# Model independent approach to galaxy clustering

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## Standard perturbation theory

Dark matter: non relativistic, non interacting, subject to gravitational interaction

$$\frac{\partial \delta(\mathbf{x}, \tau)}{\partial \tau} + \nabla_{\mathbf{x}} \Big[ \Big( 1 + \delta(\mathbf{x}, \tau) \Big) \, \mathbf{v}(\mathbf{x}, \tau) \Big] = 0 \quad \text{Continuity equation}$$

$$\frac{\partial \mathbf{v}(\mathbf{x},\tau)}{\partial \tau} + \mathcal{H}(\tau)\mathbf{v}(\mathbf{x},\tau) + \left(\mathbf{v}\cdot\nabla_{\mathbf{x}}\right)\mathbf{v}(\mathbf{x},\tau) + \nabla_{\mathbf{x}}\phi(\mathbf{x},\tau) = \frac{1}{\rho}\nabla_{\mathbf{x}}\tau$$

Euler equation

$$\nabla_{\mathbf{x}}^2 \phi(\mathbf{x}, \tau) = \frac{3}{2} \Omega_m(\tau) \mathcal{H}^2(\tau) \delta(\mathbf{x}, \tau) \quad \text{Poisson equation}$$

## Standard perturbation theory Observational effects

From dark matter to galaxies: bias, redshift space, small scale effects

$$\delta(\mathbf{k}, \tau) \longrightarrow \delta_g^{\mathrm{rs}}(\mathbf{k}, \tau) = \mathscr{F}[\delta, \partial_i \partial_j \phi, \dots](\mathbf{k}, \tau)$$

Perturbative solution

$$\delta_g^{\text{rs},(n)}(\mathbf{k}) = \int \frac{d^3 q_1}{(2\pi)^3} \dots \int \frac{d^3 q_n}{(2\pi)^3} \delta_D(\mathbf{k} - \mathbf{q}_{1...n}) Z_n\left(\mathbf{q}_1, \dots, \mathbf{q}_n\right) \delta^{(1)}(\mathbf{q}_1) \dots \delta^{(1)}(\mathbf{q}_n)$$

$$\delta_g^{\rm rs}(\mathbf{k}) = \delta_g^{\rm rs, PT}(\mathbf{k}) + \delta_g^{\rm rs, SN}(\mathbf{k}) + \delta_g^{\rm rs, CT}(\mathbf{k})$$

## Standard perturbation theory Results



D'Amico et al. 1909.05271 Ivanov et al. 1909.05277 Colas et al. 1909.07951 Chen et al. 2110.05530





## Standard perturbation theory Results



SPT



## Role of symmetries in cosmology



## Role of symmetries in cosmology

Equivalence principle

Model independent measurement using the EP in the BAO range

- MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)
- MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Constraints the perturbative kernels in a model independent way: Large Scale Structure Bootstrap!

• D'Amico G., **MM**, Pietroni M., Vernizzi F., JCAP 10 (2021)





Equivalence principle



Equivalence principle



SPT

Equivalence principle



## Consistency relations

Exact equalities among correlation functions of different order (Kehagias A., Riotto A., Nucl.Phys. 2013, Peloso M., Pietroni M., JCAP 2013, Creminelli P., Noreña J., Simonović M., Vernizzi F., JCAP 2013)

$$\langle \delta(\mathbf{x}_1, \tau_1) \dots \delta(\mathbf{x}_n, \tau_n) | \Phi_L \rangle = \langle \delta(\tilde{\mathbf{x}}_1, \tilde{\tau}_1) \dots \delta(\tilde{\mathbf{x}}_n, \tilde{\tau}_n) \rangle.$$

Effect of long modes on linear scales — Equivalence Principle (or Galilean Invariance)

$$\langle \delta_m(\mathbf{q},\tau)\delta_g(\mathbf{k}_1,\tau_1)\dots\delta_g(\mathbf{k}_n,\tau_n)\rangle'_{\mathbf{q}\to 0} = -P_L(q,\tau)\sum_{i=1}^n \frac{D_+(\tau_i)}{D(\tau)}\frac{\mathbf{q}\cdot\mathbf{k}_i}{q^2}\langle \delta_g(\mathbf{k}_1,\tau_1)\dots\delta_g(\mathbf{k}_n,\tau_n)\rangle'$$



$$\delta_g^{(n)}(\mathbf{k}) = \int \frac{d^3 q_1}{(2\pi)^3} \dots \int \frac{d^3 q_n}{(2\pi)^3} K_n\left(\mathbf{q}_1, \dots, \mathbf{q}_n\right) \delta^{(1)}(\mathbf{q}_1) \dots \delta^{(1)}(\mathbf{q}_n)$$



$$\varphi_{\mathbf{q}_{1}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{\mathbf{q}_{2}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{\mathbf{q}_{2}} \underbrace{\varphi_{\mathbf{q}_{n-1}}}_{\mathbf{q}_{n-1}} \underbrace{\delta_{g}^{(n)}(\mathbf{k})}_{\mathbf{q}_{2}} = \int \frac{d^{3}q_{1}}{(2\pi)^{3}} \dots \int \frac{d^{3}q_{n}}{(2\pi)^{3}} K_{n} \left(\mathbf{q}_{1}, \dots, \mathbf{q}_{n}\right) \delta^{(1)}(\mathbf{q}_{1}) \dots \delta^{(1)}(\mathbf{q}_{n})$$

$$\varphi_{\mathbf{q}_{1}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{\mathbf{q}_{2}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{\mathbf{q}_{n-1}} \underbrace{\varphi_{\mathbf{q}_{n-1}}}_{\mathbf{q}_{n-1}} \underbrace{\delta_{g}^{(n)}(\mathbf{k})}_{\mathbf{q}_{g}^{(n)}(\mathbf{k})} = \int \frac{d^{3}q_{1}}{(2\pi)^{3}} \dots \int \frac{d^{3}q_{n}}{(2\pi)^{3}} K_{n}\left(\mathbf{q}_{1}, \dots, \mathbf{q}_{n}\right) \delta^{(1)}(\mathbf{q}_{1}) \dots \delta^{(1)}(\mathbf{q}_{n})$$

• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)

$$\varphi_{\mathbf{q}_{1}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{q_{2}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{q_{n-1}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{q_{n-1}} \underbrace{\varphi_{\mathbf{q}_{2}}}_{q_{n-1}} \underbrace{\varphi_{\mathbf{q}_{n-1}}}_{q_{n}} \underbrace{\varphi_{\mathbf{q}_{n}}}_{q_{n}} \underbrace{\varphi_{\mathbf{q}_{n}}}_{q_{n$$

Mass and momentum conservation (only for dark matter)













• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)



• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)

- Leading Order: single momentum going  $\rightarrow 0$
- **Next-to-Leading Order**: sum of two momenta going  $\rightarrow 0$
- Next-to-Next-to-Leading Order: sum of three momenta going  $\rightarrow 0$

• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)

### Equivalence principle

Leading Order: single momentum going  $\rightarrow 0$ 

Next-to-Leading Order: sum of two momenta going  $\rightarrow 0$ 

N.<sup>*l*-1</sup>-to-Leading Order: sum of l - 1momenta going  $\rightarrow 0$ 

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N.<sup>*l*-1</sup>-to-Leading Order: sum of l - 1momenta going  $\rightarrow 0$ 

$$\lim_{\mathbf{q}_{1},\dots,\mathbf{q}_{m}\to 0} K_{n}(\mathbf{q}_{1},\dots,\mathbf{q}_{m},\mathbf{q}_{m+1},\dots,\mathbf{q}_{n}) = \frac{\mathbf{q}_{1}\cdot\mathbf{Q}_{n,m}}{q_{1}^{2}}\dots\frac{\mathbf{q}_{m}\cdot\mathbf{Q}_{n,m}}{q_{m}^{2}}K_{n-m}(\mathbf{q}_{m+1},\dots,\mathbf{q}_{n}) + O\left(\left(\frac{1}{q}\right)^{m-1}\right)$$

• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)

### Equivalence principle

Leading Order: single momentum going  $\rightarrow 0$ 

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N.<sup>*l*-1</sup>-to-Leading Order: sum of l - 1 momenta going  $\rightarrow 0$ 

$$\lim_{\mathbf{q}_1 + \mathbf{q}_2 \to 0} K_n(\mathbf{q}_1, \mathbf{q}_2, \mathbf{q}_3, \dots, \mathbf{q}_n) \supset \frac{\mathbf{k} \cdot \mathbf{q}_{12}}{q_{12}^2} K_{n-2}(\mathbf{q}_3, \dots, \mathbf{q}_n) \int^{\eta} d\eta' f_+(\eta') \frac{D_+(\eta')^2}{D_+(\eta)^2} G_2(\mathbf{q}_1, \mathbf{q}_2; \eta')$$

• D'Amico G., MM, Pietroni M., Vernizzi F., JCAP 10 (2021)

### Equivalence principle

Leading Order: single momentum going  $\rightarrow 0$ 

**Next-to-Leading Order**: sum of two momenta going  $\rightarrow 0$ 

N.<sup>*l*-1</sup>-to-Leading Order: sum of *l* momenta going  $\rightarrow 0$ 

$$\lim_{\mathbf{Q}_{l,0}\to 0} K_n(\mathbf{q}_1,\dots,\mathbf{q}_l,\mathbf{q}_{l+1},\dots,\mathbf{q}_n) \supset \frac{\mathbf{k}\cdot\mathbf{Q}_{l,0}}{Q_{l,0}^2} \int^{\eta} d\eta' f_+(\eta') \left(\frac{D_+(\eta')}{D_+(\eta)}\right)^l G_l(\mathbf{q}_1,\dots,\mathbf{q}_l;\eta') K_{n-l}(\mathbf{q}_{l+1},\dots,\mathbf{q}_n;\eta)$$

Kernel at second order

 $K_1(\mathbf{q_1}) = c_0$ 

 $K_2(\mathbf{q_1}, \mathbf{q_2}) = c_1 + c_\beta \beta(\mathbf{q_1}, \mathbf{q_2}) + c_\gamma \gamma(\mathbf{q_1}, \mathbf{q_2})$ 

$$\beta(\mathbf{q}_1, \mathbf{q}_2) = \frac{|\mathbf{q}_1 + \mathbf{q}_2|^2 \mathbf{q}_1 \cdot \mathbf{q}_2}{2q_1^2 q_2^2}$$
$$\gamma(\mathbf{q}_1, \mathbf{q}_2) = 1 - \left(\frac{\mathbf{q}_1 \cdot \mathbf{q}_2}{q_1 q_2}\right)^2$$
$$\alpha_a(\mathbf{q}_1, \mathbf{q}_2) = \frac{\mathbf{q}_1 \cdot \mathbf{q}_2}{q_1^2} - \frac{\mathbf{q}_1 \cdot \mathbf{q}_2}{q_2^2}$$

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$$\{c_0, c_1, c_\beta, c_\gamma\}$$



Kernel at second order

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 $\{c_0, c_1, c_\beta, c_\gamma\}$ 

Leading Order



Kernel at second order

 $K_1(\mathbf{q_1}) = c_0$ 

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 $\{c_0, c_1, c_\beta, c_\gamma\}$ 

Leading Order

Only 3 parameters left! (tracers)



Kernel at second order

 $K_1(\mathbf{q_1}) = c_0$ 

 $K_2(\mathbf{q_1}, \mathbf{q_2}) = c_1 + c_\beta \beta(\mathbf{q_1}, \mathbf{q_2}) + c_\gamma \gamma(\mathbf{q_1}, \mathbf{q_2})$ 

 $\{c_0, c_1, c_\beta, c_\gamma\}$ 

Only 3 parameters left! (tracers)

$$\beta(\mathbf{q}_1, \mathbf{q}_2) = \frac{|\mathbf{q}_1 + \mathbf{q}_2|^2 \mathbf{q}_1 \cdot \mathbf{q}_2}{2q_1^2 q_2^2}$$
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Leading Order Mass+momentum conservation (matter)

Kernel at second order

 $K_1(\mathbf{q_1}) = c_0$ 

 $K_2(\mathbf{q_1}, \mathbf{q_2}) = c_1 + c_\beta \beta(\mathbf{q_1}, \mathbf{q_2}) + c_\gamma \gamma(\mathbf{q_1}, \mathbf{q_2})$ 

 $\{C_0, C_1, C_\beta, C_\gamma\}$ 

Only 3 parameters left! (tracers)

$$\beta(\mathbf{q}_1, \mathbf{q}_2) = \frac{|\mathbf{q}_1 + \mathbf{q}_2|^2 \mathbf{q}_1 \cdot \mathbf{q}_2}{2q_1^2 q_2^2}$$
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Leading Order Mass+momentum conservation (matter)

Only 1 parameter left! (matter)

Kernel at third order

 $K_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}) = c_{2} + c_{\gamma 1} \gamma(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\gamma 2} \gamma(\mathbf{q}_{12}, \mathbf{q}_{3}) + c_{\beta 1} \beta(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\beta 2} \beta(\mathbf{q}_{12}, \mathbf{q}_{3})$ 

 $+c_{\gamma\gamma}\gamma(\mathbf{q}_{1},\mathbf{q}_{2})\gamma(\mathbf{q}_{12},\mathbf{q}_{3})+c_{\beta\beta}\beta(\mathbf{q}_{1},\mathbf{q}_{2})\beta(\mathbf{q}_{12},\mathbf{q}_{3})+c_{\gamma\beta}\gamma(\mathbf{q}_{1},\mathbf{q}_{2})\beta(\mathbf{q}_{12},\mathbf{q}_{3})$ 

$$+c_{\beta\gamma}\beta(\mathbf{q}_1,\mathbf{q}_2)\gamma(\mathbf{q}_{12},\mathbf{q}_3)+\left(c_{\alpha}+c_{\gamma\alpha}\gamma(\mathbf{q}_1,\mathbf{q}_2)+c_{\beta\alpha}\beta(\mathbf{q}_1,\mathbf{q}_2)\right)\alpha_a(\mathbf{q}_{12},\mathbf{q}_3)$$

Kernel at third order

 $K_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}) = c_{2} + c_{\gamma 1} \gamma(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\gamma 2} \gamma(\mathbf{q}_{12}, \mathbf{q}_{3}) + c_{\beta 1} \beta(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\beta 2} \beta(\mathbf{q}_{12}, \mathbf{q}_{3})$ 

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$$+c_{\beta\gamma}\beta(\mathbf{q}_1,\mathbf{q}_2)\gamma(\mathbf{q}_{12},\mathbf{q}_3) + \left(c_{\alpha} + c_{\gamma\alpha}\gamma(\mathbf{q}_1,\mathbf{q}_2) + c_{\beta\alpha}\beta(\mathbf{q}_1,\mathbf{q}_2)\right)\alpha_a(\mathbf{q}_{12},\mathbf{q}_3)$$

 $\{c_2, c_{\gamma 1}, c_{\gamma 2}, c_{\beta 1}, c_{\beta 2}, c_{\gamma \gamma}, c_{\beta \beta}, c_{\gamma \beta}, c_{\beta \gamma}, c_{\alpha}, c_{\gamma \alpha}, c_{\beta \alpha}\}$ 

Kernel at third order

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Leading Order

Kernel at third order

 $K_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}) = c_{2} + c_{\gamma 1} \gamma(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\gamma 2} \gamma(\mathbf{q}_{12}, \mathbf{q}_{3}) + c_{\beta 1} \beta(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\beta 2} \beta(\mathbf{q}_{12}, \mathbf{q}_{3})$ 

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Leading Order

Next-to-Leading Order

Equivalence principle

Kernel at third order

 $K_{3}(\mathbf{q}_{1}, \mathbf{q}_{2}, \mathbf{q}_{3}) = c_{2} + c_{\gamma 1} \gamma(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\gamma 2} \gamma(\mathbf{q}_{12}, \mathbf{q}_{3}) + c_{\beta 1} \beta(\mathbf{q}_{1}, \mathbf{q}_{2}) + c_{\beta 2} \beta(\mathbf{q}_{12}, \mathbf{q}_{3})$ 

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Leading Order Next-to-Leading Order Only 4 parameters left! (tracers)

Kernel at third order

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Leading Order Next-to-Leading Order Mass+mom. conservation Only 4 parameters left! (tracers)

Kernel at third order

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Leading OrderNext-to-Leading OrderMass+mom. conservationOnly 4 parameters left!Only 2 parameters left!(tracers)(matter)

General time dependence

Bootstrap expansion: valid for every scaleindependent model

$$-\frac{k^2}{\mathcal{H}^2}\phi(\mathbf{k},\eta) = \frac{3}{2}\Omega_m(\eta)\mu(\eta)\delta(\mathbf{k},\eta) + S(\mathbf{k},\eta)$$

ACDM (exact time dep., Donath Y., Senatore L., 2020) nDGP (Dvali et al. 2000, Schmidt F. 2009) k-mouflage (Babichev et al. 2009, Brax and Valaegias 2014) JBD (Brans C., Dicke R., 1961)

D'Amico G., MM, Piga L., Pietroni M., Vernizzi F., Wright B., 2209.XXXX

## Matter and velocity kernels



In the 1-loop PS the only time dependent functions are  $a_{\gamma}, d_{\gamma}, d_{\gamma a} \equiv d_{\gamma \alpha} - d_{\gamma \gamma}/2$ 

Constrain (possible) deviations from ΛCDM with a model-independent approach

SPT



## LSS Bootstrap & PNG

Consistency relations

$$\lim_{k \to 0} \frac{B_{\delta}(k, q, |\mathbf{k} + \mathbf{q}|; \tau, \tau', \tau'')}{P_{\delta}(q; \tau', \tau'')P_{\delta}(k; \tau, \tau)} = -\frac{\mathbf{q} \cdot \mathbf{k}}{k^2} \frac{D_{+}(\tau') - D_{+}(\tau'')}{D_{+}(\tau)} + \frac{6f_{NL}\Omega_{m,0}H_0^2}{k^2T(k)} \frac{D_{+}(\tau_0)}{D_{+}(\tau)} + O\left(k^0, f_{NL}^2\right)$$

Peloso M., Pietroni M., JCAP 05 (2013) 031 Goldstein S., et al., 2209.06228

Equivalence principle

## LSS Bootstrap & PNG



SPT

Equivalence principle

## Conclusions

- Importance of Equivalence Principle in Cosmology
- Perturbative EP: LSS bootstrap
- Easily generalized to NG IC
- First step for a novel approach to Galaxy Clustering

## Thanks for your attention



## In this thesis

### Equivalence principle

Model independent measurement using the EP in the BAO range

- MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)
- MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Constraints the perturbative kernels in a model independent way: Large Scale Structure Bootstrap!

• D'Amico G., **MM**, Pietroni M., Vernizzi F., JCAP 10 (2021)

Bootstrap

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Squeezed limit of **Bispectrum** (real space)

$$\lim_{q/k \to 0} B_{\alpha\beta\gamma}(q, k_+, k_-; t_\alpha, t_\beta, t_\gamma) \simeq \frac{\mathbf{k} \cdot \mathbf{q}}{q^2} P_{\alpha m}(q; t_\alpha) \left[ \frac{D(t_\beta)}{D(t_\alpha)} P_{\beta\gamma}(k_-; t_\beta, t_\gamma) - \frac{D(t_\gamma)}{D(t_\alpha)} P_{\beta\gamma}(k_+; t_\beta, t_\gamma) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$

In configuration space

$$\propto \mu^2 \frac{d \log \xi(r)}{d \log r}$$

Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)



Equivalence principle

CR and BAO

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

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$$\lim_{q/k\to 0} B_{\alpha\beta\gamma}(q,k_+,k_-;t_\alpha,t_\beta,t_\gamma) \simeq \frac{\mathbf{k}\cdot\mathbf{q}}{q^2} P_{\alpha m}(q;t_\alpha) \left[ \frac{D(t_\beta)}{D(t_\alpha)} P_{\beta\gamma}(k_-;t_\beta,t_\gamma) - \frac{D(t_\gamma)}{D(t_\alpha)} P_{\beta\gamma}(k_+;t_\beta,t_\gamma) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$

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Equal-time squeezed limit

Equivalence principle

CR and BAO

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Squeezed limit of **Bispectrum** (real space)

$$\lim_{q/k\to 0} B_{\alpha\beta\gamma}(q,k_+,k_-;t_\alpha,t_\beta,t_\gamma) \simeq \frac{\mathbf{k}\cdot\mathbf{q}}{q^2} P_{\alpha m}(q;t_\alpha) \left[ \frac{D(t_\beta)}{D(t_\alpha)} P_{\beta\gamma}(k_-;t_\beta,t_\gamma) - \frac{D(t_\gamma)}{D(t_\alpha)} P_{\beta\gamma}(k_+;t_\beta,t_\gamma) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$

Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)

CR and BAO

In configuration space

$$\propto \mu^2 \frac{d \log \xi(r)}{d \log r}$$

$$\mu = \hat{\mathbf{k}} \cdot \hat{\mathbf{q}}$$

$$r_{BAO}$$

$$r_{BAO}$$

$$r_{BAO}$$
(2015)

Bootstrap

Equal-time squeezed limit

$$\lim_{q/k \to 0} \frac{B_{\alpha\alpha\alpha}(q, k_+ \cdot k_-)}{P_{\alpha\alpha}(q)P_{\alpha\alpha}(k)} = -\frac{\mu^2}{b_{\alpha}(q)} \frac{d\log P_{\alpha\alpha}(k)}{d\log k} + O\left(\left(\frac{q}{k}\right)^0\right)$$

Unchanged by / nonlinearities!

Beyond ACDM

SPT

Baldauf T. et al., Phys.Rev.D (2015)

MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)
MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Squeezed limit of **Bispectrum** (real space)

Beyond ACDM

SPT

$$\lim_{q|k\to 0} B_{\alpha\beta\gamma}(q, k_{+}, k_{-}; t_{\alpha}, t_{\beta}, t_{\gamma}) \approx \frac{\mathbf{k} \cdot \mathbf{q}}{q^{2}} P_{\alpha m}(q; t_{\alpha}) \left[ \frac{D(t_{\beta})}{D(t_{\alpha})} P_{\beta\gamma}(k_{-}; t_{\beta}, t_{\gamma}) - \frac{D(t_{\gamma})}{D(t_{\alpha})} P_{\beta\gamma}(k_{+}; t_{\beta}, t_{\gamma}) \right] + O\left(\left(\frac{q}{k}\right)^{0}\right)$$

$$Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)$$

$$\mu = \hat{\mathbf{k}} \cdot \hat{\mathbf{q}}$$

$$\lim_{q|k\to 0} \frac{B_{\alpha\alpha\alpha}(q, k_{+}, k_{-})}{P_{\alpha\alpha}(q)P_{\alpha\alpha}(k)} = \left(\frac{\mu^{2}}{b_{\alpha}(q)}\right) \frac{d \log P_{\alpha\alpha}(k)}{d \log k} + O\left(\left(\frac{q}{k}\right)^{0}\right)$$

$$Baldauf T. et al., Phys.Rev.D (2015)$$

$$In presence of a scale like the BAO, the oscillating part of the derivative is enhanced by a ~ k r_{BAO} factor, we can isolate it to verify CR and measure bias$$

$$SPT \quad \text{Equivalence principle} \qquad CR and BAO \qquad \text{Bootstrap}$$

Equivalence principle

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Squeezed limit of **Bispectrum** (real space)

$$\lim_{q/k\to 0} B_{\alpha\beta\gamma}(q,k_+,k_-;t_\alpha,t_\beta,t_\gamma) \simeq \frac{\mathbf{k}\cdot\mathbf{q}}{q^2} P_{\alpha m}(q;t_\alpha) \left[ \frac{D(t_\beta)}{D(t_\alpha)} P_{\beta\gamma}(k_-;t_\beta,t_\gamma) - \frac{D(t_\gamma)}{D(t_\alpha)} P_{\beta\gamma}(k_+;t_\beta,t_\gamma) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$

Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)



Bootstrap

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

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Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)

$$\mu = \hat{\mathbf{k}} \cdot \hat{\mathbf{q}}$$

### Equal-time squeezed limit



Bootstrap

#### Equivalence principle

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

Squeezed limit of **Bispectrum** (real space)

$$\lim_{q/k \to 0} B_{\alpha\beta\gamma}(q, k_+, k_-; t_\alpha, t_\beta, t_\gamma) \simeq \frac{\mathbf{k} \cdot \mathbf{q}}{q^2} P_{\alpha m}(q; t_\alpha) \left[ \frac{D(t_\beta)}{D(t_\alpha)} P_{\beta\gamma}(k_-; t_\beta, t_\gamma) - \frac{D(t_\gamma)}{D(t_\alpha)} P_{\beta\gamma}(k_+; t_\beta, t_\gamma) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$

Peloso M. and Pietroni M., JCAP (2013) Kehagias A. and Riotto A., Nucl. Phys. (2013)

$$\mu = \hat{\mathbf{k}} \cdot \hat{\mathbf{q}}$$

Equal-time squeezed limit

$$\lim_{q/k \to 0} \frac{B_{\alpha\alpha\alpha}(q, k_+, k_-)}{P_{\alpha\alpha}(q)P_{\alpha\alpha}(k)} = -\frac{\mu^2}{b_{\alpha}(q)} \frac{d\log P_{\alpha\alpha}(k)}{d\log k} + O\left(\left(\frac{q}{k}\right)^0\right)$$

Unchanged by / nonlinearities!

Bootstrap

• MM, Nishimichi T., Pietroni M., Phys. Rev. D100 (2019)

• MM, Nishimichi T., Pietroni M., JCAP 07 (2020)

 $b_{\alpha}(q)$ 

Squeezed limit of **Bispectrum** (real space)

$$\lim_{q/k\to 0} B_{\alpha\beta\gamma}(q,k_+,k_-;t_{\alpha},t_{\beta},t_{\gamma}) \simeq \frac{\mathbf{k}\cdot\mathbf{q}}{q^2} P_{\alpha m}(q;t_{\alpha}) \left[ \frac{D(t_{\beta})}{D(t_{\alpha})} P_{\beta\gamma}(k_-;t_{\beta},t_{\gamma}) - \frac{D(t_{\gamma})}{D(t_{\alpha})} P_{\beta\gamma}(k_+;t_{\beta},t_{\gamma}) \right] + O\left(\left(\frac{q}{k}\right)^0\right)$$
Peloso M. and Pietroni M., JCAP (2013)  
Kehagias A. and Riotto A., Nucl. Phys. (2013)

In configure in presence of a scale like the BAO, the oscillating part of the Equal-time derivative is enhanced by a  $\sim k r_{BAO}$  factor, we can isolate it to verify CR and measure bias  $\lim \frac{B_{aaa}(q,h_{\mp},h_{\mp})}{2} = \frac{1}{2} \frac{1}{2}$ 

 $q/k \rightarrow 0$   $P_{\alpha\alpha}(q)P_{\alpha\alpha}(k)$ Unchanged by

nonlinearities!

Baldauf T. et al., Phys.Rev.D (2015)

 $d \log k$ 

CR and BAO

N-body simulations: real space w/ biased tracers, MM+ (2019)



#### CR and BAO

Bootstrap

Equivalence principle

SPT

N-body simulations: redshift space w/ biased tracers, MM+ (2020)

### Multipoles + Kaiser

$$\lim_{qR \to 0} \frac{B_{l}^{(0)}(q,k)}{P_{l}^{(0)}(q)P_{l}^{(0)}(k)} = - \begin{bmatrix} \frac{1}{3b_{l}} + \frac{b_{l}-1}{9b_{l}} \beta_{l} \frac{1 + \frac{3}{3}\beta_{l}}{1 + \frac{2}{3}\beta_{l} + \frac{1}{5}\beta_{l}^{2}} \end{bmatrix} \frac{\partial \log P_{l}^{(0)}(k)}{\partial \log k} + \dots \qquad 12$$

$$\lim_{qR \to 0} \frac{B_{l}^{(l_{2}-2)}(q,k)}{P_{l}^{(0)}(q)P_{l}^{(0)}(k)} = -\frac{2\beta_{l}}{45b_{l}} \frac{2 + b_{l}(5 + 3\beta_{l})}{1 + \frac{2}{3}\beta_{l} + \frac{1}{5}\beta_{l}^{2}} \frac{\partial \log P_{l}^{(0)}(k)}{\partial \log k} + \dots \qquad 0.8$$

$$\lim_{k \to 0} \frac{P^{(2)}(k)}{P_{l}^{(0)}(k)} = \frac{4\beta_{l}}{21} \frac{7 + 3\beta_{l}}{1 + \frac{2}{3}\beta_{l} + \frac{1}{5}\beta_{l}^{2}} \qquad 0.4$$

$$\text{Angles} \qquad 0.2$$

$$\text{Growth rate} \quad f \equiv \frac{d \log D}{d \log a} \qquad \mu \equiv \hat{\mathbf{k}} \cdot \hat{\mathbf{q}} \qquad 0.4$$

$$\beta_{l} \equiv \frac{f}{b_{l}} \qquad \mu_{q} \equiv \hat{\mathbf{q}} \cdot \hat{\mathbf{z}} \qquad 0.4$$

CR and BAO

## N-body simulations: redshift space w/ biased tracers, MM+ (2020)



$M_{ m min} = 10^{13} h^{-1} M_{\odot} \qquad z = 0$					
$q_{max} (h/Mpc)$	$b_h$	$b_h^{fid}$	$f = \beta_h b_h$	$f^{fid}$	
0.020	$1.58_{-0.13}^{+0.15}$	1.47	$0.57\substack{+0.06\\-0.06}$	0.528	
0.030	$1.53_{-0.08}^{+0.08}$	1.47	$0.55\substack{+0.04 \\ -0.04}$	0.538	
$M_{\rm min} = 10^{13} h^{-1} M_{\odot} \qquad z = 1$					
$q_{max} (h/Mpc)$	$b_h$	$b_h^{fid}$	$f = \beta_h b_h$	$f^{fid}$	
0.020	$2.85_{-0.32}^{+0.39}$	2.686	$0.93^{+0.14}_{-0.14}$	0.877	
0.030	$2.58^{+0.17}_{-0.16}$	2.686	$0.82^{+0.07}_{-0.07}$	0.877	
$M_{\rm min} = 10^{14} h^{-1} M_{\odot} \qquad z = 0$					
$q_{max}$ (h/Mpc)	$b_h$	$b_h^{fid}$	$f = \beta_h b_h$	$f^{fid}$	
0.020	$3.40^{+0.83}_{-0.61}$	2.446	$0.73_{-0.18}^{+0.18}$	0.528	
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Model independent measurement of f(z) and  $b_t$  at 10%!

Bootstrap

SPT

Equivalence principle

## N-body simulations: redshift space w/ biased tracers, MM+ (2020)



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C	D	Т
$\circ$	T	T

#### Equivalence principle

#### CR and BAO

## N-body simulations: redshift space w/ biased tracers, MM+ (2020)



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Model independent measurement of f(z) and  $b_t$  at 10%!

q<sub>max</sub>=0.020 h/Mpc

Halo

z=0

1.8

 $M_{min} = 10^{13} M_{\odot} h^{-1}$ 

2,0

2,2

q<sub>max</sub>=0.030 h/Mpc

#### SPT

#### Equivalence principle

CR and BAO