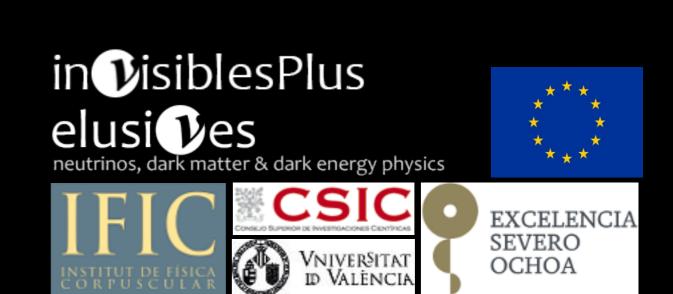
Searching for Neutrinos through Cosmology

Olga Mena (Spăin)





Today's menu

Antipasto

- The classical cosmic pizza
- The neutrino slice
- Neutrino decoupling in the early universe
- Other cosmic pizzas for tasting

Main course

- Number of neutrinos and Big-Bang Nucleosynthesis
- Number of neutrinos and Cosmic Microwave Background Radiation
- Neutrino masses and Cosmic Microwave Background Radiation
- Neutrino masses and structure formation in the universe
- The Dark Justice League game: Σm_{ν} versus w(z)

Doggy Bag

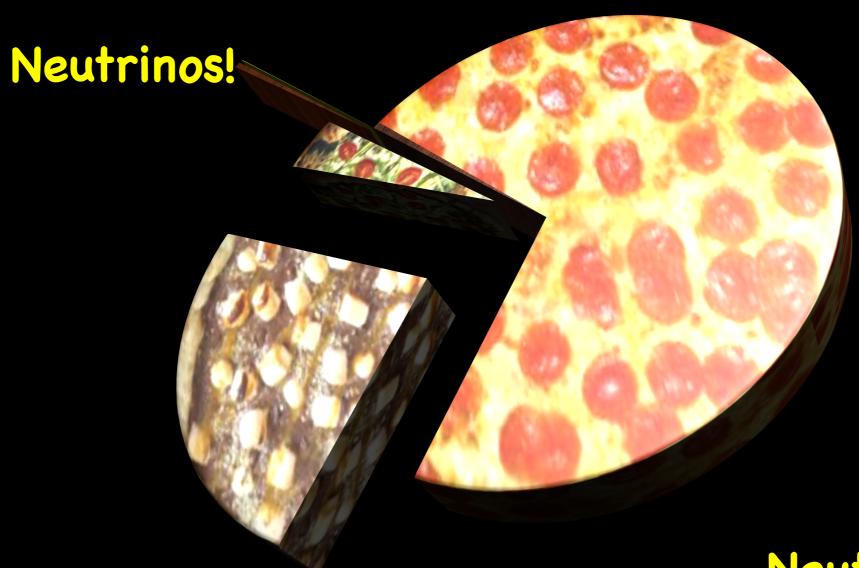
Take home messages

ACDM Pizza

 Ω_i



How large is the MASSIVE neutrino contribution?



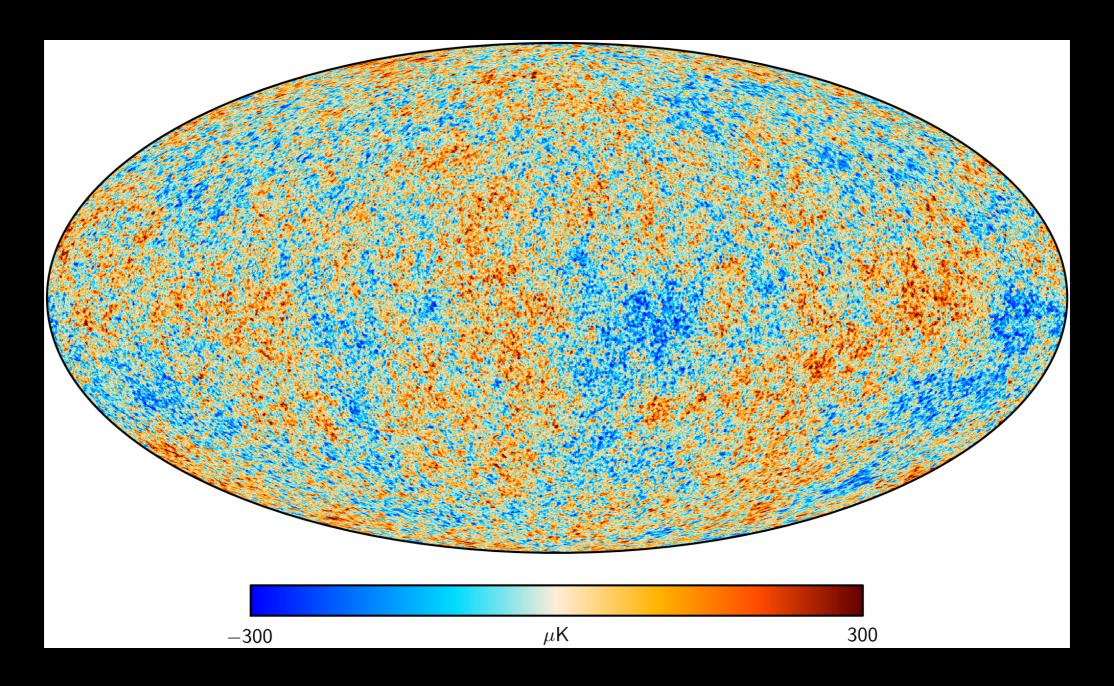
Cosmology tells us

 $\Omega_{\nu} \lesssim 0.0024 95\% \text{ CL}$

Neutrino oscillations tell us

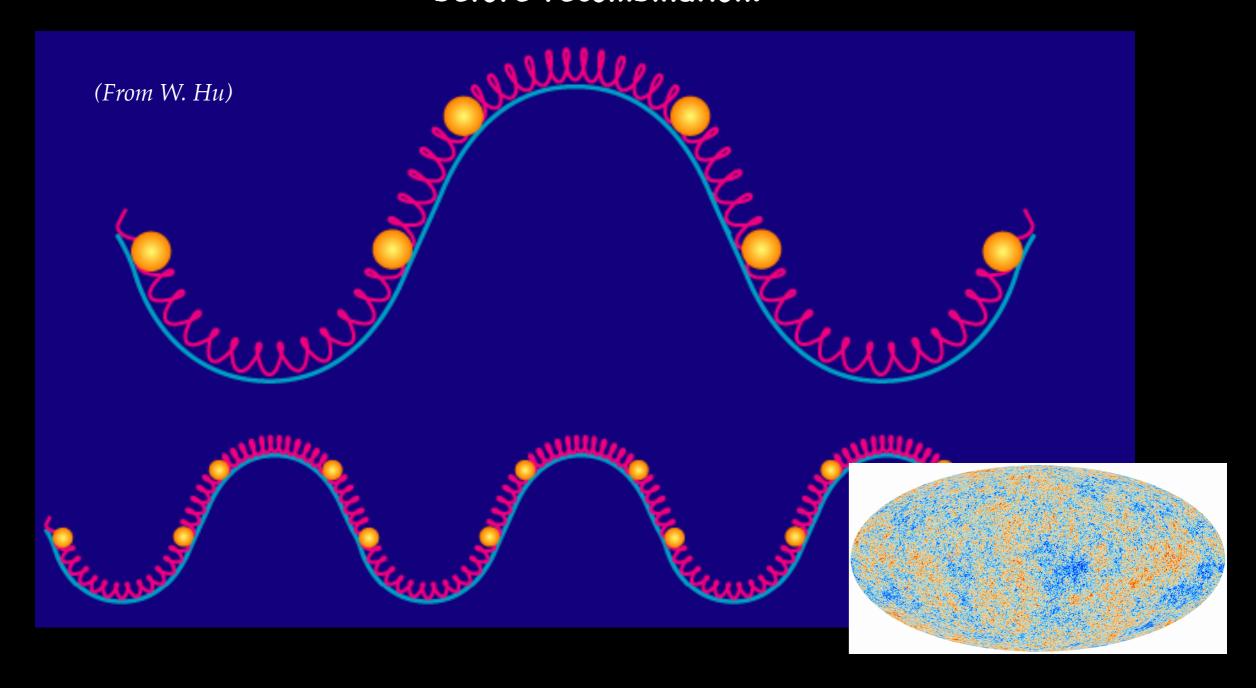
 $\Omega_{\nu} \gtrsim 0.0012 \ 95\% \ {\rm CL}$

I guess all you know about dark matter/what about dark radiation? But radiation is visible, and it has a mean T= 2.725 K!



This map is just telling us how the CMB temperature fluctuations vary with the angular size of patches in the sky...

The CMB fluctuations are due to the acoustic oscillations in the baryon-photon fluid before recombination.

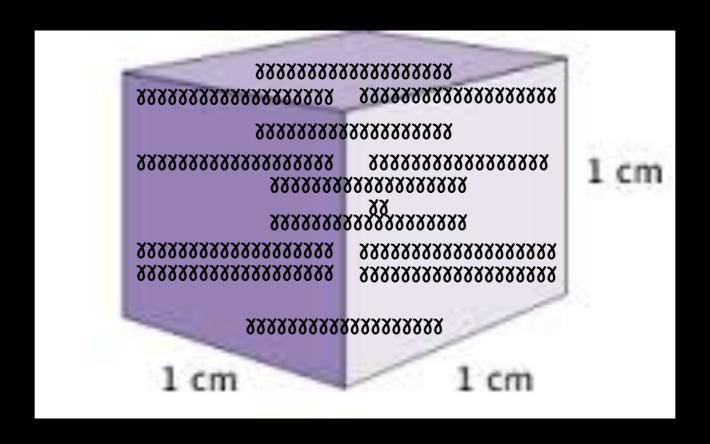


Potential wells \longrightarrow High density \longrightarrow COLD SPOTS in CMB maps Potential hills \longrightarrow Low density \longrightarrow HOT SPOTS in CMB maps

I guess all you know about dark matter/what about dark radiation?

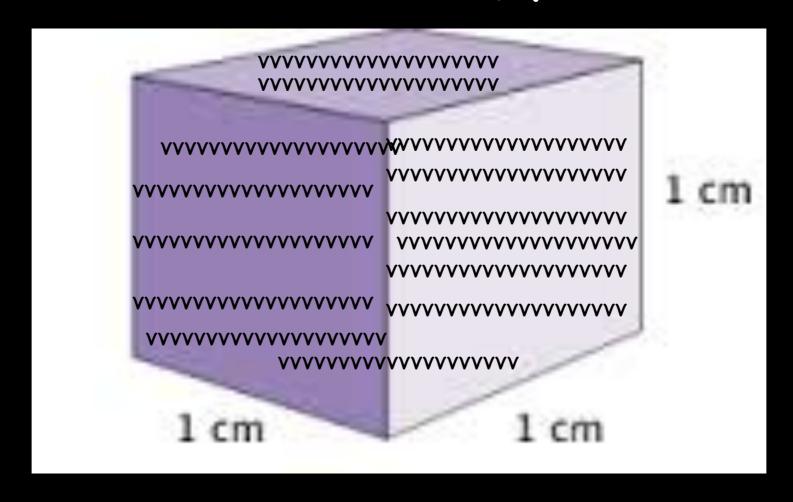
But radiation is visible!

410 photons/cm³



According to standard cosmology, there is a cosmic neutrino background, equivalent to the CMB photon background, albeit slightly colder $T \approx 1.94 \text{ K}$

340 neutrinos/cm³

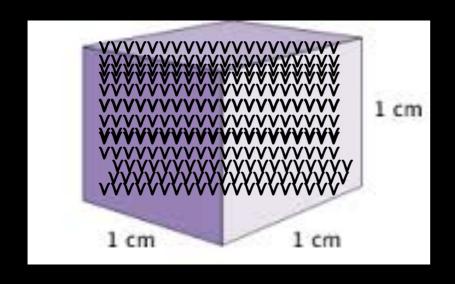


This cosmic relic neutrino background has never been detected directly.

The universe is filled with a dense flux of "relic neutrinos" created in the Big Bang.

This makes neutrinos the most abundant KNOWN form of...

340 neutrinos/cm³



HOT dark matter!

According to standard cosmology, there are three active Dirac or Majorana neutrinos, which decouple from the thermal bath when their scattering rate is smaller than the expansion rate of the universe:

$$\Gamma_{\nu} \lesssim H$$

• Neutrinos only interact via weak interactions, with a rate:

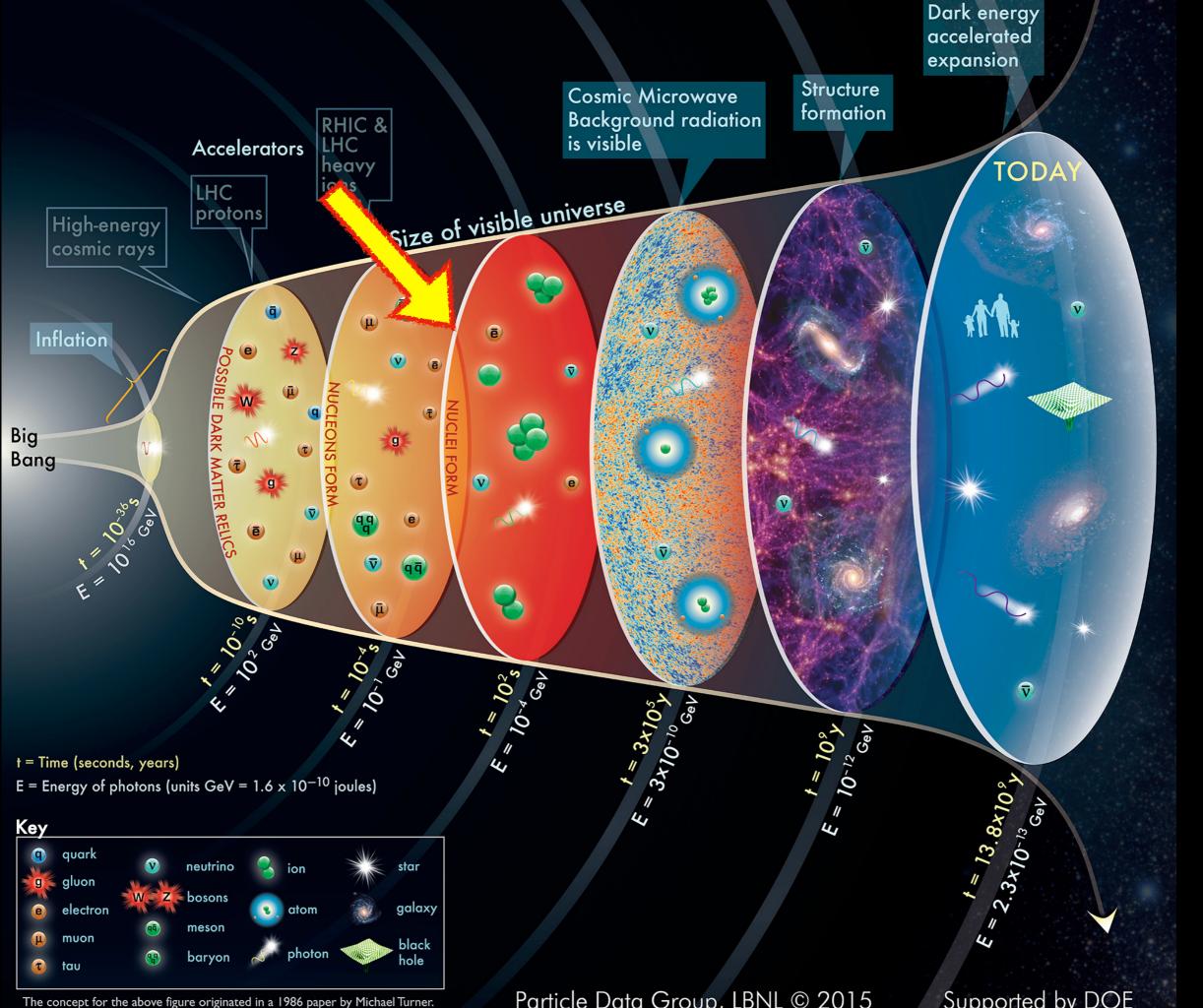
$$\Gamma_{\nu} = n\sigma v \simeq T^3 G_F^2 T^2 \sim G_F^2 T^5$$

• While the expansion rate of the universe is given by the Hubble factor:

$$H^2 = \frac{8\pi G}{3}\rho \sim T^4/m_{pl}^2$$

$$\Gamma_{\nu}/H \sim \left(\frac{T}{1 \text{ MeV}}\right)^3$$

• Therefore neutrinos decouple from the thermal bath around 1 MeV.



Particle Data Group, IBNI © 2015

Event	Time	Redshift	Temperature	
Baryogenesis	?	?	?	
EW phase transition	$2 \times 10^{-11} s$	10^{15}	100 GeV	
QCD phase transition	$2 \times 10^{-5} s$	10^{12}	150 MeV	
Neutrino decoupling	1s	6×10^9	1 MeV	
Electron-positron annihilation	6s	2×10^9	500 keV	
Big bang nucleosynthesis	3min	4×10^8	100keV	
Matter-radiation equality	$6 \times 10^4 yrs$	3400	.75eV	
Recombination	$2.6 - 3.8 \times 10^5 yrs$	1100-1400	.2633eV	
CMB	$3.8 \times 10^5 yrs$	1100	.26eV	

They do not inherit any of the energy associated to e+ e- annihilations, being colder than photons:

$$T_{\nu 0} = \left(\frac{4}{11}\right)^{1/3} T_0 = 1.945 \text{ K} \sim 1.697 \times 10^{-4} \text{ eV}$$

If these neutrinos are massive, their energy density, at T<<m is

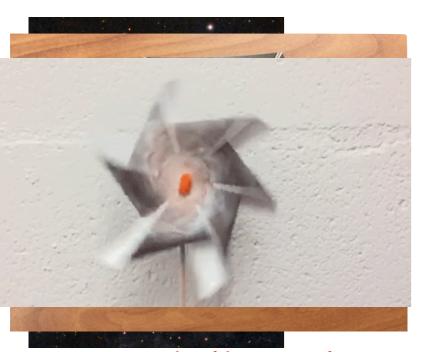
$$\rho_{\nu} = m_{\nu} n_{\nu}$$
 $n_{\nu_i}(T_{\nu 0}) \approx 56 \text{ cm}^{-3}$ $\Omega_{\nu} h^2 = \frac{\sum m_{\nu}}{93 \text{ eV}}$

Then, demanding that massive neutrinos do not over-close the universe, $\sum m_
u \lesssim 45~{
m eV}$

Their thermal motion is:
$$\langle v_{\rm thermal} \rangle \simeq 81(1+z) \left(\frac{{\rm eV}}{m_{\nu}}\right) {\rm ~km~s^{-1}}$$

For a 1 eV neutrino, thermal motion is comparable to the typical velocity dispersion of a galaxy.

For dwarf galaxies, the velocity dispersion is smaller, 10 km/s





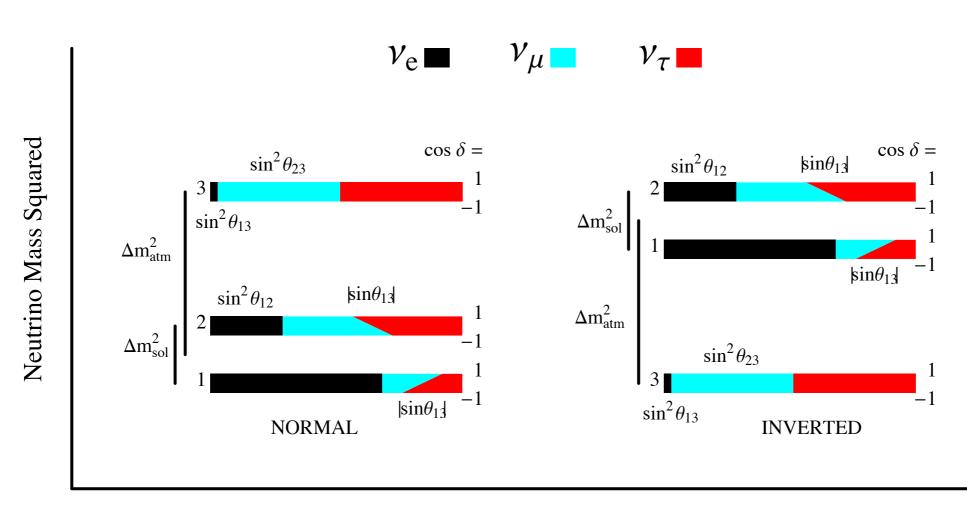
Too much thermal energy to be squeezed into small volumes to form the smaller structures we observe today! 13

According to neutrino oscillation physics, we know that there are at lest two Dirac or Majorana massive neutrinos:

$$\Delta m_{12}^2 = (7.05 - 8.14) \times 10^{-5} \text{eV}^2$$

$$\Delta m_{13}^2 = (2.41 - 2.60) \times 10^{-3} \text{eV}^2$$

$$\Delta m_{13}^2 = -(2.31 - 2.51) \times 10^{-3} \text{eV}^2$$



Fractional Flavor Content varying $\cos \delta$

(Mena, Parke, PRD'04)

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We are sure then that two neutrinos have a mass above:

$$\sqrt{\Delta m_{12}^2} \simeq 0.008 \text{ eV}$$

and that at least one of these neutrinos has a mass larger than

$$\sqrt{|\Delta m_{13}^2|} \simeq 0.05 \text{ eV}$$

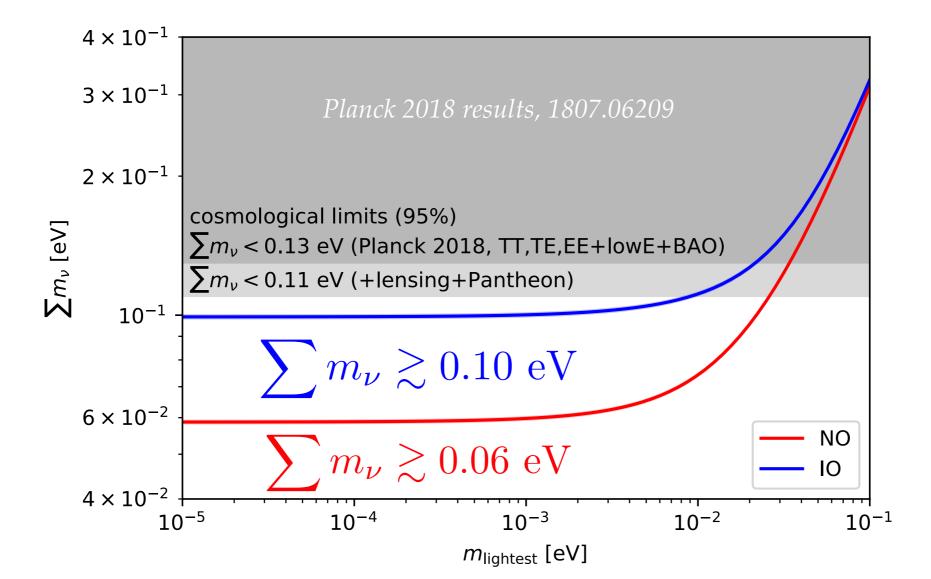
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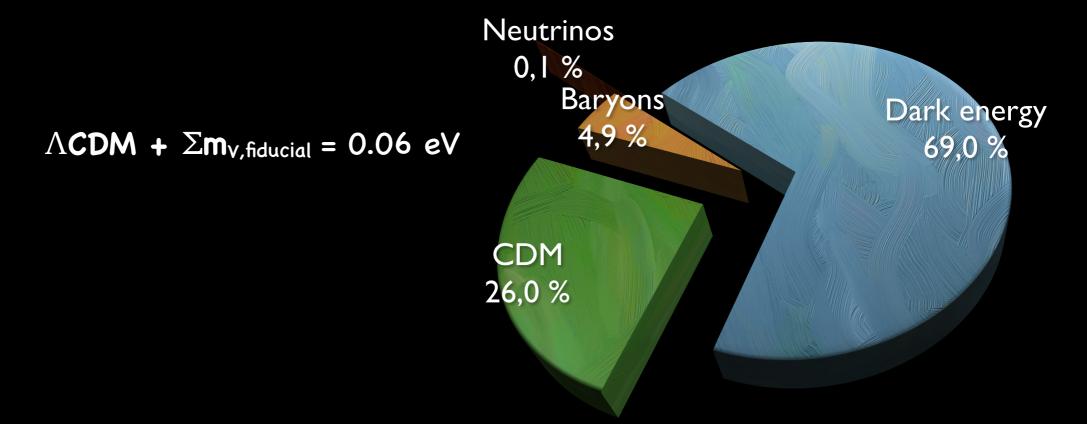
which translates into a lower bound on the total neutrino mass, depending on the ordering:

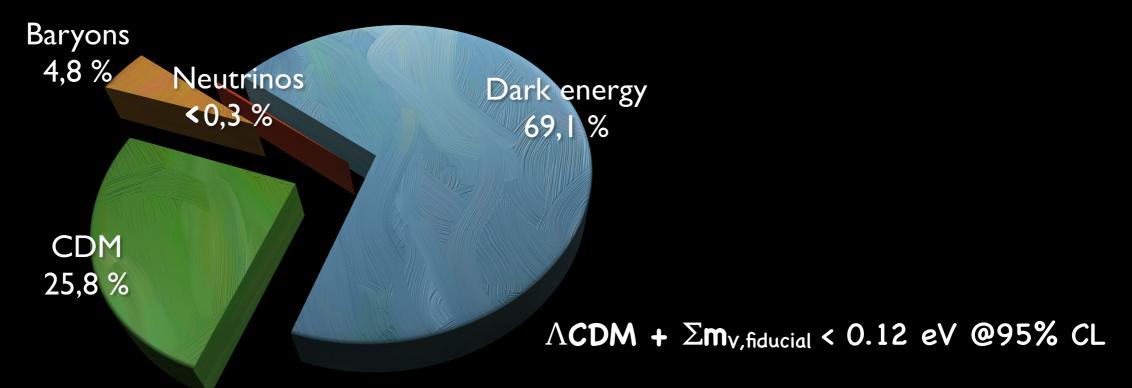


Planck 2018 Cosmic Pizzas

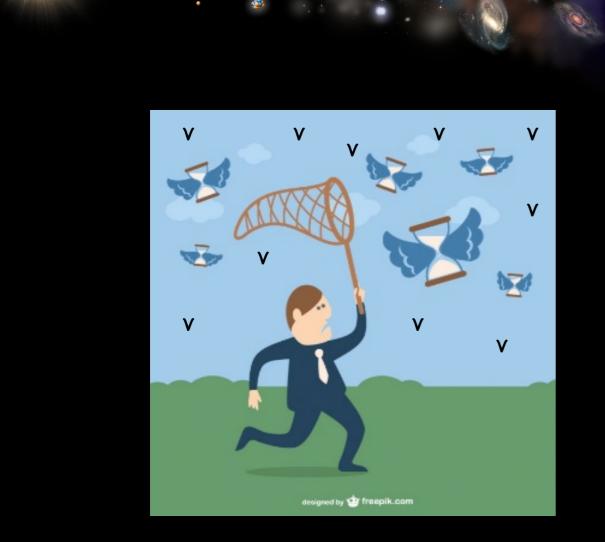


Planck 2018 Cosmic Pizzas





To hunt and bound their abundances and their masses, we need to look at several epochs in our universe's evolution:



and it's as challenging as trapping time!

Are you willing to join me in the Cosmo neutrino hunting trip?



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Doggy Bag

Take home messages

Number of neutrinos: Neff

The total radiation in the universe can be written as:

$$\Omega_r h^2 = \left(1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}}\right) \Omega_{\gamma} h^2$$

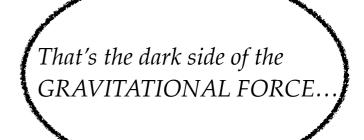
 N_{eff} = 3.046 standard scenario: electron, muon and tau neutrinos

Neff < 3.046 (less neutrinos): Neutrino decays?

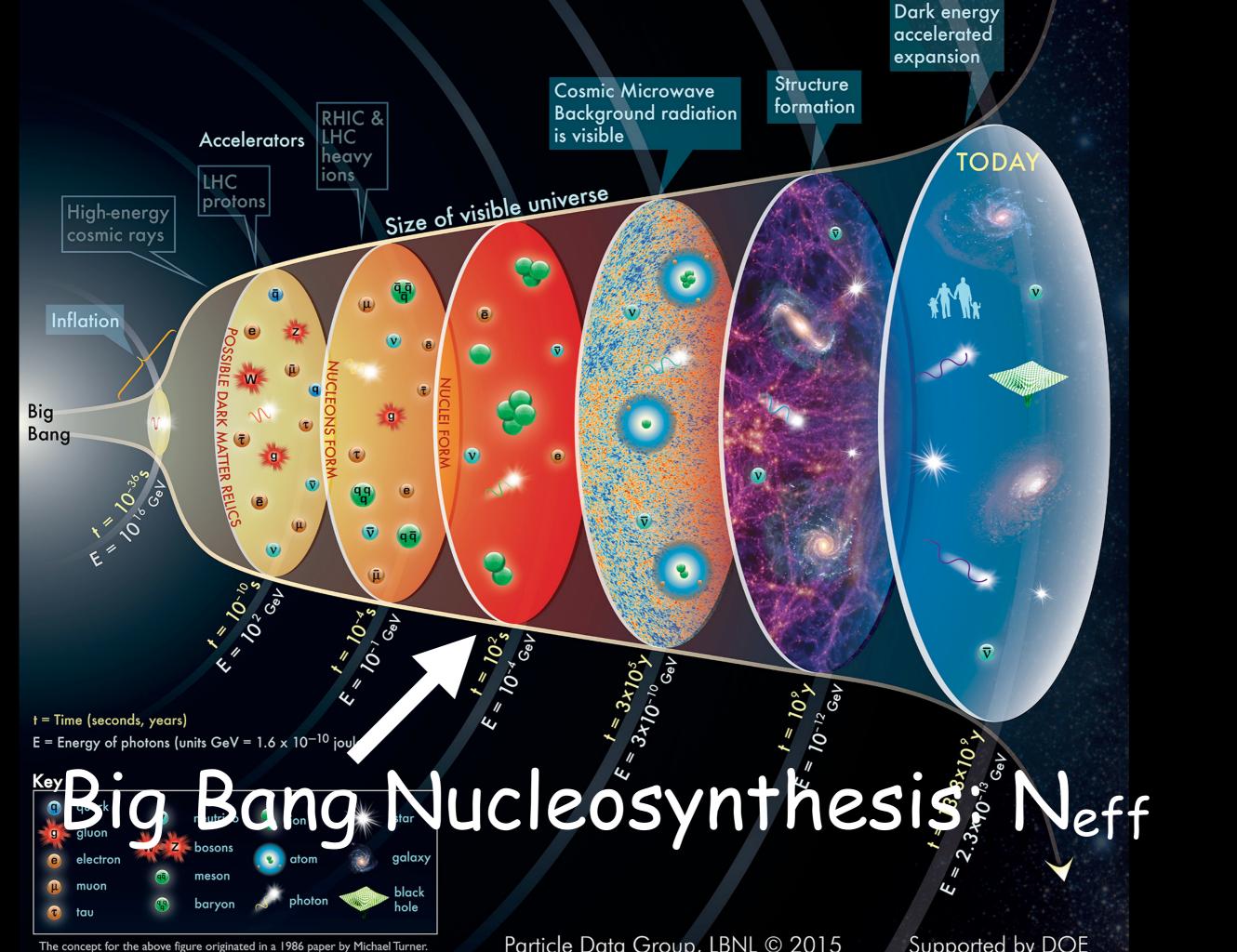
N_{eff} > 3.046 (more neutrinos): Sterile neutrino species ?

But....if they are sterile, and do not interact with other particles, how cosmologists measure them?









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Big Bang Nucleosynthesis: Neff

BBN theory predicts the abundances of D, 3 He 4 He and 7 Li which are fixed by $t \approx 180$ s. They are observed at late times: low metallicity sites with little evolution are "ideal".

Big Bang Nucleosynthesis: Neff

Neff changes the freeze out temperature of weak interactions:

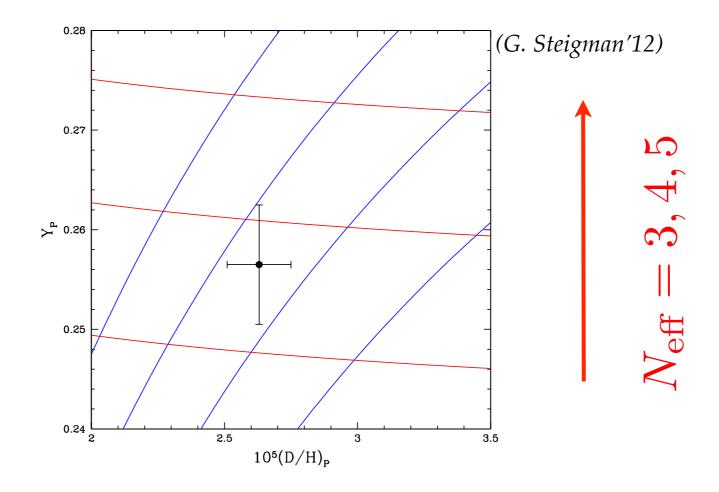
$$\Gamma_{n\leftrightarrow p}\sim H$$

MORE NEUTRINOS:

Higher Neff: larger expansion rate & freeze out temperature, MORE HELIUM 4

$$n/p \simeq e^{-\frac{m_n - m_p}{T_{freeze}}}$$

$$Y_p = \frac{2(n/p)}{1 + n/p}$$



Today's menu

Antipasto

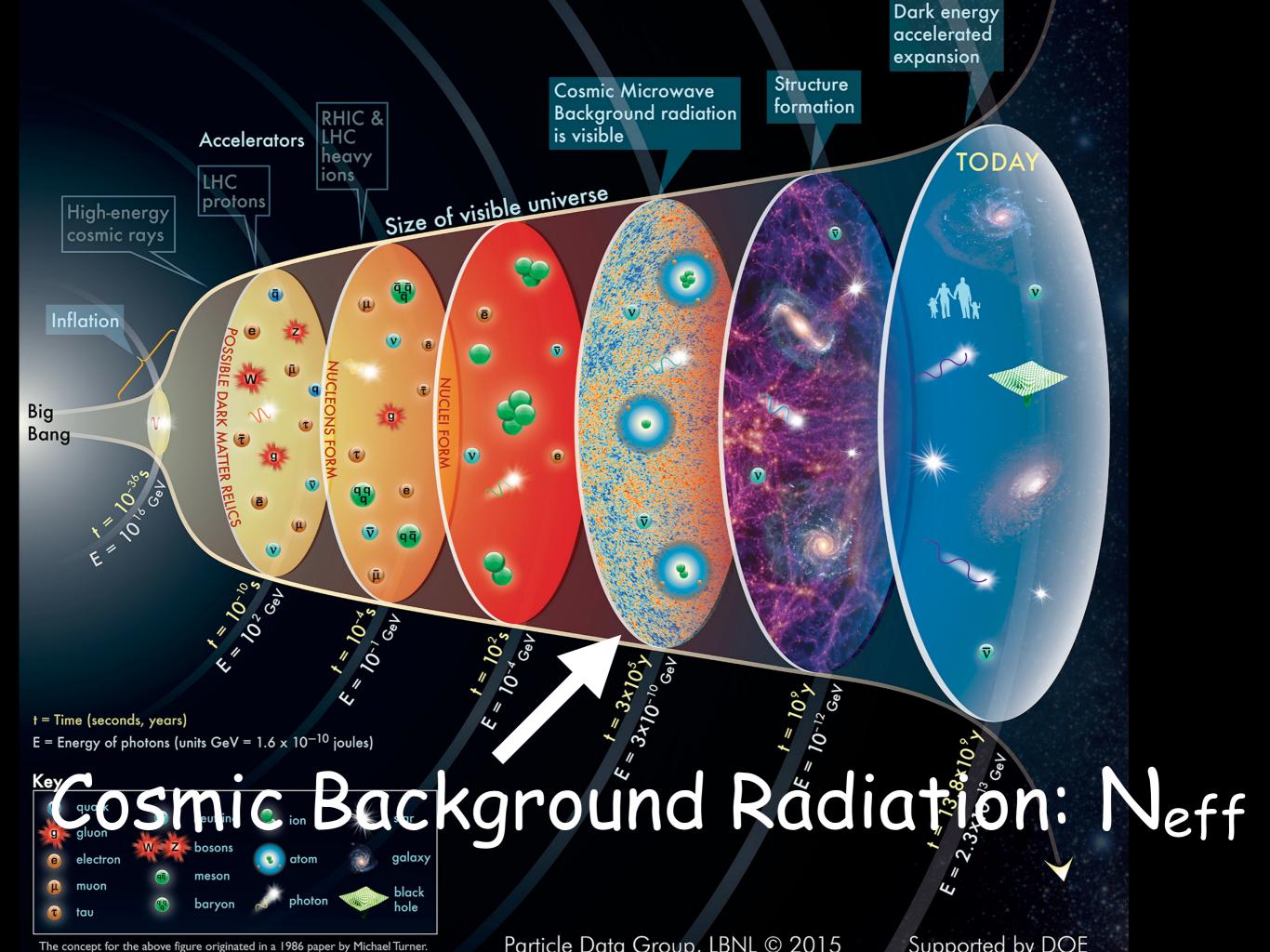
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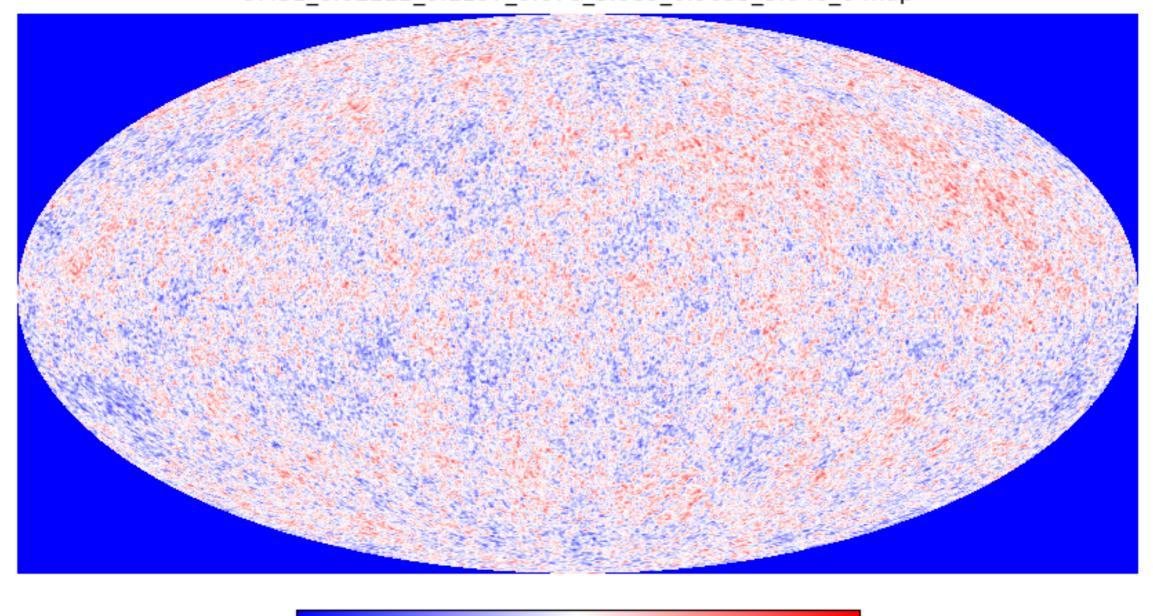
Take home messages



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Also known as "photon decoupling", as photons started freely travel through the universe without interacting with matter and the CMB is "frozen"

 $67.31_0.02222_0.1197_0.078_3.089_0.9655_3.046_0\ map$





Spherical harmonics decomposition:

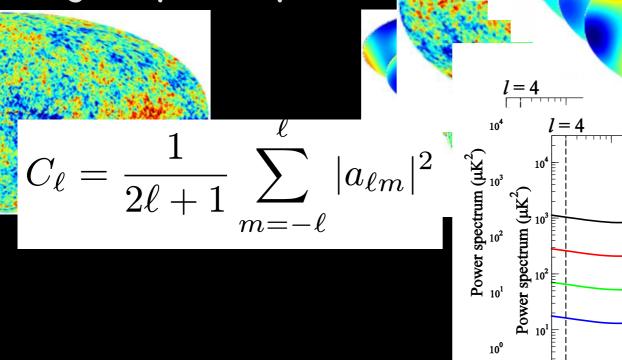
With expansion coefficients:

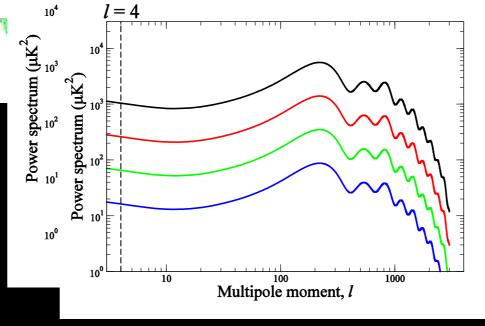
$$T(\hat{n}) = \sum_{\ell=0}^{\ell_{\text{max}}} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\hat{n})$$

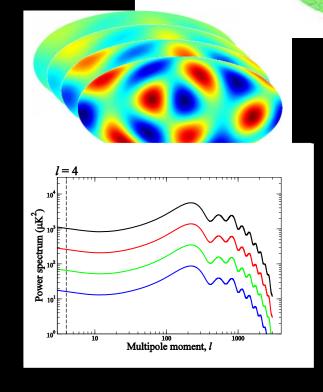
$$a_{\ell m} = \int_{4\pi} T(\hat{n}) Y_{\ell m}^*(\hat{n}) d\Omega$$

The angular power spectrum

es the amplitude as a function of the

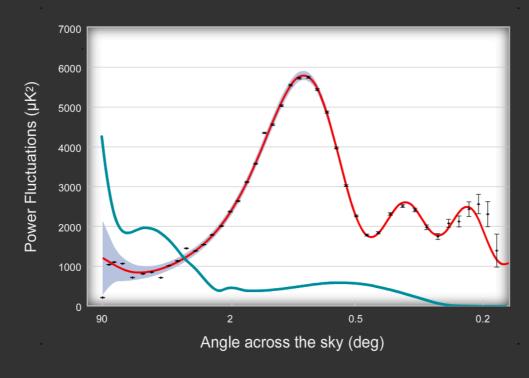






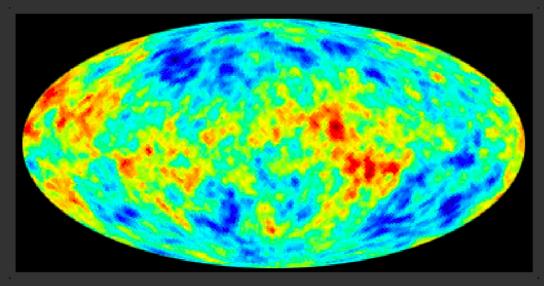
CMB Analyzer





Pie Chart: Graphically shows the composition of your universe. The wedges compare the amount of each component; the size of the pie compares the total composition (matter + dark matter + dark energy) with the critical density (black circle).

- A universe at critical density is geometrically flat and probably infinite.
- A universe can have more or less than the critical density.
- Flatness the term we use for closeness to critical density.



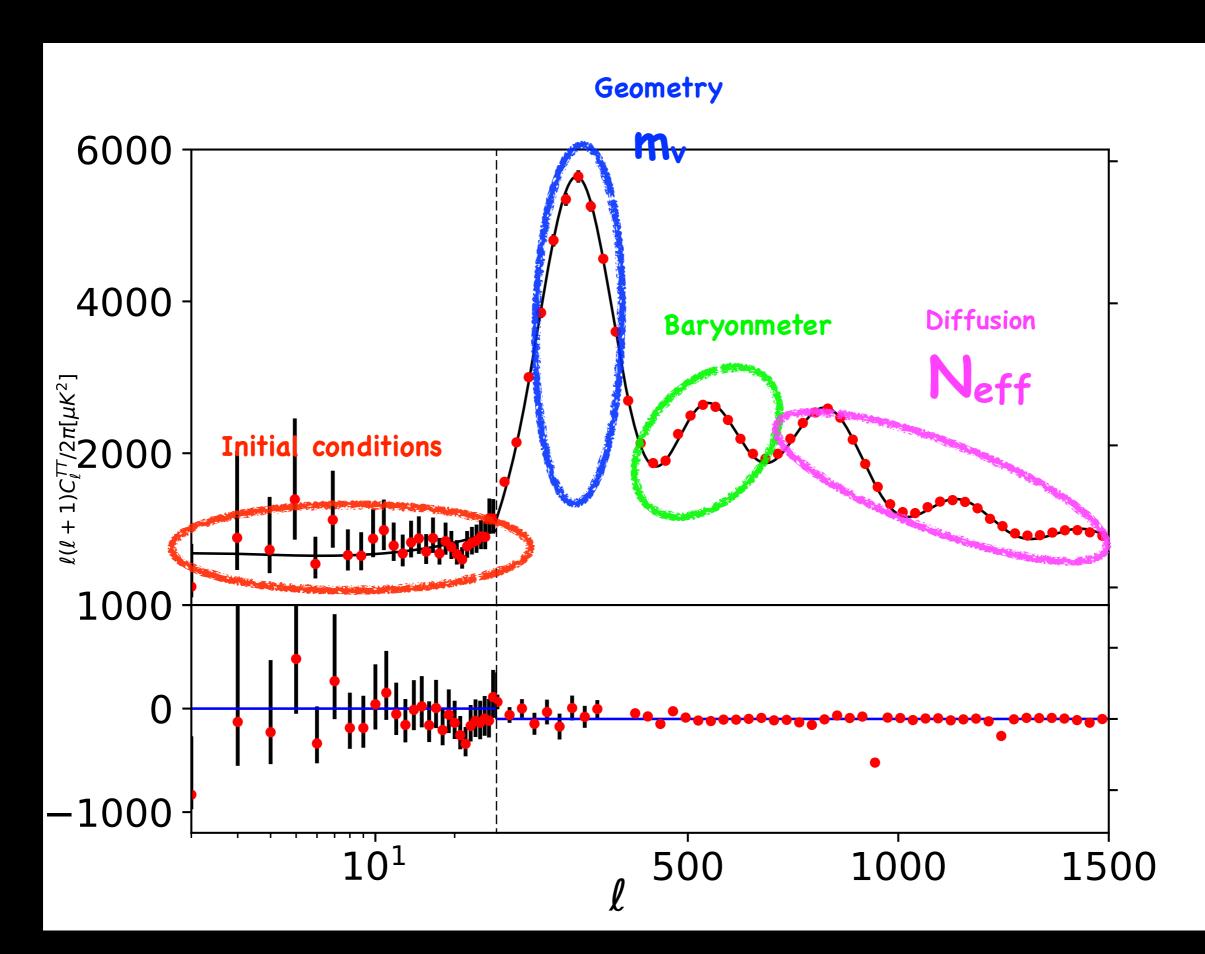
Age: 6.9 billion years

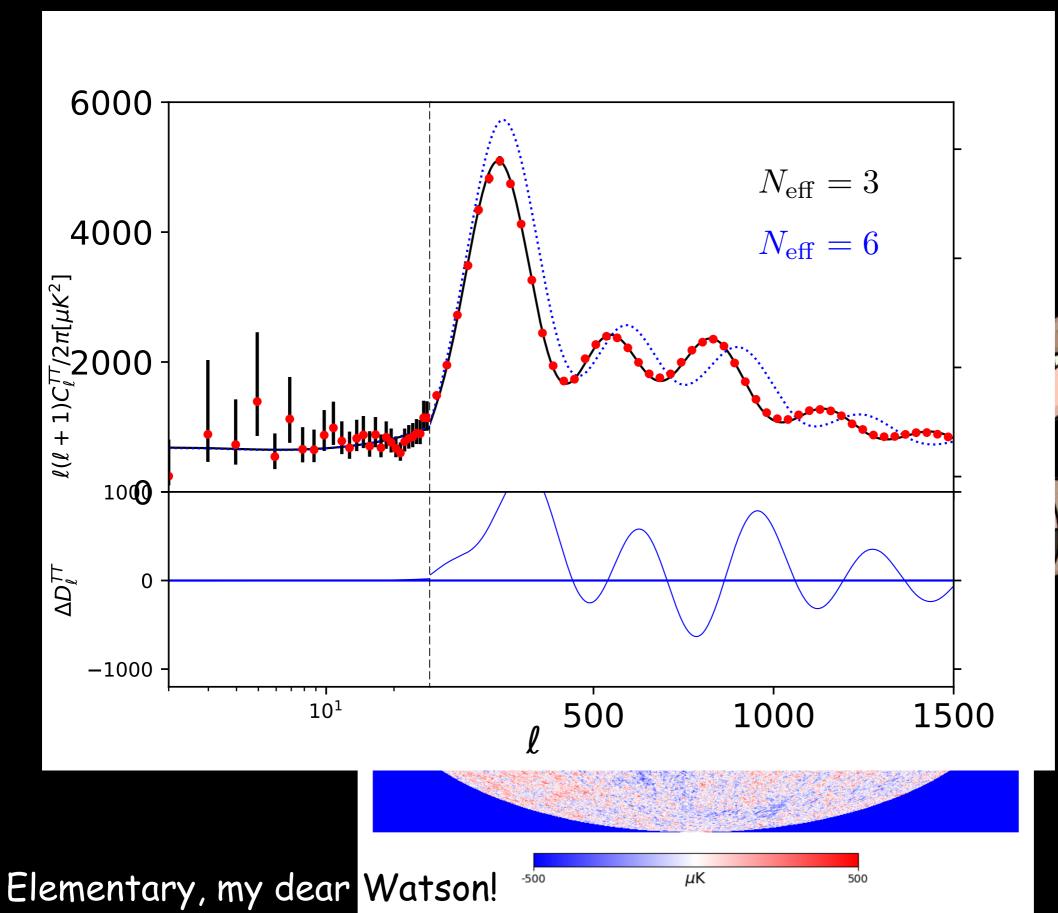
Flatness: 1.7

ANSWER

RESET

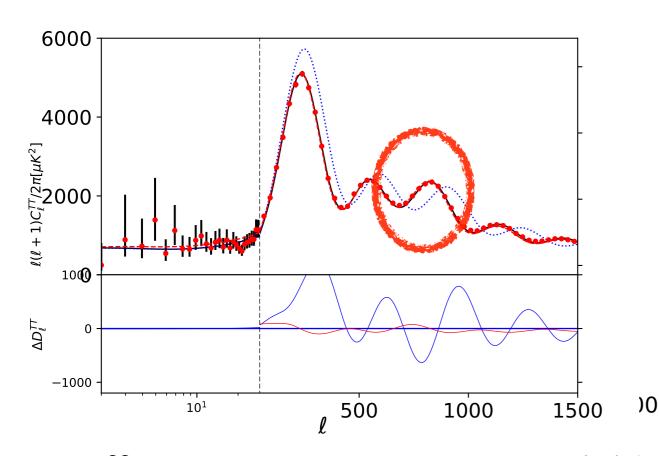
CMB: a lot to learn about....







$$N_{\text{eff}} = 6$$
 $N_{\text{eff}} = 3$ $N_{\text{eff}} = 6$



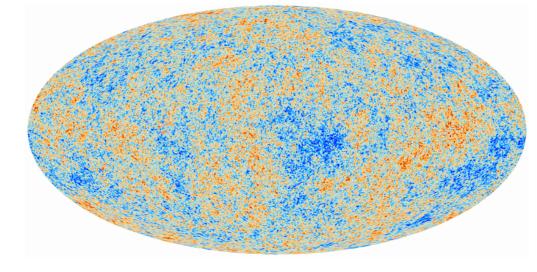
$$(\omega_b, \omega_m, h, A_s, n_s, \tau, N_{\text{eff}})$$



It is elementary, Sherlock Holmes!

Only effect at I<1000 that can not be mimicked by others: anisotropic stress, around 3rd peak







Neutrinos are free-streaming particles propagating at the speed of light, faster than the sound speed in the photon fluid, suppressing the oscillation amplitude of CMB modes that entered the horizon in the radiation epoch.

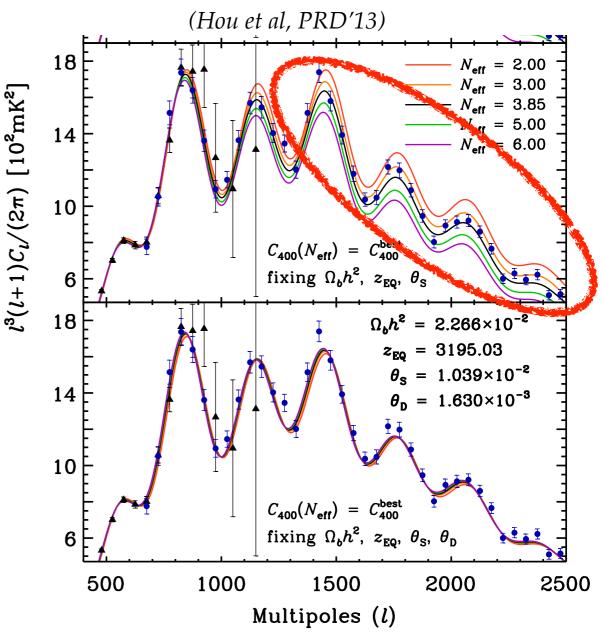
@Cosmic Microwave Background in the damping tail, measured by SPT, ACT & Planck:

Higher Neff will increase the expansion rate AND the damping at high multipoles.





$$r_d^2 \propto \int_0^{a_\star} \frac{da}{a^3 \sigma_T n_e H}$$





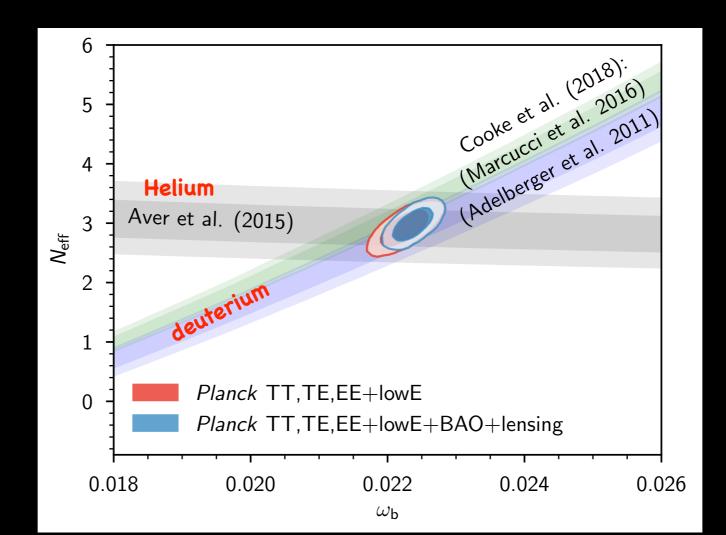
Planck 2018 CMB temperature polarization and lensing potential data:

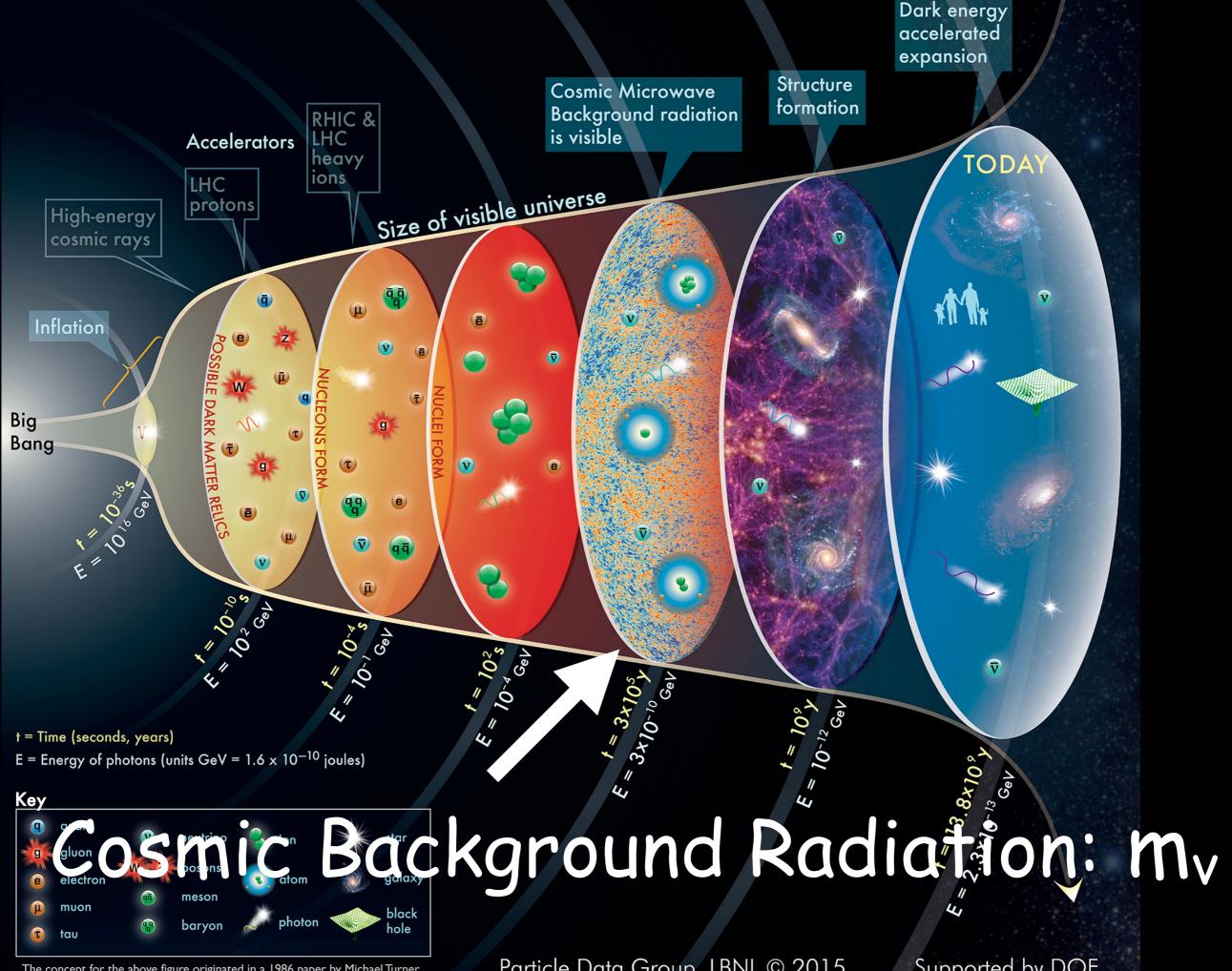
$$N_{\text{eff}} = 2.89^{+0.36}_{-0.38} 95\% \text{CL}$$

If we add large scale structure information in the BAO shape form:

$$N_{\text{eff}} = 2.99^{+0.34}_{-0.33} 95\% \text{CL}$$

Perfectly consistent with BBN estimates:





The concept for the above figure originated in a 1986 paper by Michael Turner.

Particle Data Group, IBNI © 2015

Supported by DOF

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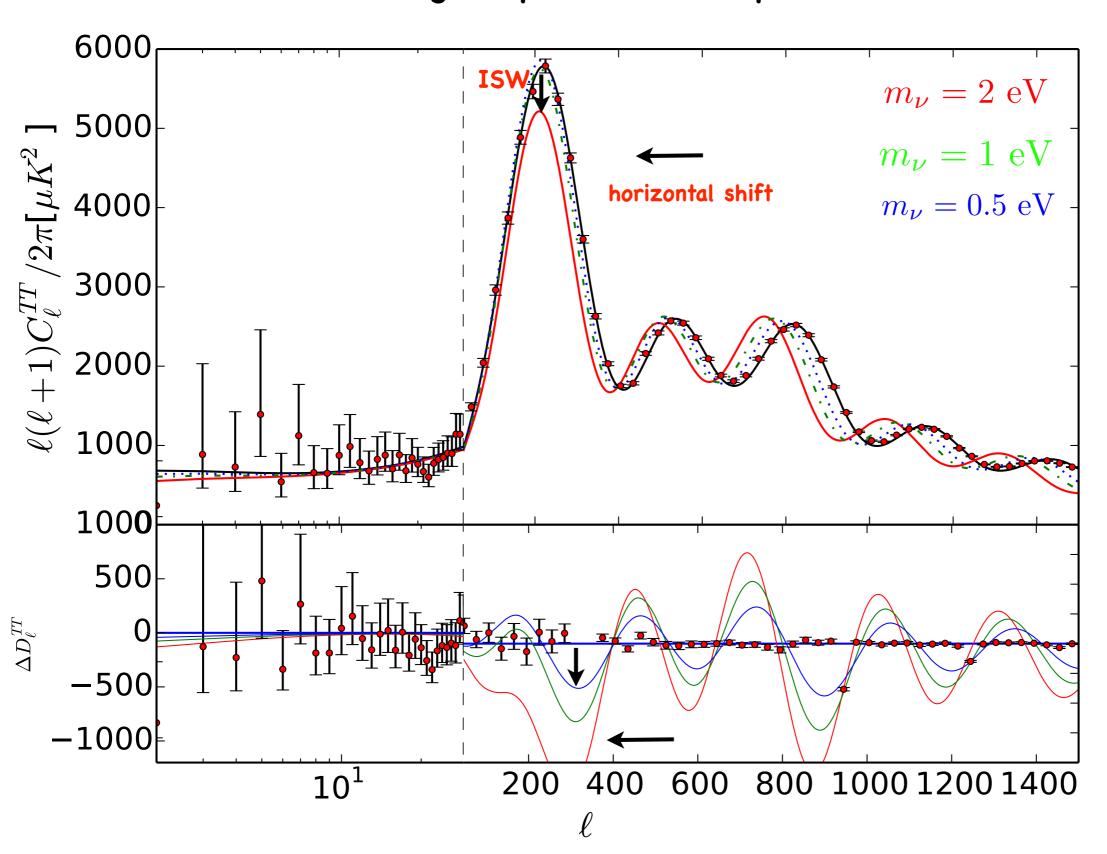
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Doggy Bag

Take home messages

CMB: Σm_v

@ CMB: Early Integrated Sachs Wolfe effect (ISW). Shift in the angular position of the peaks.

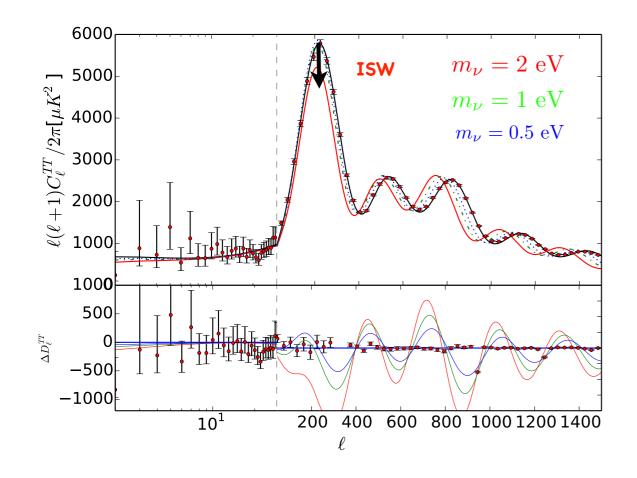


CMB: Σm_v

@ CMB: Early Integrated Sachs Wolfe effect (ISW)

$$\Theta(\hat{n}) = \frac{\delta T}{T}(\hat{n}) \simeq \Theta_0 + \Psi + \hat{n}(\hat{v}_e - v) + \int \dot{\Psi} + \dot{\Phi} d\eta$$

In matter domination, the gravitational potential is constant: NO ISW effect! The transition from the relativistic to the non relativistic neutrino regime gets imprinted in the decays of the gravitational potentials near the recombination period, contributing to the ISW effect!



This early ISW effect leads to a depletion of:

$$\frac{\Delta C_{\ell}}{C_{\ell}} = -(\sum m_{\nu}/0.1 \text{ eV})\%$$

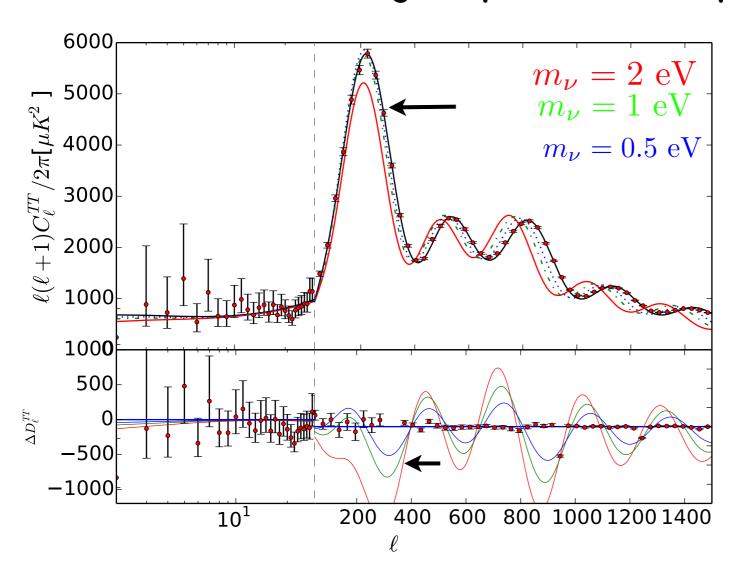
on multipoles:

$$20 < \ell < 200$$

CMB: Σm_v

@ CMB: Early Integrated Sachs Wolfe effect (ISW).

Shift in the angular position of the peaks.



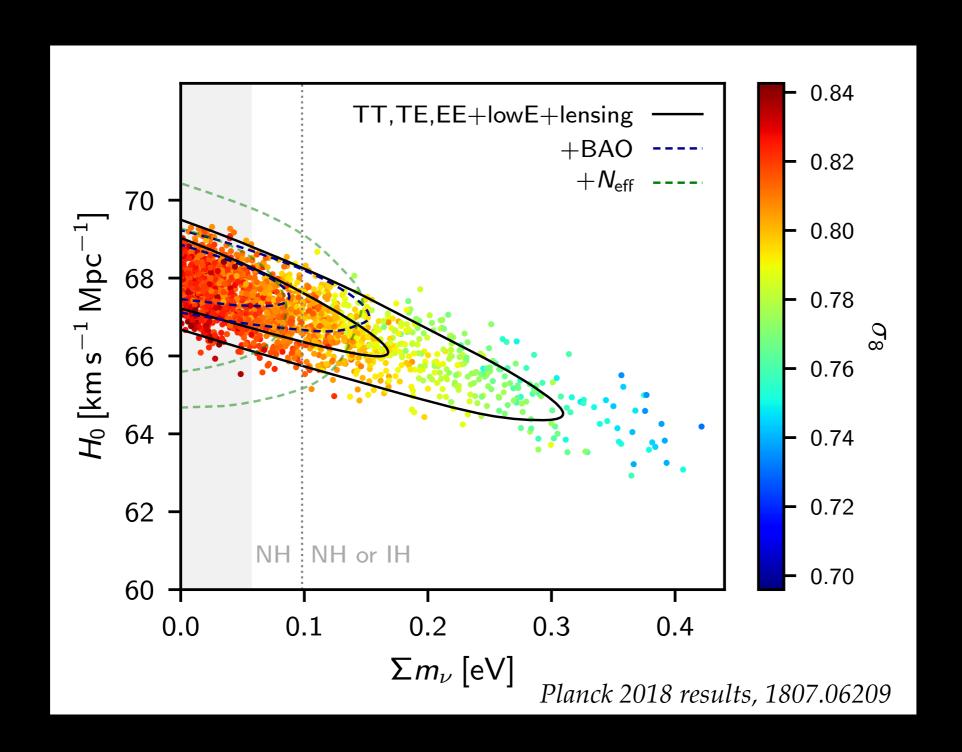
$$\theta_s = \frac{r_s}{D_A}$$

$$\Upsilon_s = \int_0^{t(z_d)} c_s (1+z) dt = \frac{2}{3k_{\text{eq}}} \sqrt{\frac{6}{R_{\text{eq}}}} \ln \frac{\sqrt{1 + R_d} + \sqrt{R_d + R_{\text{eq}}}}{1 + \sqrt{R_{\text{eq}}}}$$

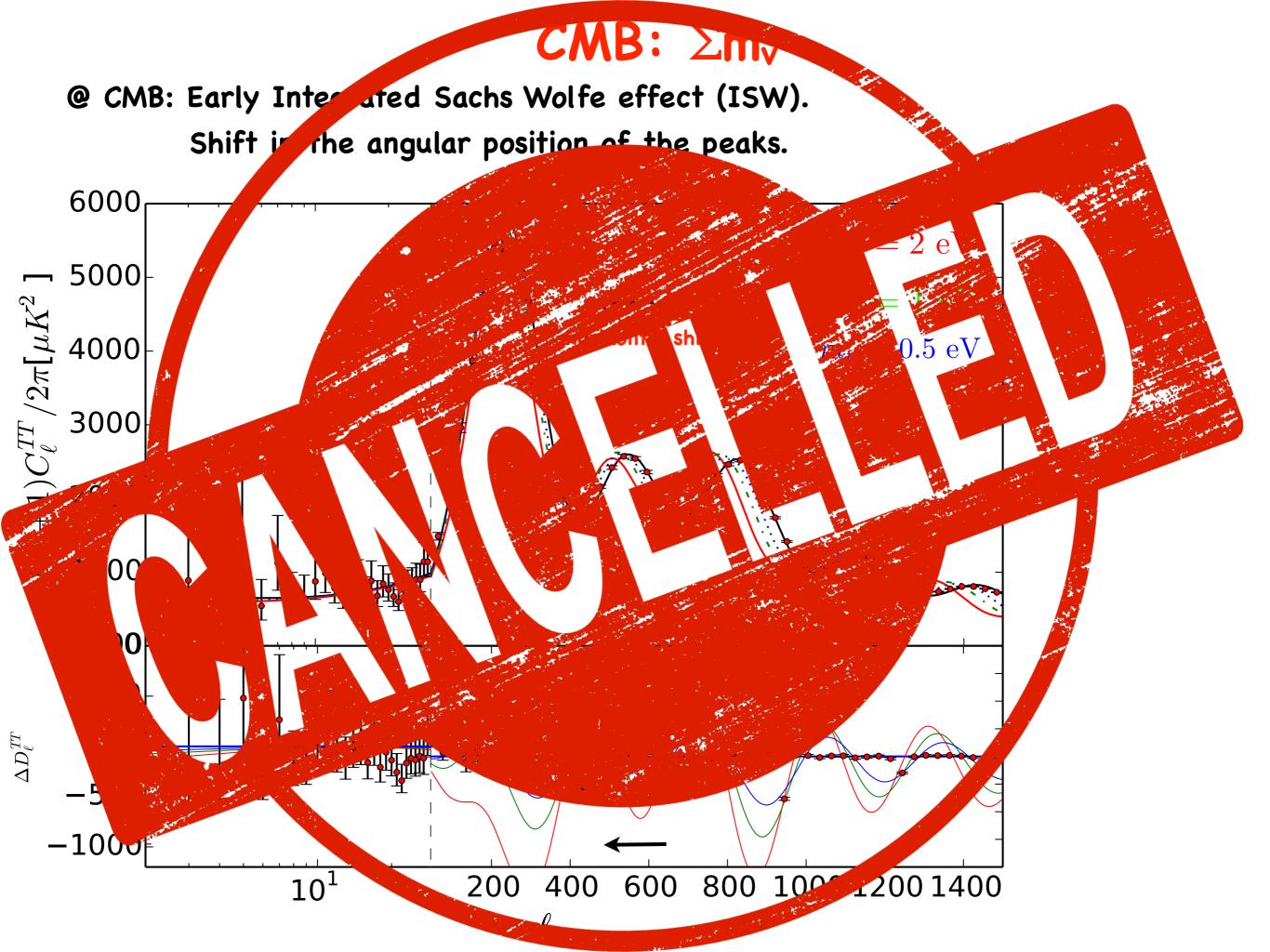
$$D_A = \int_0^{z_{rec}} \frac{dz}{H(z)}$$

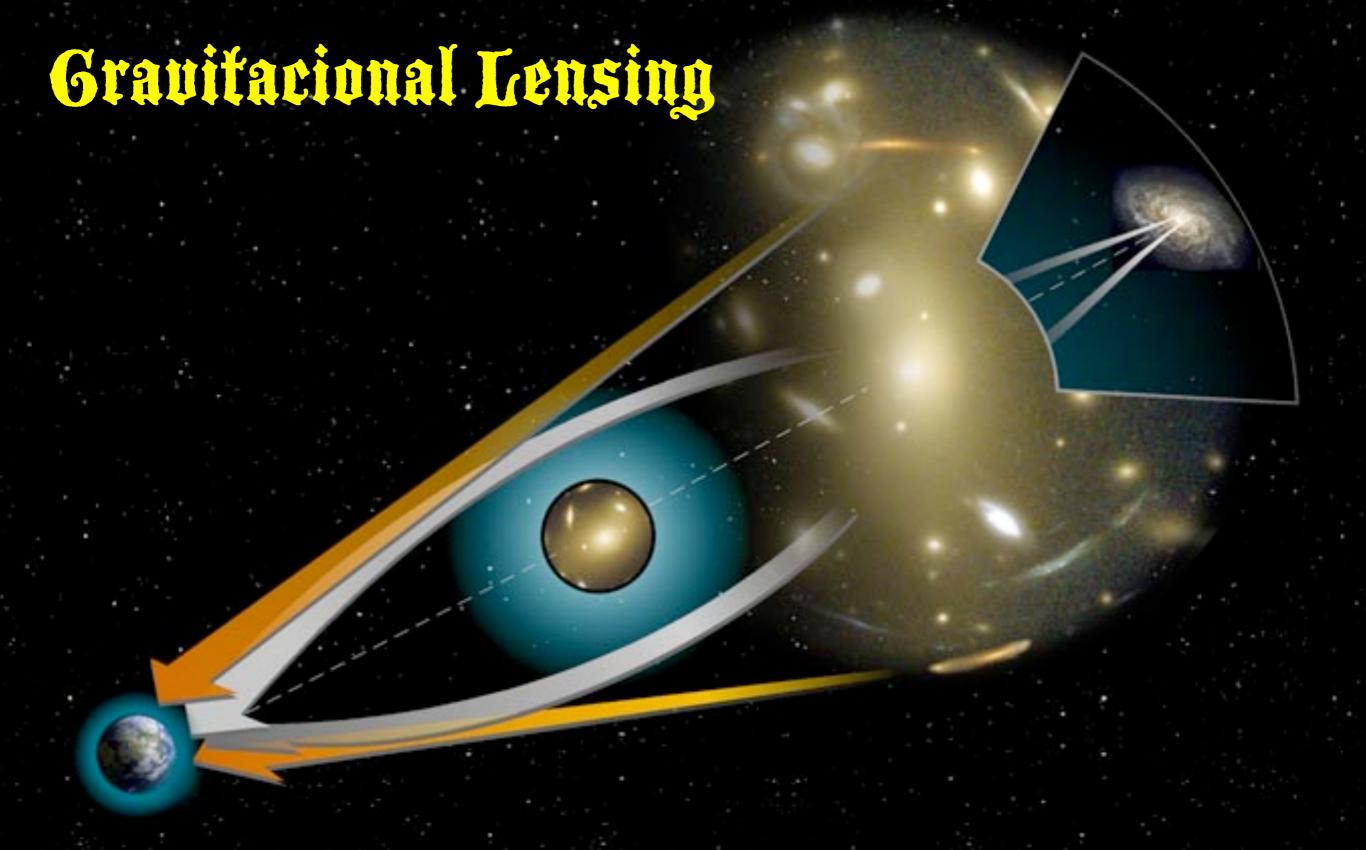
The higher the neutrino mass, the lower the angular diameter distance. Peaks shift to lower multipoles. But this effect can be compensated with a lower Hubble constant:

Strong degeneracy between $\sum m_{\nu}$ and the Hubble constant H₀!



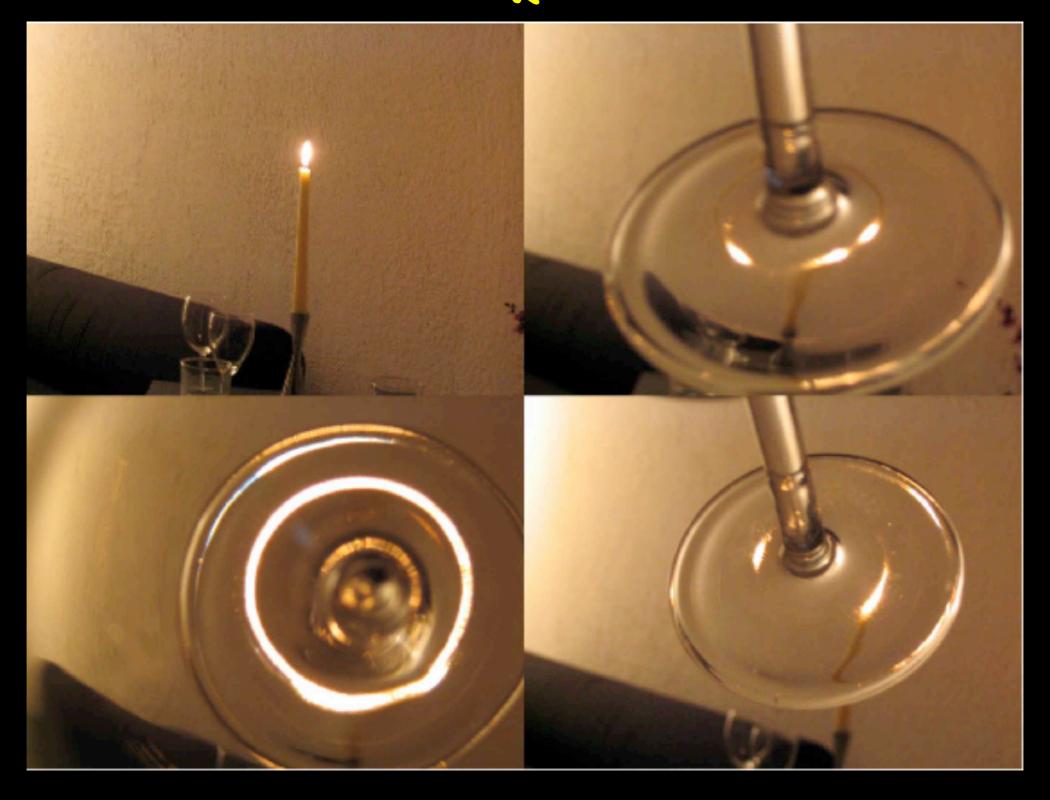
Strong degeneracy between Σm_{ν} and the Hubble constant H₀!





Einstein's relativity predicts that the presence of a massive body will curve space time, distorting the light trajectory. The shape of the background objects will change/multiplied by the presence of intervening galaxies.

Gravitacional Lensing

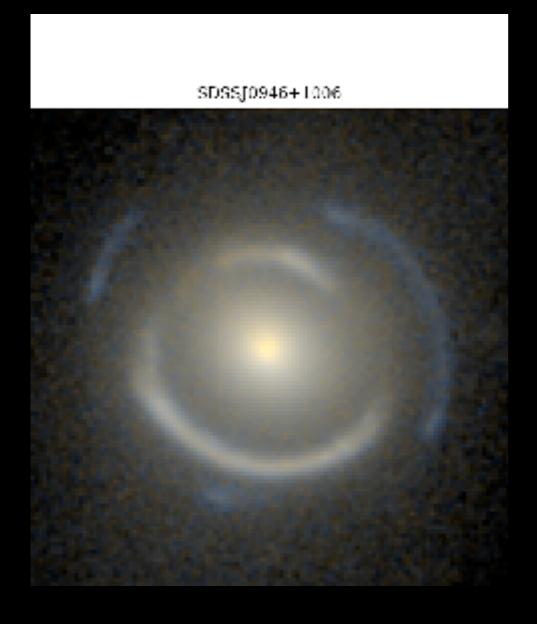


Lensing Galaxy



This movie shows a spiral galaxy acting as a lense of a background quasar (Quasi-stellar radio source) moving behind the galaxy. When the alignment source-lens-observer is perfect, we see the formation of the Einstein ring!

Gravitacional Lensing



Double Einstein ring! 3 perfectly aligned galaxies (probably less than 100 cases in all the universe, and we have observed one!)

CMB Lensing: Σm_v

Lensing remaps the CMB fluctuations: $\Theta_{lensed}(\hat{n}) = \Theta(\hat{n} + \nabla \phi(\hat{n}))$

$$\Theta_{\text{lensed}}(\hat{n}) = \Theta(\hat{n} + \nabla \phi(\hat{n}))$$

Lensing potential ϕ is a measure of the integrated mass distribution back to the last

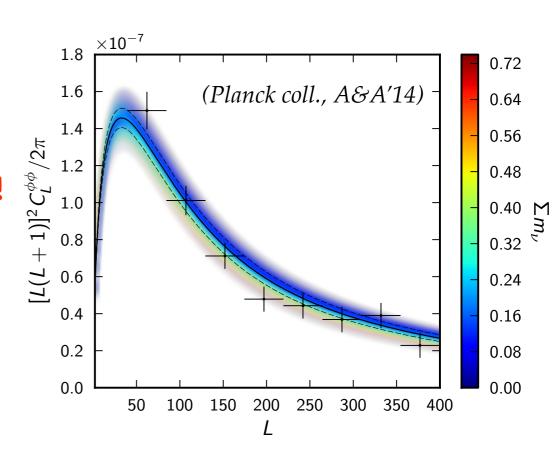
scattering surface

$$\phi(\hat{n}) = -2 \int_0^{z_{rec}} \frac{dz}{H(z)} \Psi(z, D(z) \hat{n}) \left(\frac{D(z_{rec}) - D(z)}{D(z_{rec})D(z)} \right)$$
 Geometry Matter distribution

$$C_{L}^{\phi\phi} = \frac{8\pi^{2}}{L^{3}} \int_{0}^{z_{rec}} \frac{dz}{H(z)} D(z) \left(\frac{D(z_{rec}) - D(z)}{D(z_{rec})D(z)} \right)^{2} P_{\Psi}(z, k = L/D(z))$$

Neutrino free-streaming implies less clustering on small scales, reducing therefore CMB lensing!

(Kaplinghat et al PRL'03, Lesgourgues et al, PRD'06)





Planck TTTEEE+lowT+lowE+lensing

$$\sum m_{\nu} < 0.24 \text{ eV } 95\%\text{CL}$$

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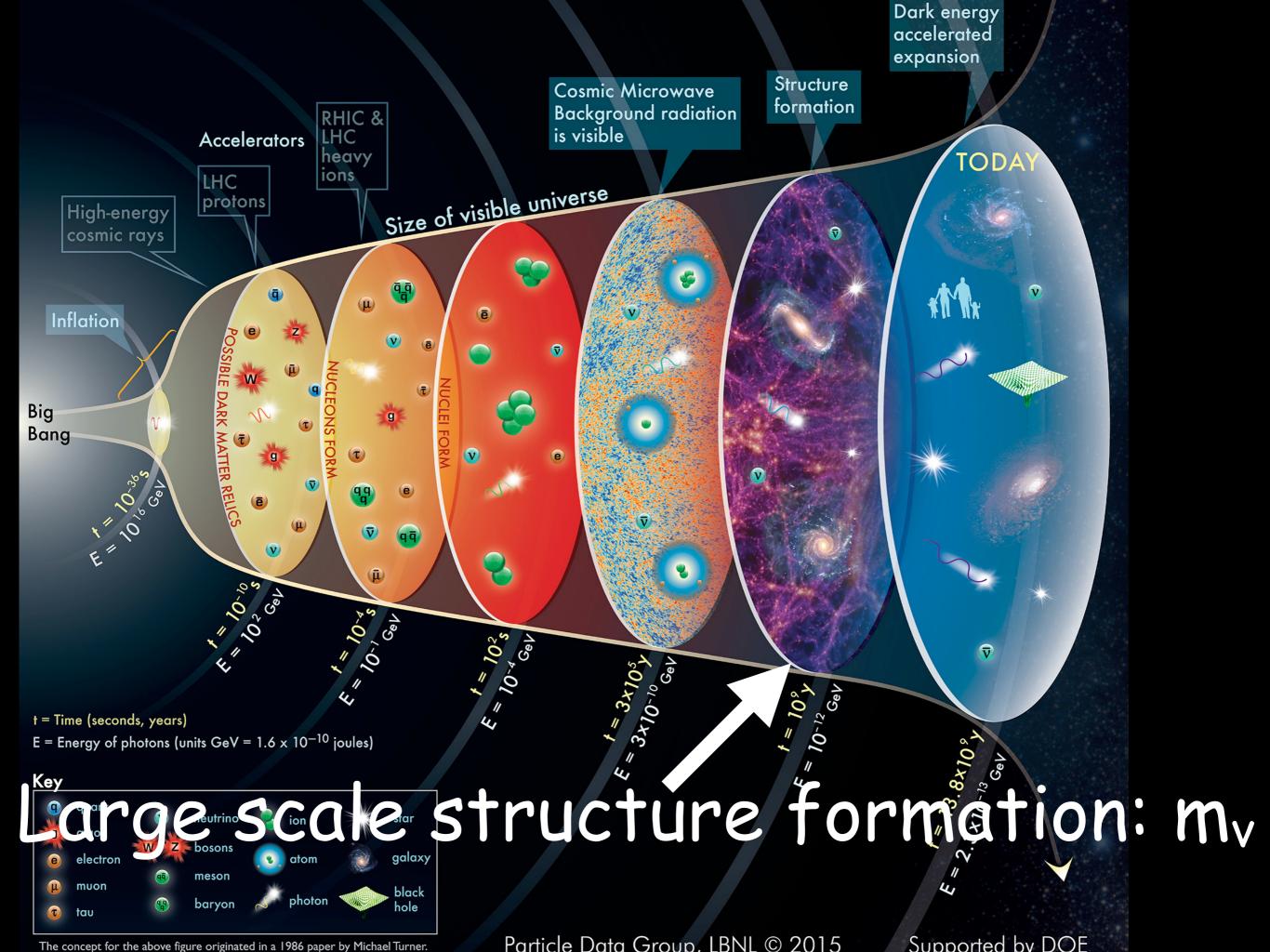
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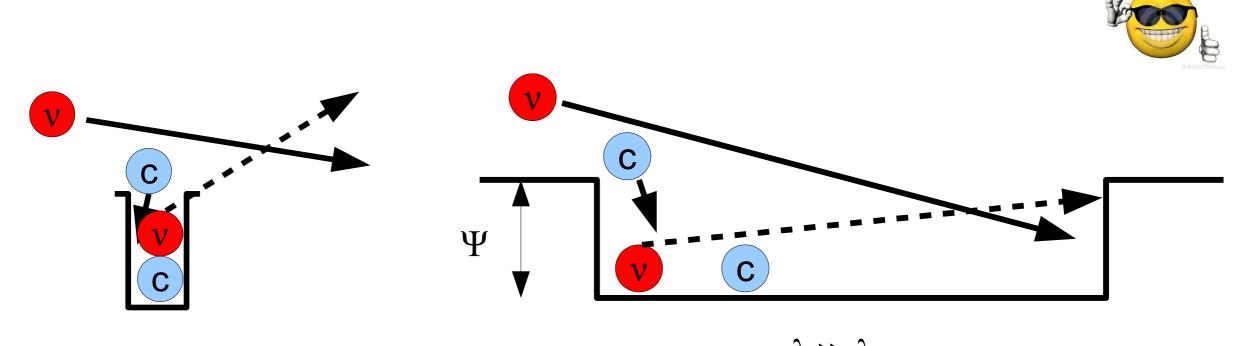


Large scale structure: m_V

Neutrino masses suppress structure formation on scales larger than their free streaming scale when they turn non relativistic. (Bond et al PRL'80)

Neutrinos with eV or sub-eV masses are HOT relics with LARGE thermal velocities!

Cold dark matter instead has zero velocity and therefore it clusters at any scale!



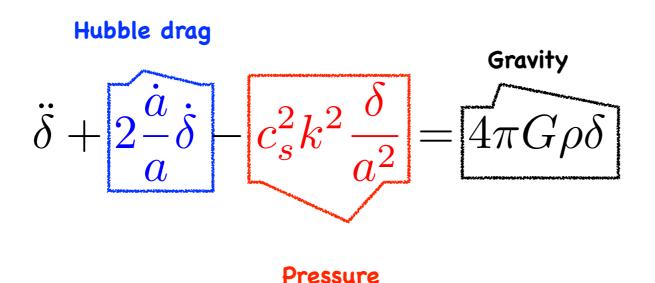


$$\lambda \ll^{\lambda} \widetilde{\mathcal{A}}_{fs,\nu}^{\lambda_{Fs}} \to k \gg k_{fs,\nu}$$

$$\lambda \gg \lambda_{fs,\nu}^{\lambda_{FS}} \to k \ll k_{fs,\nu}$$

Large scale structure: m_V

Growth equation for a single uncoupled fluid, linear regime, with constant sound speed:



Jeans scale:

$$k_J \equiv \sqrt{\frac{4\pi G\rho}{c_s^2(1+z)^2}}$$

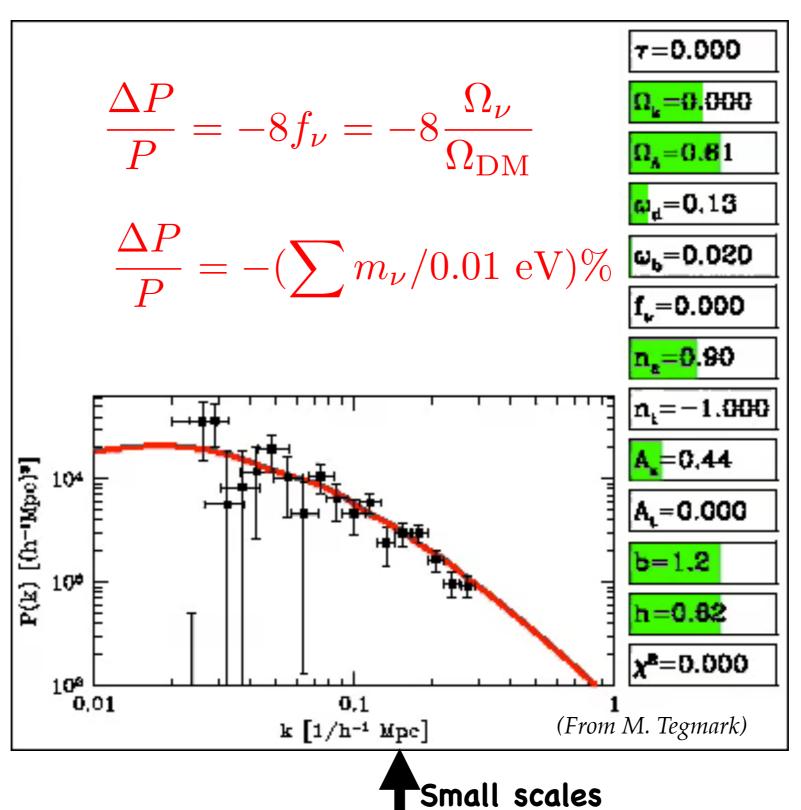
k>k_J no growth can occur k<k_J density perturbations growth

Neutrino free streaming scale:

$$k_{fs,\nu}(z) \equiv \sqrt{\frac{3}{2}} \frac{H(z)}{(1+z)\sigma_{v,\nu}(z)}$$

Large scale structure: m_V

Matter power spectrum suppression:

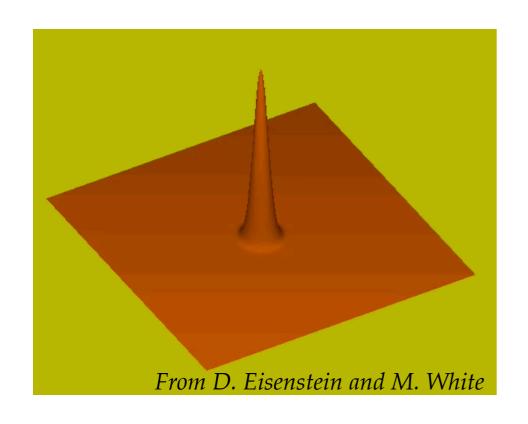


Baryon Acoustic Oscillations

Photons and baryons in the early universe behave as a tightly coupled fluid, resembling acoustic waves, generated as the baryon-photon fluid is attracted and falls onto the overdensities:

$$\ddot{\delta} + [\mathbf{Pressure - Gravity}]\delta = 0$$
 $\delta(\vec{x}) \equiv \frac{\rho(\vec{x}) - \bar{\rho}(\vec{x})}{\bar{\rho}(\vec{x})}$

The time when the baryons are "released" from the drag of the photons is known as the drag epoch. From then on photons expand freely while the acoustic waves "freeze in" the baryons at a scale given by the size of the horizon at the drag epoch:



$$r_s = 147.09 \pm 0.26 \text{ Mpc}$$

(Planck 2018 results)

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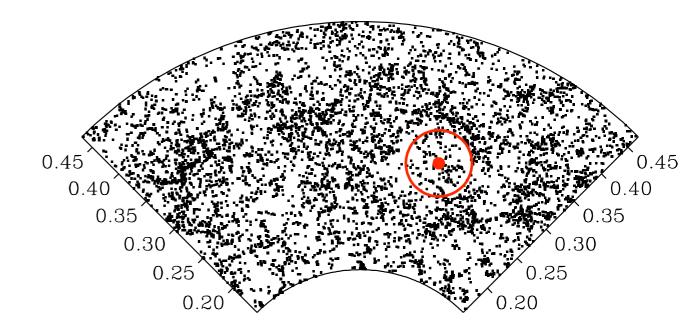
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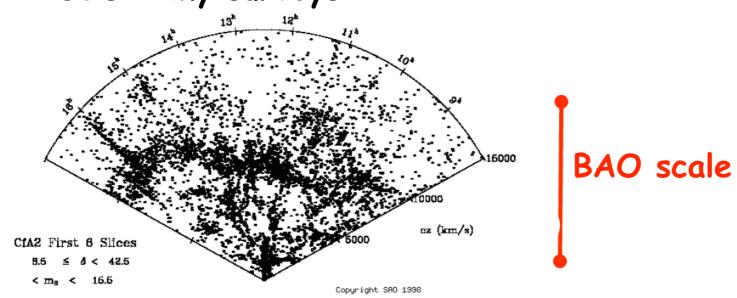
(Planck 2018 results)

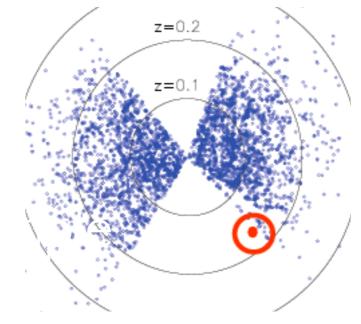
There should be a small excess in the two-point galaxy correlation function around 150 Mpc!

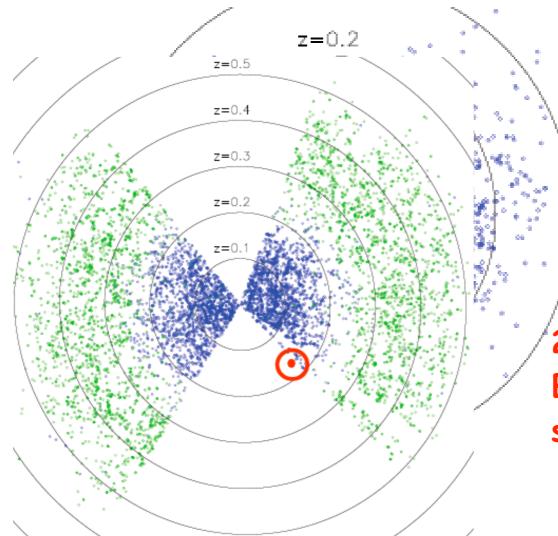


80's: Tiny surveys

Baryon Acoustic Oscillations







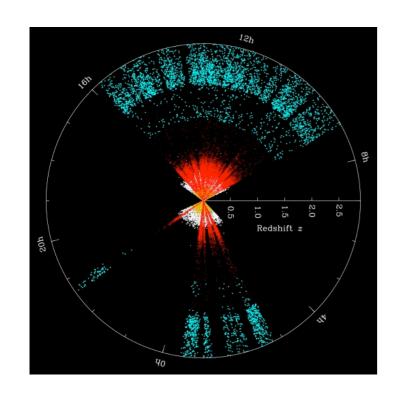
2000: Main galaxies @SDSS. Big number, but small volume

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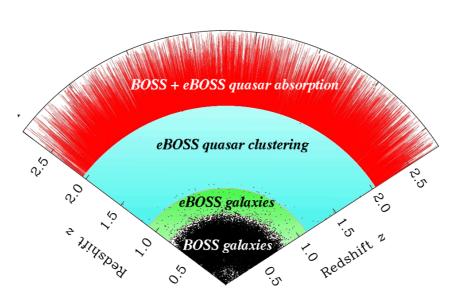
Baryon Acoustic Oscillations

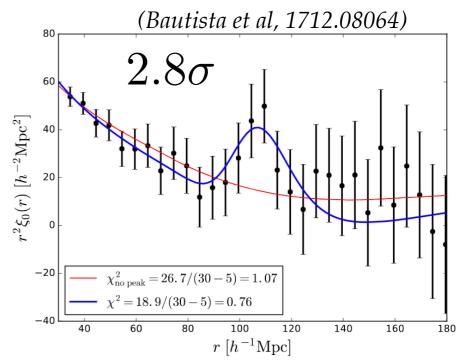
2009-2014= SDSS III

2014 - 2020= SDSS IV eBOSS



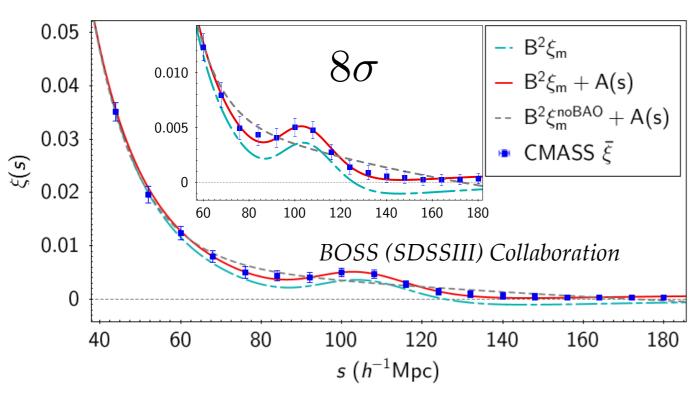
0.015

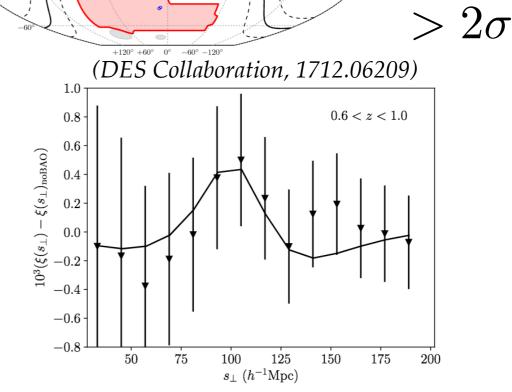




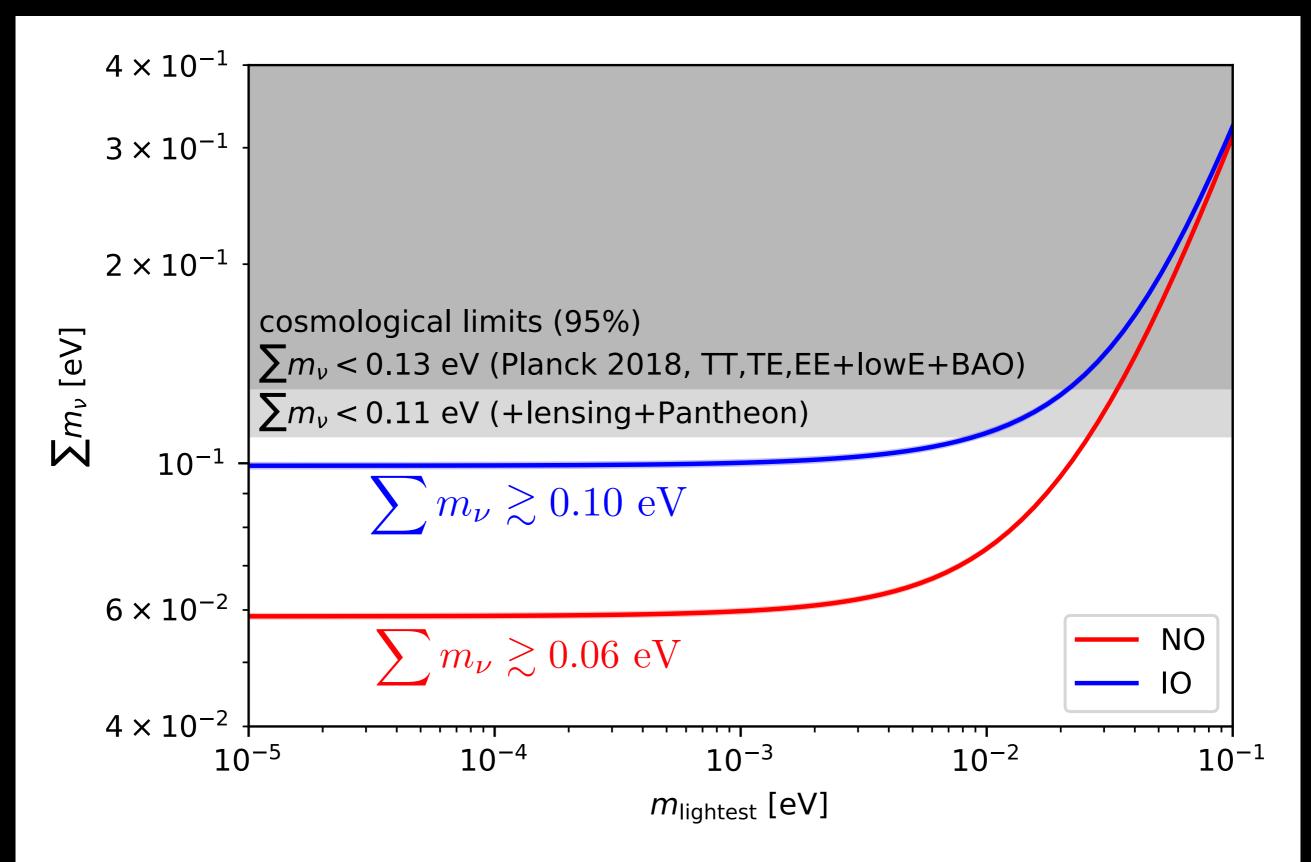
2013

2018= DES











Planck TTTEEE+lowT+lowE+lensing

$$\sum m_{\nu} < 0.24 \text{ eV } 95\% \text{CL}$$

+ BAO

$$\sum m_{\nu} < 0.12 \text{ eV } 95\% \text{CL}$$

+ BAO + SNIa

$$\sum m_{\nu} < 0.11 \text{ eV } 95\%\text{CL}$$

 $+ BAO + SNIa + H_0=73.45 \pm 1.66 \text{ km/s/Mpc}$

Riess et al, APJ'18

$$\sum m_{\nu} < 0.0970 \text{ eV } 95\% \text{CL}$$



Planck TTTEEE+lowT+lowE+lensing

$$\sum m_{\nu} < 0.26 \text{ eV}$$
 $N_{\text{eff}} = 2.90 \pm 0.37 95\% \text{CL}$

+ BAO

$$\sum m_{\nu} < 0.12 \text{ eV}$$
 $N_{\text{eff}} = 2.96^{+0.34}_{-0.33} 95\% \text{CL}$

+ BAO + SNIa

$$\sum m_{\nu} < 0.11 \text{ eV}$$
 $N_{\text{eff}} = 2.98^{+0.35}_{-0.33} 95\% \text{CL}$

+ BAO + SNIa + H₀=73.45 ±1.66 km/s/Mpc + w + nrun

$$\sum m_{\nu} < 0.16 \text{eV}$$
 $N_{\text{eff}} = 3.11^{+0.38}_{+0.38}$ 95%CL

Today's menu

Antipasto

- The classical cosmic pizza
- The neutrino slice
- Neutrino decoupling in the early universe
- Other cosmic pizzas for tasting

Main course

- Number of neutrinos and Big-Bang Nucleosynthesis
- Number of neutrinos and Cosmic Microwave Background Radiation
- Neutrino masses and Cosmic Microwave Background Radiation
- Neutrino masses and structure formation in the universe

Doggy Bag

Take home messages

The 3-Neutrino representative,





The Dark JUSTICE GAME:

The 3-Neutrino representative Σm_{ν}

versus

the Dark energy equation of state w(z)

GAME RULES:

- 1. Choose your favourite cosmological model for dark energy
- 2. Derive cosmological bounds on Σm_{ν} within that model, discarding neutrino oscillations (i.e. using just with the prior Σm_{ν} >0)
- 3. Are cosmological bounds consistent with oscillation data?

YES!

GREAT! You just won the game!
Your model is not ruled out (yet!)
Go to 1.

NO!

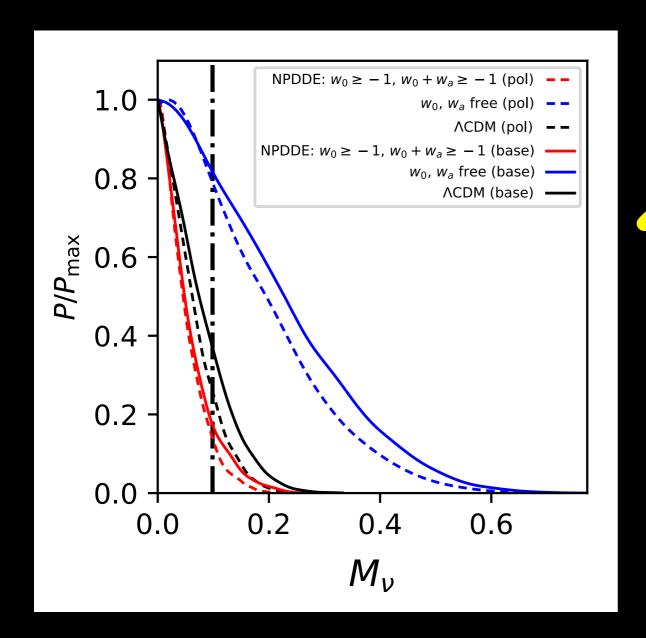
(i.e. Σm_{ν} <0.06 eV or Σm_{ν} <0.1 eV)

Write a paper
GAME OVER....after referral process

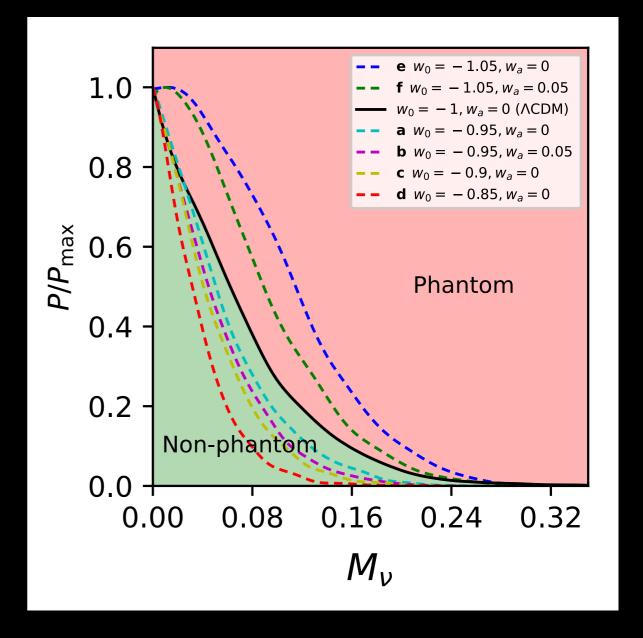
The Dark Justice Game

$$w(z) = w_0 + w_a \frac{z}{1+z}$$

Chevalier & Polarsky'01 & Linder'03



$w_0 \ge -1$ $w_0 + w_a \ge -1$ NON-PHANTOM REGION



Vagnozzi et al PRD'18
Choudury & Choubey, JCAP'18

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- Neutrino masses and structure formation in the universe
- The Dark Justice League game: Σm_{ν} versus w(z)

Doggy Bag

→ Take home messages

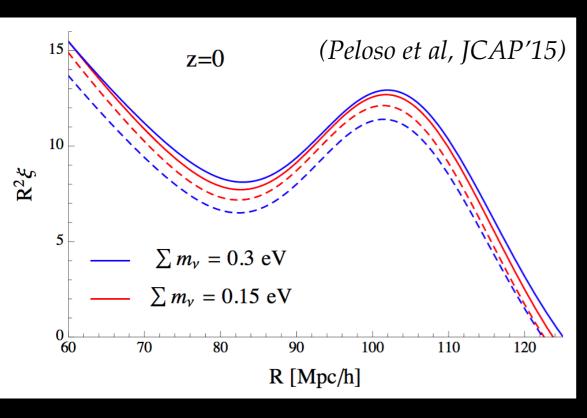
The "Take Home" messages

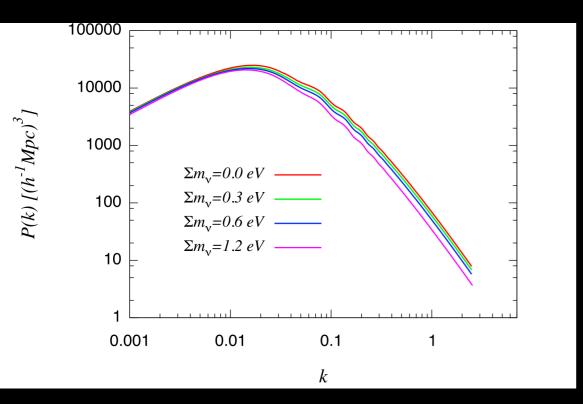


- v masses & abundances leave key signatures in cosmological observables.
- NO hints so far for neutrino masses or extra dark radiation species!
- N_{eff} @BBN: Light element abundances (⁴He) abundances.
- N_{eff} @CMB: damping tail
- N_{eff} = 2.99 +0.34_{-0.33}, (95% CL) from 2018 Planck TTTEEE+lensing, perfectly consistent with BBN.
- Cosmology provides currently the tightest bounds to neutrino masses.
- v masses@CMB: Early ISW, gravitational lensing
- v masses@LSS: Free streaming reduces clustering at small scales
- Σm,<0.12 eV (95%CL) from 2018 Planck TTTEEE+lensing plus BAO data
 </p>
- For non-phantom (physical) dynamical dark energy, ∑m, bounds get tighter!



What we know: Massive neutrinos cosmological signatures





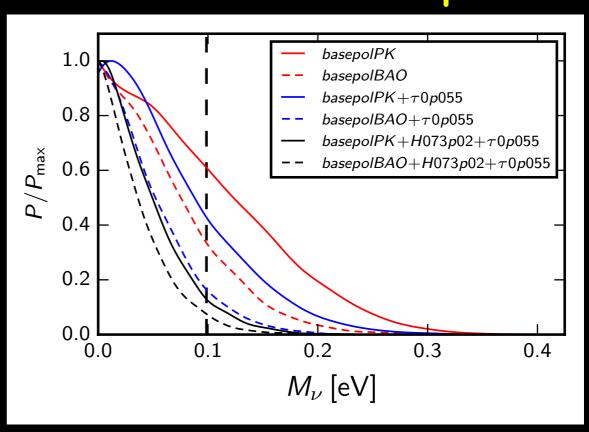
Large scale structure measurements can be interpreted either in the geometrical or shape forms

2 point correlation function

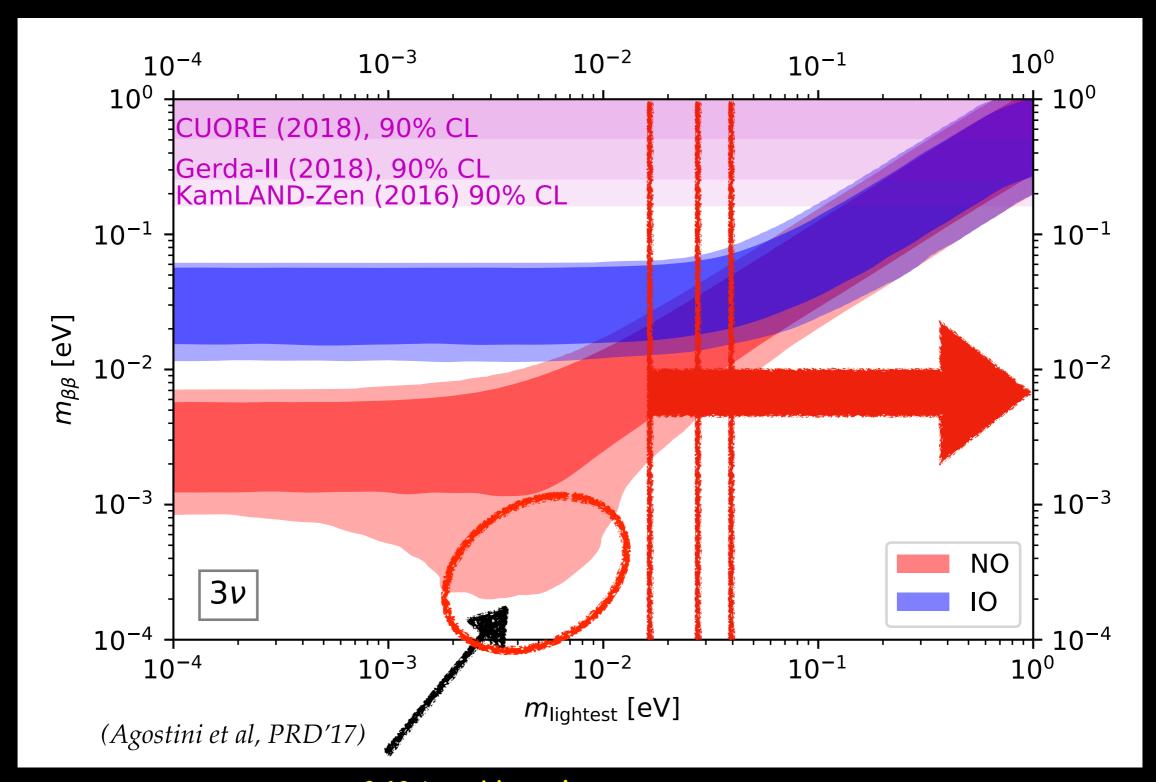
Fourier Transform

Matter power spectrum

BAO information more powerful

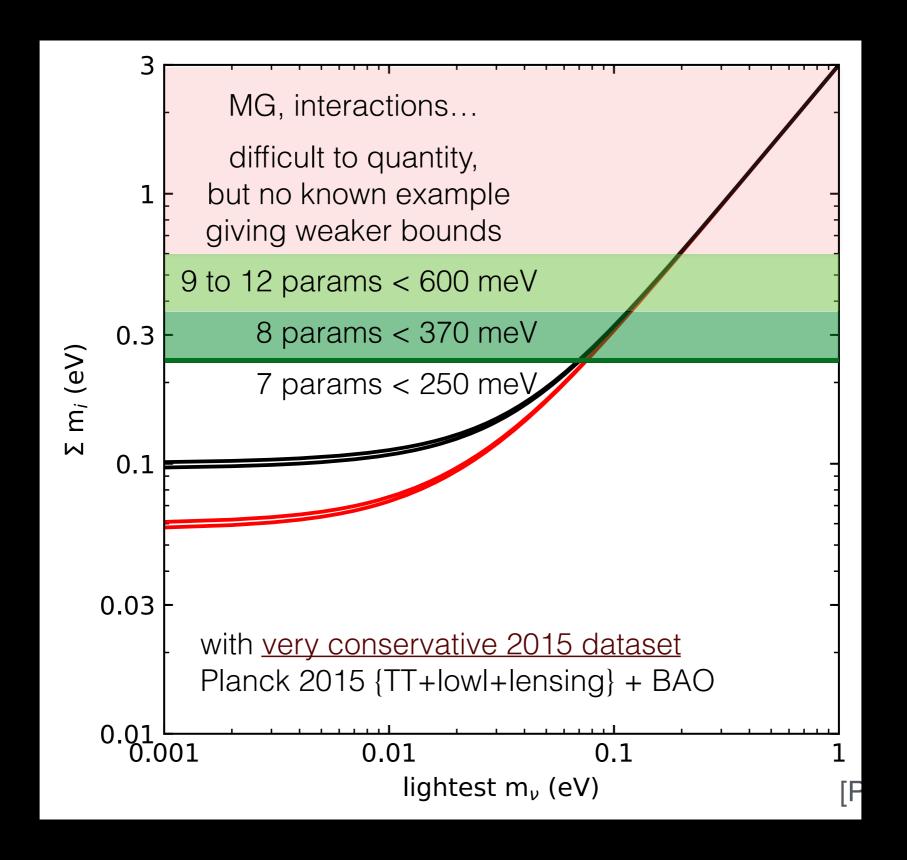


What we know



 $m_{\beta\beta}$ < 2 10⁻⁴ would require some fine tuning in the Majorana phases

What we know



BACKUP SLIDES

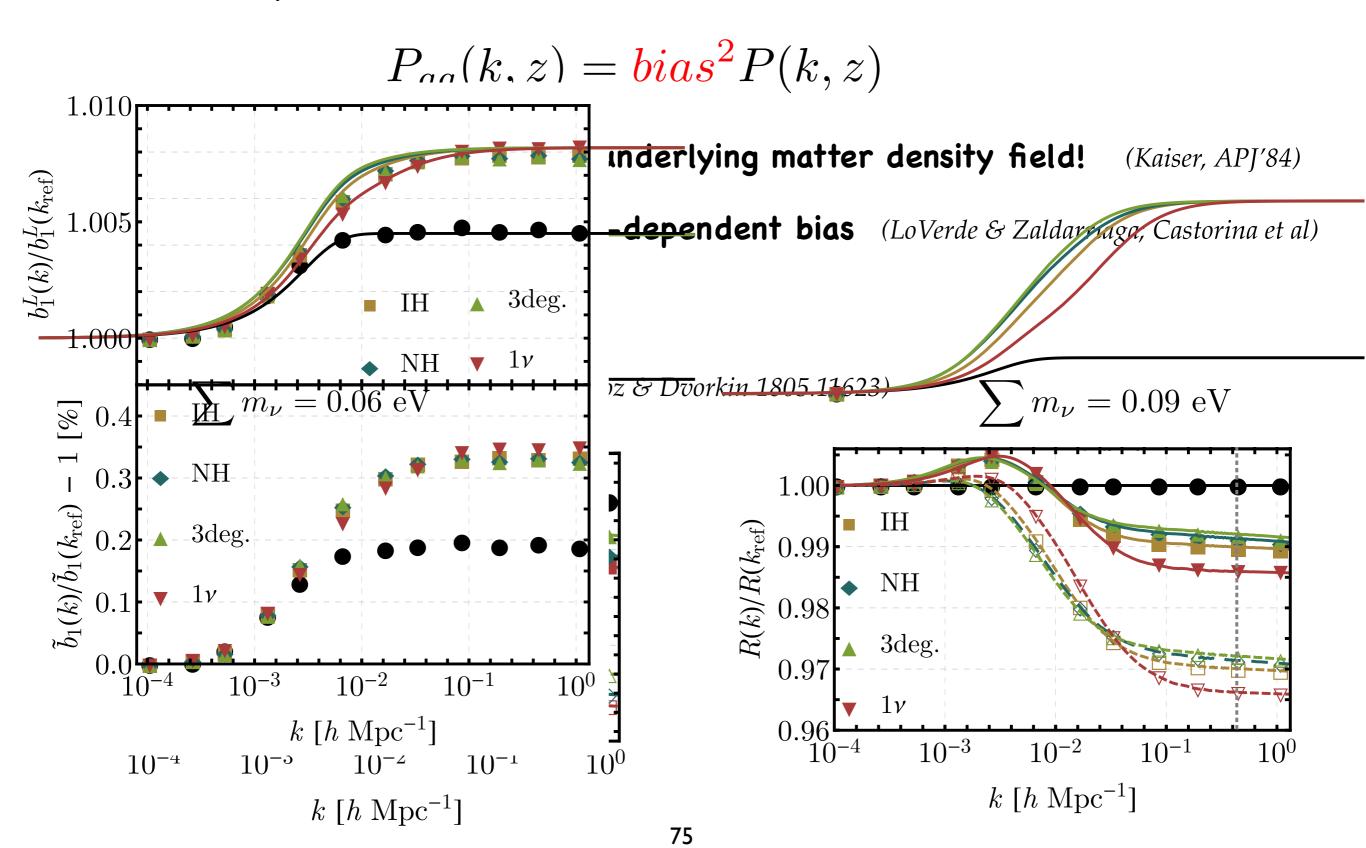






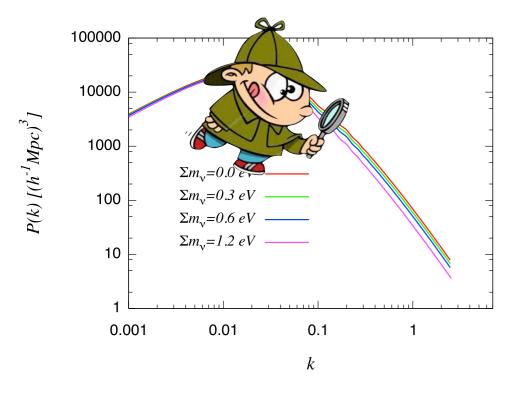
What we know: Massive neutrinos cosmological signatures

@LSS: Caveats, BIAS!

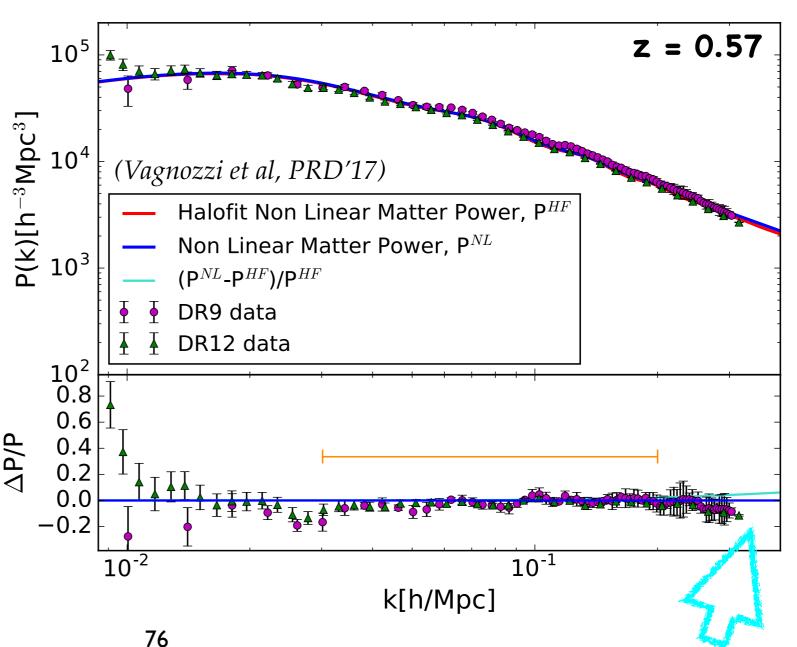


What we know: Massive neutrinos cosmological signatures

@LSS: Caveats, NON-LINEARITIES

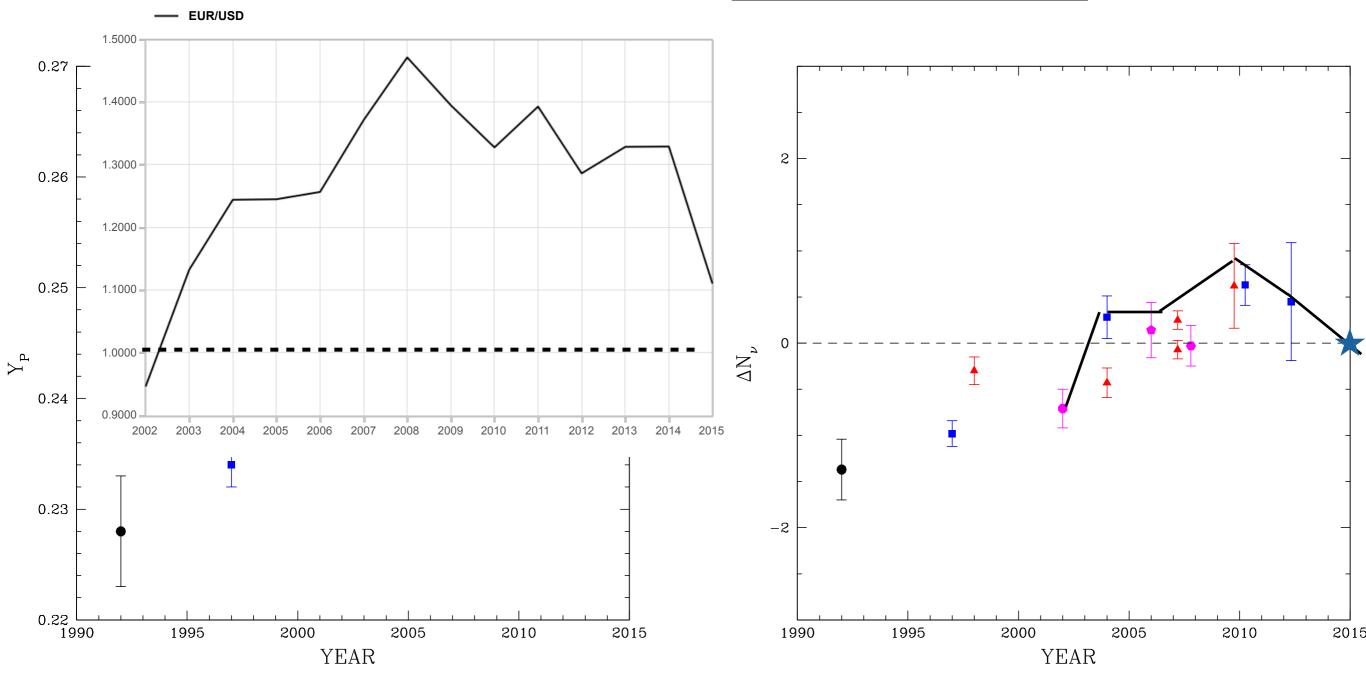


Beyond a given scale k_{nl} , linear perturbation theory breaks down!



Big Bang Nucleosynthesis: Neff

Chronology, over some years, of the published observational determinations of Y_P : the extracted value of ΔN_{eff} mirrors the helium fraction Y_P .



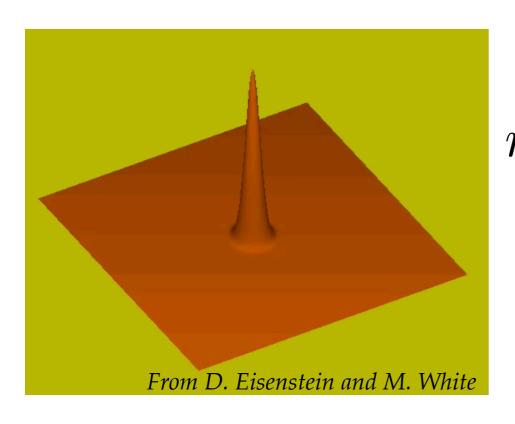
(G. Steigman'12)

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$$R \equiv 3\rho_b/4\rho_\gamma$$

$$\gamma_s = \int_0^{t(z_d)} c_s (1+z) dt = \frac{2}{3k_{\text{eq}}} \sqrt{\frac{6}{R_{\text{eq}}}} \ln \frac{\sqrt{1+R_d} + \sqrt{R_d + R_{\text{eq}}}}{1+\sqrt{R_{\text{eq}}}}$$

$$r_s = 147.09 \pm 0.26 \; \mathrm{Mpc}$$
 (Planck 2018 results)

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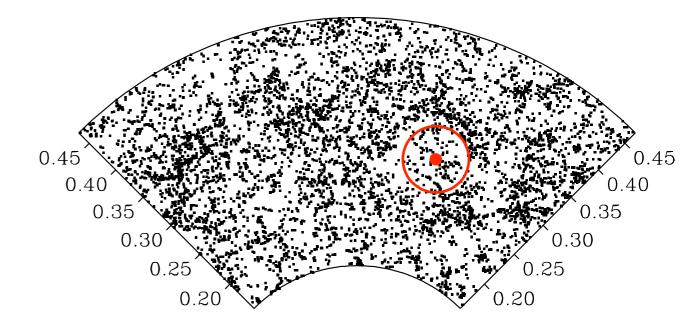
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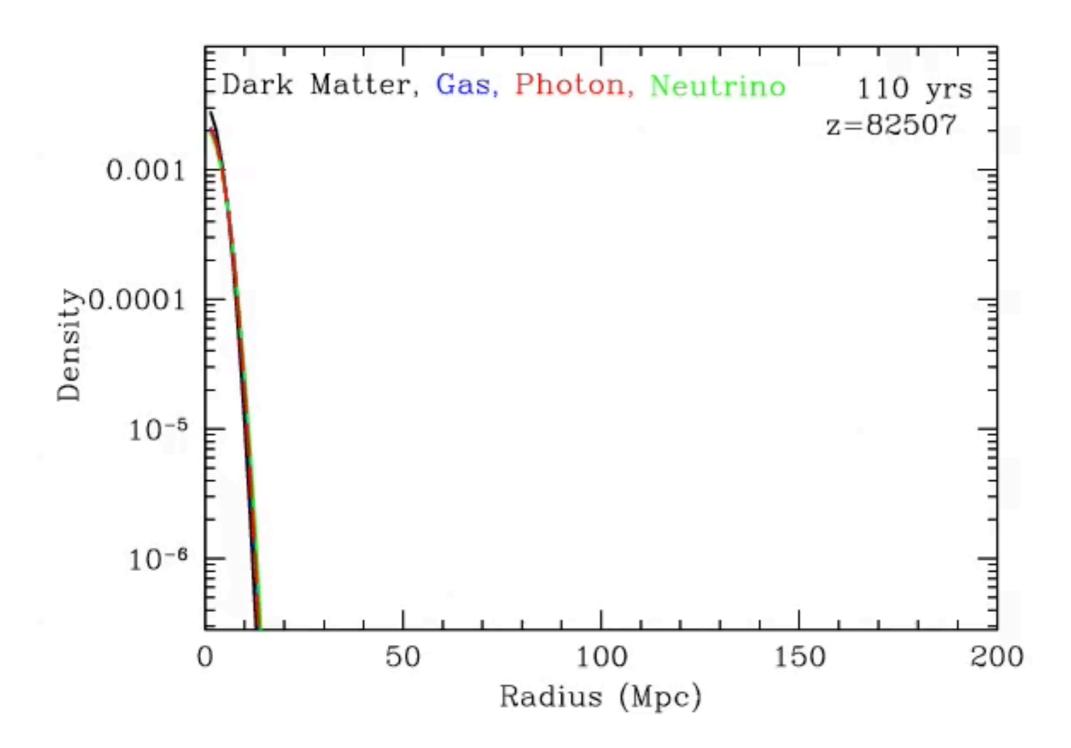
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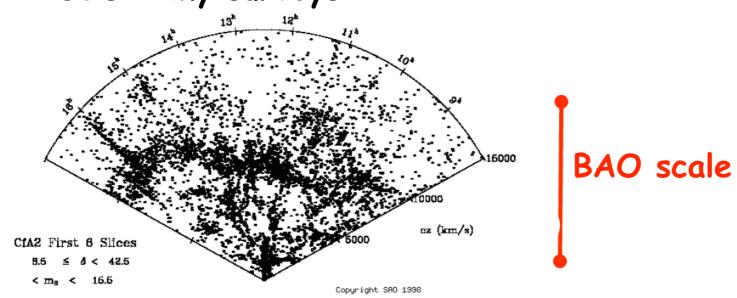
Baryon Acoustic Oscillations

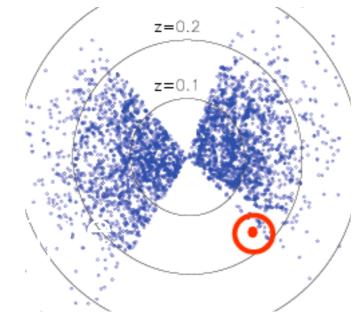


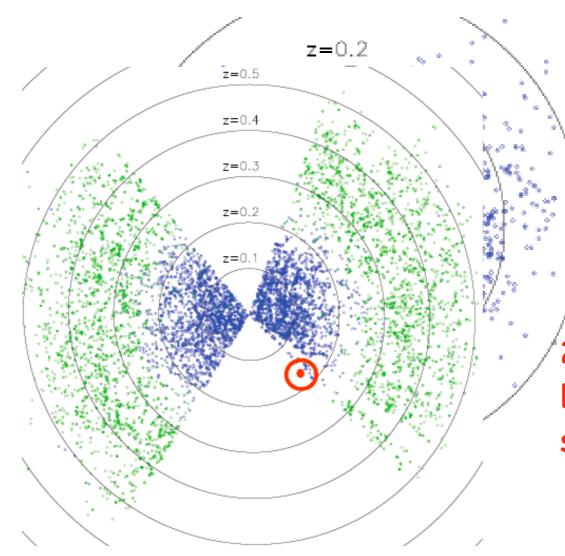
From D. Eisenstein and M. White

80's: Tiny surveys

Baryon Acoustic Oscillations





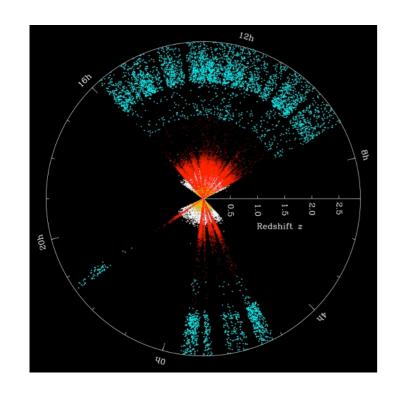


2000: Main galaxies @SDSS. Big number, but small volume

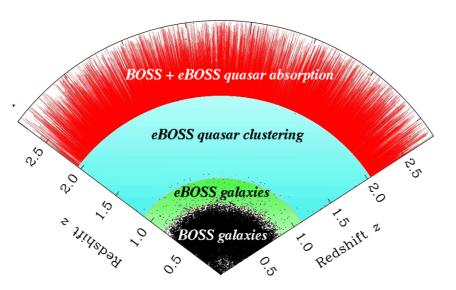
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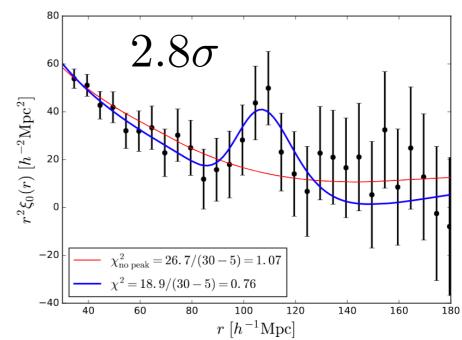
2009-2014= SDSS III

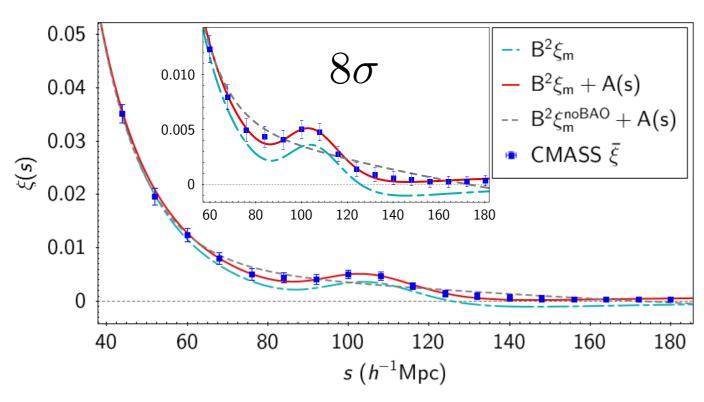
Baryon Acoustic Oscillations 2014 - 2020= SDSS IV eBOSS

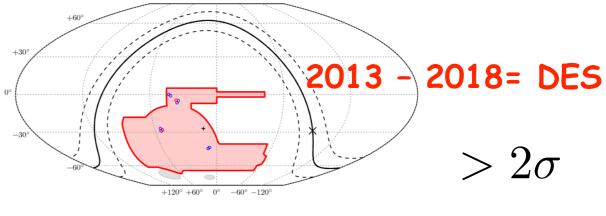


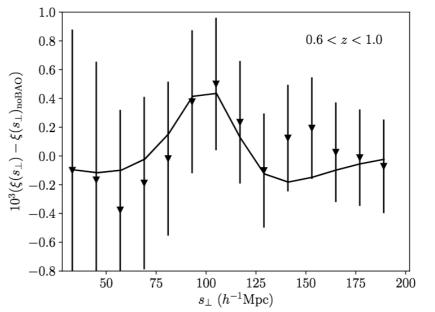
0.015



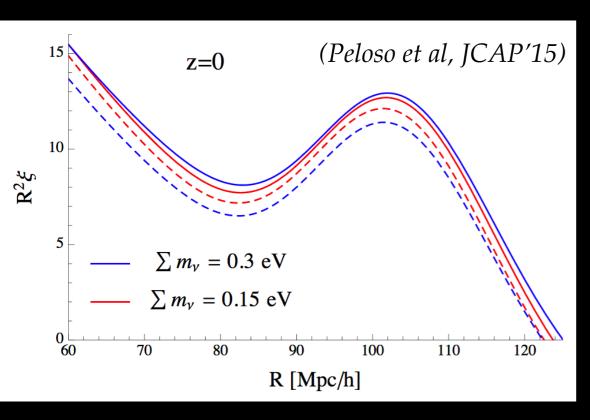








Large scale structure and mv



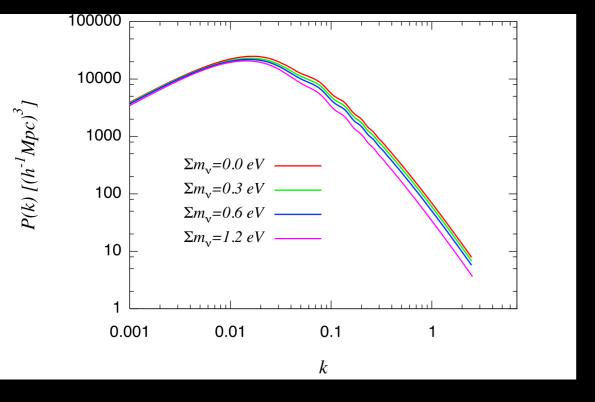
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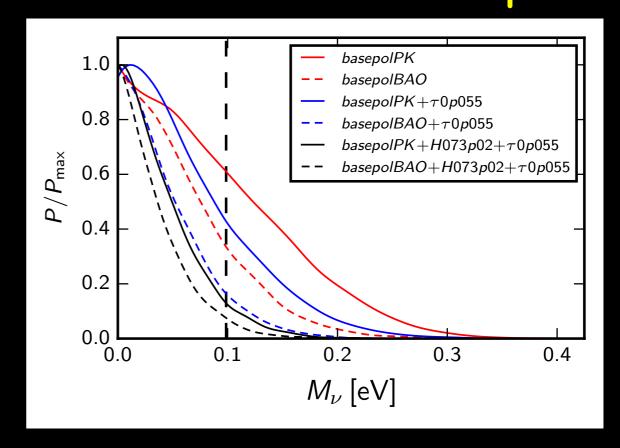
2 point correlation function

Fourier Transform

Matter power spectrum

BAO information still more powerful





One can look for (sterile) neutrinos in something not so

shiny and bright....

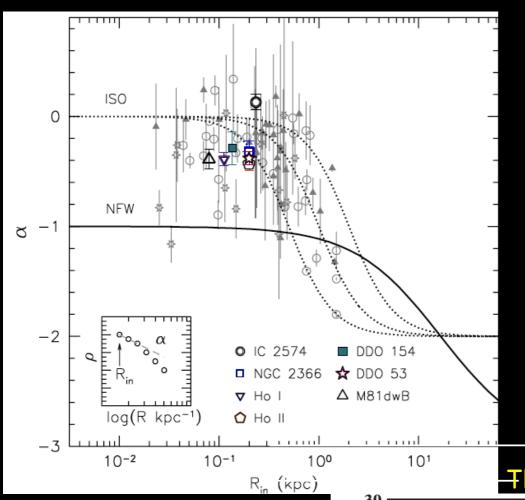
COLD dark matter **HOT** dark matter CDM WDM WARM dark matter

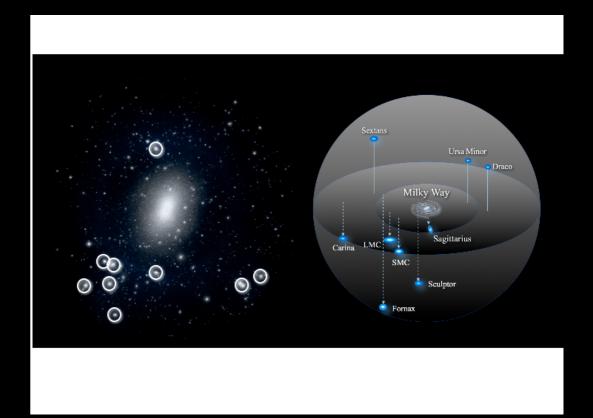
Small scale crisis of ACDM@galactic and sub-galactic scales

Why do these subhalos not

18 km/s

12 km/s





Missing satellite problem

The predicted satellite population far exceeds the observed one

Core/Cusp problem

Observations seem to indicate 25 n approximately constant galaxies (core), while cosmological simulations indicate 4 km

15

10

tter density in the inner parts of vertaw-like behaviou(T844) problem

Massive dark subhalos are too dense to match data. Expected 10 subhalos in the Milky Way with v >30 km/s, only 3 known!

(Boylan et al, MNRAS'11)

A controversial unidentified line has been detected at with a significance > 3σ in two independent samples of X-ray clusters with XMM-Newton.

It is independently seen by the same group in the Perseus Cluster with Chandra data.

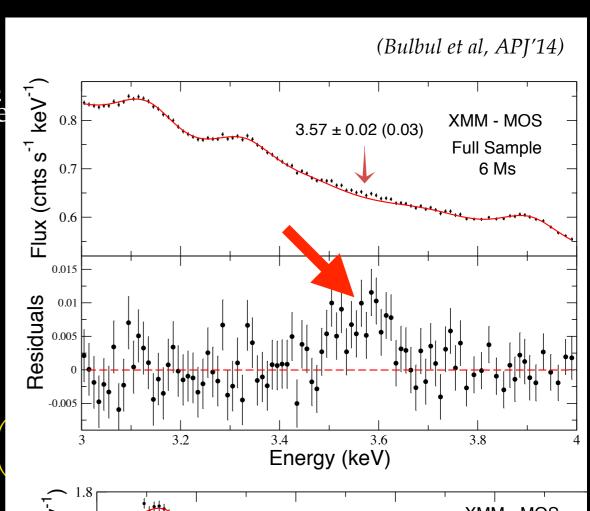
(Bulbul et al, APJ'14)

An independent group finds a line at the same energy toward Andromeda and Perseus with XMM-Newton, with a combined statistical evidence of 4.40. (Boyarsky et al, PRL'14)

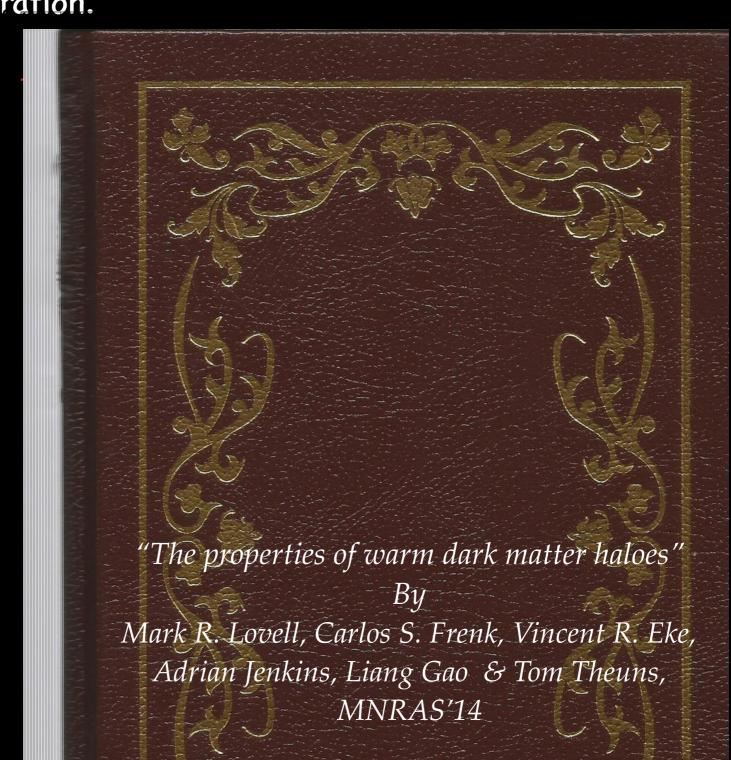
$$u_s
ightarrow
u_{\alpha} + \gamma
ightharpoonup_{v_2}
ightharpoonup_{v_1}
ightharpoonup_{v_1}$$

$$m_s = 2E = 7.1 keV$$
 $\Gamma_{\gamma = 1.62 \times 10^{-28} \text{ s}^{-1}}$





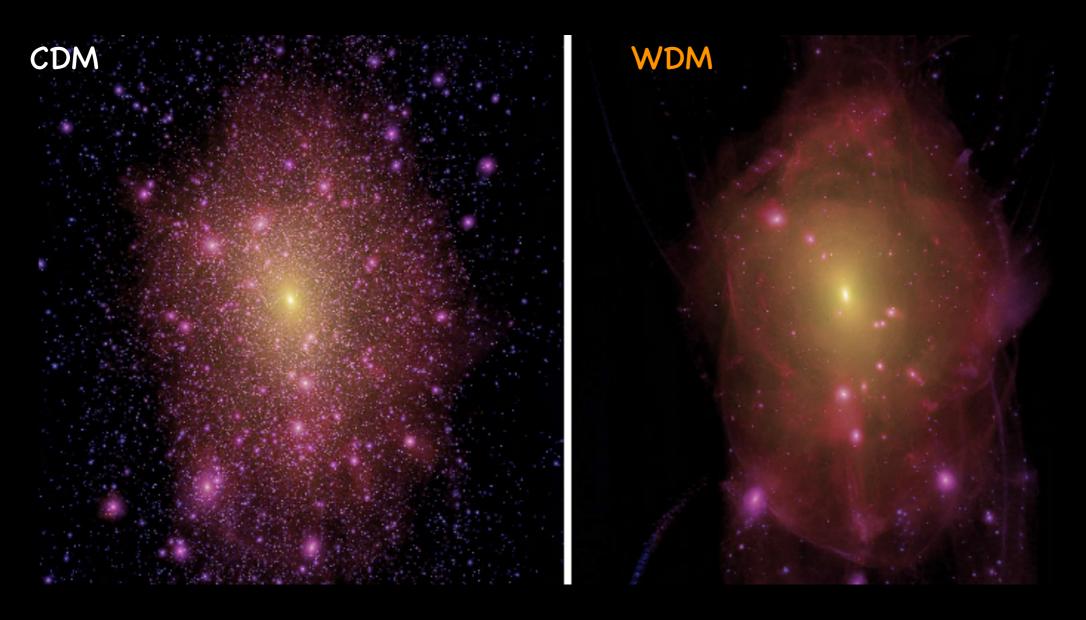
WDM leads to an identical large scale structure pattern than CDM, but very different subhaloes abundance, structure and dynamics: the free streaming of a keV sterile neutrino will reduce power at the small scales, delaying structure formation and lowering the haloes concentration.



WDM leads to an identical large scale structure pattern than CDM, but very different subhaloes abundance, structure and dynamics: the free streaming of a keV sterile neutrino will reduce power at the small scales, delaying structure formation and lowering the haloes concentration.



WDM could reconcile theory with observations!



"The Haloes of Bright Satellite Galaxies in a Warm Dark Matter Universe", Mark R. Lovell, Vincent R. Eke, Carlos S. Frenk, Liang Gao, Adrian Jenkii Jie Wang, D.M. White, Alexey Boyarsky & Oleg Ruchayskiy MNRAS'12

"The properties of warm dark matter haloes", Mark R. Lovell, Carlos 🗫 renk, Vincent R. Eke, Adrian Jenkins, Liang Gao & Tom Theuns, MNRAS'14

Backup material

The entropy density is:

$$s \equiv \frac{\rho + p}{T}$$

¿How are related the photon and the neutrino temperatures?

- · Electron positron annihilation takes place AFTER neutrino decoupling.
- In an expanding universe the entropy density pero comoving volume is conserved:
- Boson's entropy contribution:
- Fermion's entropy contribution:

$$2\pi^2 T^3 / 45$$
$$7/8 \times 2\pi^2 T^3 / 45$$

•Before electron/positron annihilation= electrons (g=2), positrons (g=2), neutrinos (3), antineutrinos (3) and photons (g=2) therefore:

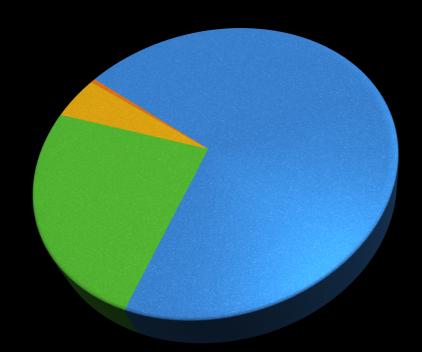
$$s(a_1) = 2\pi^2 T_1^3 / 45(2 + 7/8(2 + 2 + 3 + 3))$$

After, only neutrinos, antineutrinos and photons but at different temperature!

$$s(a_2) = 2\pi^2/45(2T_{\gamma}^3 + 7/8(3+3)T_{\nu}^3)$$

$$s(a_1)a_1^3 = s(a_2)a_2^3 \qquad a_1T_1 = a_2T_{\nu} \qquad \left(\frac{T_{\nu}}{T_{\gamma}}\right) = \left(\frac{4}{11}\right)^{1/3}$$

Summary's Dessert Cosmic Cake



Sterile neutrino masses and abundances leave key signatures in cosmological observables.

Abundances (N_{eff}): via Cosmic Microwave Background measurements damping tail & Big Bang Nucleosynthesis light element abundances, N_{eff} <3.41 (@95%CL).

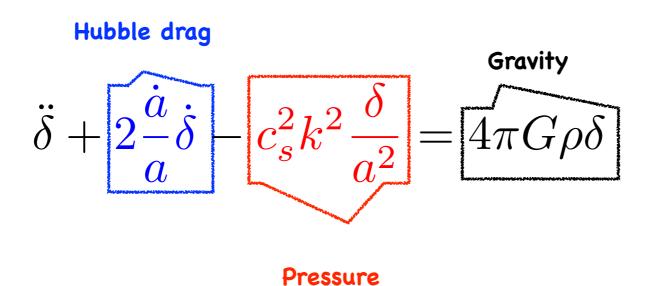
Masses: neutrino free-streaming nature induce a small-scale suppression matter power spectrum, $m_{\nu,eff}$ < 0.7 eV (< 0.0000014 m_e) (@95%CL).

No current cosmological evidence for sterile neutrinos or dark radiation.

Warm Dark Matter sterile keV neutrinos: alleviate/solve mostly all the small scale problems of the standard Cold Dark Matter scenario, still a very interesting option!

Massive neutrinos cosmological signatures: m_V

Growth equation for a single uncoupled fluid, linear regime, with constant sound speed:



Jeans scale:

$$k_J \equiv \sqrt{\frac{4\pi G\rho}{c_s^2(1+z)^2}}$$

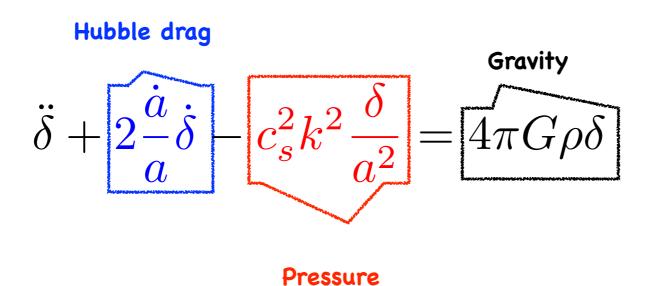
k>k_J no growth can occur k<k_J density perturbations growth

Neutrino free streaming scale:

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Neff dark radiation species cosmological signatures

N_{eff} changes the freeze out temperature of weak interactions:

$$\Gamma_{n \leftrightarrow p} \sim H$$

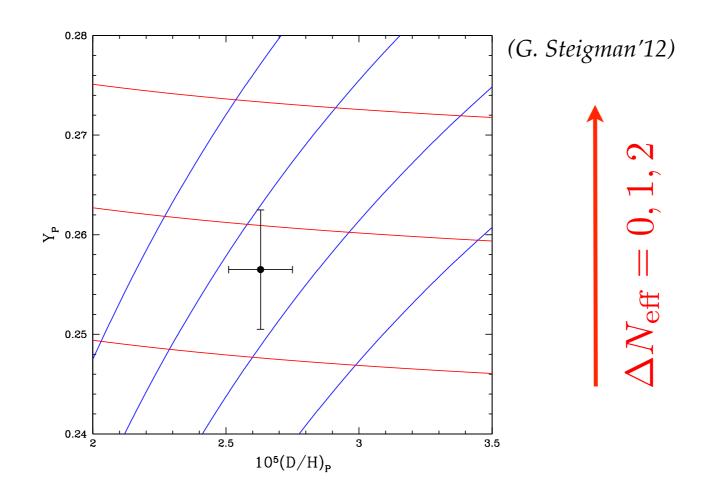
$$\Omega_r h^2 = \left(1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}}\right) \Omega_\gamma h^2$$

 $\Omega_r h^2 = \left(1+\tfrac{7}{8}\left(\tfrac{4}{11}\right)^{4/3} N_{\rm eff}\right) \Omega_\gamma h^2$ Higher Neff, Higher expansion rate, higher freeze out temperature, higher ⁴He

fraction:

$$n/p \simeq e^{-\frac{m_n - m_p}{T_{freeze}}}$$

$$Y_p = \frac{2(n/p)}{1 + n/p}$$



• The entropy density is:

$$s \equiv \frac{\rho + p}{T}$$

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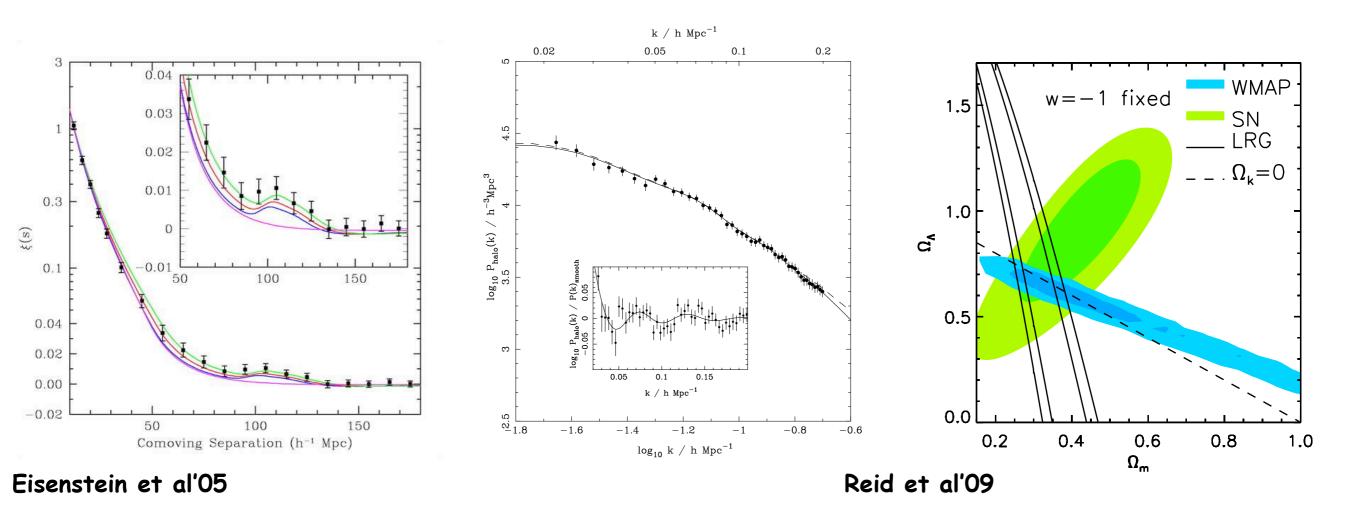
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Los catálogos de galaxias miden la función de correlación:

$$\xi(\vec{r}) \equiv \langle \delta(\vec{x}) \delta(\vec{x} + \vec{r}) \rangle_{\text{Volume}} \qquad \langle \tilde{\delta}(\vec{k}) \tilde{\delta}(\vec{k}') \rangle_{\text{Volume}} = (2\pi)^3 P(k) \delta^3(\vec{k} - \vec{k}')$$

$$\delta(\vec{x}) \equiv \frac{\rho(\vec{x}) - \bar{\rho}(\vec{x})}{\bar{\rho}(\vec{x})} \qquad \qquad \tilde{\delta}(\vec{k}) \equiv \int d^3 \vec{r} \ e^{i\vec{k}\vec{r}} \ \delta(\vec{r})$$



SSDS 2005: Primera detección de la señal BAO (3.4s) (47000 LRGs, 4000 deg 2 , z=0.35) SDSS II 2009: 110 000 LRGs, 8000 deg 2 , z=0.35.