

Classical de Sitter solutions and the swampland

David ANDRIOT

CERN, Geneva, Switzerland

Based on [arXiv:1609.00385](#) (with J. Blåbäck), [1710.08886](#)
[arXiv:1806.10999](#), [1807.09698](#)

Vistas over the Swampland
20/09/2018, IFT-Madrid, Spain

Introduction

Stringy de Sitter

Constraints
class. de Sitter

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 \Rightarrow constraints on cosmological models.
Still, many models $\checkmark \Rightarrow$ theoretical constraints?

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Requirement: can the model be derived/is compatible with a
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“Modern” formulation: is the cosmological model in the
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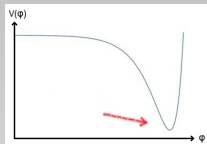
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Relation to de Sitter solutions (having 4d de Sitter space-time):
“mainstream” models (\checkmark with observ.) exhibit de Sitter
solutions:

- present/future universe
- end-point of inflation
- \sim inflation phase



(Meta)stable or slightly unstable solutions.

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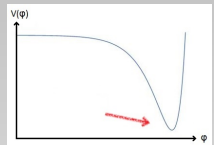
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\Rightarrow Get de Sitter solutions from quantum gravity?

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Idea behind recent “**de Sitter swampland criterion**”

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For a 4d theory of minimally coupled scalars ϕ_i ($M_4 = 1$)

$$\mathcal{S} = \int d^4x \sqrt{|g_4|} (\mathcal{R}_4 + \text{kin. terms} - V(\phi_i))$$

solutions as extrema of potential: $\partial_{\phi_i} V|_0 = 0$, $\mathcal{R}_4 = 2V|_0$
⇒ de Sitter solutions: $\Lambda_4 = \frac{1}{2}V|_0 > 0$.

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Criterion: if NOT in the swampland, one has:

$$|\nabla V| \geq cV$$

with $c > 0$, $|\nabla V| = \sqrt{g^{ij} \partial_{\phi_i} V \partial_{\phi_j} V}$.

Extremum: $|\nabla V|_0 = 0 \Rightarrow V|_0 \leq 0$: no de Sitter solution.

Conclusion:

- Requirement (can model be derived from quant. gravity?) is fruitful, constraining.
- Problem of de Sitter solutions in string theory is central.

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Plan:

- Comments on de Sitter solutions in string theory and the swampland
- Constraints on classical de Sitter solutions

De Sitter solutions in string theory

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Recent review:

[U. H. Danielsson, T. Van Riet, \[arXiv:1804.01120\]](#)

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Complicated interplay between quantum gravity (**10d** supergravity/string theory) and cosmological model (**4d** low energy effective theory)

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Stringy constructions to get de Sitter:

- “4d approaches”

- “10d approach”

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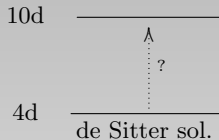
- “4d approaches”: (at some point) include an ingredient at 4d level \Rightarrow controlled realisation in string theory?

Prime example: KKLT: include 4d non-perturbative contributions, 4d potential term for \overline{D}_3 ...

Non-geometric: many different non-geometric fluxes \Rightarrow realise as a non-trivial stringy background?

NS Bianchi identities not satisfied: exotic sources?
(some 4d non-geometric fluxes lifted to an understood 10d realisation)

- “10d approach”



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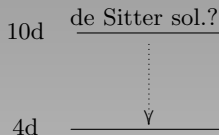
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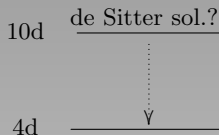
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No such solution in heterotic string.

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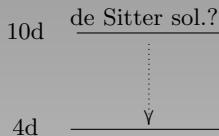
S. R. Green, E. J. Martinec, C. Quigley, S. Sethi [arXiv:1110.0545]

In type IIA/B: not excluded but very constrained

Intrinsically difficult

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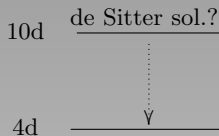
In type IIA/B: not excluded but very constrained

Intrinsically difficult: e.g. parallel D_6/O_6 :

$$\begin{aligned} \Lambda_4 = & -\frac{e^{2A}}{8} \left(*_{\perp} H|_{\perp} + e^{\phi} F_0 \right)^2 - \frac{e^{2A}}{4} \left(e^{4A} *_{\perp} de^{-4A} - e^{\phi} F_2^{(0)} \right)^2 \\ & - \frac{e^{2A}}{8} \sum_{a_{\parallel}} \left(*_{\perp} de^{a_{\parallel}}|_{\perp} - e^{\phi} (\iota_{a_{\parallel}} F_2^{(1)}) \right)^2 \\ & - \frac{e^{2A+2\phi}}{8} \left(2(F_4^{(0)})^2 + 2(F_4^{(1)})^2 + (F_4^{(2)})^2 + (F_6)^2 \right) \\ & + \frac{e^{2A}}{8} \left(-2\mathcal{R}_{\parallel} - 2\mathcal{R}_{\parallel}^{\perp} + (H^{(2)})^2 + 2(H^{(3)})^2 \right) \end{aligned}$$

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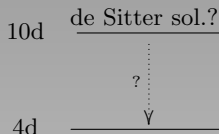
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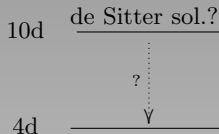
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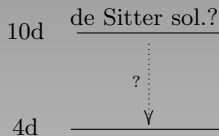
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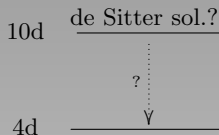
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3 points in favor of de Sitter swampland criterion.

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Criticism on these solutions (smeared O-planes, Romans mass, flux quantization/large volume/small coupling, etc.)

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\hookrightarrow refine de Sitter swampland criterion

$\exists b_i \in \mathbb{R}, c_i \in \mathbb{R}_+$ such that

$$V + \sum_i b_i \phi_i \partial_{\phi_i} V + \sum_i c_i \phi_i^2 \partial_{\phi_i}^2 V \leq 0$$

$$\text{Solution: } V|_0 + \sum_i c_i (\phi_i^2 \partial_{\phi_i}^2 V)|_0 \leq 0$$

\Rightarrow **no stable de Sitter solution**, tachyonic de Sitter sol. \checkmark

D. Andriot [arXiv:1806.10999]

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Replace $\partial_{\phi_i} V$ term by a power of $|\nabla V|$.

Single field and $V > 0$: refined criterion becomes:

$$\sqrt{\epsilon_V} - a \eta_V \geq c, \quad \text{with } a \geq 0, c > 0$$

D. Andriot [arXiv:1806.10999]

\Rightarrow checks? Cosmological implications?

S. K. Garg, C. Krishnan [arXiv:1807.05193]

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Replace $\partial_{\phi_i} V$ term by a power of $|\nabla V|$.

Single field and $V > 0$: refined criterion becomes:

$$\sqrt{\epsilon_V} - a \eta_V \geq c, \quad \text{with } a \geq 0, c > 0$$

D. Andriot [arXiv:1806.10999]

\Rightarrow checks? Cosmological implications?

S. K. Garg, C. Krishnan [arXiv:1807.05193]

Interesting general statements.

The whole situation now requires more precise studies.

Less example based and more **analytical** results,
understanding.

\Rightarrow constraints on 10d classical de Sitter solutions

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Constraints on 10d classical de Sitter solutions

D. Andriot [[arXiv:1807.09698](https://arxiv.org/abs/1807.09698)]

Reasoning and known results

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M. P. Hertzberg, S. Kachru, W. Taylor, M. Tegmark [arXiv:0711.2512],

E. Silverstein [arXiv:0712.1196]

$$\begin{array}{l|l} \mathcal{R}_4 = \dots > 0 & \\ \partial_\rho V|_0 = 0 & \\ \partial_\tau V|_0 = 0 & \end{array}$$

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Combine 3 equations, get constraints

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Existence of sol.: necessary ingredients (parallel D_p/O_p)

T. Wrase, M. Zagermann [arXiv:1003.0029], G. Shiu, Y. Sumitomo [arXiv:1107.2925]

D. Andriot, J. Blåbäck, [arXiv:1609.00385]

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	A de Sitter solution requires $T_{10} > 0$ and	
$p = \dots$	$\mathcal{R}_6 \geq 0$	$\mathcal{R}_6 < 0$
3	F_1, H	nothing
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5		F_1
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Considering on top the 10d sourced Bianchi identity

$$dF_{8-p} - H \wedge F_{6-p} = \varepsilon_p \frac{T_{10}}{p+1} \text{vol}_\perp$$

\Rightarrow more constraints, more ingredients needed

J. Blåbäck, U. H. Danielsson, D. Junghans, T. Van Riet, T. Wrase, M. Zagermann

[arXiv:1009.1877], D. Andriot, J. Blåbäck, [arXiv:1609.00385]

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Analogous results derived for intersecting D_p/O_p :
slightly less constraining \leftrightarrow solutions known.

G. Shiu, Y. Sumitomo [arXiv:1107.2925], D. Andriot [arXiv:1710.08886]

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3 equations + sourced Bianchi identity

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$p = 4$: Different Bianchi identity: $dF_{4-p} = dF_0 = 0$.

F_0 is a scalar $\Rightarrow F_0$ constant.

But $F_0 \rightarrow -F_0$ under O_4 projection $\Rightarrow F_0 = 0$

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Introduce **new scalar** field σ : distinguishes \parallel, \perp dir. of D_p/O_p

U. H. Danielsson, G. Shiu, T. Van Riet, T. Wrase [arXiv:1212.5178]

Gives new constraints through

$$\partial_\sigma V|_0 = 0 \quad | \quad |$$

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Reproduces constraints obtained in 10d + new ones

$$\begin{aligned}
 V(\rho, \tau, \sigma) = & -\tau^{-2} \left(\rho^{-1} \mathcal{R}_6(\sigma) - \frac{1}{2} \rho^{-3} \sum_n \sigma^{-An-B(3-n)} |H^{(n)}|^2 \right) \\
 & - g_s \tau^{-3} \rho^{\frac{p-6}{2}} \sigma^{B\frac{p-9}{2}} \frac{T_{10}}{p+1} \\
 & + \frac{1}{2} g_s^2 \left(\tau^{-4} \sum_{q=0}^4 \rho^{3-q} \sum_n \sigma^{-An-B(q-n)} |F_q^{(n)}|^2 - \tau^4 \rho^3 |F_6|^2 \right. \\
 & + \frac{1}{2} \sum_n (\tau^{-4} \rho^{-2} \sigma^{-An-B(5-n)} |F_5^{(n)}|^2 \\
 & \left. - \tau^4 \rho^2 \sigma^{-An-B(1-n)} |(*_6 F_5)^{(n)}|^2 \right)
 \end{aligned}$$

$$\begin{aligned}
 \text{with } \mathcal{R}_6(\sigma) = & \sigma^{-B} \left(\mathcal{R}_\perp + \delta^{ab} \partial_{a_\perp} f^{c_\perp}_{c_\perp b_\perp} + \mathcal{R}_\perp^\parallel + |f^\parallel_{\perp\perp}|^2 \right) \\
 & + \sigma^{-A} \left(\mathcal{R}_\parallel + \delta^{ab} \partial_{a_\parallel} f^{c_\parallel}_{c_\parallel b_\parallel} + \mathcal{R}_\parallel^\perp + |f^\perp_{\parallel\parallel}|^2 \right) \\
 & - \frac{1}{2} \sigma^{-2A+B} |f^\perp_{\parallel\parallel}|^2 - \frac{1}{2} \sigma^{-2B+A} |f^\parallel_{\perp\perp}|^2
 \end{aligned}$$

and $A = p - 9$, $B = p - 3$

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On **group manifolds** (with constant fluxes): constraints in
terms of $\lambda = -\frac{\delta^{cd} f^{b\perp}{}_{a\parallel} c_{\perp} f^{a\parallel}{}_{b\perp} d_{\perp}}{\frac{1}{2} \delta^{ab} \delta^{cd} \delta_{ij} f^{i\parallel}{}_{a\perp} c_{\perp} f^{j\parallel}{}_{b\perp} d_{\perp}}$:

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No classical 10d de Sitter solution (with parallel D_p/O_p) for $\lambda \leq 0$ or $\lambda \geq 1$

\Leftrightarrow no de Sitter solution on nilmanifold, semi-simple group manifold, some solvmanifolds (in standard basis).

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Stability island:

a de Sitter solution would have $\partial_{\rho}^2 V|_0 > 0$, $\partial_{\tau}^2 V|_0 > 0$, $\partial_{\sigma}^2 V|_0 > 0$ for

$$0 < \lambda < \frac{1}{17} \quad \text{with } p = 6$$

$$0 < \lambda < \frac{1}{10} \quad \text{with } p = 4, 5$$

On **group manifolds** (with constant fluxes): constraints in terms of $\lambda = -\frac{\delta^{cd} f^{b\perp}{}_{a\parallel} c_{\perp} f^{a\parallel}{}_{b\perp} d_{\perp}}{\frac{1}{2} \delta^{ab} \delta^{cd} \delta_{ij} f^i{}_{a\perp} c_{\perp} f^j{}_{b\perp} d_{\perp}}$:

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\Rightarrow Tiny region of parameter space where

possible stable de Sitter solutions (with parallel D_p/O_p)

\hookrightarrow explore!

Remark on the Bianchi identity and the swampland

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Remark on the Bianchi identity and the swampland

$p = 3$, no flux, ρ, τ

$$\mathcal{S} = \int d^4x \sqrt{|g_4|} (\mathcal{R}_4 + \text{kin. terms} - V)$$

with
$$V(\rho, \tau) = -\tau^{-2} \rho^{-1} \mathcal{R}_6 - \tau^{-3} \rho^{-\frac{3}{2}} g_s \frac{T_{10}}{4}$$

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3 equations of motion \Rightarrow de Sitter solution with

$$\mathcal{R}_4 = -\frac{2}{3} \mathcal{R}_6 = g_s \frac{T_{10}}{4} > 0$$

\Rightarrow landscape of 4d theories (\mathcal{R}_6, T_{10}) , some with de Sitter solutions.

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Bianchi identity: 10d origin, added by hand in 4d:
“quantum gravity origin stamp”:

$$dF_{8-p} - H \wedge F_{6-p} = \varepsilon_p \frac{T_{10}}{p+1} \text{vol}_\perp$$

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No flux $\Rightarrow T_{10} = 0$

Remark on the Bianchi identity and the swampland

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$$dF_{8-p} - H \wedge F_{6-p} = \varepsilon_p \frac{T_{10}}{p+1} \text{vol}_\perp$$

No flux $\Rightarrow T_{10} = 0 \Rightarrow$ no de Sitter solution.

Remark on the Bianchi identity and the swampland

$p = 3$, no flux, ρ, τ

$$\mathcal{S} = \int d^4x \sqrt{|g_4|} (\mathcal{R}_4 + \text{kin. terms} - V)$$

$$\text{with } V(\rho, \tau) = -\tau^{-2} \rho^{-1} \mathcal{R}_6 - \tau^{-3} \rho^{-\frac{3}{2}} g_s \frac{T_{10}}{4}$$

3 equations of motion \Rightarrow de Sitter solution with

$$\mathcal{R}_4 = -\frac{2}{3} \mathcal{R}_6 = g_s \frac{T_{10}}{4} > 0$$

\Rightarrow landscape of 4d theories (\mathcal{R}_6, T_{10}) , some with de Sitter solutions.

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Similar with $\lambda \leq 0$: inconclusive without Bianchi identity, exclusion de Sitter with it.

Summary

David
ANDRIOT

- What cosmological model is (not) in the swampland?
- \leftrightarrow Get, in a controlled manner, de Sitter vacua from string theory?
- de Sitter swampland criterion, or a refined version?

- Difficulties/constraints on classical 10d de Sitter solutions
- Explore a stability island for parallel D_p/O_p
- Explore more intersecting D_p/O_p

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Thank you for your attention!

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