

Asli Abdullahi^{*} IFT NuTs workshop, 7 June 2022



*in collaboration with

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Muon g-2 experiment seems to confirm existing tension with SM



https://mod.fnal.gov/mod/stillphotos/2017/0100/17-0188-19

B. Abi, et al. Phys. Rev. Lett. 126, 141801 (2021)



Combined BNL and FNAL result at 4.2σ deviation from the SM theory prediction



Possible that the discrepancy is purely SM but exciting to consider otherwise!

Among the possibilities is the existence of a dark photon (DP)

$$\mathscr{L}_{\rm SM} - \frac{\varepsilon}{2c_{\rm W}} F_{\mu\nu} X^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + \frac{m_X^2}{2} X_\mu X^\mu$$

DP provides a positive contribution to Δa_{μ}



 $\Delta a_{\mu} \propto \int_{0}^{1} dz \frac{2zm_{\mu}^{2}(1-z)^{2}}{m_{\mu}^{2}(1-z)^{2} + zm_{\tau \tau}^{2}}$

M. Pospelov, Phys. Rev. D80, 095002

New interactions from gauging accidental U(1) symmetries of SM, e.g.



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Or comprise a new and secluded dark sector, e.g. a U(1)_X



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Fabbrichesi, et al. 10.1007/978-3-030-62519-1

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D. Banerjee et al, PRL 123, 121801 (2019)

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G.Mohlabeng, 10.1103/PhysRevD.99.115001

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Evades visible DP searches with missing energy

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- Evades visible DP searches with **missing energy**
 - Evades invisible DP searches with visible final states

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1) Does the kinetically-mixed semi-visible DP remain a viable solution to g-2?

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- 1) Does the kinetically-mixed semi-visible DP remain a viable solution to g-2?
 - 2) Which models does it point us toward?

Rest of this talk

• Semi-visible DP models

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- Inelastic DM (iDM)
- Dirac iDM (i2DM)
- Heavy Neutral Leptons

• Semi-visible DP models

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- Heavy Neutral Leptons
- Constraints and recasts

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 - Inelastic DM (iDM)
 - Dirac iDM (i2DM)
 - Heavy Neutral Leptons
- Constraints and recasts
 - BaBar and NA64
- Future searches + Conclusions

Vector Portal Inelastic Dark Matter

• Introduced as a solution to DAMA anomaly D.Smith, N. Wiener, Phys. Rev. D 64, 043502

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G. Giudice et al. Phys. Lett. B780, 543–552
A. Berlin, F. Kling, Phys. Rev. D99, 015021
G. Mohlabeng, Phys. Rev. D99, 115001
J. Eby et al. J. High Energ. Phys. 2019, 115
Y.D. Tsai et. al Phys. Rev. Lett. 126 (2021) 181801
M. Duerr et al. JHEP 02 (2020) 039

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... and many many more

In addition to DP Lagrangian

$$\mathscr{L} = \mathscr{L}_{\rm SM} - \frac{\varepsilon}{2c_{\rm W}} F_{\mu\nu} X^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_X X_\mu \mathcal{J}_X^\mu + \frac{m_X^2}{2} X_\mu X^\mu$$

introduce a pair of Weyl fermions,



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After SSB of U(1)', generate Majorana masses

$$\mathscr{L}_{\mathrm{mass}} = -rac{1}{2} egin{pmatrix} \mu_L & \overline{\chi_R^c} \ m_D & \mu_R \end{pmatrix} egin{pmatrix} \mu_L & m_D \ m_D & \mu_R \end{pmatrix} egin{pmatrix} \chi_L^c \ \chi_R \end{pmatrix} + \mathrm{h.c.}$$

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In the mass basis,

$$\mathscr{L}_{\mathrm{mass}} = rac{m_1}{2} \overline{\psi_1} \psi_1 + rac{m_2}{2} \overline{\psi_2} \psi_2 + \mathrm{h.c.}$$

$$m_{1,2} = rac{\mu_L + \mu_R}{2} \pm \sqrt{M_D^2 + \left(rac{\mu_R - \mu_L}{2}
ight)^2}$$
 and $\Delta = rac{m_2 - m_1}{m_1}$

Dark sector current reads

$$\begin{aligned} \mathcal{J}_X^{\mu} &= \frac{Q_L c_{\theta}^2 - Q_R s_{\theta}^2}{2} \overline{\psi_2} \gamma^{\mu} \gamma^5 \psi_2 + \frac{Q_L s_{\theta}^2 - Q_R c_{\theta}^2}{2} \overline{\psi_1} \gamma^{\mu} \gamma^5 \psi_1 \\ &+ i \frac{Q_L + Q_R}{2} \sin 2\theta \cos \delta \, \overline{\psi_2} \gamma^{\mu} \psi_1 + \frac{Q_L + Q_R}{2} \sin 2\theta \sin \delta \, \overline{\psi_2} \gamma^{\mu} \gamma^5 \psi_1, \text{ with } \tan 2\theta = m_D / (\mu_L - \mu_R). \end{aligned}$$

In the vector-like limit, $Q_L = Q_R$, and for pseudo-Dirac states

 $Z' \sim \sqrt{\frac{\psi_1}{2}}$ Dark photon couples **purely off-diagonally**, with diagonal couplings suppressed by a factor $\mu_L - \mu_R$



Z' with diagonal couplings suppressed by a factor $\mu_L - \mu_R$

- Light state constitute the DM \rightarrow relic abundance sourced by co-annihilations
- Heavy state decays to DM and light SM fermions •



Constraints on iDM

In the g-2 preferred region, strongest constraints come from NA64 and BaBar searches for invisible dark photon decays



D. Banerjee et al, PRL 123, 121801 (2019)

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BaBar search for initial state radiation (ISR) in the process $e^+e^- \rightarrow \gamma Z'$, accompanied by missing energy



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DP assumed to have 100% invisible branching ratio

BaBar search for initial state radiation (ISR) in the process $e^+e^- \rightarrow \gamma Z''$, accompanied by missing energy



DP assumed to have 100% invisible branching ratio

In the case of a semi-visible DP (e.g. in iDM), photon accompanied by **displaced**, **visible lepton tracks**




BaBar monophoton search

In the case of a semi-visible DP (e.g. in iDM), photon accompanied by **displaced**, **visible lepton tracks**



If leptons sufficiently energetic, events are vetoed leading to significant relaxation of the bound

Dedicated monophoton trigger lines:



(IFR) or Muon system

Low-mass DP (high-energy monophoton):

- One cluster in ECAL with $E \ge 2 \text{ GeV}$
- NO e+e- tracks in DCH

High-mass DP (low-energy monophoton):

- One cluster in ECAL with $E \ge 1$ GeV
- NO e+e- tracks in DCH

Photons accepted in polar angle range, $|\cos \theta_{\gamma}^*| < 0.6$

BaBar monophoton search

<u>Veto criteria:</u> For decays that occur in instrumented regions of the detector (SVT, DCH, ECAL, IFR)



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Require

1. Energy of lepton tracks must exceed BaBar detection threshold for charged lepton tracks

If $\theta_{sep} > 10^{\circ}$, then cut: $E_{\pm} > 100 \text{ MeV}$

If $\theta_{sep} < 10^{\circ}$, then cut: $E_+ + E_- > 100 \text{ MeV}$

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If $\theta_{sep} < 10^{\circ}$, then cut: $E_+ + E_- > 100 \text{ MeV}$

2. The polar angles are sufficiently wide that the electrons do not escape along the beam pipeline cut: $17^{\circ} < \theta_{pol.} < 142^{\circ}$

NA64 performed a missing energy search with DP production in dark bremsstrahlung, $e^-N \rightarrow e^-NZ'$



...performing two distinct analyses for the DP model

1) Invisible DP





D. Banerjee et al, PRL 123, 121801 (2019)

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1) invisible search

→ Apply *trigger* conditions on energy deposited in pre-shower ECAL: $E_{tot} \in [0.2, 80]$ GeV



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Selection criteria

 ECAL: Total energy deposited < 50 GeV, i.e. total missing energy > 50 GeV

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→ Apply *trigger* conditions on energy deposited in pre-shower ECAL: $E_{tot} \in [0.2, 80]$ GeV

Selection criteria

- ECAL: Total energy deposited < 50 GeV, i.e. total missing energy > 50 GeV
- 2. Veto counter: < 10 MeV deposited in veto
- 3. HCAL: < 1 GeV deposited in HCALs

2) semi-visible DP in iDM



Cazzaniga, C. et al, Eur. Phys. J. C 81 (2021) 10, 959

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2) semi-visible DP in iDM



Cazzaniga, C. et al, Eur. Phys. J. C 81 (2021) 10, 959

Constraint from semi-visible search **not** sensitive to prompt decays

• Efficiency to distinguish collimated e+eshowers from primary electron beam in ECAL is challenging 2) semi-visible DP in iDM



Selection criteria

- 1. HCAL: < 1 GeV deposited in HCALs
- 2. All decays beyond HCALs selected (long-lived fermions)

Note: In the semi-visible search, all decays that occur prior to HCAL are vetoed regardless of energy

"The largest limitation to probe the missing region comes from the increasingly short decay time of χ_2 , which makes the chance of detection vanishingly small." - Eur. Phys. J. C 81 (2021) 10, 959

Treat invisible search as "sensitivity" rather than a constraint

iDM Recast

iDM RECAST

We consider a benchmark model of iDM inspired by recent work on this topic in relation to g-2



iDM: $\Delta_{21} = 0.4$, r = 0.33, $\alpha_D = 0.1$

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We consider a benchmark model of iDM inspired by recent work on this topic in relation to g-2

10-2 iDM: $\Delta_{21} = 0.4$, r = 0.33, $\alpha_D = 0.1$ Ψ 10⁻¹ **EWPO NA64** (semi-visible) 10^{-2} BaBar M 10⁻⁵ 10^{-1} mA [GeV 10-3 model indep: ω 10^{-2} NA64 (invisible) NuCal relic density E137 (5:37(81) iDM model for (g-2) is disfavoured at 20 5×10^{-1} 10^{-4} 10^{-1} 10^{0} 10^{1} 10^{-} $m_{Z'}/\text{GeV}$

G. Mohlabeng, Phys. Rev. D99, 115001

10-

 $\Delta=0.4\;m_{\chi1}$, $m_{A'}=3\;m_{\chi1}$, $\alpha_D=0.1$

iDM RECAST

ω

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iDM: $\Delta_{21} = 0.4$, r = 0.33, $\alpha_D = 0.1$ **EWPO NA64** 10^{-2} BaBar M 10 mA [Ge] nodel indep. 10^{-2} Dashed = only energy cut Solid = energy + pipeline cut 10-3 NA64 (invisible) 10^{-1} NuCal Ψ Kinetic mixing, ϵ -01 relic density siscatter E137 5×10^{-1} 10^{-4} 10^{-1} 10^{0} 10^{1} 10^{-10} $m_{Z'}/\text{GeV}$ 10^{-3} 0.05 0.100.150.200.00BaBar constraint on kinetic mixing for Mz' = 1.25 e^{\pm} energy threshold/GeV

GeV as function of e+e- energy threshold cut

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0.30

Majorana DM

HNL BP 1 HNL BP 2

0.25

pseudo-Dirac DM

 $(g - 2)_{\mu}$ central value

G. Mohlabeng, Phys. Rev. D99, 115001

10-

10-2

 $\Delta = 0.4 \ m_{\chi 1}$, $m_{A'} = 3 \ m_{\chi 1}$, $\alpha_D = 0.1$

Vector Portal *Dirac* Inelastic Dark Matter • We would like to consider a model in which we couple diagonally to the heavy fermion state, while suppressing the diagonal coupling to the DM, e.g.



Introduce the four-component fields $\eta = \eta_L + \eta_R$ (sterile), $\chi = \chi_L + \chi_R$ (charged)

in addition to standard DP Lagrangian, we have

$$\begin{split} \mathscr{L} \supset \overline{\chi_L} (i \partial \!\!\!/ - g_X Q_L \not\!\!\!X) \chi_L + \overline{\chi_R} (i \partial \!\!\!/ - g_X Q_R \not\!\!\!X) \chi_R + \overline{\eta} i \partial \!\!\!/ \eta \\ - \left\{ M_1 \overline{\eta} \eta + \frac{1}{2} \mu_L \overline{\eta_L} \eta_L^c + \frac{1}{2} \mu_R \overline{\eta_R} \eta_R^c + M_2 \overline{\chi} \chi + y_L \Phi \overline{\chi_L} \eta_R + y_R \Phi \overline{\chi_R} \eta_L + \text{h.c.} \right\} \end{split}$$

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If we assume Majorana masses of sterile state are small, we have

$$\mathscr{L}_{\mathrm{mass}} = -\frac{1}{2} \begin{pmatrix} \overline{\eta_L} & \overline{\chi_L} \end{pmatrix} \begin{pmatrix} M_1 & M \\ & \\ M & M_2 \end{pmatrix} \begin{pmatrix} \eta_R \\ \chi_R \end{pmatrix} + \mathrm{h.c.}$$

$$\mathscr{L}_{\mathrm{mass}} = -rac{1}{2} egin{pmatrix} M_1 & \overline{\chi_L} \end{pmatrix} egin{pmatrix} M_1 & M \ & M \ & M \end{pmatrix} egin{pmatrix} \eta_R \ & \chi_R \end{pmatrix} + \mathrm{h.c.}$$

If we also also take $M \ll M_{1,2}$, the mass basis is very **nearly aligned** with the flavour basis

Perform a nearly diagonal rotation by $\theta \approx M/(M_2 - M_1)$ to find the mass states

$$\mathcal{J}_X^{\mu} = Q(s_{\theta}^2 \overline{\psi_1} \gamma^{\mu} \psi_1 - s_{\theta} c_{\theta} \overline{\psi_2} \gamma^{\mu} \psi_1 + c_{\theta}^2 \overline{\psi_2} \psi_2)$$

i2DM RECAST

For comparison, keep most parameters the same as iDM benchmark model



Full 3σ g-2 preferred region opens up for DP masses between 600 MeV - 2 GeV

Main features:

- Relaxation occurs at larger DP masses
- NA64 semi-visible constraint stronger

Dirac iDM RECAST

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Important caveat!

• Benchmark model *may not* give the dark matter relic abundance

(see arXiv:2201.08409 for more detail about Dirac iDM story)

Heavy Neutral Leptons

 Similar to i2DM scenario, with the sterile state now coupling to SM neutrinos as in the Type 1 seesaw

$$\mathcal{L} \supset \overline{\chi_L} (i \partial \!\!\!/ - g_X Q X) \chi_L + \overline{\chi_R} (i \partial \!\!\!/ - g_X Q X) \chi_R + \overline{\eta} i \partial \!\!\!/ \eta - \{ y_\nu \overline{L} \tilde{H} \eta^c + y_R \overline{\eta} \chi_R \Phi^* + y_L \overline{\eta} \chi_L^c \Phi + \frac{1}{2} M_N \overline{\eta} \eta^c + M_X \overline{\chi} \chi + \text{ h.c.} \}$$

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Dark sector Yukawa allows for mass mixing between the dark fermions and SM neutrinos

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Dark sector Yukawa allows for mass mixing between the dark fermions and SM neutrinos



Can have light neutrino upscattering to heavy states

Heavy Neutral Leptons

This phenomenological signature has been extensively studied for MiniBooNE



1.00

1.0

0.50 0.75

1.2

- Increased particle content allows for rapid, cascade decays in the detector
- multi-lepton final states can vastly increase the visibility of the dark photon decays

$$\mathcal{J}_X^\mu pprox V_{ij}\overline{\psi_i}\gamma^\mu P_L\psi_j$$
 with $V_{ij} \equiv U_{\chi_L i}^\dagger U_{\chi_L j} - U_{\chi_R i}^\dagger U_{\chi_R j}$



• Increased particle content allows for rapid, cascade decays in the detector

- $N_{3} \qquad N_{2} \qquad N_{1} \qquad \dots \qquad N_{1} \qquad \bigvee \qquad f^{-} \qquad f^{-} \qquad f^{+} \qquad f^{+$
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$$\mathcal{J}_X^{\mu} \approx V_{ij} \overline{\psi_i} \gamma^{\mu} P_L \psi_j$$
 with $V_{ij} \equiv U_{\chi_L i}^{\dagger} U_{\chi_L j} - U_{\chi_R i}^{\dagger} U_{\chi_R j}$

Parameter space relevant for MiniBooNE predicts semi-visible dark photon states!

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$$\mathcal{J}_X^\mu pprox V_{ij} \overline{\psi_i} \gamma^\mu P_L \psi_j$$
 with $V_{ij} \equiv U^\dagger_{\chi_L i} U_{\chi_L j} - U^\dagger_{\chi_R i} U_{\chi_R j}$

Benchmark models:

- 1) theory parameters consistent with a MB explanation
- 2) generic parameters for $g-2 \rightarrow$ designed to maximise visibility of decay signal

1) MiniBooNE benchmark model



Small opening between 1.2 - 2.8 GeV

Main features:

 Double bump feature due to two unstable fermions, BR(Z' → 32) ~ 96%

Important!

- Benchmark consistent with constraints on the light neutrino masses and the MiniBooNE excess → highly constrained parameter space
- Compatibility of g-2 and MiniBooNE requires further study

Heavy Neutral Leptons

2) g-2 benchmark model



Significant relaxation between 150 MeV - 3 GeV

Main features:

- Double bump feature due to two unstable fermions, BR(Z' → 32) ~ 99%
- A prompt decay search from NA64 may exclude part of this region

Testing the allowed parameter space of the models presented will be an important task for currently running and upcoming experiments

- NA64 prompt regime for semi-visible dark photon search
 - **Work-in-progress** understanding NA64 sensitivity to multiple showers in the ECAL
- Monophoton searches at Belle II and BES III will provide stronger constraints on the invisible DP and, consequently, the semi-visible DP
- Displaced vertex search at Belle II
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 - Potential connection to DM

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Thanks for listening!