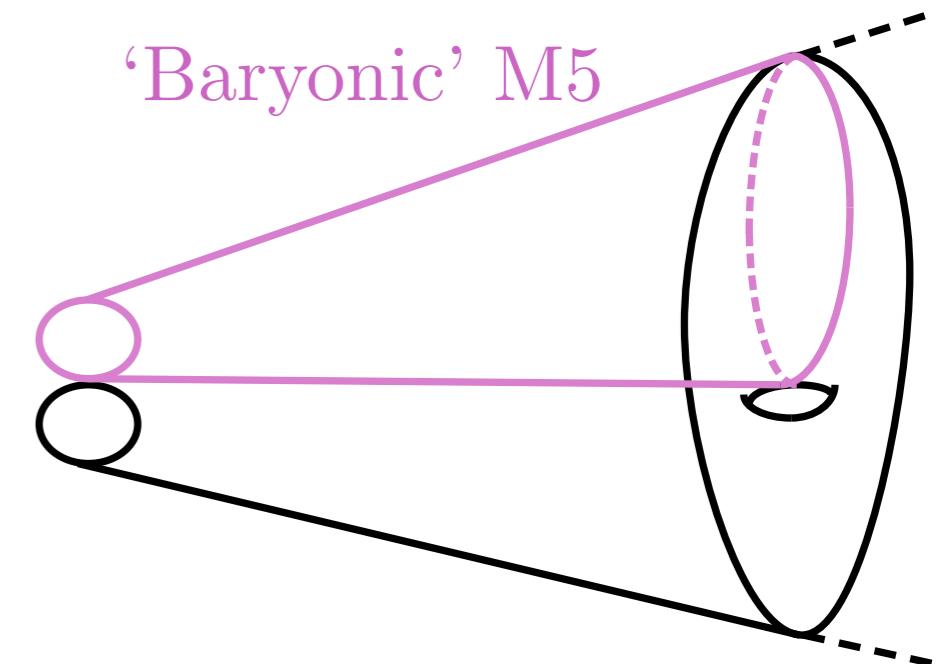
e.g.  $\Delta_{e^{-2\pi\rho}} = \frac{N}{3}$



[Klebanov-Murugan '07]  
[Benishti-Rodríguez Gómez-Sparks '10]

$\langle B^N \rangle \sim e^{-2\pi\rho}$  have correct  $\Delta = \frac{1}{3}N$

[Jafferis, Klebanov,  
Pufu, Safdi '11]

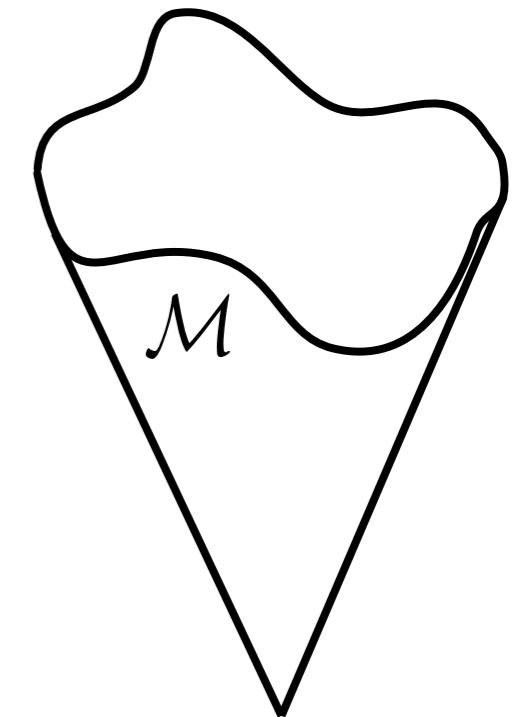


- $\Delta[K_M^{\text{flux}}] = 1$

(non-linearly) realised  
superconformal symmetry!

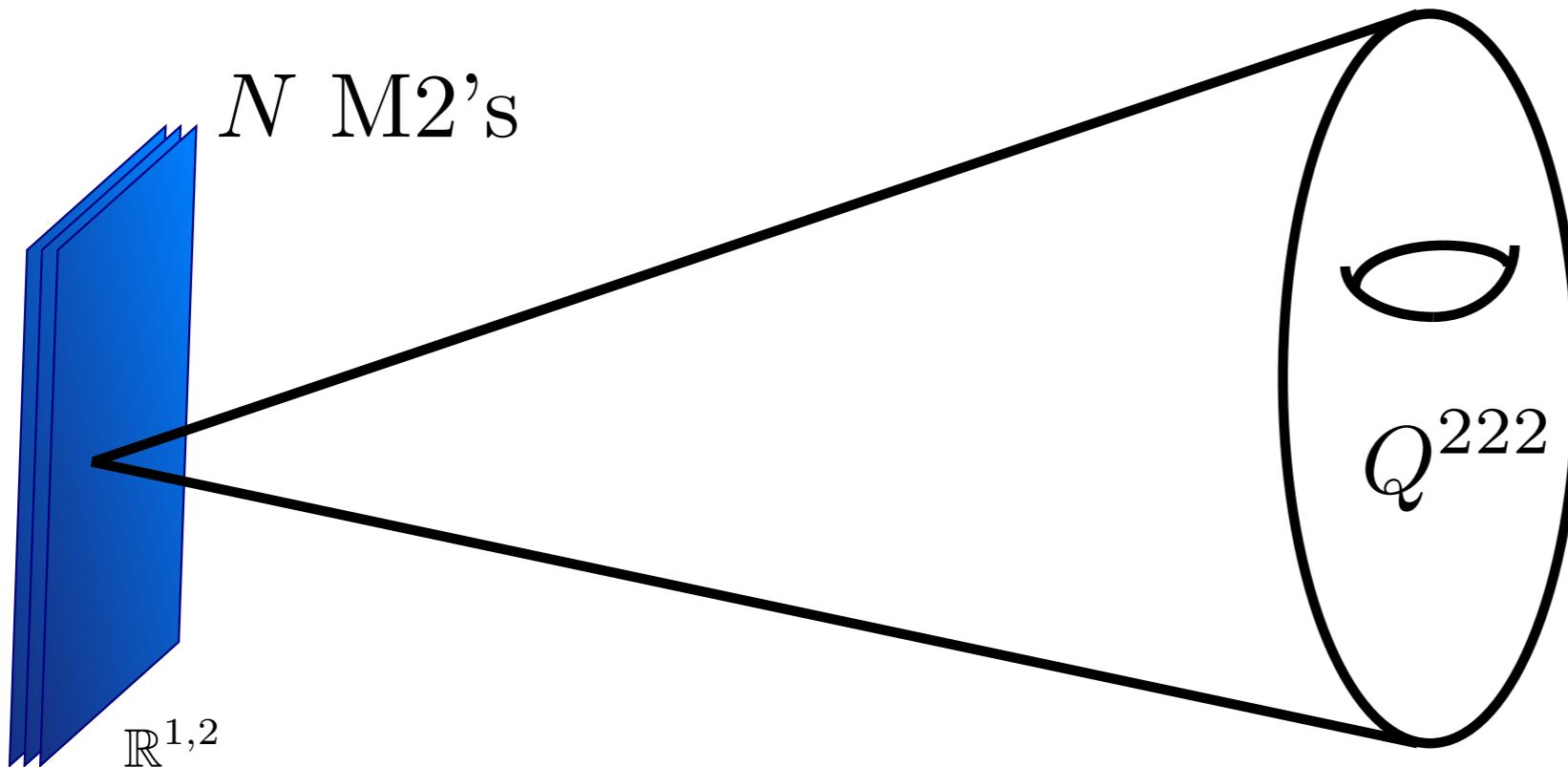


conical Kähler metric  
over  $\mathcal{M}$

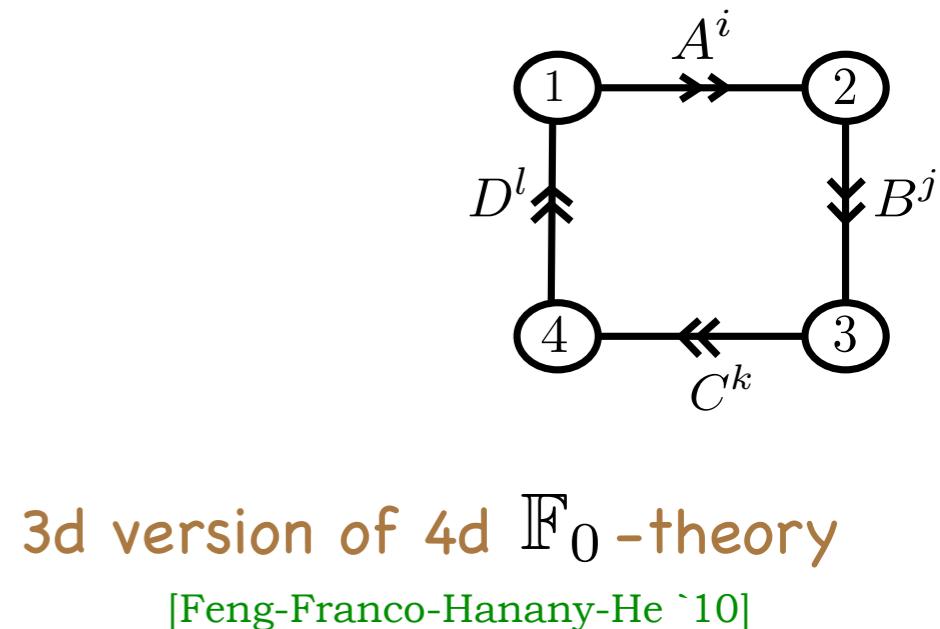


# Non-perturbative (de)stabilisation in $\text{AdS}_4/\text{CFT}_3$

# The $Q^{222}$ model



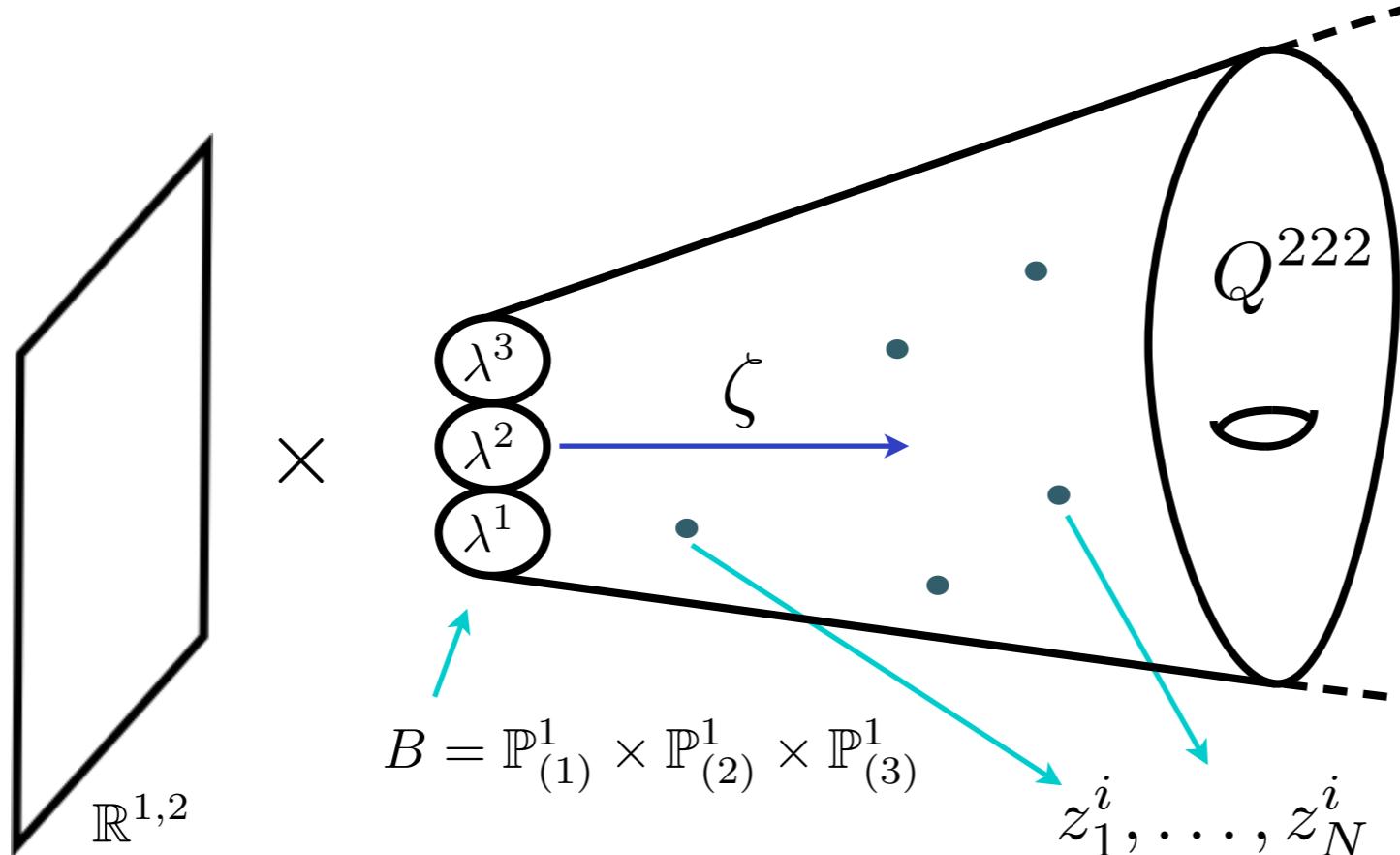
- ✿ UV quiver gauge theory:
- \* gauge group:  $\frac{U(N)_1 \times U(N)_2 \times U(N)_3 \times U(N+2)_4}{U(N)_{B,1} \times U(N)_{B,2}}$
- \* superpotential:  $W = \epsilon_{ik} \epsilon_{jl} \text{Tr}(A^i B^j C^k D^l)$



'Seiberg dual' of 3d quiver proposed in:

[Closset-Cremonesi '12]

# HEFT of the $Q^{222}$ model



$$X = \mathcal{K}_B$$

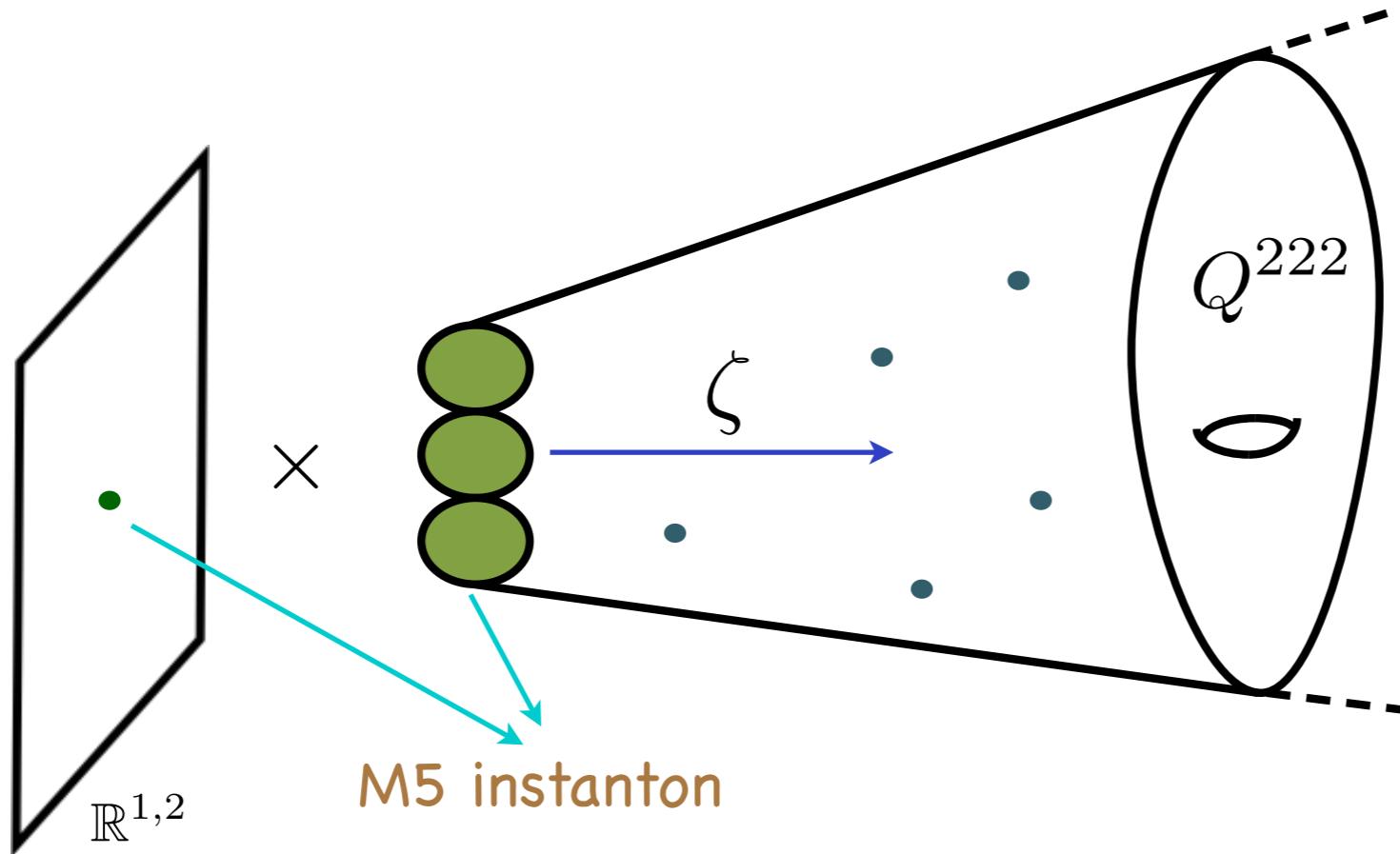
cf [Cvetic-Gibbons-Lu-Pope '01]  
 [Benishti-Rodríguez Gómez-Sparks '10]

$$z^i = (\zeta, \lambda^1, \lambda^2, \lambda^3)$$

Moduli	chiral fields	
M2-brane positons	$z_I^i = (\zeta_I, \lambda_I^1, \lambda_I^2, \lambda_I^3)$	$I = 1, \dots, N$
volumes $\mathbb{P}^1$ 's + $C_6$ -axions	$\rho_1, \rho_2, \rho_3$	

- Explicit perturbative HEFT:  $K_{\mathcal{M}}(\rho, \bar{\rho}, z_I, \bar{z}_I)$

# HEFT of the $Q^{222}$ model



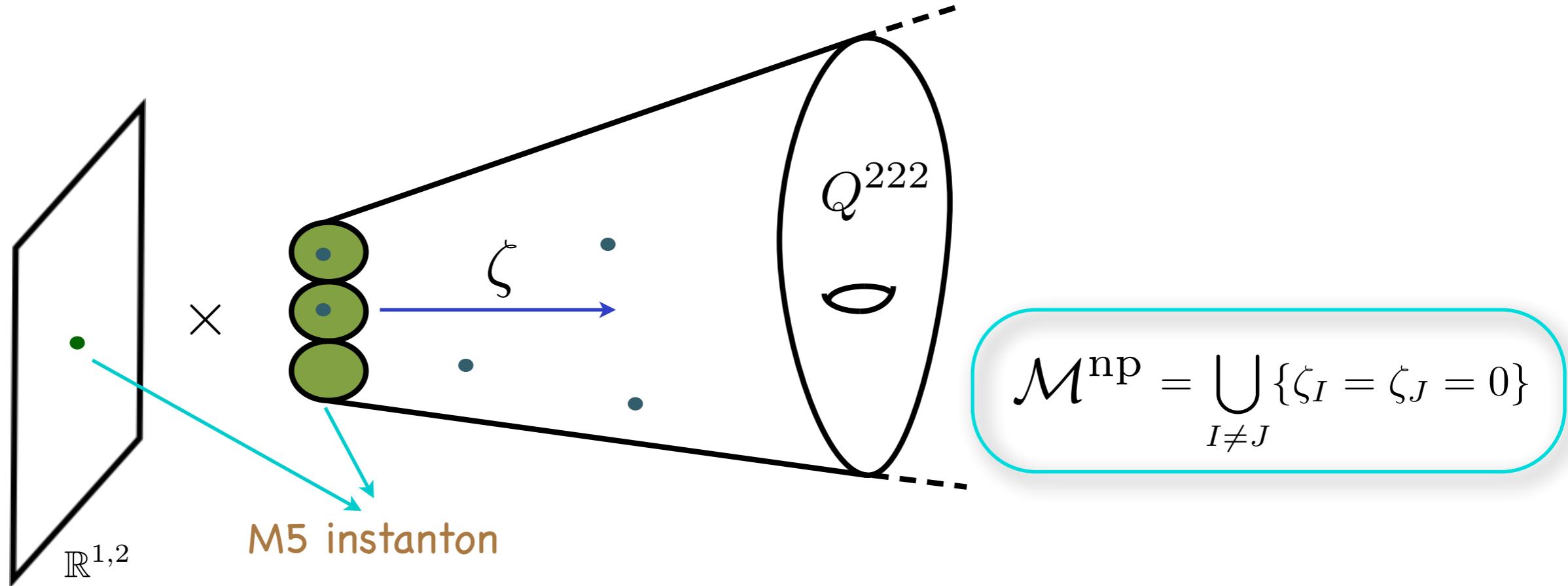
see also [Benishti-Rodríguez Gómez-Sparks `10]

- M5-brane instanton generates a **non-perturbative superpotential**

$$W_{\text{np}} = e^{4\pi(\rho_1 + \rho_2 + \rho_3)} \prod_{I=1}^N \zeta_I$$

cf. [Witten `96] [Ganor `96]  
[Katz-Vafa `96] [Diacunescu-Gukov `98]  
[Baumann-Dymarsky-Klebanov-Maldacena-McAllister-Murugan]

# HEFT of the $Q^{222}$ model



see also [Benishti-Rodríguez Gómez-Sparks '10]

- M5-brane instanton generates a **non-perturbative superpotential**

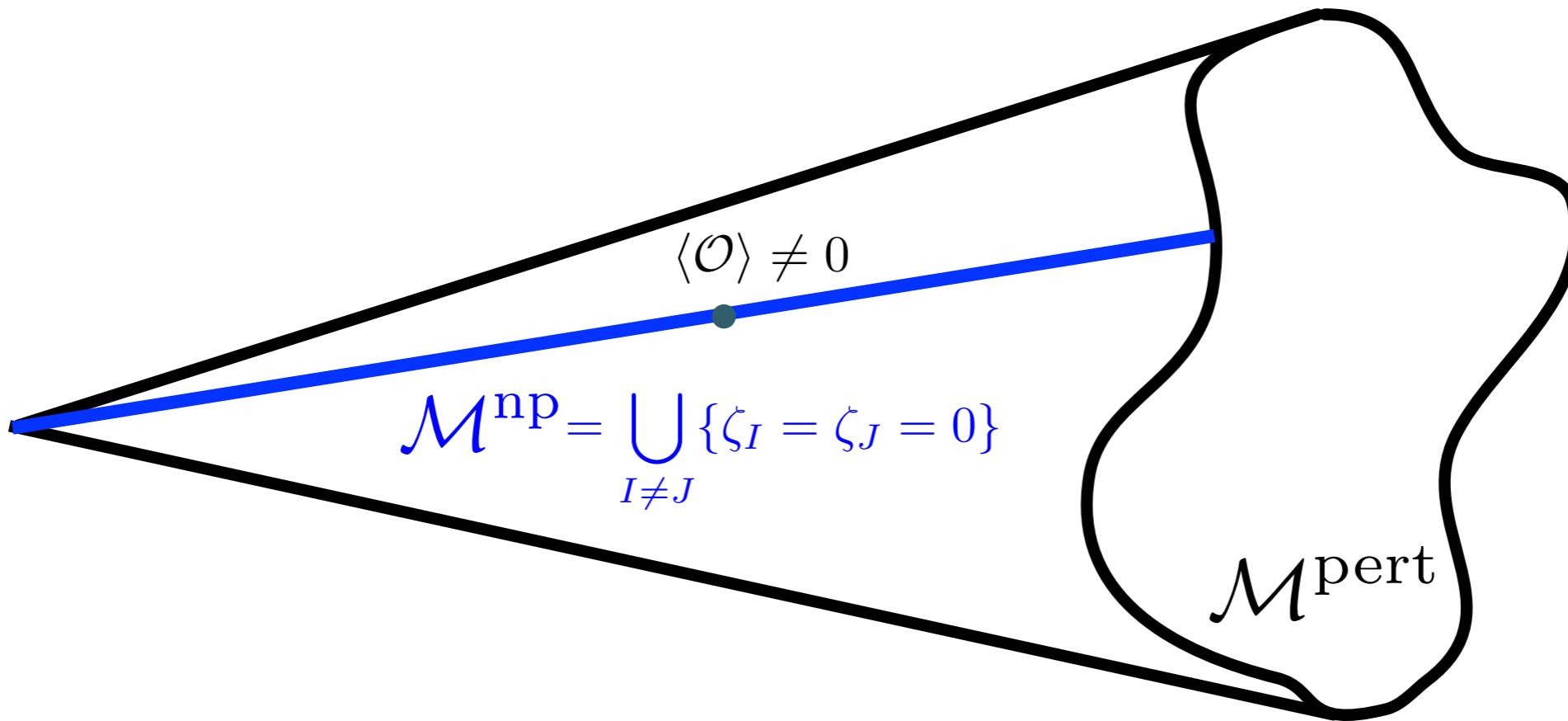
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 [Baumann-Dymarsky-Klebanov-Maldacena-McAllister-Murugan]

- Strong coupling counterpart of weakly coupled SU(2) instanton

[Affleck, Harvey, Witten '82] [De Boer-Hori-Oz '97] [Aharony-Hanany-Intriligator-Seiberg-Strassler '97] ...

# Complete HEFT & moduli space



- \*  $K(\rho, \bar{\rho}, z, \bar{z})$  **perturbative in**  $1/N \ll 1$

- \*  $W_{\text{np}} = e^{4\pi(\rho_1 + \rho_2 + \rho_3)} \prod_{I=1}^N \zeta_I$  **non-perturbative in**  $1/N \ll 1 \sim O(e^{-cN})$



$$V_{\text{eff}} = K^{A\bar{B}} \partial_A W_{\text{np}} \partial_{\bar{B}} \bar{W}_{\text{np}} \neq 0$$

# Summary

- ✿ M-theory derivation of **holographic EFT's** for 3d  $\mathcal{N} = 2$  CFT's
  - \* natural parametrization of bulk+brane moduli
  - \* general formulas for the Kähler potential of bulk+brane moduli
  - \* incorporation of  $G_4$ -fluxes and M5-instantons
- ✿ Explicit formulas for  $Q^{111}, Q^{222}, Y^{12}(\mathbb{P}^2)$  from explicit CY metric.  
**Existence of more efficient computational technics???**

*Thanks*