The Swampland, duality, and de Sitter entropy

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To appear, with Gary Shiu, Hiroshi Ooguri, Cumrun Vafa

Vistas over the Swampland Madrid, Sep 2018 Many string theory settings satisfy a relation of the form

$$\left|\underline{\nabla}V\right| > c V \qquad \qquad c \sim \mathcal{O}(1)$$

See talks by Vafa, ...

It was suggested that this may be a Swampland criterion: the dS conjecture [Obied, Ooguri, Spodyneiko, Vafa '18]

Could there be general physics underlying this behaviour?

A different proposal is the Swampland Distance Conjecture

For supersymmetric settings (at least 8 supercharges) this can be stated most sharply: approaching any infinite distances locus in moduli space there is an infinite tower of states which becomes exponentially light

$$m_{tower} \sim e^{-a \phi}$$
 for $\phi \to \infty$

[Ooguri, Vafa '06]

For at least 8 supercharges, it is one of the most supported of the Swampland conjectures

[Cecotti '15; Grimm, EP, Valenzuela '18; Lee, Lerche, Weigand '18; Grimm, Li, EP '18 (To Appear)] See talks by Valenzuela, Weigand

It can be tied to a Scalar version of the Weak Gravity Conjecture

[EP '17; Grimm, EP, Valenzuela '18]

The distance conjecture has a stronger version. The Refined SDC:

- The exponential behaviour appears at $\phi \sim M_p$
- The exponent $a \sim \mathcal{O}(1)$
- The conjecture holds for fields with a potential $V(\phi)$

[Baume, EP '16; Klaewer, EP '16]

(Implicit to some degree in [Ooguri, Vafa '06])

Has passed some non-trivial tests, but still many open questions

[Baume, EP '16; Klaewer, EP '16; Valenzuela '16; EP '17; Blumenhagen, Valenzuela, Wolf '17; Hebecker, Henkenjohann, Witkowski '17; Cicoli, Ciupke, Mayrhofer, Shukla '18; Blumenhagen, Klaewer, Schlechter, Wolf '18]

See talk by Hebecker

The underlying motivation for the distance conjecture is duality

At large distance there is a dual description in terms of the light states



Couplings in String Theory are scalar fields $(g_s, t, u, g_I, f_I, ...)$

Weak Coupling
$$g o 0$$
 () Large distance $\phi o \infty$

Examples in string theory: <u>Type IIB on Calabi-Yau</u>

• The conifold:

[Strominger '95]

CY Moduli space



• Gopakumar-Vafa invariants (mirror):

[Gopakumar, Vafa '98]

Geometric invariants



1-loop integral of D2-D0 states

• Moduli space at infinite distance:

[Grimm, EP, Valenzuela '18] See talk by Valenzuela

 $g_{\phi\phi}(\partial\phi)^{2}$ $Im N_{IJ} F^{I,\mu\nu}F^{J}_{\mu\nu}$



Infinite tower of D3 branes

[See also Harlow '15; Heidenreich, Reece, Rudelius '17+'18]

We interpret the Swampland Distance Conjecture as:

Any weakly coupled region in string theory should have a dual description in terms of the tower of light states

Propose that can apply this to the potential $V(oldsymbol{\phi})$

Some motivation: In N = 2 a potential can only be induced by gauging

$$Dz^{i} = (\partial z^{i} + k_{I}^{i}(z)A^{I}), \qquad V \sim k(z)$$

If the gauge fields are emergent then so must be the potential

How can we make a general connection between the tower and $V(\phi)$?

de Sitter space has a finite horizon for an observer, of radius R

It is natural to associate to this an entropy

[Gibbons, Hawking '77; Banks '00]

$$S_{dS} = Log \dim \mathcal{H} = R^2$$

We will consider de Sitter vacua or slow-rolling quasi-de Sitter

$$\partial_{\phi} V \sim \alpha V \qquad \alpha \ll 1$$

And identify

$$S_{dS}(\phi) \sim H^{-2} \sim \frac{1}{V(\phi)}$$

<u>Note</u>: we can not say too much away from these regimes. But if $\alpha \sim O(1)$ then anyway the dS conjecture is satisfied.

On the other hand, the distance conjecture says we get many light states



N is the number of fields below some scale Λ_{eff} which we will define

$$\widetilde{N} \sim Exp[a\phi]$$
 $N \sim Exp[b\phi]$ $a, b \sim \mathcal{O}(1)$

The entropy associated to these light states is a function of N and R

$$S_{tower}(N,R)$$

The cutoff Λ_{eff} is then defined as the scale above which the states no longer contribute significantly to the entropy

Since $N, R \gg 1$, we expect $S_{tower}(N, R)$ to be dominated by a single term

$$S_{tower}(N,R) \sim N^p R^q$$

The duality proposal implies that this should account for the dS entropy

$$N^p R^q \sim R^2$$



We therefore recover the de Sitter conjecture.

This is a generalisation of a Dine-Seiberg type argument to arbitrary couplings [Dine, Seiberg '85]

<u>Note</u>: Even if the tower only accounts for a fraction of the total entropy, still there are no de Sitter vacua, but could have $c \ll 1$

Recall our assumptions:

- The SDC holds for potentials
- We are in a weakly-coupled regime where the tower is a dual description
- We are in a quasi-de Sitter setting (accelerating expansion with horizon)

⇒ The de Sitter conjecture

We are not sure how to determine p, q, Λ_{eff}

The problem amounts to counting microstates in the dual theory

We will later parameterise this ignorance, but first let us try..

Recall, in de Sitter space have 3 types of states:

- QFT states localised within the bulk of de Sitter space
- Black holes
- States localised on the Horizon

One proposal for counting the microstates is to consider all localised QFT states within a box of order the de Sitter radius R

Consider a relativistic particle in a box of size R

We want to calculate the entropy associated to it by counting its momentum modes up to a maximum value k_{max}

$$S_1 \sim (k_{max}R)^3$$

The maximum energy associated to these modes is

$$E_1 \sim k_{max} (k_{max} R)^3$$

We count only localised QFT states, so not black holes, therefore $E_1 < R$

$$k_{max} < R^{-\frac{1}{2}} \qquad S_1 < R^{\frac{3}{2}}$$

[Page '81; Banks '05]

Now we consider N species of such particles

We need these to be relativistic to contribute significantly to the entropy, so we have

$$\Lambda_{eff} \sim k_{max}$$

Can deduce the number of species needed by simple thermodynamics

To maximise entropy set temperature in each species equal $(k_{max} \rightarrow T)$

$$S_N \sim N(TR)^3$$
 $E_N \sim NT(TR)^3$

Not forming black holes implies

$$T \leq \frac{1}{\frac{1}{N^{\frac{1}{4}}R^{\frac{1}{2}}}} \qquad S_N \sim N^{\frac{1}{4}}R^{\frac{3}{2}}$$

We therefore find that we need

$$N \sim R^2$$
 $T \sim \frac{1}{R}$ $S_1 \sim 1$

The low temperature and entropy per species means at borderline of thermodynamics, but can explicitly check by counting microstates



While the de Sitter conjecture is not sensitive to the microstate counting scheme, the cosmology is

How would these bounds apply to our universe?

 $R \sim 10^{60} \qquad \qquad c \lesssim 0.6$

[Agrawal, Obied, Steinhardt, Vafa '18]

We assume an equally-spaced tower of states $m_n \sim n m$

$$N \sim R^{\frac{2-q}{p}} \qquad \qquad \Lambda_{eff}^{(1)} \equiv 1 \qquad \qquad \Lambda_{eff}^{(2)} \equiv \frac{1}{\sqrt{N}} \sim R^{\frac{q-2}{2p}}$$
$$m^{(1)} \equiv \frac{\Lambda_{eff}^{(1)}}{N} \sim R^{\frac{q-2}{p}} \qquad \qquad m^{(2)} \equiv \frac{\Lambda_{eff}^{(2)}}{N} \sim R^{\frac{3(q-2)}{2p}}$$

• **QFT counting:** $q = \frac{3}{2}, p = \frac{1}{4}$

$$N \sim 10^{120}$$
 $\Lambda_{eff} \sim \Lambda_{eff}^{(2)} \sim 10^{-60}$ $m \sim 10^{-180}$

 $\Lambda_{eff}^{(2)} \sim 10^{-7}$

• Scenario 1: $q = \frac{3}{2}, p = 1$

$$N \sim 10^{30}$$
 $\Lambda_{eff}^{(2)} \sim 10^{-15} \sim TeV$

$$m^{(1)} \sim 10^{-30} \sim m_{\nu}$$
$$m^{(2)} \sim 10^{-45} \sim 10^{-15} m_{\nu}$$

• Scenario 2:
$$q = \frac{7}{4}$$
, $p = 1$

 $N \sim 10^{15}$

 $m^{(1)} \sim 10^{-15} \sim TeV$ $m^{(2)} \sim 10^{-22} \sim MeV$

• **Decoupling scenario**: q = 2, p = p

Summary

- Proposed that Weakly-coupled regions in string theory admit a dual description, including the potential, by the tower of light states of the Swampland Distance Conjecture
- In a quasi-de Sitter setting, the duality applied to the entropy leads to the Swampland de Sitter conjecture

• We have calculated an example approach to counting microstates

• For that approach, we find that our universe would not be in the weakly coupled regime. For other microstate counting recover varied cosmology.

Thank You