

# Are Early Extremely Supermassive Black Holes Consistent with $\Lambda$ CDM?



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2022-09-20

# “community” view

SMBH challenging for  $\Lambda$ CDM

hard to grow BHs to  $10^9+ M_{\odot}$  by  $z > 6$

- ... this (QSO) is only the latest and most extreme of a growing number of known giant BHs at early times whose rapid growth, within the (somewhat squishy) constraint of the Eddington limit, is difficult to understand. ...
- The point worth making is this: Such objects are so rare that any attempt to find a “natural” explanation is probably wrong. **If the suggested process that makes these objects is not extremely unusual, it is probably the wrong process.** Kormendy 2013
- rare-processes are a good way to probe models at their limits!

# Modifications to a “base model” $\Lambda$ CDM which might more easily produce early SMBHs?

- mechanism: early universe production of compact (black hole?) “seeds”
  - enhanced (super horizon) inhomogeneities on small scales
  - non-gaussianity enhancing rare small scale inhomogeneities
- mechanism: enhance growth rate of BHs with “new” physics
  - typically involves dark matter accretion enhanced by self-interactions

## No modifications needed?

- baryonic “direct formation” of say,  $10^4 M_{\odot}$ , BH seeds
- growth rate of BHs is short enough to produce observed BHs



# Black Hole Growth by Accretion

- Eddington Limit
  - for accretion radiation Thomson scattering repulsion < gravitational attraction

$$\frac{L_{\text{bol}}}{M} < \frac{L_{\text{Edd}}}{M} = \frac{4\pi G m_p}{\sigma_T \left(1 - \frac{1}{2} Y_{\text{He}}\right)} = 3.72 \times 10^4 \frac{L_{\odot}}{M_{\odot}}$$

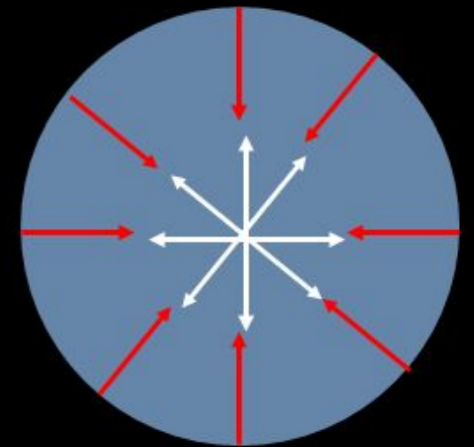
- PARAMETERS

- Eddington ratio:  $\lambda_{\text{Edd}} = \frac{L_{\text{bol}}}{L_{\text{Edd}}}$

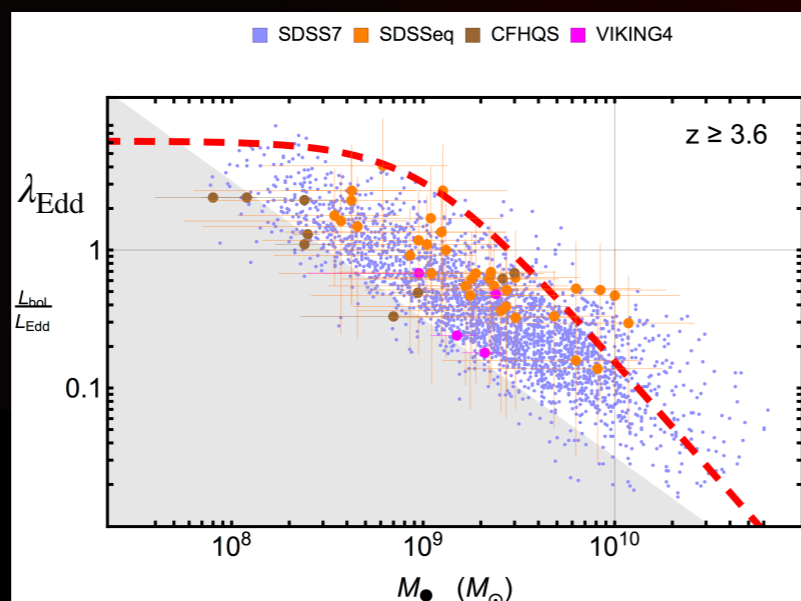
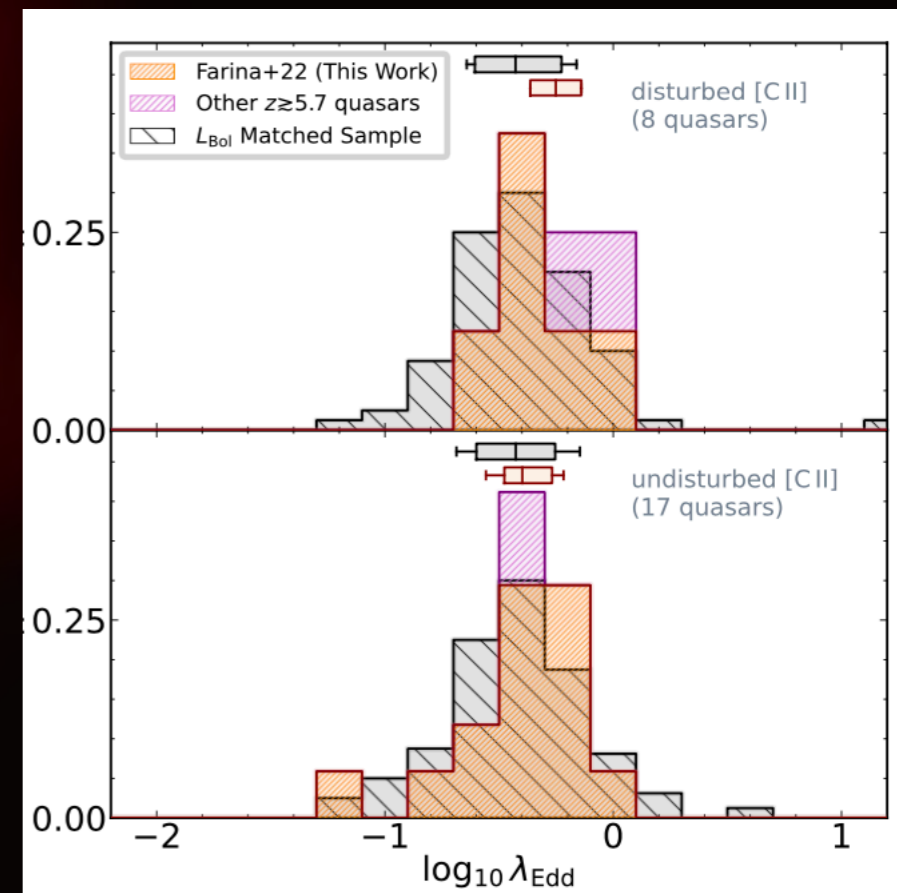
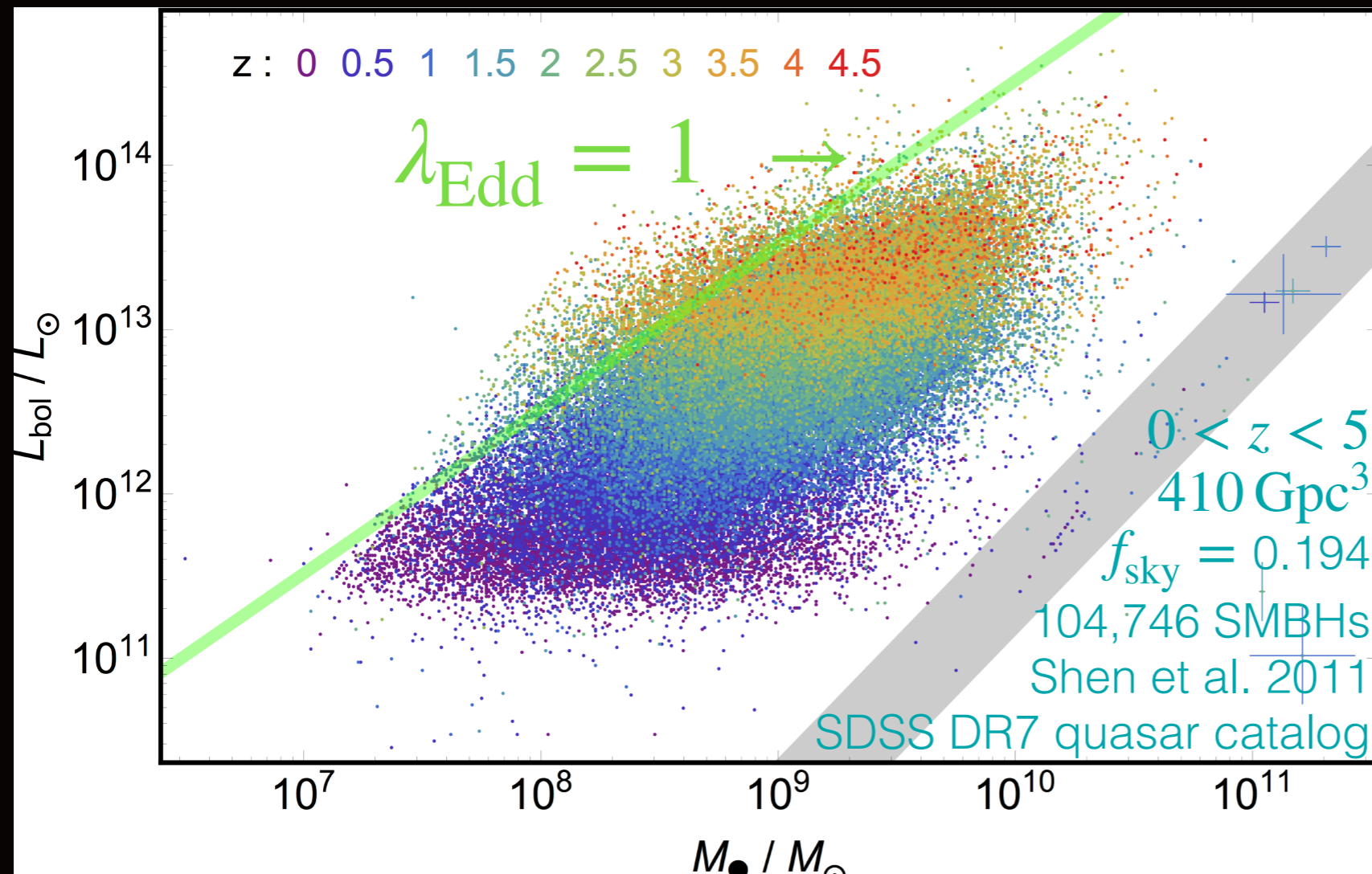
- radiative efficiency:  $\eta_{\text{rad}} = \frac{L_{\text{bol}}}{\dot{M}_{\text{gas}} c^2}$   $\epsilon_{\text{rad}} = \frac{L_{\text{bol}}}{\dot{M}_{\text{BH}} c^2} = \frac{\eta_{\text{rad}}}{1 - \eta_{\text{rad}}}$

- Salpeter growth timescale:  $\tau_{\text{Salp}} = \frac{M_{\text{BH}}}{\dot{M}_{\text{BH}}} = \frac{39.6 \text{ Myr}}{\lambda_{\text{Edd}}} \frac{\epsilon_{\text{rad}}}{0.1}$

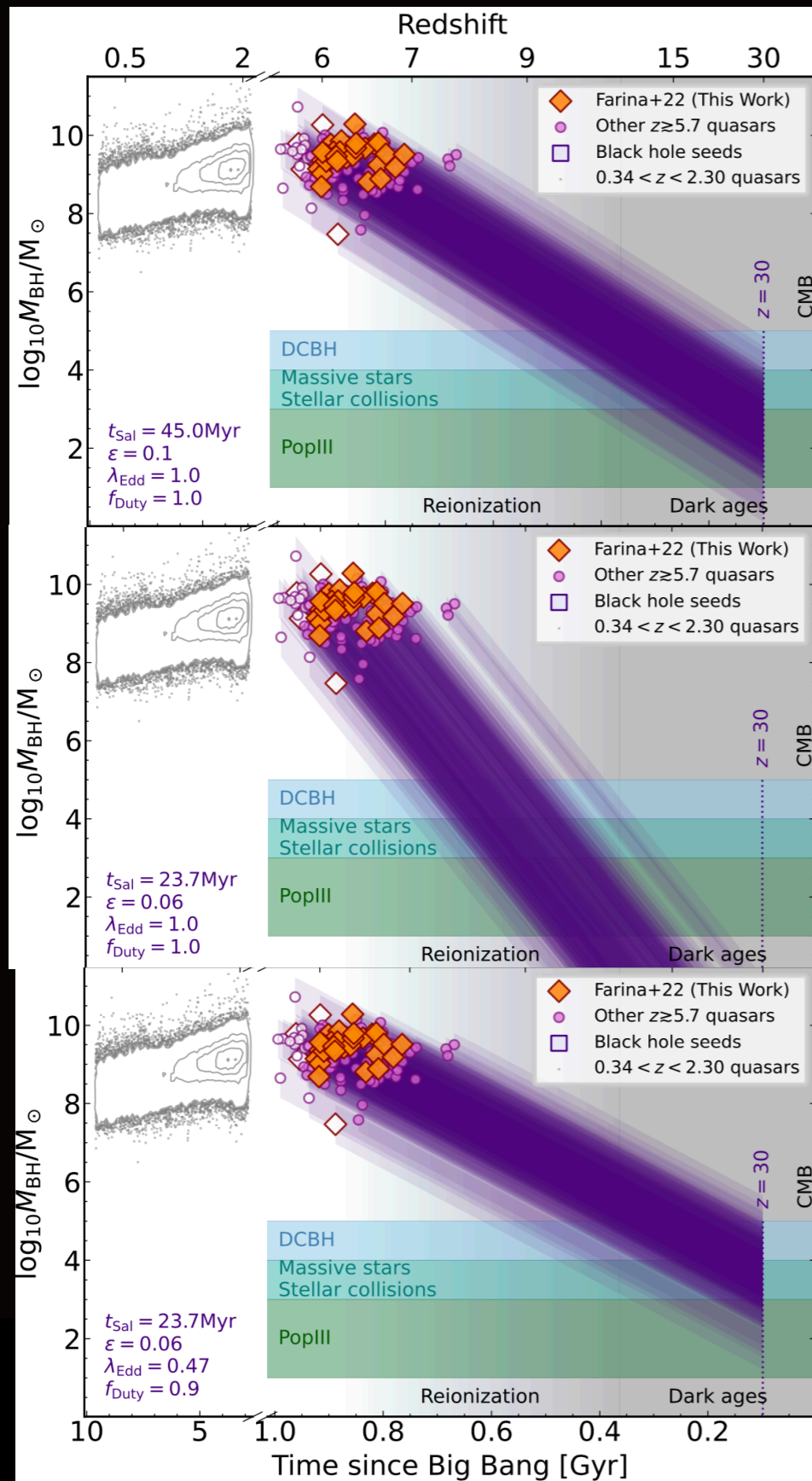
- N.B. one can measure  $L_{\text{bol}}$ ,  $M_{\text{BH}}$ ,  $\lambda_{\text{Edd}}$  directly in quasars - but not  $\epsilon_{\text{rad}}$



# $\lambda_{\text{Edd}} < 1$ is typical



# Are Early SMBHs really problematic?



← canonical parameter

← decrease efficiency by 40%

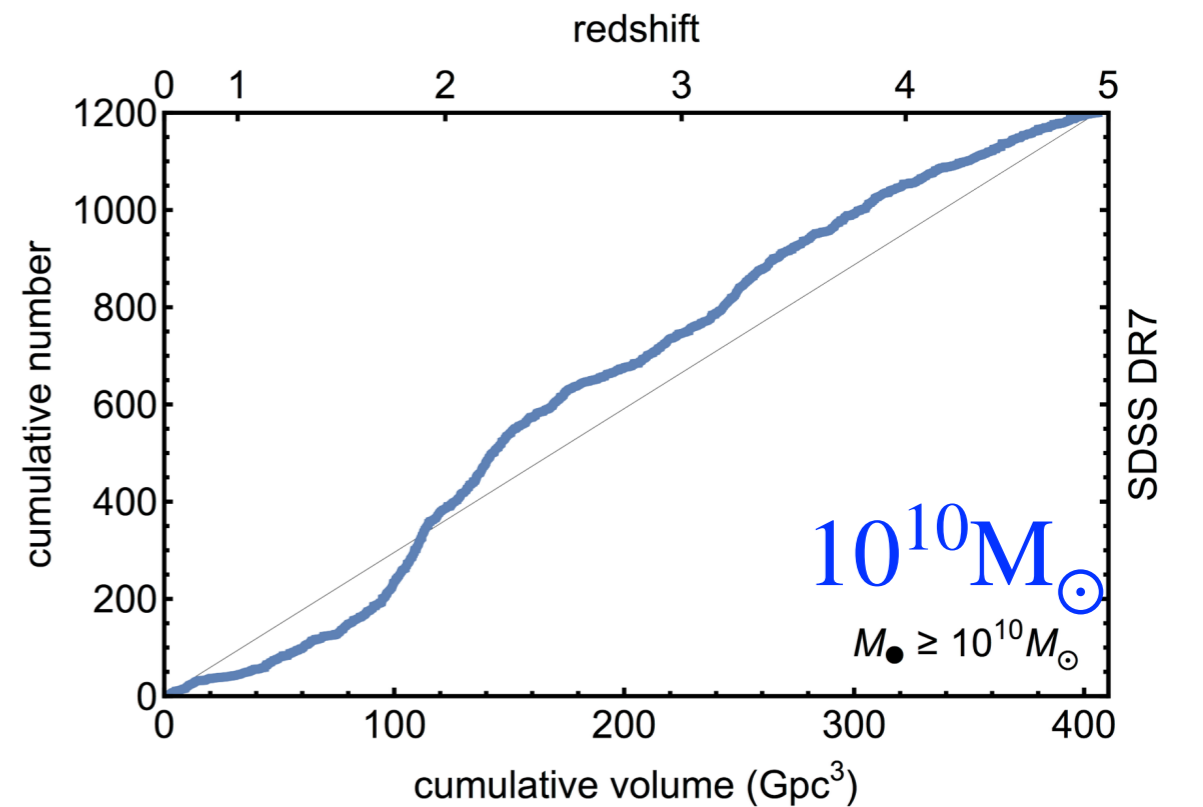
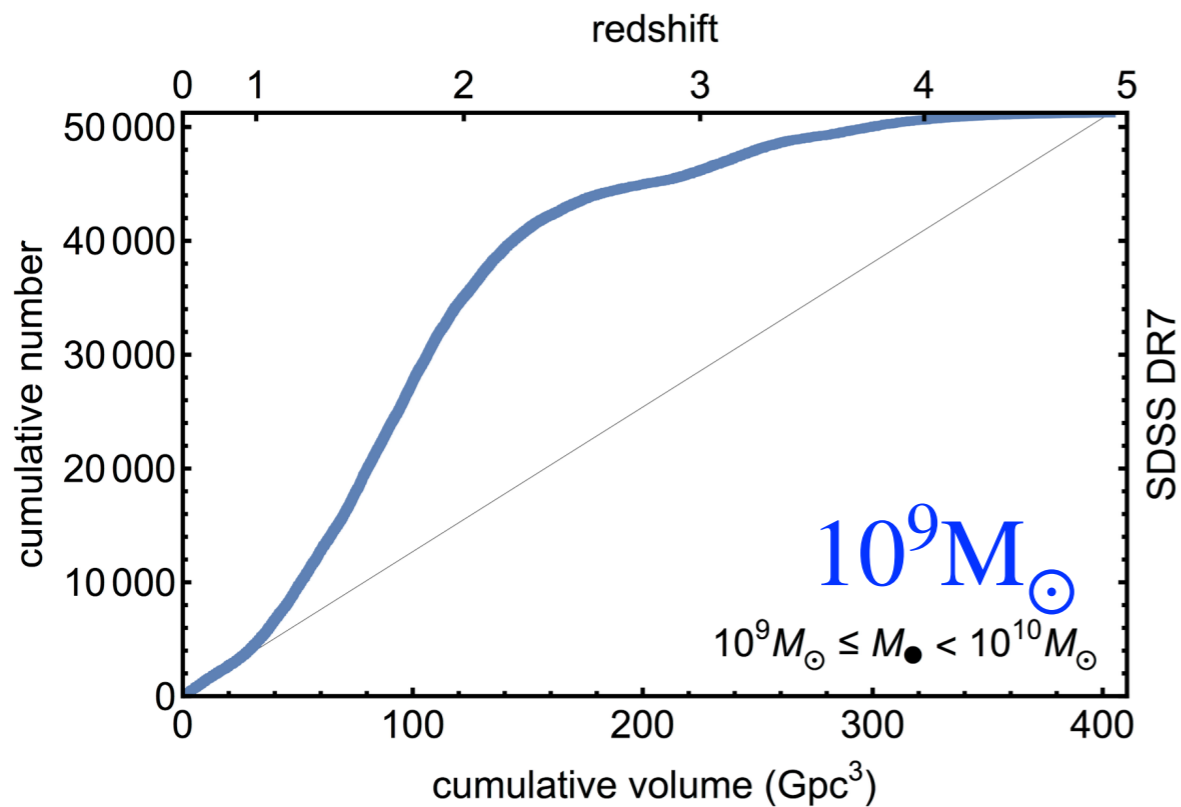
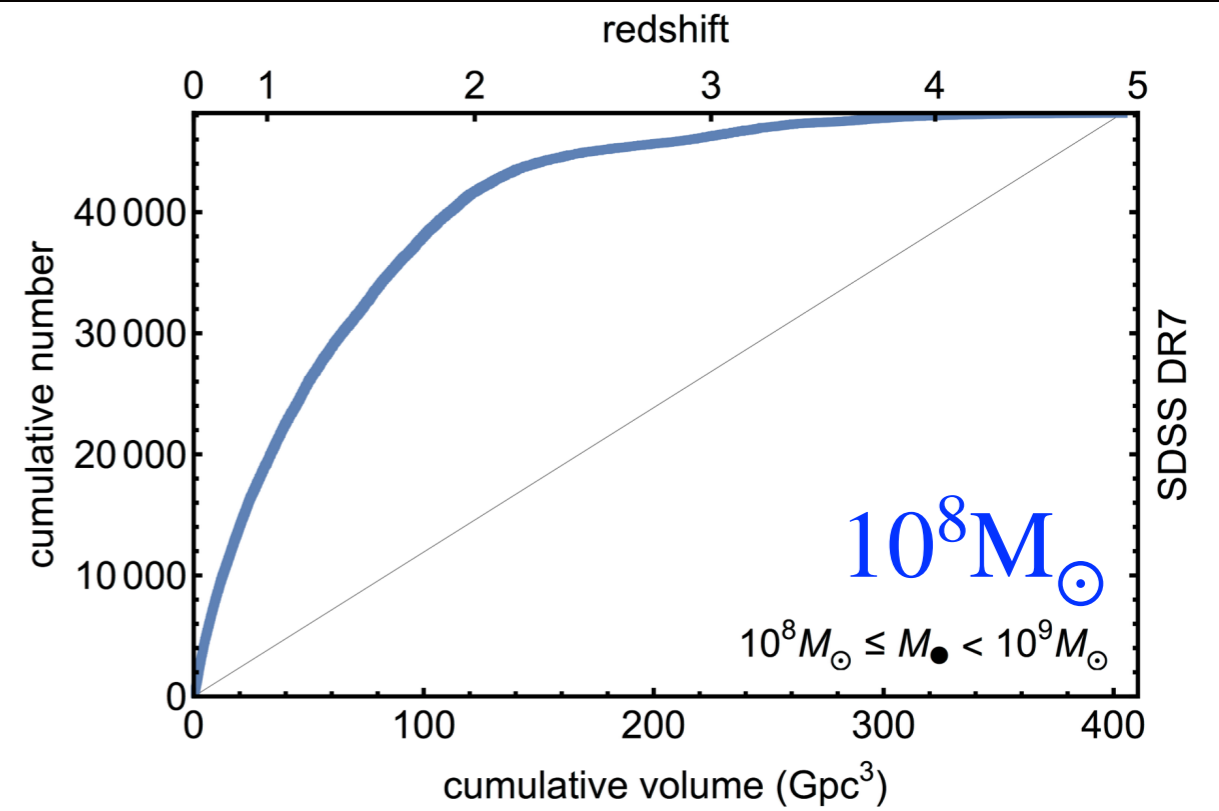
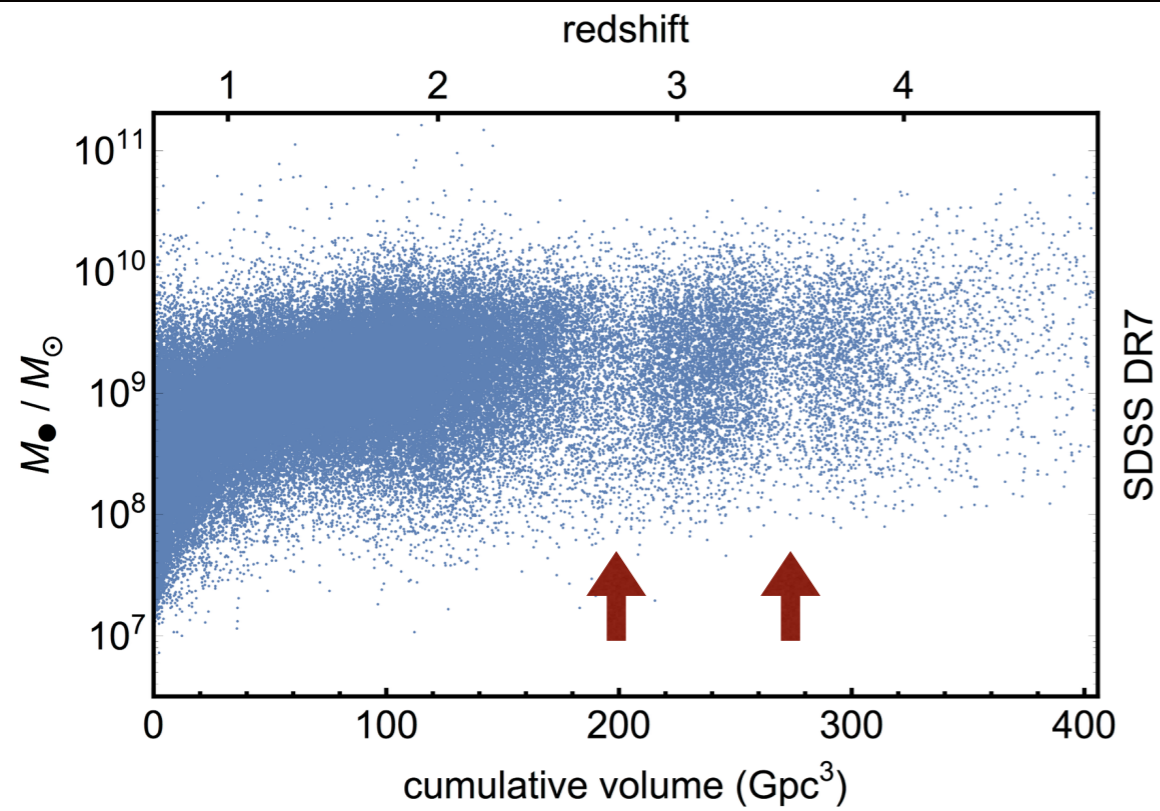
← more “realistic”  $\lambda_{\text{Edd}}$ , duty cycle

Individual QSO growth argument inconclusive as to whether this is a problem for the base  $\Lambda$ CDM model

due to uncertainty as to reasonable values for  $\tau_{\text{Salp}}$

can learn more about  $\tau_{\text{Salp}}$  from QSO population dynamics

# Slow evolution of most massive luminous BHs



SDSS DR7



# BH Population Dynamics

- Quasar distribution in luminosity redshift give clues for BH timescales

- consider (toy) non-stochastic accretion only evolution:  $\dot{M}[M]$

- $\Psi[M, t] \equiv \frac{dn_{\text{co}}}{d \ln M}$  cosmic time  $\frac{dt}{dz} = \frac{1}{(1+z)H[z]}$

- $\frac{\partial \Psi[M, t]}{\partial t} = -M \frac{\partial}{\partial M} \left( \frac{\dot{M}[M]}{M} \Psi[M, t] \right) = -M \frac{\partial}{\partial M} \left( \frac{\Psi[M, t]}{\tau[M]} \right)$

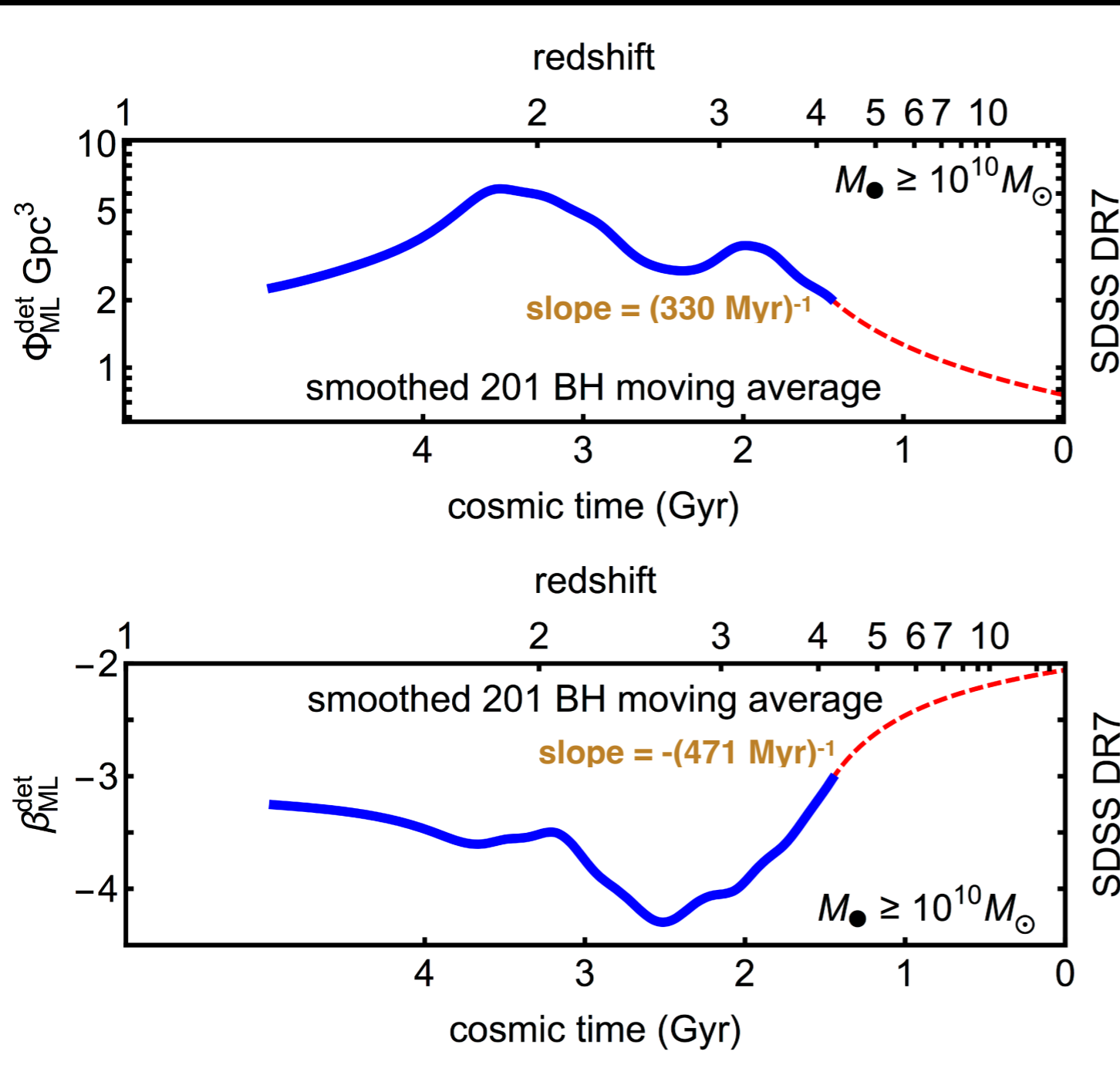
- closed form solution w/ initial condition  $\Psi[M, t_0] = \Psi_0[M]$

- $\Psi[M, t] = \frac{\tau[M] \Psi_0[\tilde{M}[M_1, \tilde{T}[M_1, M] - (t - t_0)]]}{\tau[\tilde{M}[M_1, \tilde{T}[M_1, M] - (t - t_0)]}$ ,  $M_1$  arbitrary

- $\tilde{T}[M_1, M_2] \equiv \int_{M_1}^{M_2} \frac{dM}{M} \tau[M]$  with inverse  $\tilde{M}[M_1, \tilde{T}[M_1, M_2]] = M_2$

\*for compactness subscripts are dropped here

# $t = 0$ extrapolation of $M_{\text{BH}} > 10^{10} M_{\odot}$ population



$$\beta \equiv \frac{\partial \ln \Psi}{\partial \ln M_{\text{BH}}} \rightarrow -2.05$$

$$\tau_{\text{Salp}} \rightarrow 805 \text{ Myr} \left( \frac{M_{\text{BH}}}{10^{10} M_{\odot}} \right)^{1.6}$$

$$\int_{10^{10} M_{\odot}}^{\infty} \frac{dM}{M} \Psi \rightarrow \frac{0.76}{\text{Gpc}^3} \left( \frac{M_{\text{BH}}}{10^{10} M_{\odot}} \right)^{-1.05}$$

- extrapolates to extremely massive extremely rare SMBHs in early universe!
- extending to smaller masses
- “point defect” type of non-Gaussianity

# QSO flickering

- SMBHs not always bright enough to be seen as QSOs or AGNs
  - sometimes  $\lambda_{\text{Edd}} \ll 1$  (for most SMBHs in late universe  $z \lesssim 1$ )
  - model: flicker on (duty cycle  $f_{\text{on}}$ ) and off (duty cycle  $1 - f_{\text{on}}$ )
  - caused by interruptions in gas supply for accretion
- large fraction of unseen SMBHs modifies population dynamics
  - $\tau_{\text{Salp}}^{\text{effective}} \rightarrow \frac{\tau_{\text{Salp}}}{f_{\text{duty}}}$   $f_{\text{duty}}$  time varying in uncontrolled way
  - contributes to slow evolution
  - hard to quantify



# QSO surveys now extend to higher $z$

- **Shen et al. 2020 (SHFAARH20)**
- 10 or 11 parameter fit of QSO luminosity function  $0 < z < 7$

$$\Phi[L, z] = \frac{2 \Phi_*[z]}{\left(\frac{L}{L_*[z]}\right)^{\gamma_1[z]} + \left(\frac{L}{L_*[z]}\right)^{\gamma_2[z]}}$$

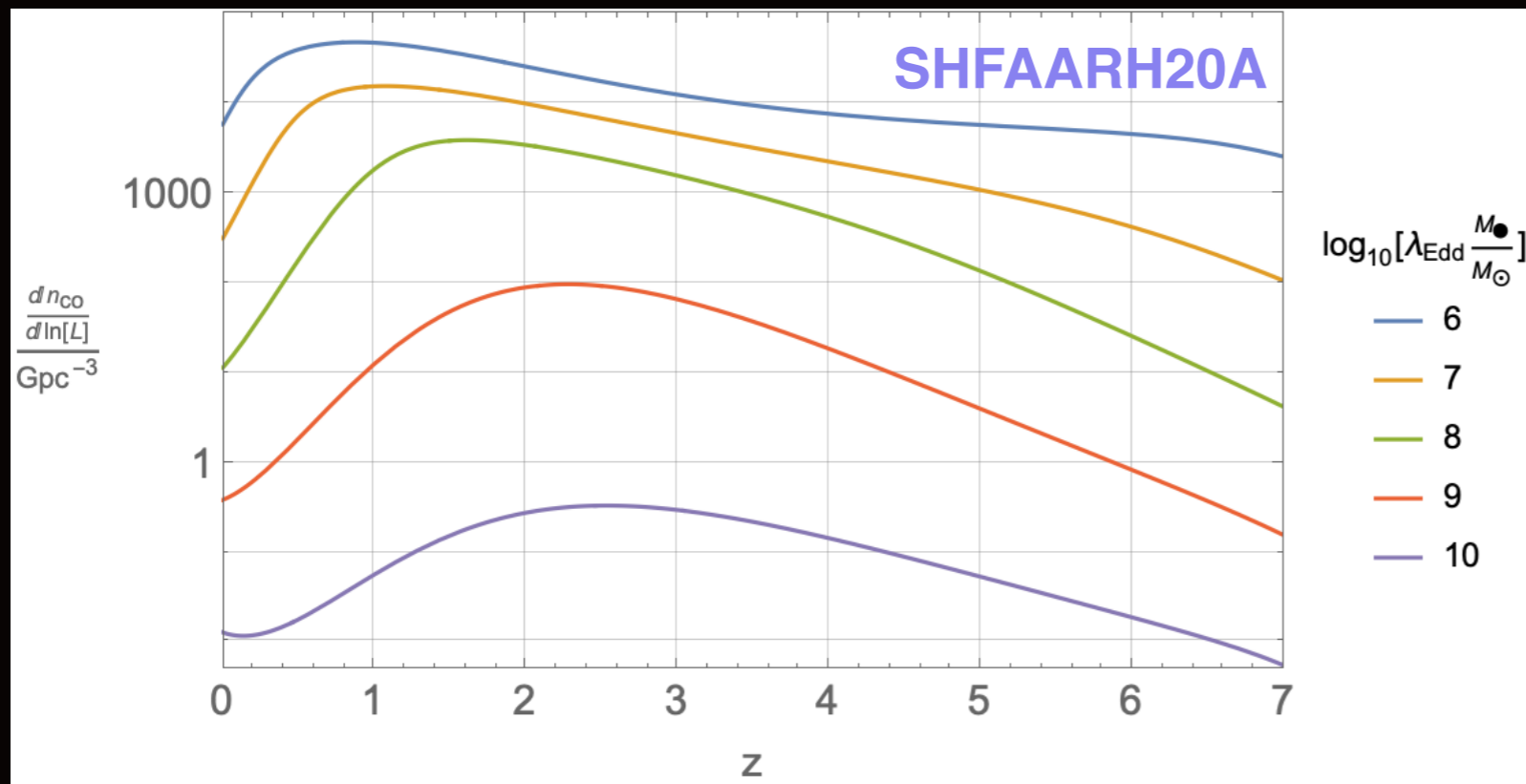
- model A:  $\gamma_1[z] = \tilde{a}_0 + \tilde{a}_1 z + \tilde{a}_2 z^2$  / model B:  $\gamma_1[z] = a_0 \left(\frac{1+z}{3}\right)^{a_1}$

$$\gamma_2[z] = \frac{2 b_0}{\left(\frac{1+z}{3}\right)^{b_1} + \left(\frac{1+z}{3}\right)^{b_2}}$$

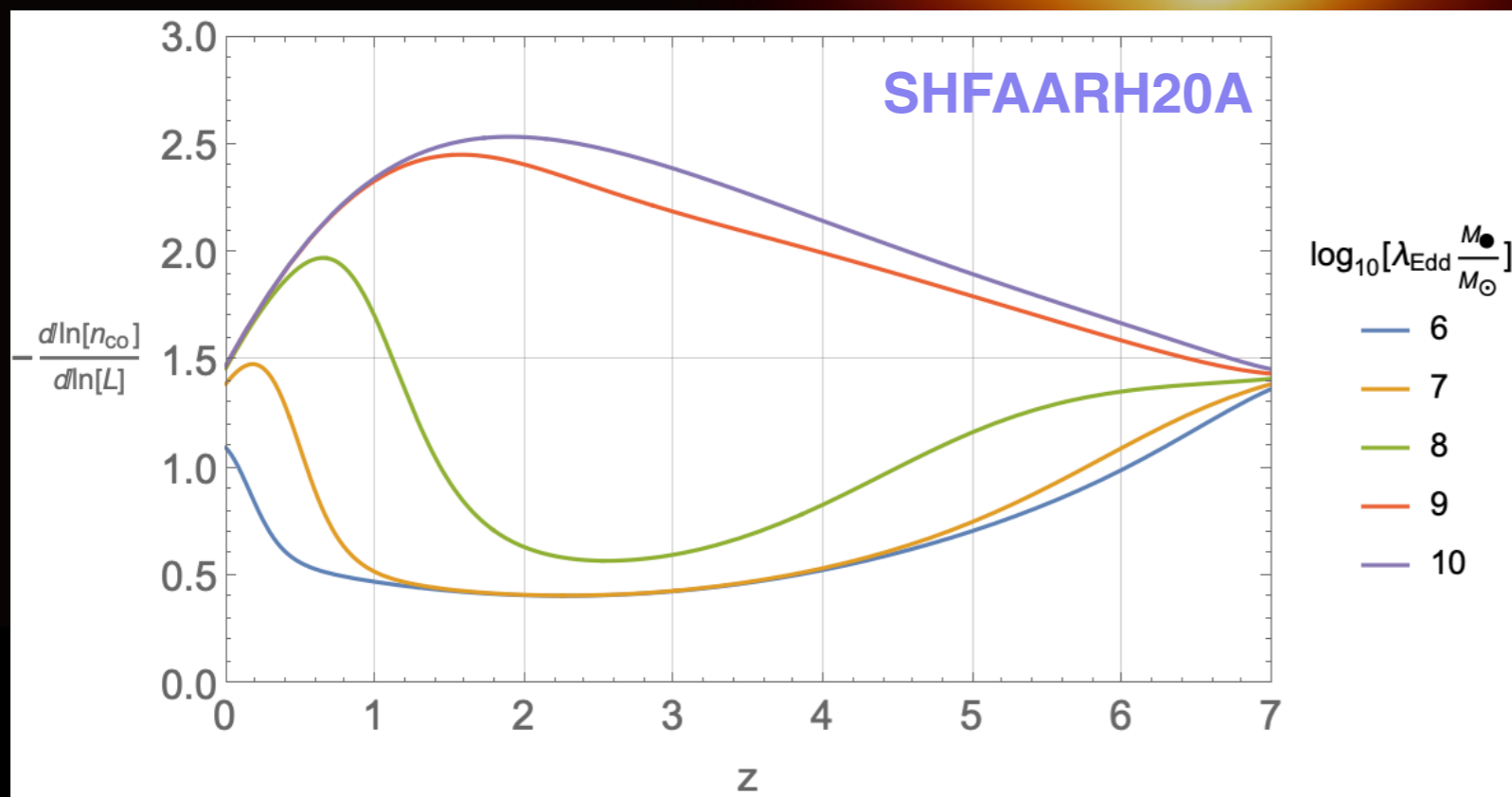
$$\log_{10} \left[ \frac{L_*[z]}{L_\odot} \right] = \frac{2 c_0}{\left(\frac{1+z}{3}\right)^{c_1} + \left(\frac{1+z}{3}\right)^{c_2}}$$

- $\log_{10} [\text{Mpc}^3 \Phi_*[z]] = \tilde{d}_0 + \tilde{d}_1 z$
- hi- $z$  functional form “biased” by lo- $z$  observations - much extrapolation

# QSO evolution speeds up at hi- $z$



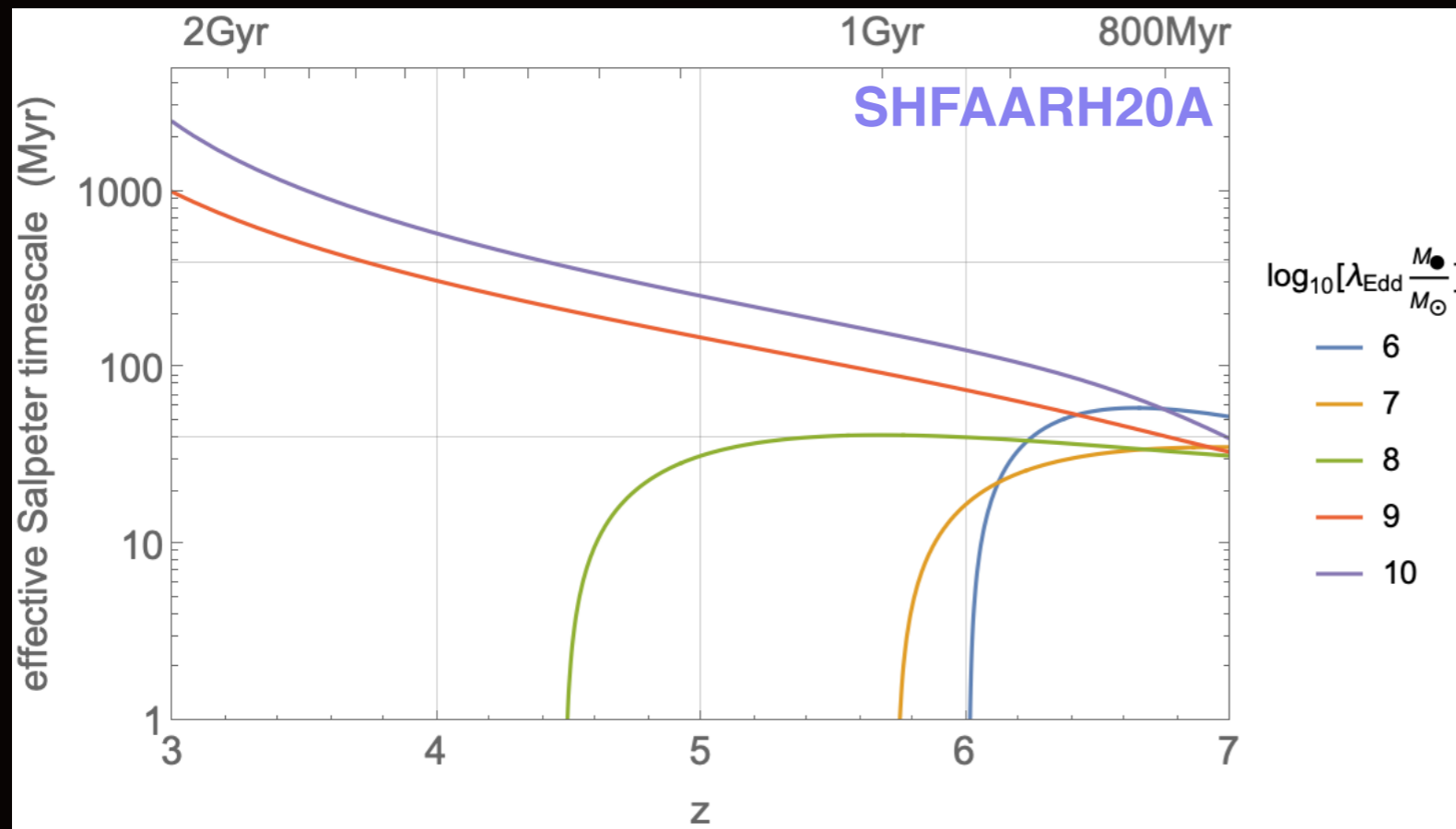
from  $z = 5$  to 7 number of  $M_{\text{BH}} \sim 10^{10} M_{\odot}$  falls by order of magnitude



luminosity function slope flattens to  $\beta \approx -1.5$  for over entire mass range

not for small mass in model B

# evolution timescale attains canonical value



$$\tau_{\text{Salp}}^{\text{effective}} \approx 40 - 50 \text{ Myr}$$

(canonical value!) at  $z = 7$  in this model fit for all  $M_{\text{BH}}$

gas supply secure for  $z > 7$  - no flickering (?)

This evolution timescale is not fast enough to easily explain the most massive early SMBHs from population III stars!

Should we (PNG) declare victory and go home?

**No!**

- evolution timescale may continue to decrease
- SHFAARH20 fit suspect - data still too sparse

# my take

- there is still room for very large black hole seeds from the early universe
- however, whether the existence of early high mass SMBHs is problematic in any sense has not been empirically established
- quasar observations are only now getting into a redshift regime where we might or might not see fast exponential growth of SMBH populations
- while  $\epsilon_{\text{rad}} \approx 0.1$  is well motivated by accretion disk theory how certain are we that in very rare cases that it might, sustainably, be 30-40% smaller
  - in rare regions of very low angular momentum gas?
  - can PNG make this more common?