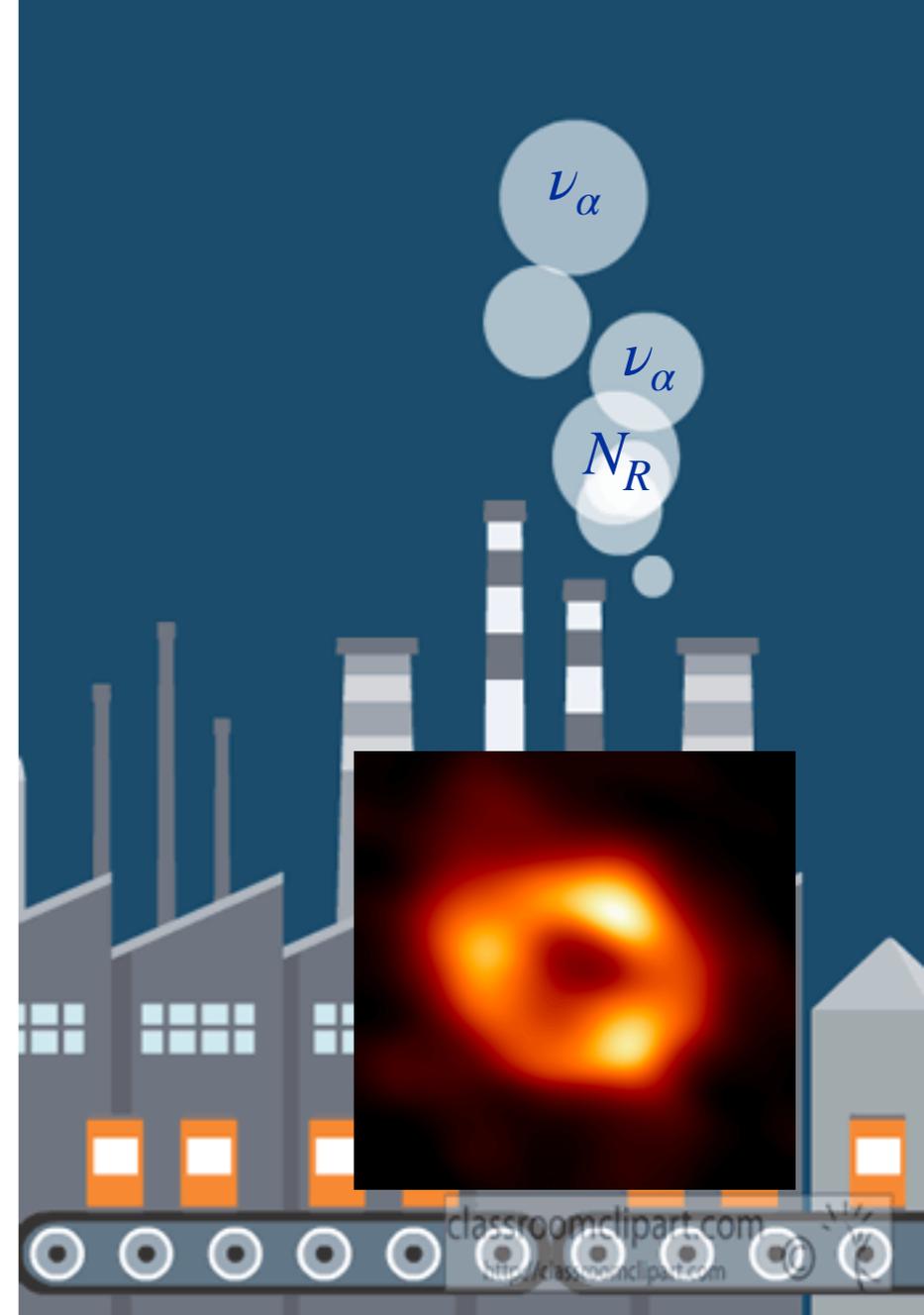


# Primordial Black Holes and Leptogenesis: An Unexpected Interplay

Yuber F. Perez-Gonzalez

IFT Seminar/NuTs Workshop  
May 23rd, 2022



EHT Collaboration



# *What's the impact of having evaporating Primordial Black Holes for Leptogenesis models?*

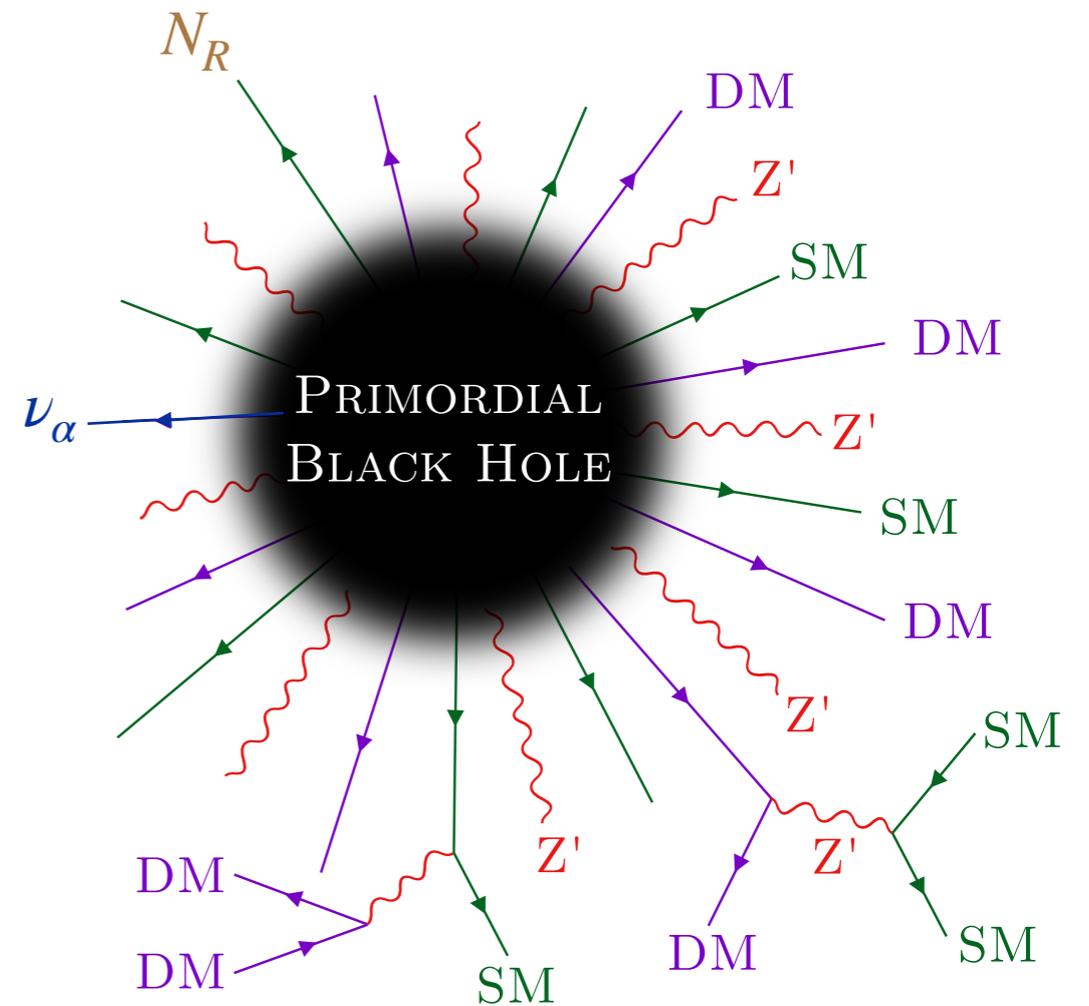
## ❖ Introduction:

- PBHs — Formation/Evaporation
- Thermal Leptogenesis

## ❖ Interplay between PBHs $\leftrightarrow$ Leptogenesis

## ❖ Testing these scenarios?

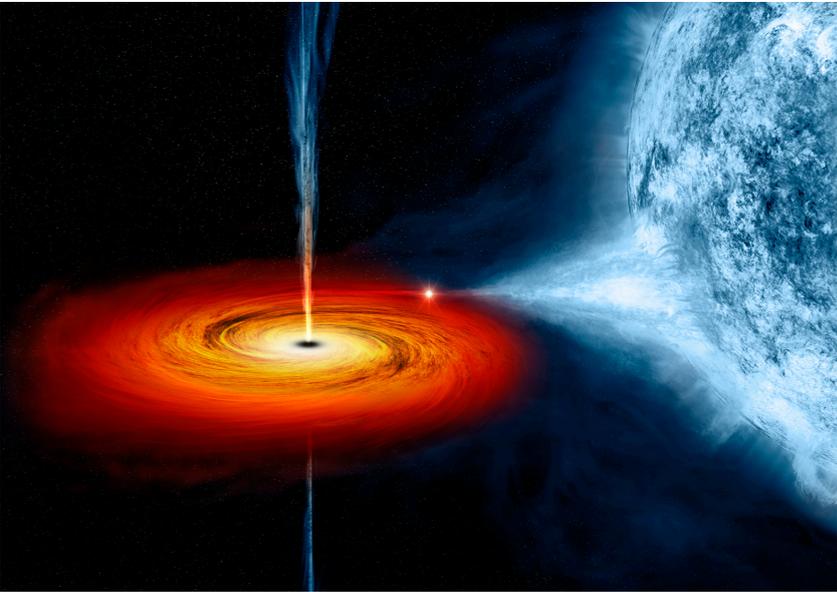
## ❖ In progress



# Primordial Black Holes (PBH)

Astrophysical Black Holes

$$M \gtrsim 3M_{\odot}$$



Formation

Lighter Black Holes

$$r_S = 2GM$$

Large densities

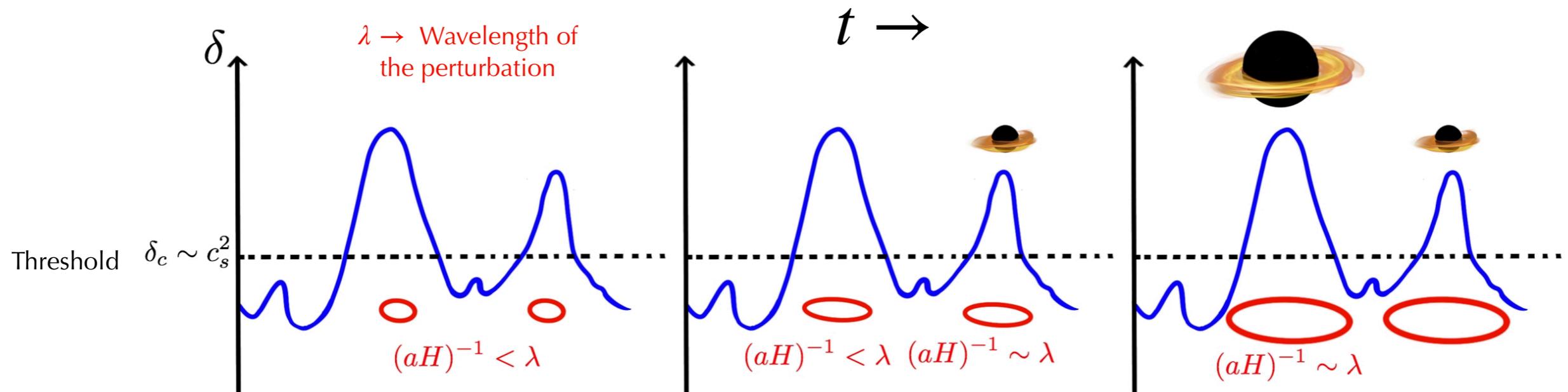
$$M_i \sim \frac{t}{G} \sim 10^{15} \left( \frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

$$\delta = \frac{\rho - \bar{\rho}}{\bar{\rho}}$$

# Formation

- ❖ Bubble collisions
- ❖ Pressure reduction
- ❖ Collapse of density fluctuations

$$M_{\text{BH},i} \propto \underbrace{\frac{4\pi}{3} \rho \frac{1}{H^3}}_{\text{Mass contained in the horizon}}$$



Taken from Villanueva-Domingo, Mena, Palomares-Ruiz 2103.12087

PBH form when the density perturbation enters the Hubble horizon

Fraction of the total energy in PBH

$$\beta = \frac{\rho_{\text{PBH}}}{\rho_{\text{tot}}}$$

Assume a monochromatic mass distribution

All PBHs with the same mass

Carr et al. 2002.12778

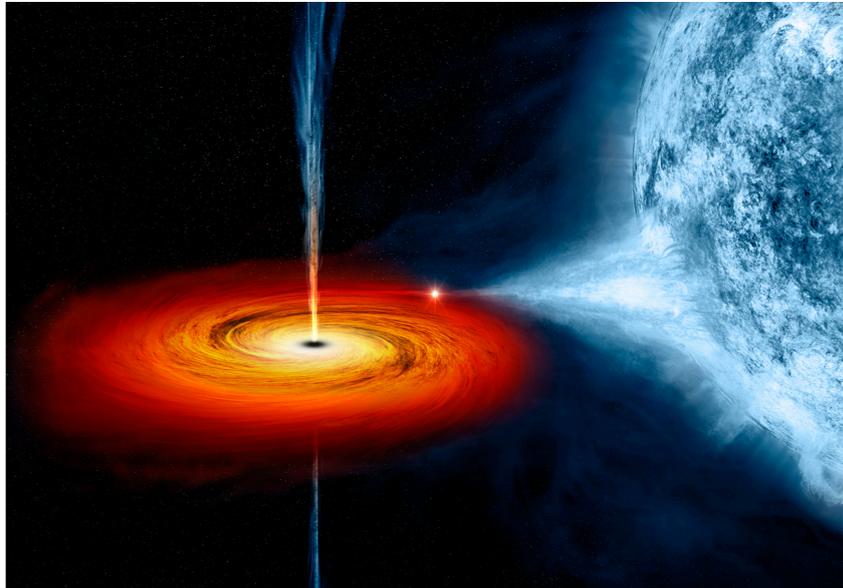
$$M_{\text{BH},i}, \beta$$

PBHs don't grow as fast as the Horizon  
Hawking, Carr & Hawking, '74

# Primordial Black Holes (PBH)

Astrophysical Black Holes

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Formation

Lighter Black Holes

Large densities

$$r_S = 2GM$$

$$M_i \sim \frac{t}{G} \sim 10^{15} \left( \frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

Assume a monochromatic mass distribution

$$M_{\text{BH},i}, \beta$$

Evaporation

Quantum effects are important

Black Holes evaporate by thermal emission

$$r_S \sim \lambda_C$$

Hawking, 1975

$$r_S \sim 0.015 \text{ fm} \left( \frac{M}{10^{13} \text{ g}} \right)$$

All degrees of freedom!

# Evaporation

Schwarzschild BH

Hawking  
Instantaneous  
Spectrum

$$\frac{d^2 N_i}{dp dt} = \frac{g_i}{2\pi^2} \frac{\sigma_{s_i}(M, p, \mu_i) p^3}{\exp[E_a(p)/T] - (-1)^{2s_i} E_a(p)}$$

Absorption  
cross section

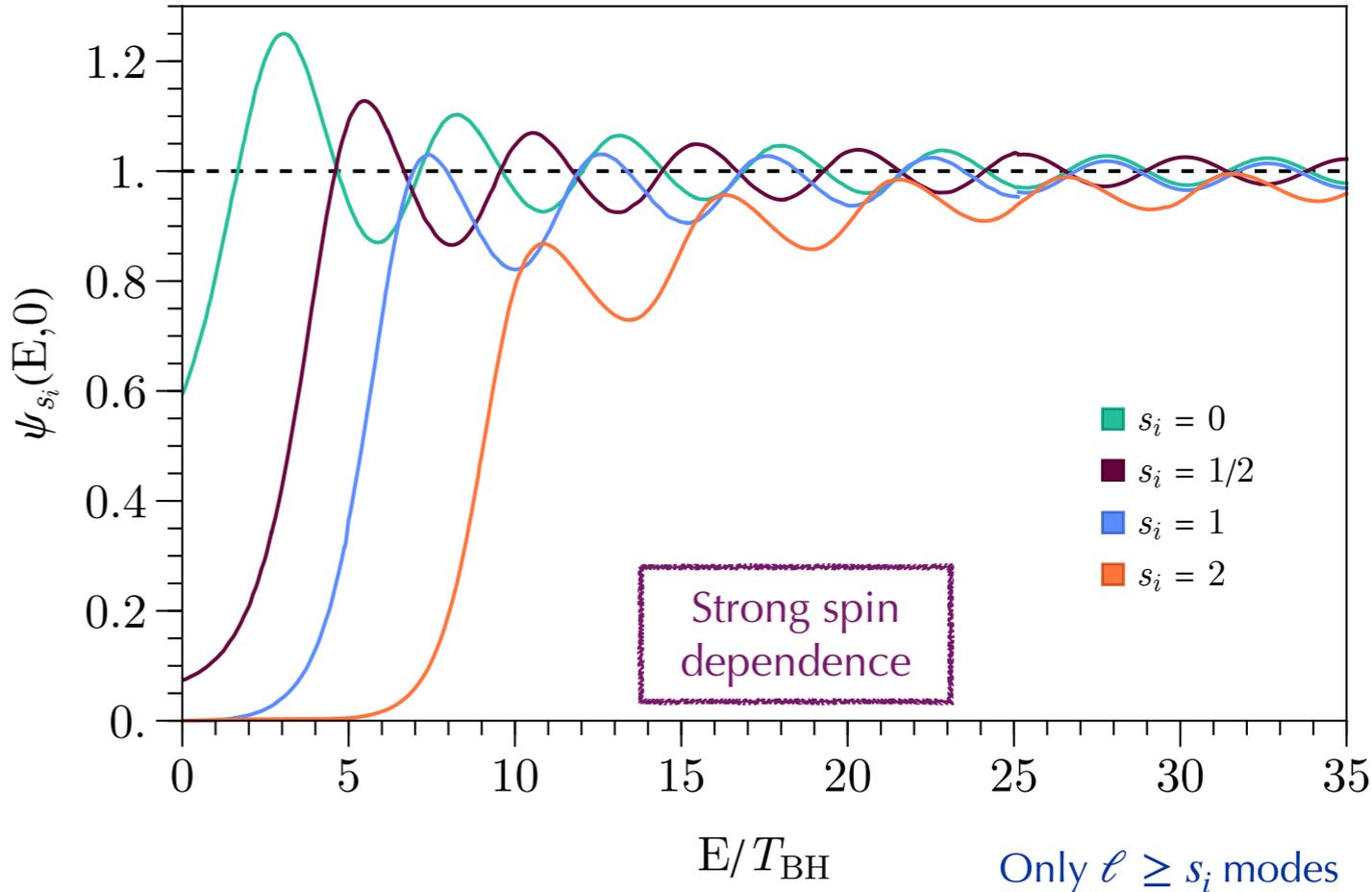
BH Temperature

$$T = \frac{\hbar c^3}{8\pi G M k}$$

$$\approx 1.06 \text{ GeV} \left( \frac{10^{13} \text{ g}}{M} \right)$$

$$\psi_{s_i}(E) \equiv \frac{\sigma_{s_i}(E)}{27\pi G^2 M_{\text{BH}}^2}$$

Reduced Absorption Cross Section



\*Hic depositum est, quod mortale fuit Isaaci Newtoni

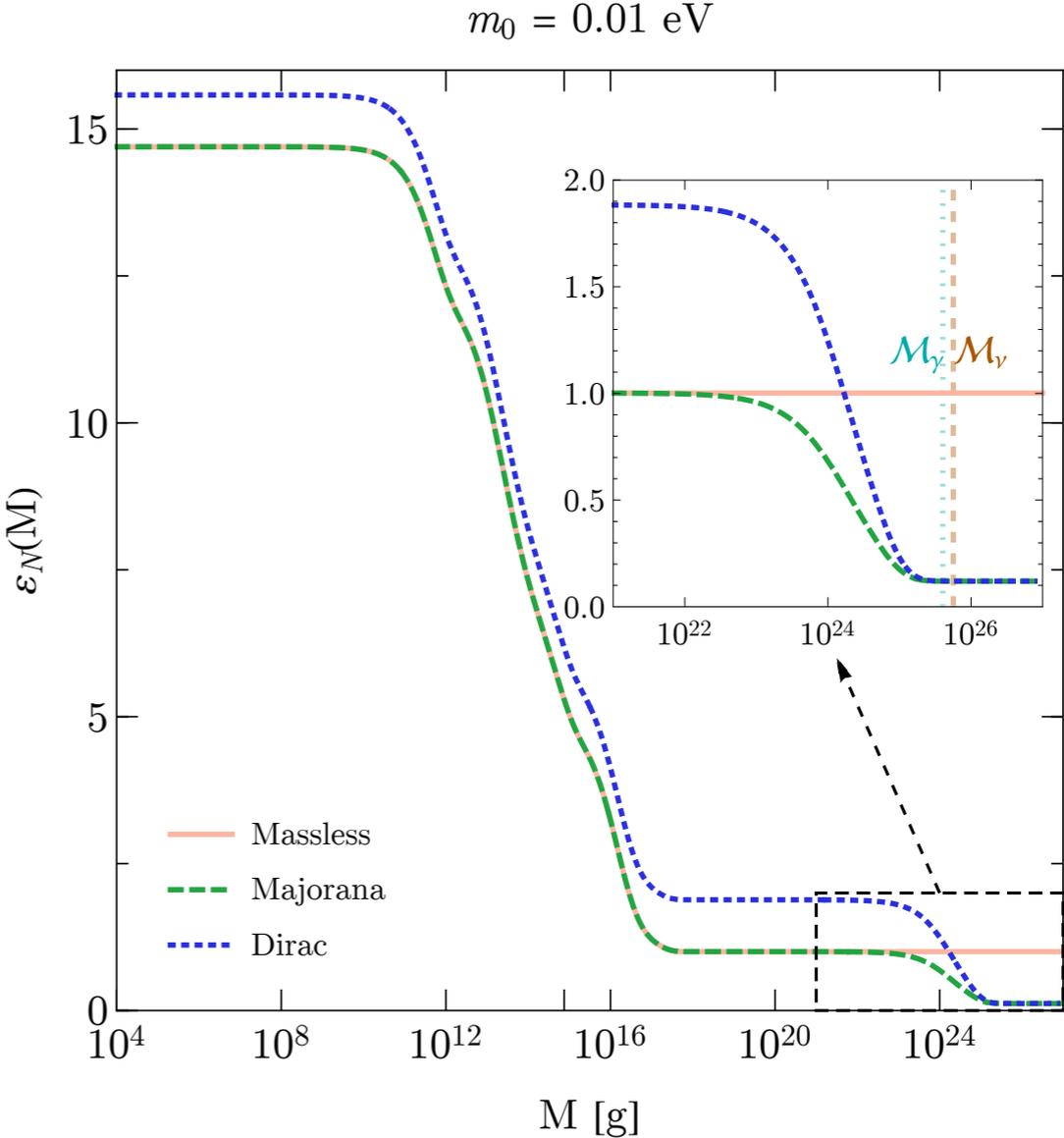
# Evaporation

BHs lose mass over time

Depends on the set of possible particles to be emitted

J. MacGibbon, 1991

$$\begin{aligned} \frac{dM_{\text{BH}}}{dt} &= - \sum_i \int_0^\infty E_i \frac{d^2 \mathcal{N}_i}{dp dt} dp \\ &= - \underbrace{\varepsilon(M_{\text{BH}})}_{\text{Evaporation function}} \frac{M_{\text{P}}^4}{M_{\text{BH}}^2} \end{aligned}$$



BH lifetime

$$\tau = \frac{1}{M_{\text{P}}^4} \int_{M_{\text{P}}}^{M_i} \frac{M^2 dM}{\varepsilon(M)}$$

Planck mass

BH with  $M_i = M_{\odot}$  :

$$\tau \approx 2 \times 10^{67} \text{ years}$$

$$\approx 10^{57} * \text{age of the Universe}$$

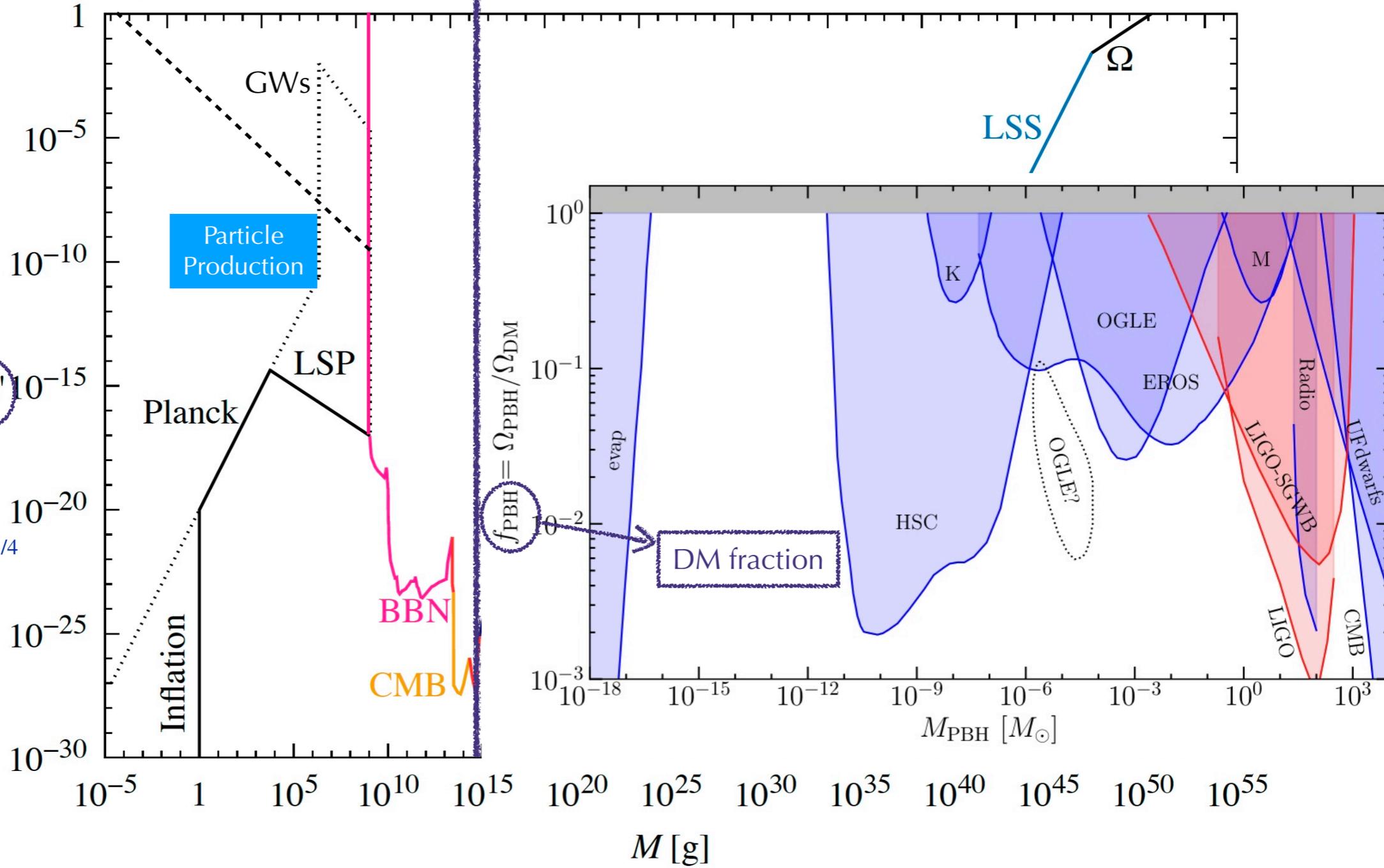
Evaporated ←  $M/M_\odot$  → (Part of) Dark Matter?

Reduced  $\beta$

$\beta'$

$$\beta' = \beta \gamma^{1/2} \left( \frac{g_{*f}}{106.75} \right)^{1/4}$$

Gravitational collapse factor



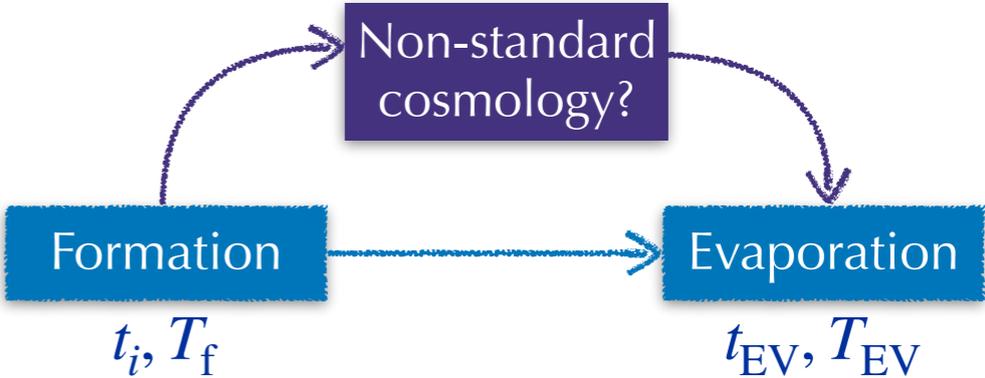
We focus on the region  $M_{\text{BH},i} \leq 10^9 \text{ g}$

Carr et al. 2002.12778  
Domènech et al. 2012.08151

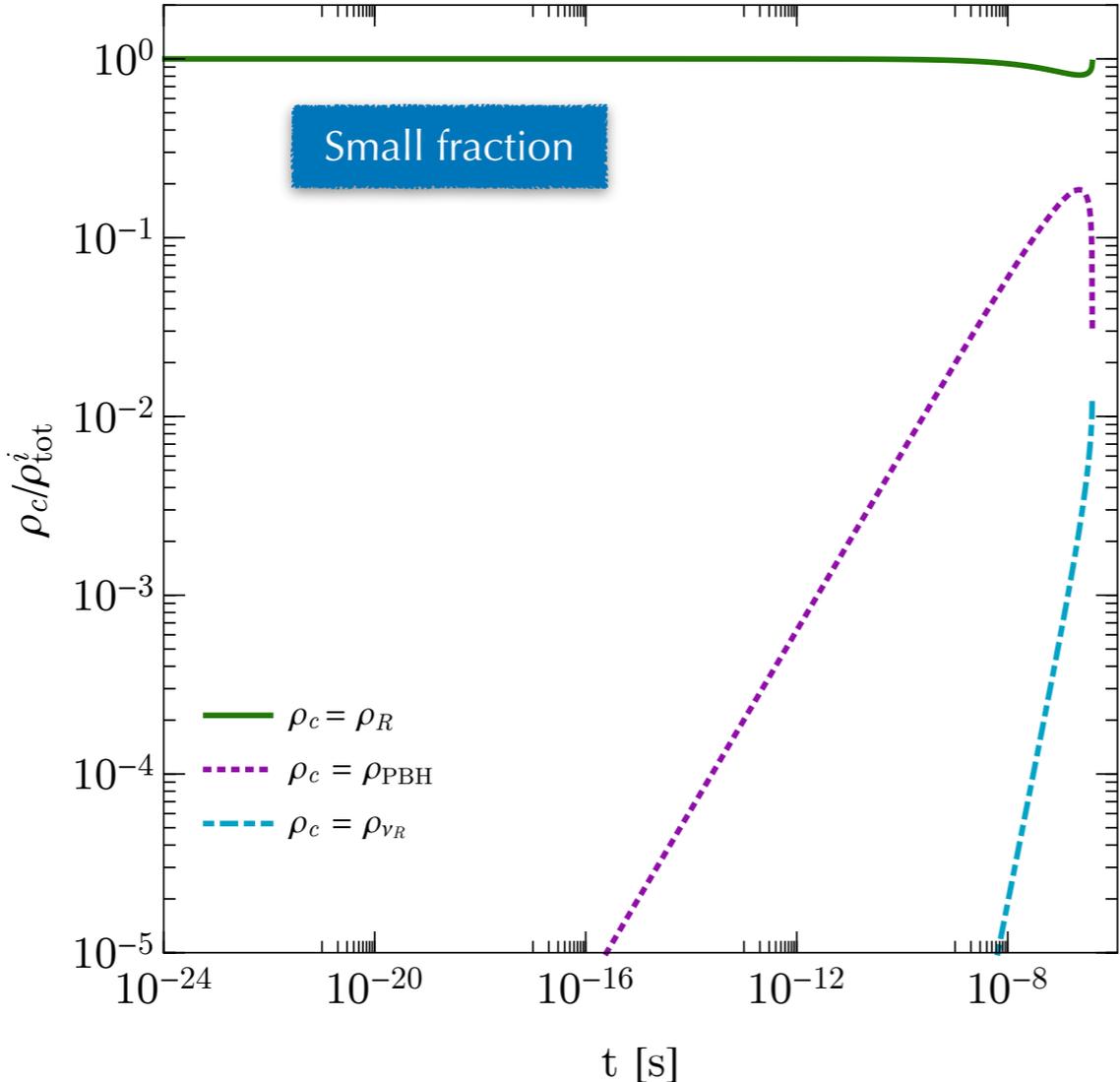
B. Kavanagh  
10.5281/zenodo.3538999

# PBHs: Rise and Fall

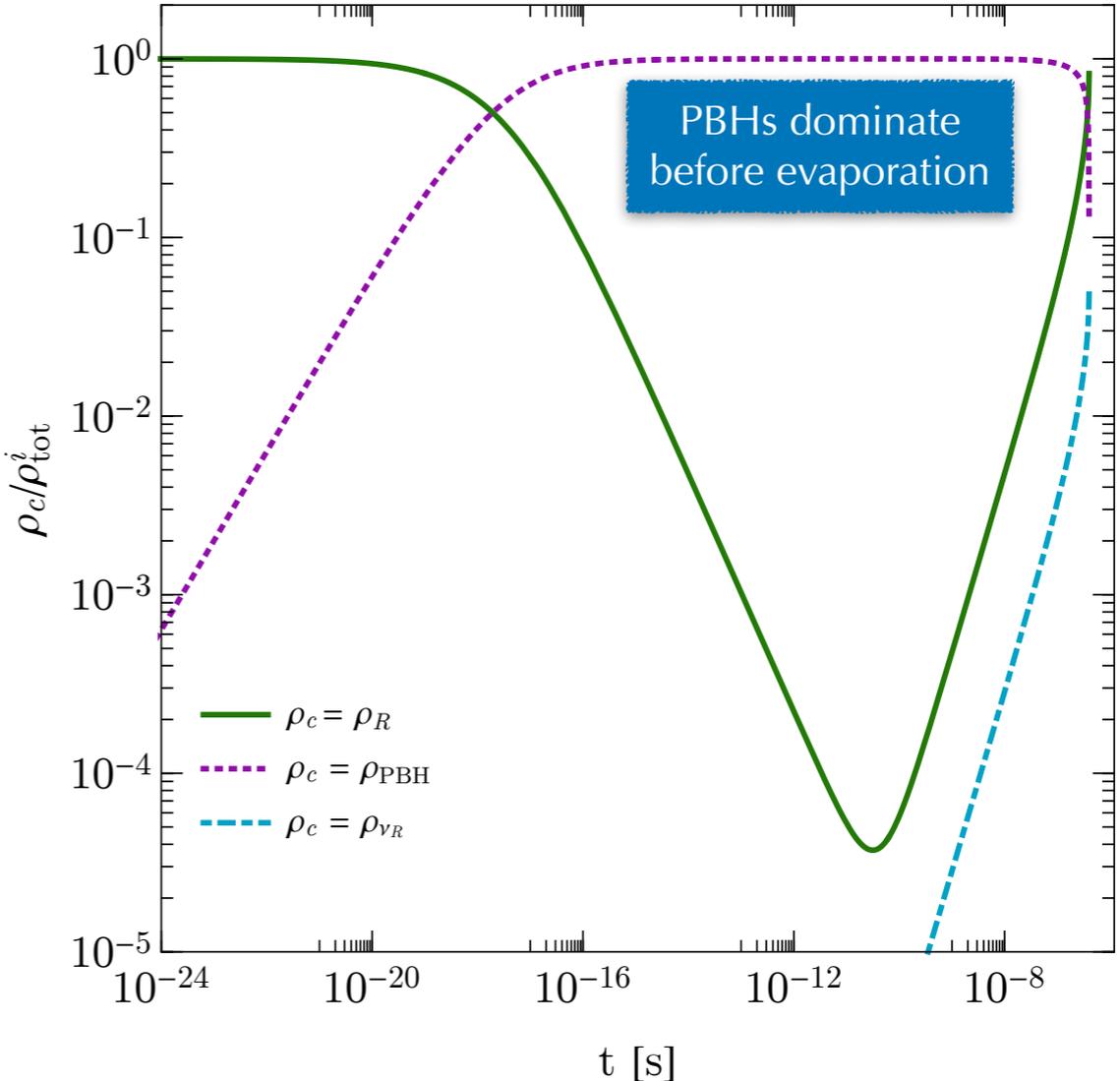
PBHs behave (almost) as matter



$$M_{\text{BH},i} = 10^7 \text{ g}, \beta = 10^{-13}$$



$$M_{\text{BH},i} = 10^7 \text{ g}, \beta = 10^{-7}$$



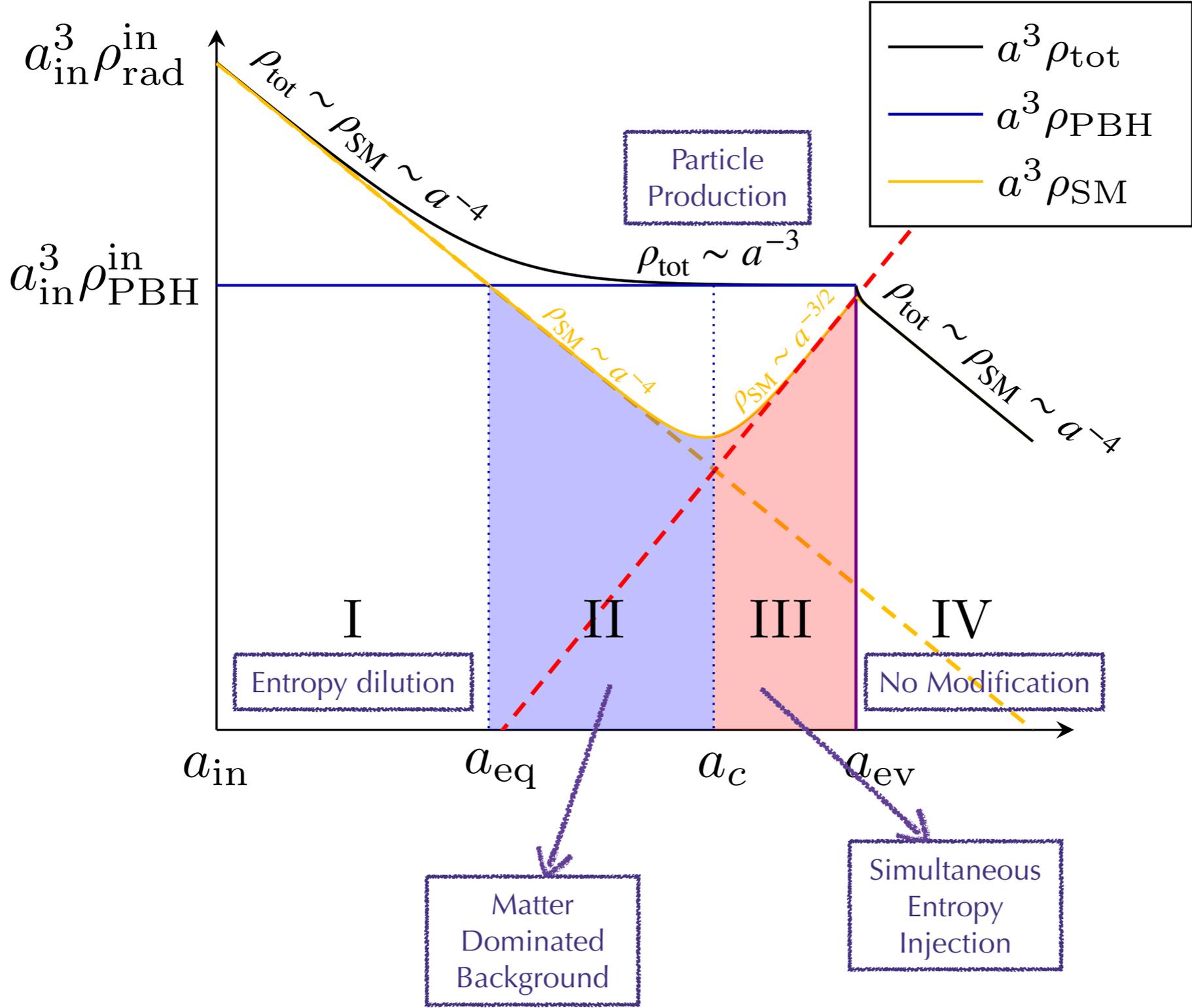
Initial fraction

$$\beta \gtrsim 2.5 \times 10^{-14} \left( \frac{g_*(T_f)}{106.75} \right)^{-\frac{1}{4}} \left( \frac{M_i}{10^8 \text{ g}} \right)^{-1} \left( \frac{\epsilon_D(M_i)}{15.35} \right)^{\frac{1}{2}}$$

Hooper et al,  
1905.01301

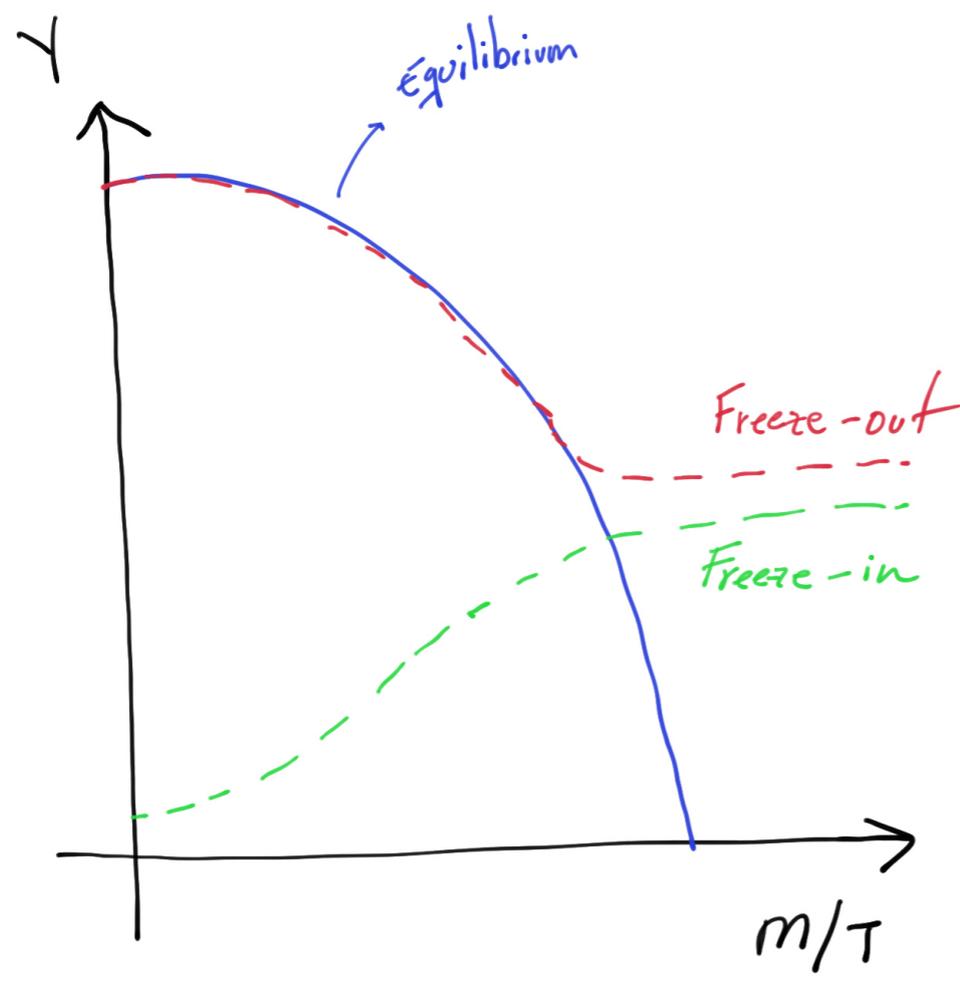
# PBHs: Rise and Fall

In a PBH dominated Early Universe



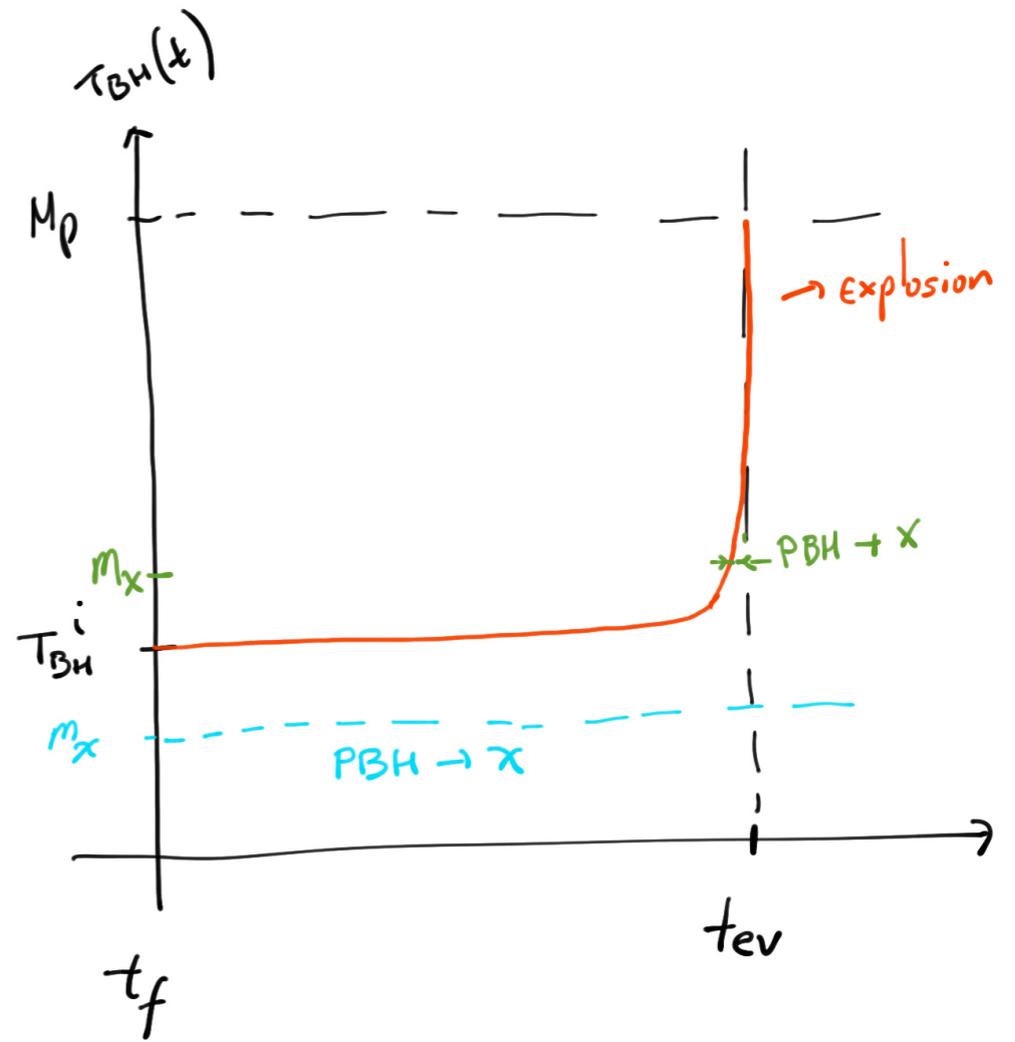
# Why is interesting the particle production via evaporation?

## (Thermal) Particle decoupling



Abundances depend on masses and interactions

## Hawking Evaporation



Abundances depend on masses and PBH parameters



- ❖ Modified Cosmology
- ❖ Particle production
- ❖ Entropy dilution

# Why is interesting the particle production via evaporation?

DM production  
(Besides PBH-DM)

- ❖ Purely Gravitationally interacting DM
- ❖ Modify Freeze-In/Freeze-out mechanisms

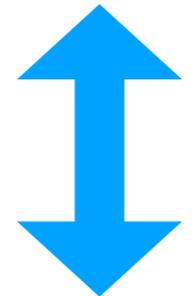
Fujita et al, 1401.1909  
 Morrison et al,1812.10606  
 Baldes et al, 2004.14773  
 Masina, 2004.04740, 2103.13825  
 Cheek, Heurtier, YFPG, Turner  
 2107.00013, 2107.00016  
 .  
 .  
 .



Baryon Asymmetry

- ❖ Modifying Baryogenesis scenarios
- ❖ **Leptogenesis scenarios**
- ❖ Producing a local asymmetry at PBH formation
- ❖ Connections with PBH-DM

Baumann, Steinhadt, Turok, 0703250  
 Yamada and Iso, 1610.02586  
 Fujita et al, 1401.1909  
 Morrison et al,1812.10606  
 García-Bellido, Carr, Clesse,  
 1904.11482  
 Hooper and Krnjaic, 2010.01134  
 .  
 .  
 .



Dark Radiation

- ❖ Production of hot gravitons
- ❖ Testable from future measurements on  $\Delta N_{\text{eff}}$

Hooper, Krnjaic, McDermott, 1905.01301  
 Lunardini, YFPG, 1910.07864  
 Masina, 2004.04740, 2103.13825  
 .  
 .  
 .

# Thermal Leptogenesis in a nutshell

Baryon asymmetry

$$\eta_{\text{CMB}} = (6.23 \pm 0.17) \times 10^{-10}$$

$$\eta_{\text{BBN}} = (6.08 \pm 0.06) \times 10^{-10}$$

Type I seesaw:

$$-\mathcal{L} \supset \frac{1}{2} M_{N_i} \bar{N}_i^c N_i + \bar{\ell}_\alpha H^* Y_{\alpha i} N_i + \text{H.c.}$$

$$m_\nu \sim \frac{Y_\nu^2 v^2}{M_N} \sim \mathcal{O}(0.1 \text{ eV})$$

Sakharov Conditions

- ❖ Baryon and Lepton number violation
- ❖ CP violation
- ❖ Departure from thermal equilibrium

$$Y_\nu = \frac{1}{v} U_{\text{PMNS}} \sqrt{m_\nu} R^T M_N^T$$

Casas, Ibarra, 2001

Boltzmann Equations

$$\frac{dn_{N_i}}{dz} = D_i (n_{N_i} - n_{N_i}^{\text{eq}})$$

$$\frac{dn_{\text{B-L}}}{dz} = \sum_i \epsilon_i D_i (n_{N_i} - n_{N_i}^{\text{eq}}) - \mathcal{W}_i n_{\text{B-L}}$$

Washout

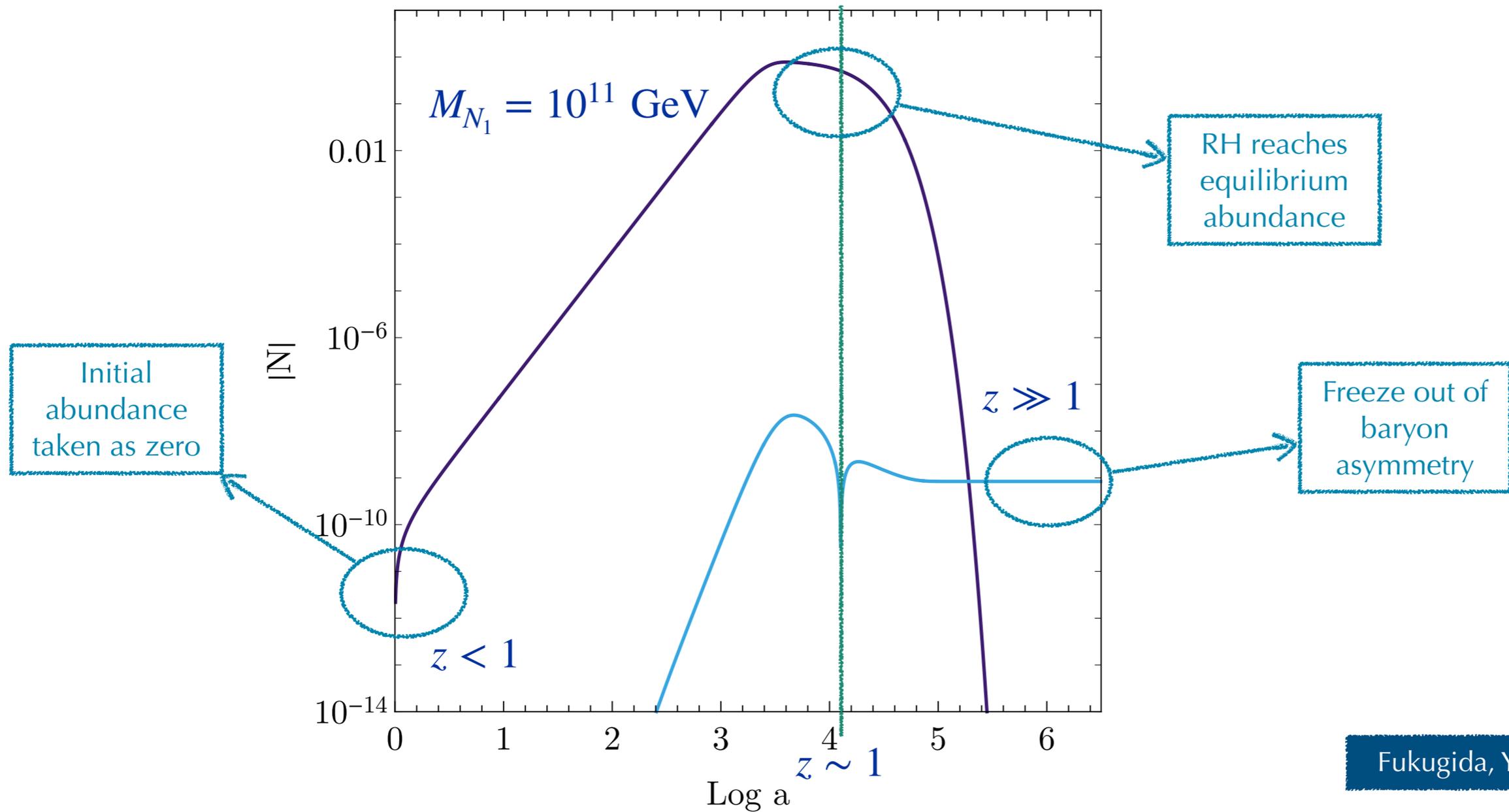
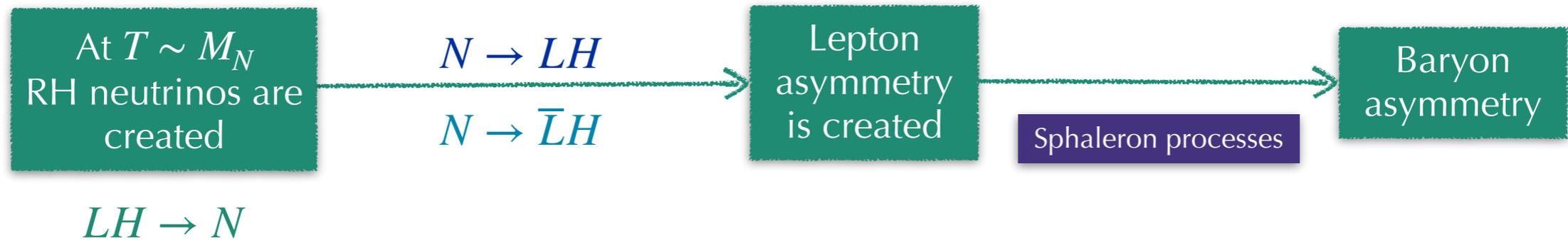
$$LH \leftrightarrow \bar{L}H$$

$$\Delta L = 2$$

Fukugida, Yanagida, '86

# Thermal Leptogenesis in a nutshell

$$z \equiv \frac{M_{N_1}}{T}$$



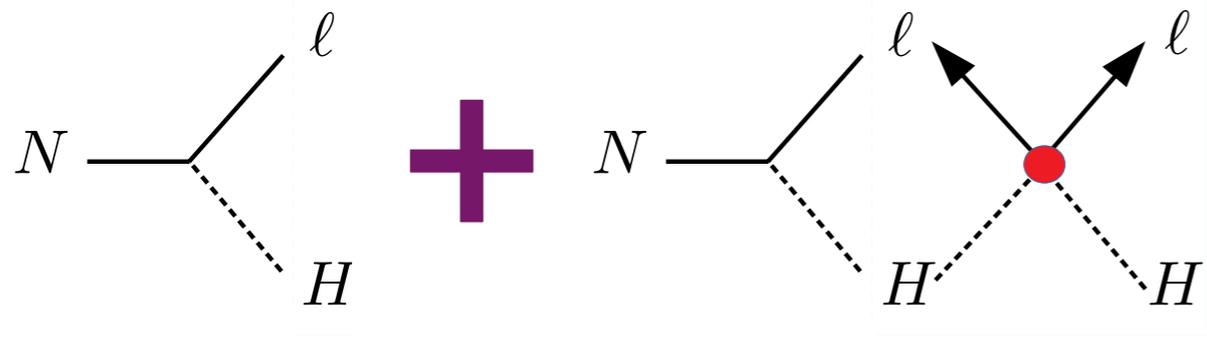
Fukugida, Yanagida, '86

# High Scale Leptogenesis

$$M_N \gtrsim 10^{12} \text{ GeV}$$

$$M_{N_1} < M_{N_2} < M_{N_3}$$

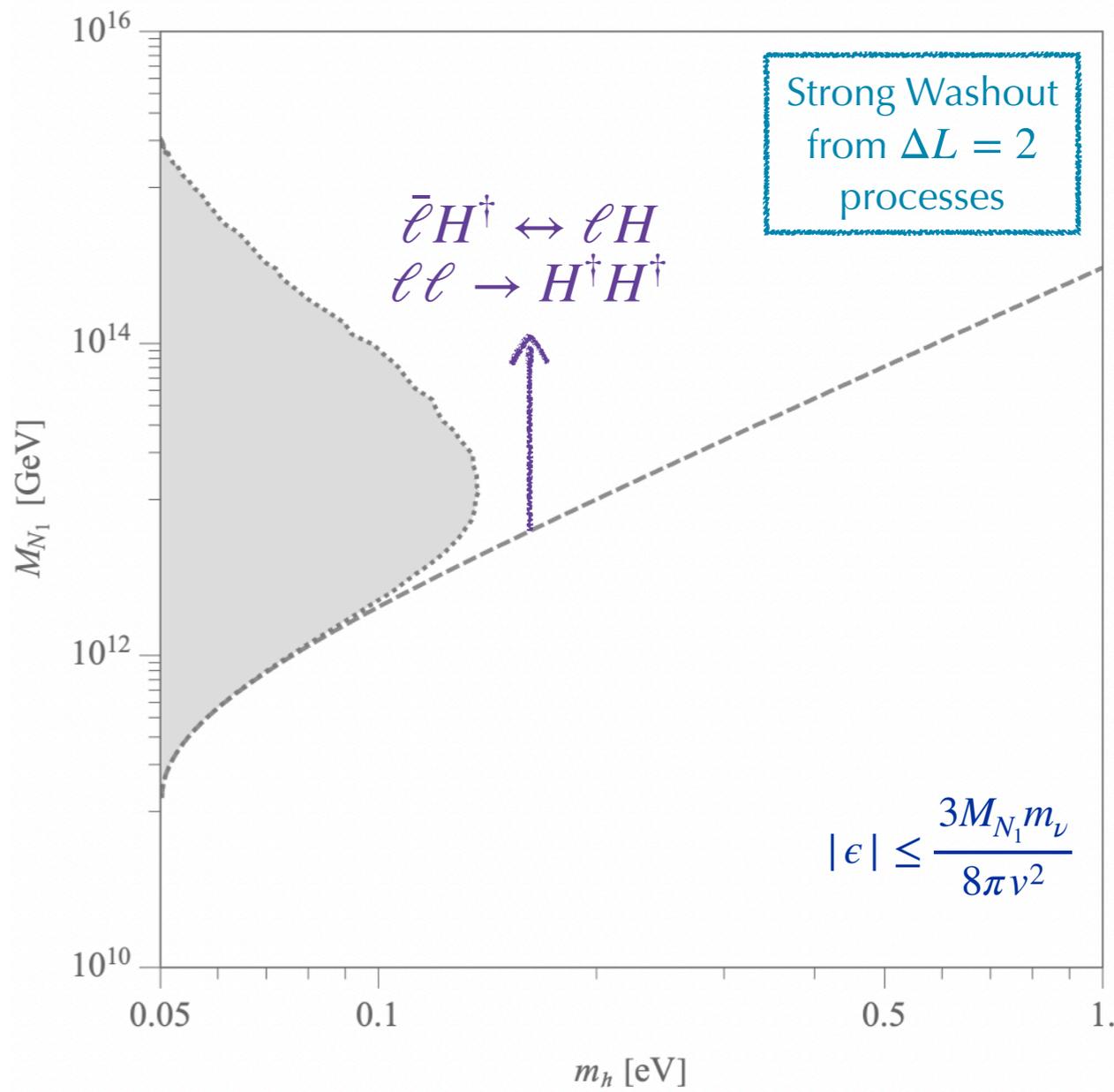
CP violation from interference



$$\frac{\epsilon}{M_N} \propto \frac{m_\nu}{v^2} \propto \sqrt{\sigma_{\Delta L=2}}$$

$\Delta L = 2$  interactions would be in equilibrium if

$$M_{N_1}^T \gtrsim 4 \left( \frac{0.1 \text{ eV}}{m_\nu} \right) 10^{12} \text{ GeV}$$



$$|\epsilon| \leq \frac{3M_{N_1}m_\nu}{8\pi v^2}$$

$m_h =$  Heaviest neutrino mass

Maximizing over Yukawa parameters

How to save HSL?

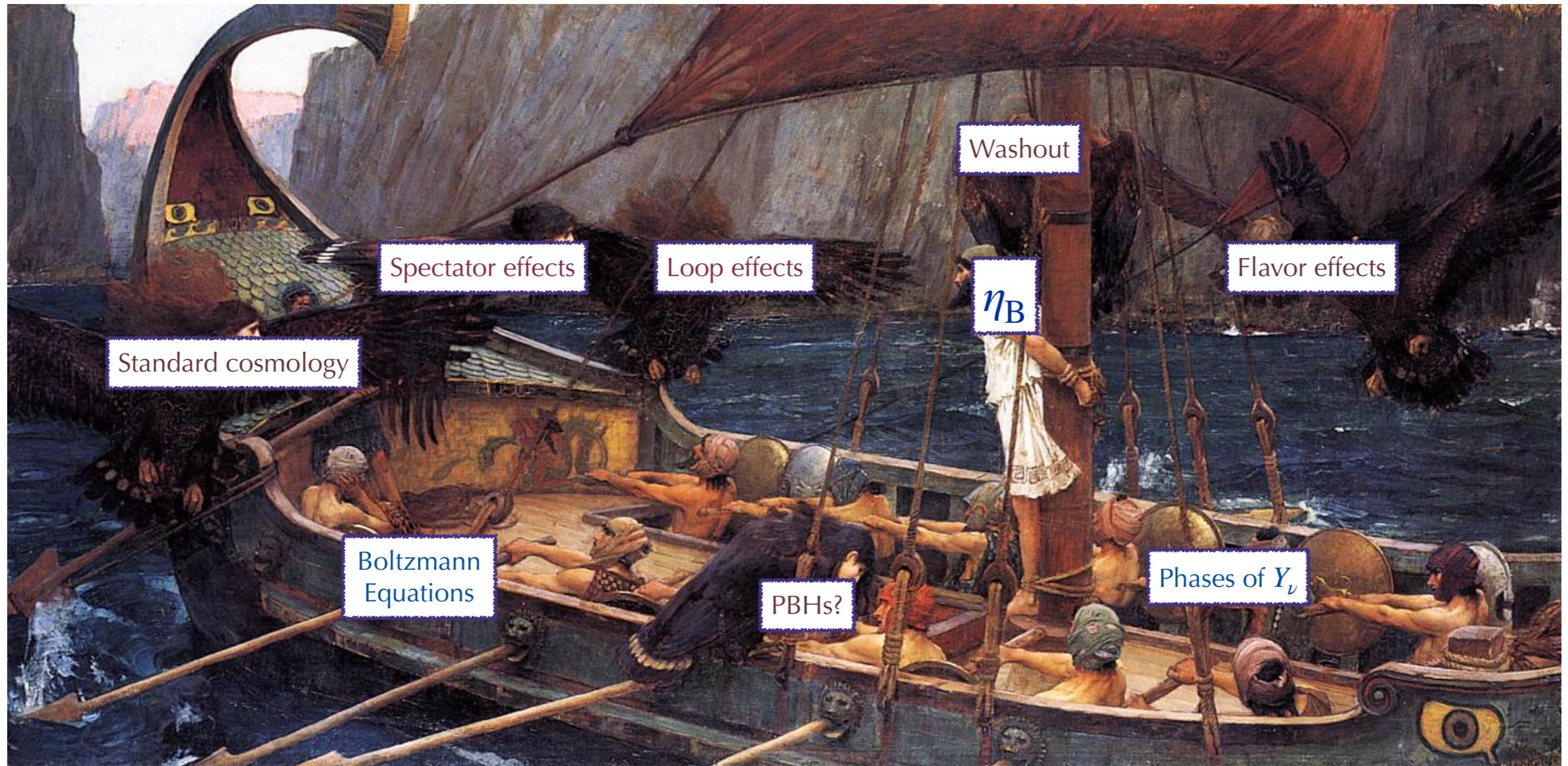
Produce RHNs after washout process have frozen out?

PBHs!

Giudice *et al.*, 2004  
Buchmuller *et al.*, 2005

*Intermezzo.*

# Universal LeptogeneSiS Equation Solver (ULYSSES)



A Granelli, K Moffat, YFPG,  
H Schulz and J Turner,  
arXiv: [2007.09150](https://arxiv.org/abs/2007.09150)

- ❖ Leptogenesis via decays and resonant leptogenesis
- ❖ Easy parallelization
- ❖ Rapid evaluation
- ❖ Multidimensional scan of the parameter space

# Interplay between PBH Evaporation and Leptogenesis

Based on:  
YFPG and Turner: 2010.03565  
Bernal, Fong, YFPG, Turner 2203.08823

# Baryogenesis from PBHs only

Assume and additional  $X$  particle producing a  $B - L$  asymmetry

$$N_X = \int_0^\tau dt \int_{M_X}^\infty dE \frac{d^2 N_X}{dt dE}$$

Total number of  $X$  particles produced

Yield of  $B - L$

$$Y_X \equiv \frac{3\beta T_0 N_X}{4M_{\text{BH0}}} \frac{s(T_{\text{ev}})}{s(\tilde{T})}$$

Entropy dilution

$$Y_X \simeq \frac{135\zeta(3)}{\pi^3 g_\star} \beta \frac{T_0}{M_{\text{BH0}}} \left(1 + \beta \frac{T_0}{T_{\text{ev}}}\right)^{-\frac{3}{4}} \times \begin{cases} \left(\frac{M_{\text{BH0}}}{M_P}\right)^2 & M_X \leq r_f T_{\text{BH0}}, \\ \frac{15\zeta(5)}{64\pi^2 \zeta(3)} \left(\frac{M_P}{M_X}\right)^2 & M_X \geq r_f T_{\text{BH0}}. \end{cases}$$

CP violation

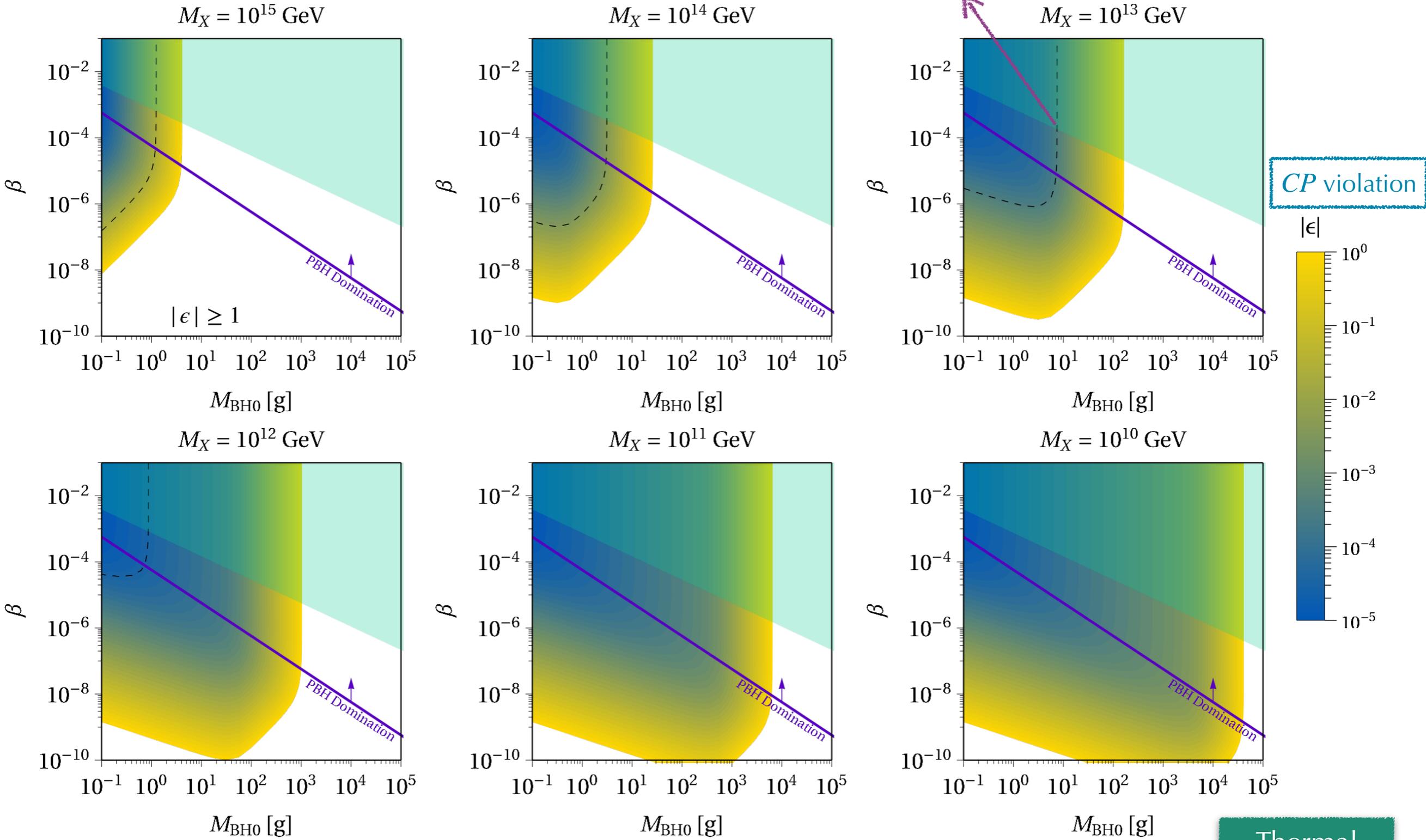
$$Y_B^{\text{max}} = \frac{30}{97} \epsilon Y_X$$

Toussiant et al, '75  
 Barrow et al, Phys. Rev. D **43**, 984  
 Baumann, Steinhardt, Turok 0703250  
 Fujita et al, 1401.1909  
 Morrison et al, 1812.10606  
 ...

$$r_f \equiv \sqrt{15 \zeta(5)/\zeta(3)}$$

# Baryogenesis from PBHs only

$$Y_B^{\max} = \frac{30}{97} \epsilon Y_X$$

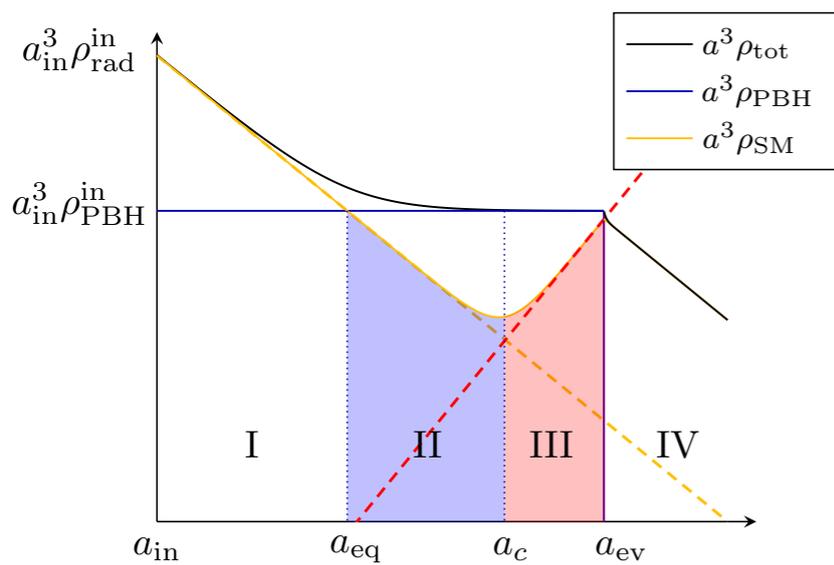
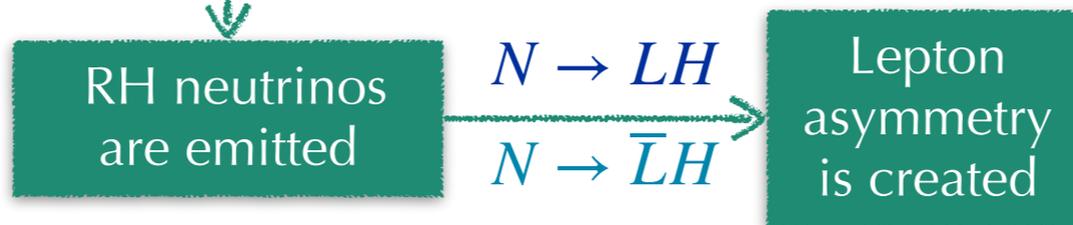
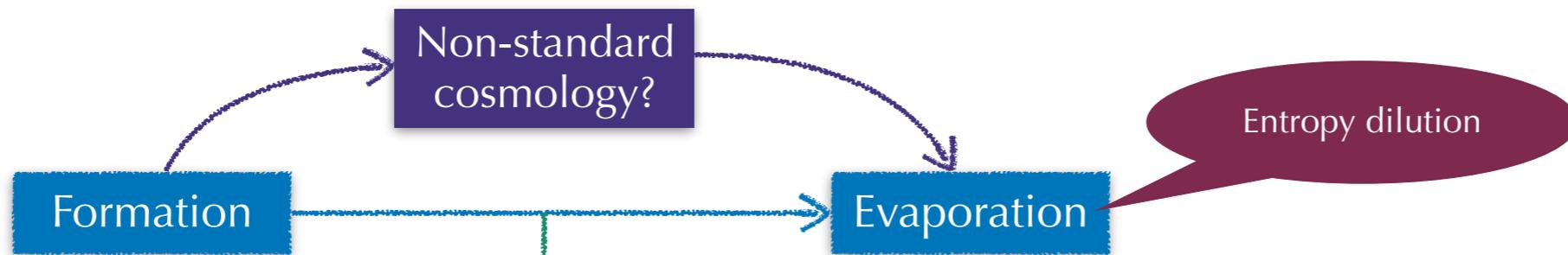
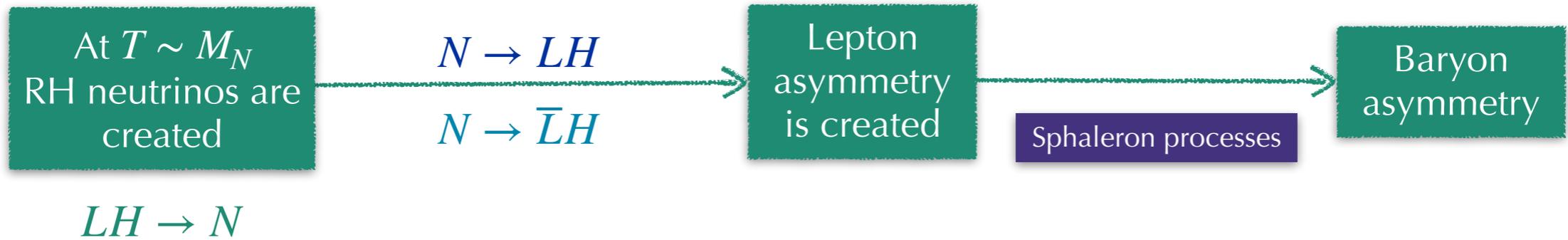


$$m_h^2 = |\Delta m_{\text{atm}}^2|$$

Tuning required for  $M \lesssim 10^{12}$  GeV

Thermal processes?

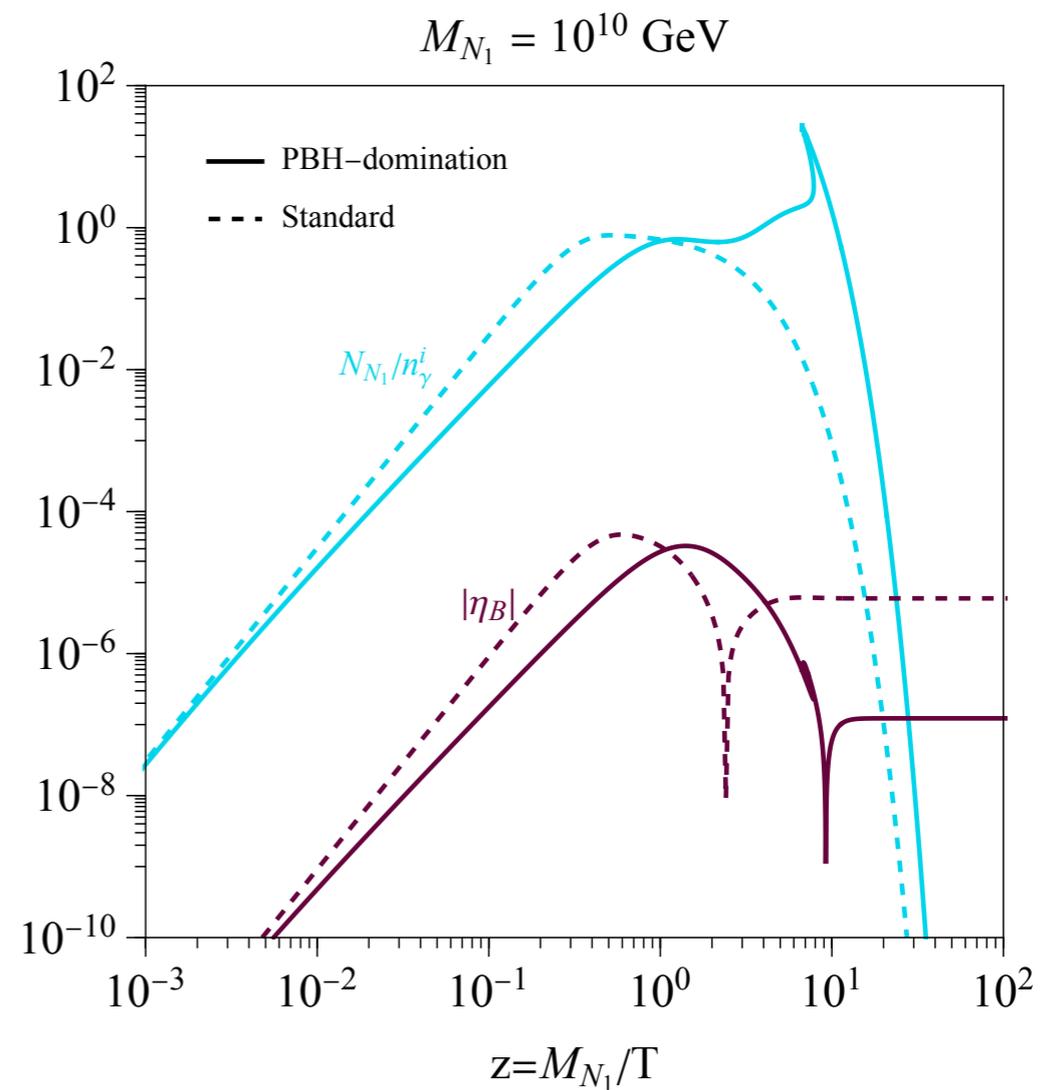
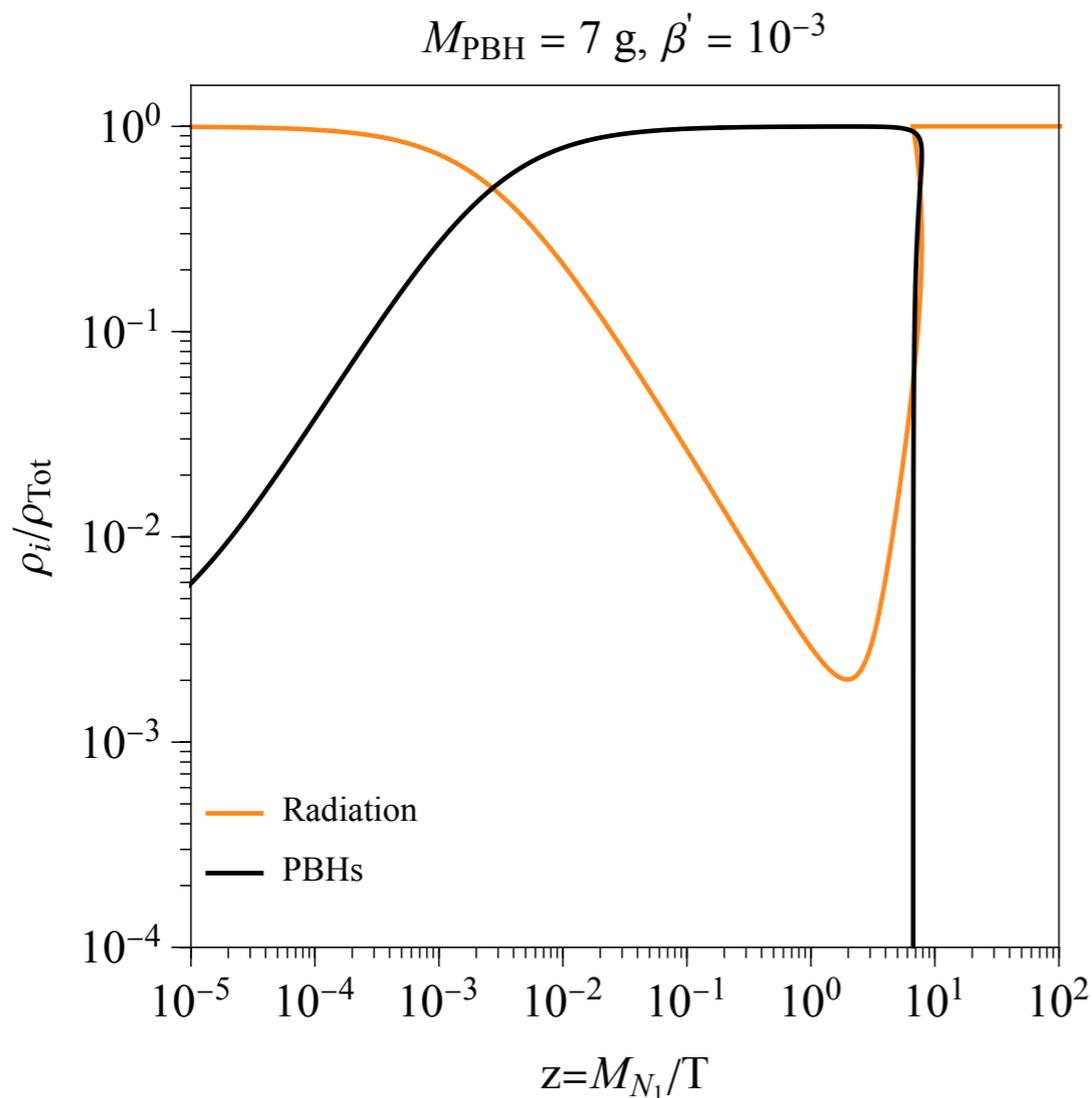
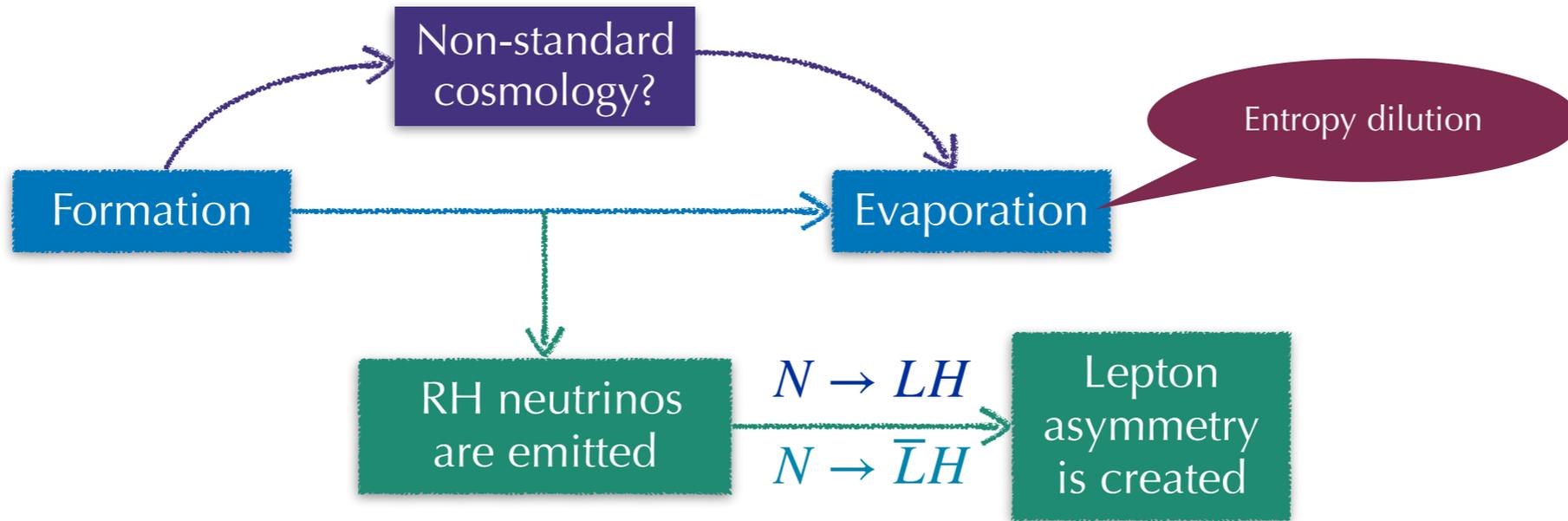
# PBH + Leptogenesis



$$z_{BH} \equiv \frac{M_{N_1}}{T_{BH}}$$

# PBH + Leptogenesis

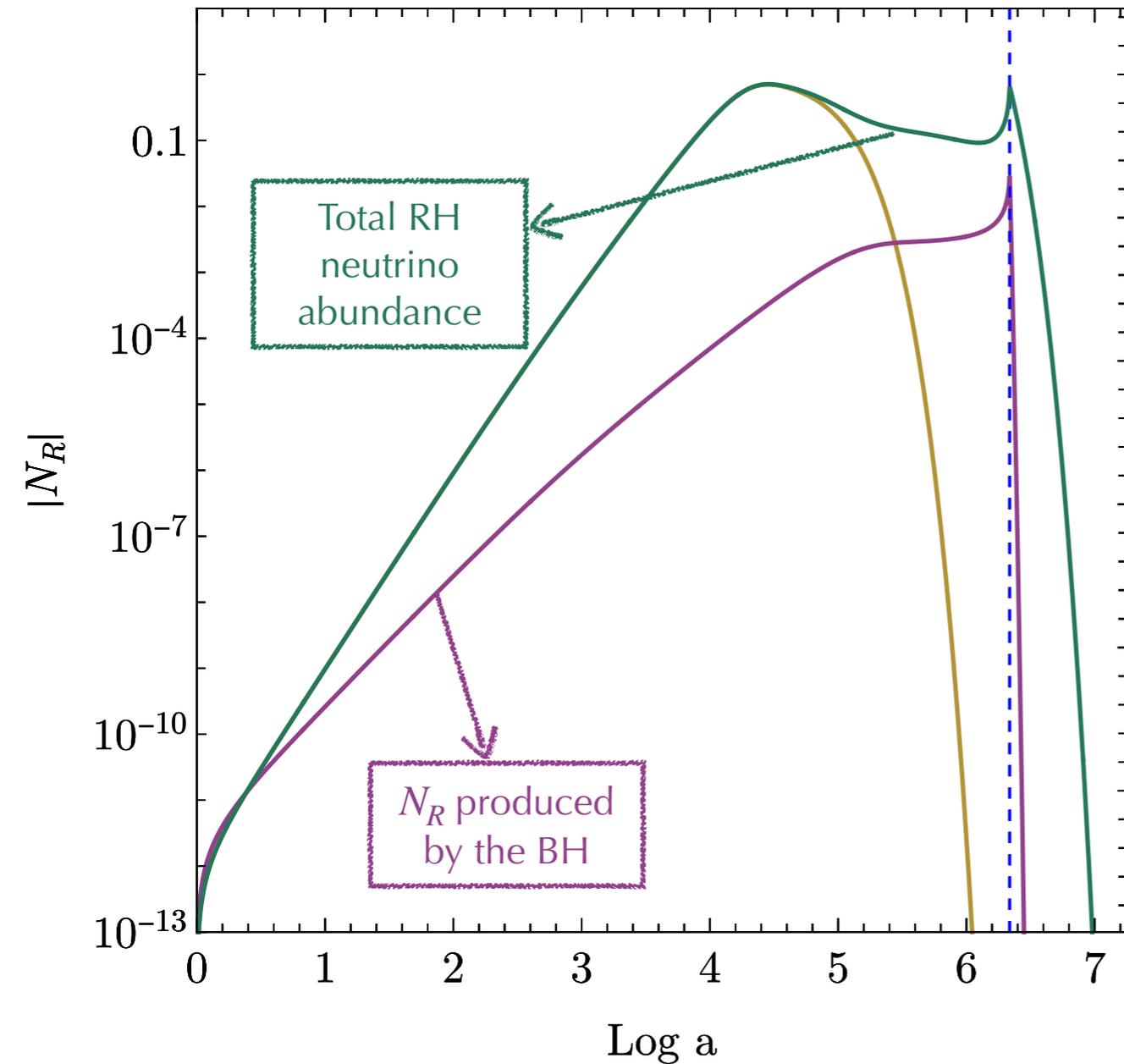
$$z_{\text{BH}} \equiv \frac{M_{N_1}}{T_{\text{BH}}}$$



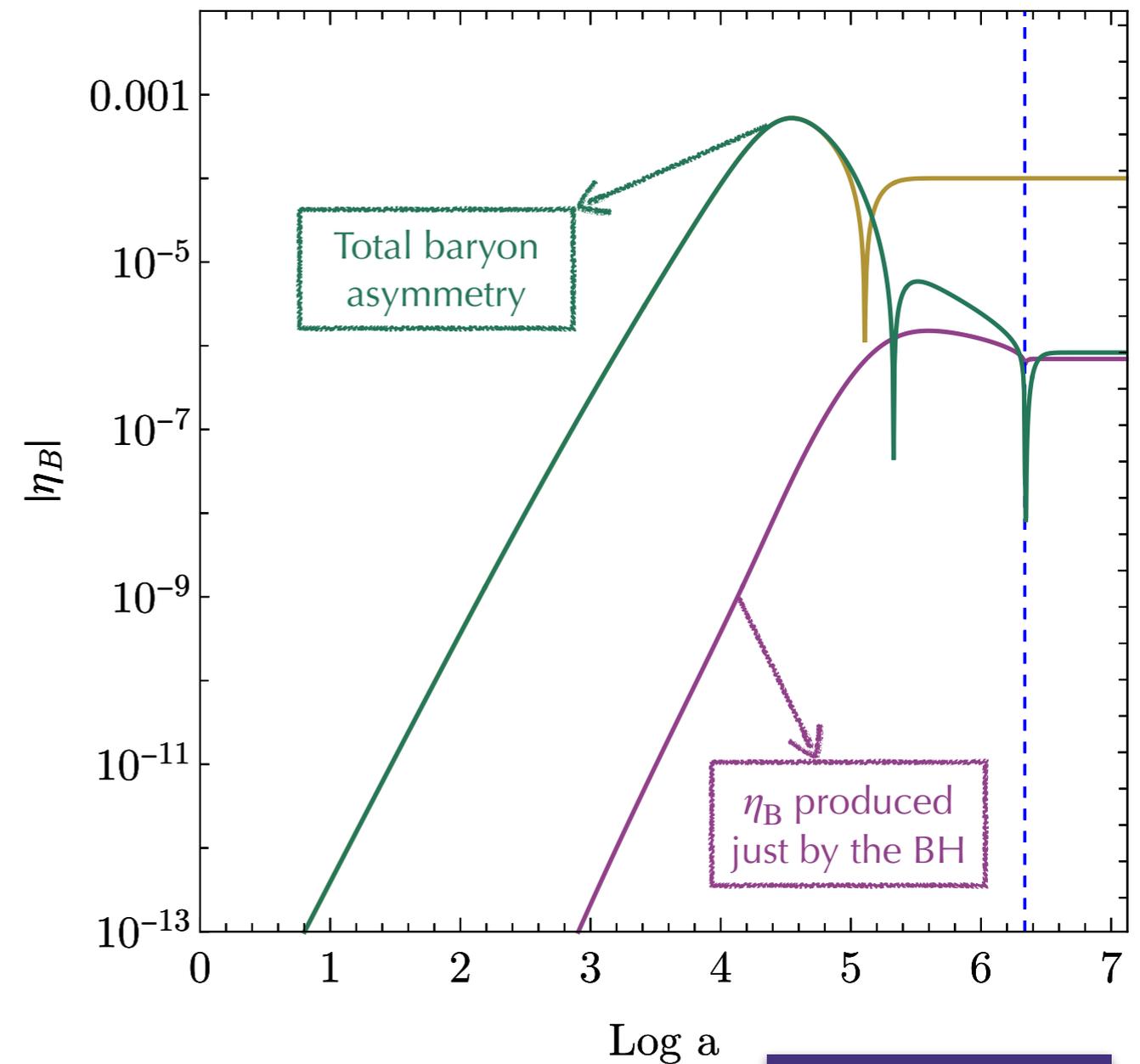
Three scenarios

- A. PBH evaporate before RHs are thermally produced (IV)
- B. Evaporation happens during thermal leptogenesis (II-III)
- C. PBH create a RH density after the thermal case (II-III)

RH neutrino abundance



Final baryon asymmetry



$$M_i = 1.7 \text{ g}$$

$$\beta'_i = 10^{-3}$$

$$M_{N_1} = 10^{11} \text{ GeV}$$

BH evaporation  
can diminish  
the asymmetry!

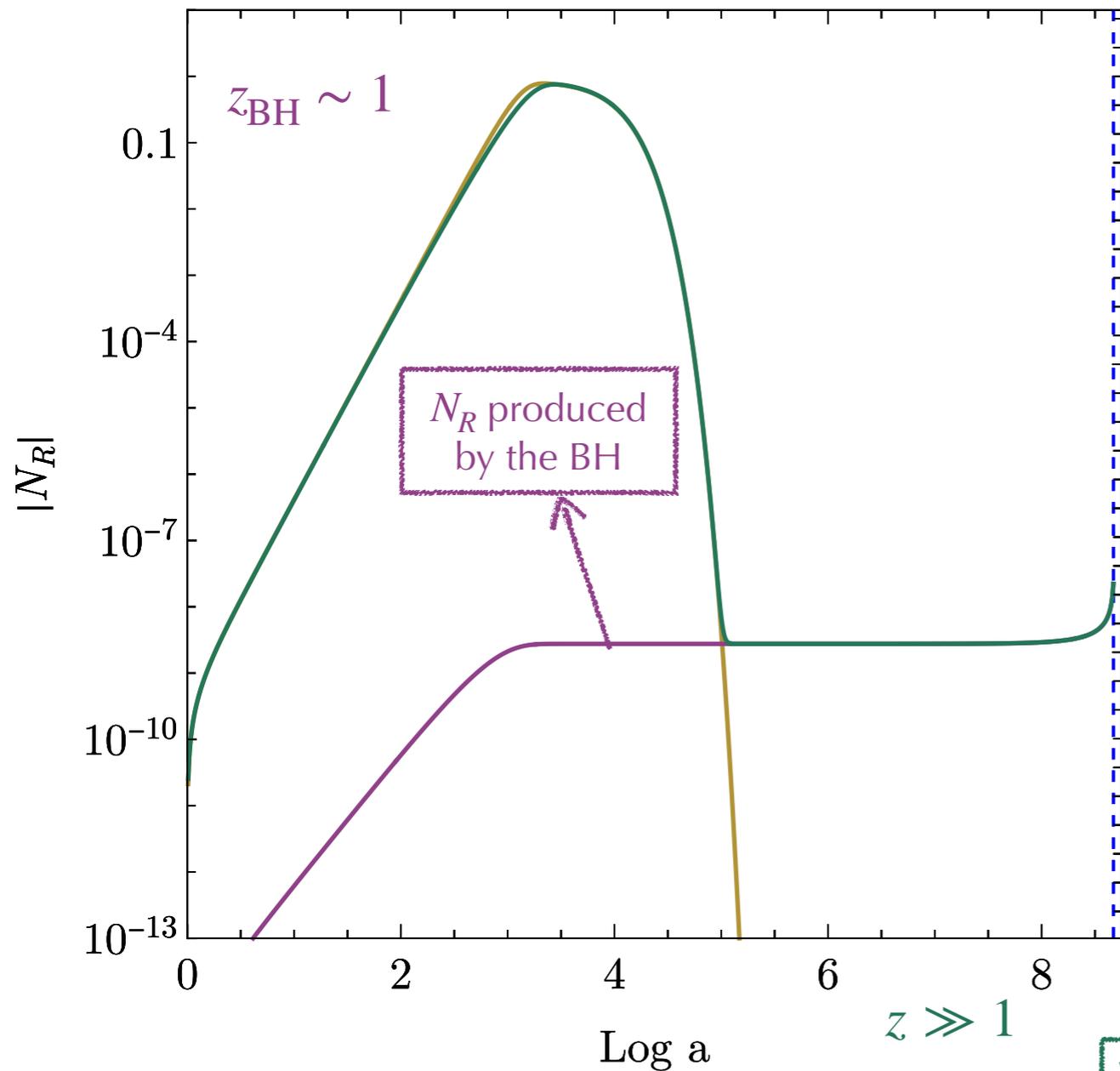
C. PBH create a RH density after the thermal case

$$M_i = 10^2 \text{ g}$$

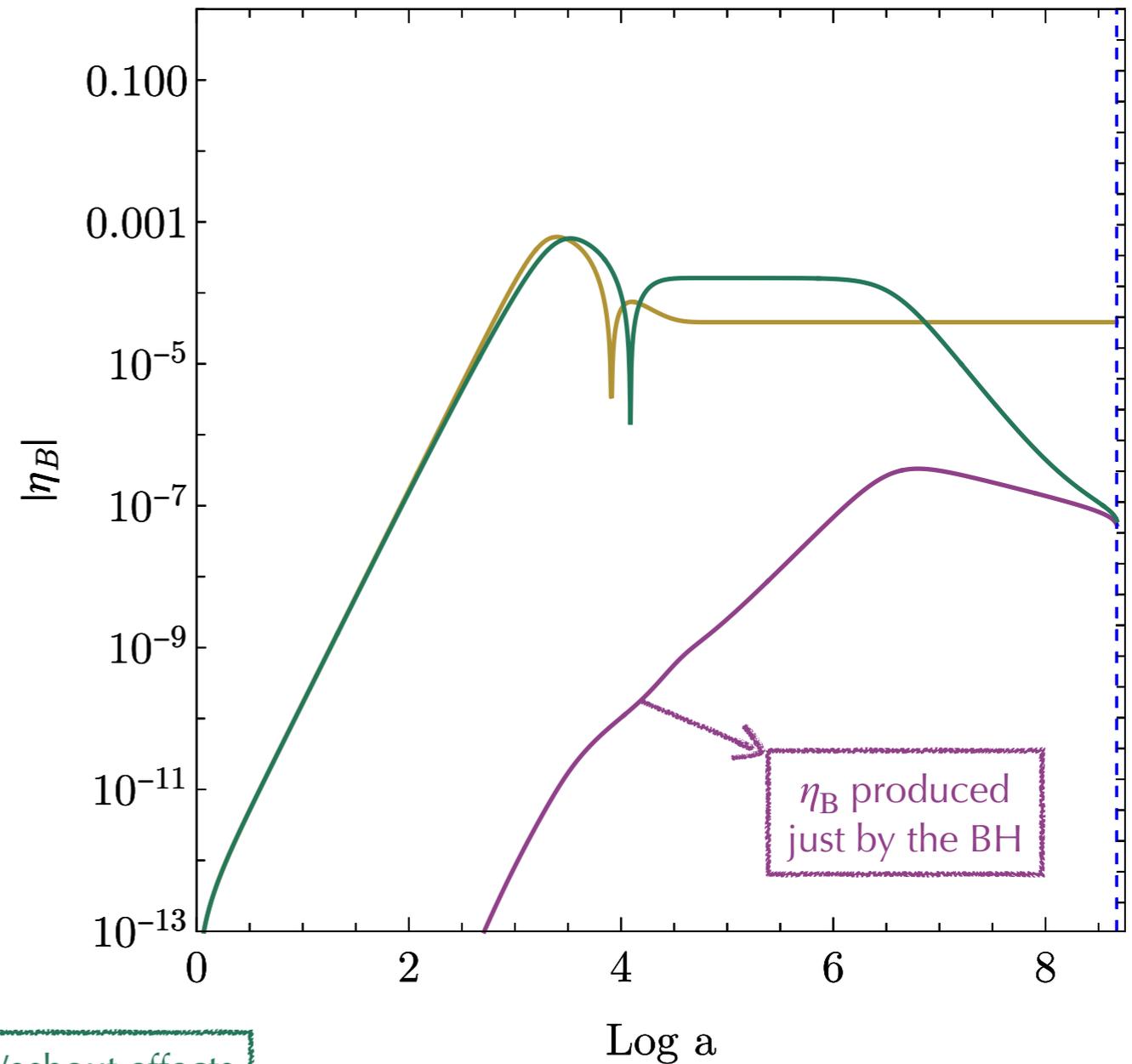
$$\beta'_i = 10^{-3}$$

How after is "after"?

RH neutrino abundance



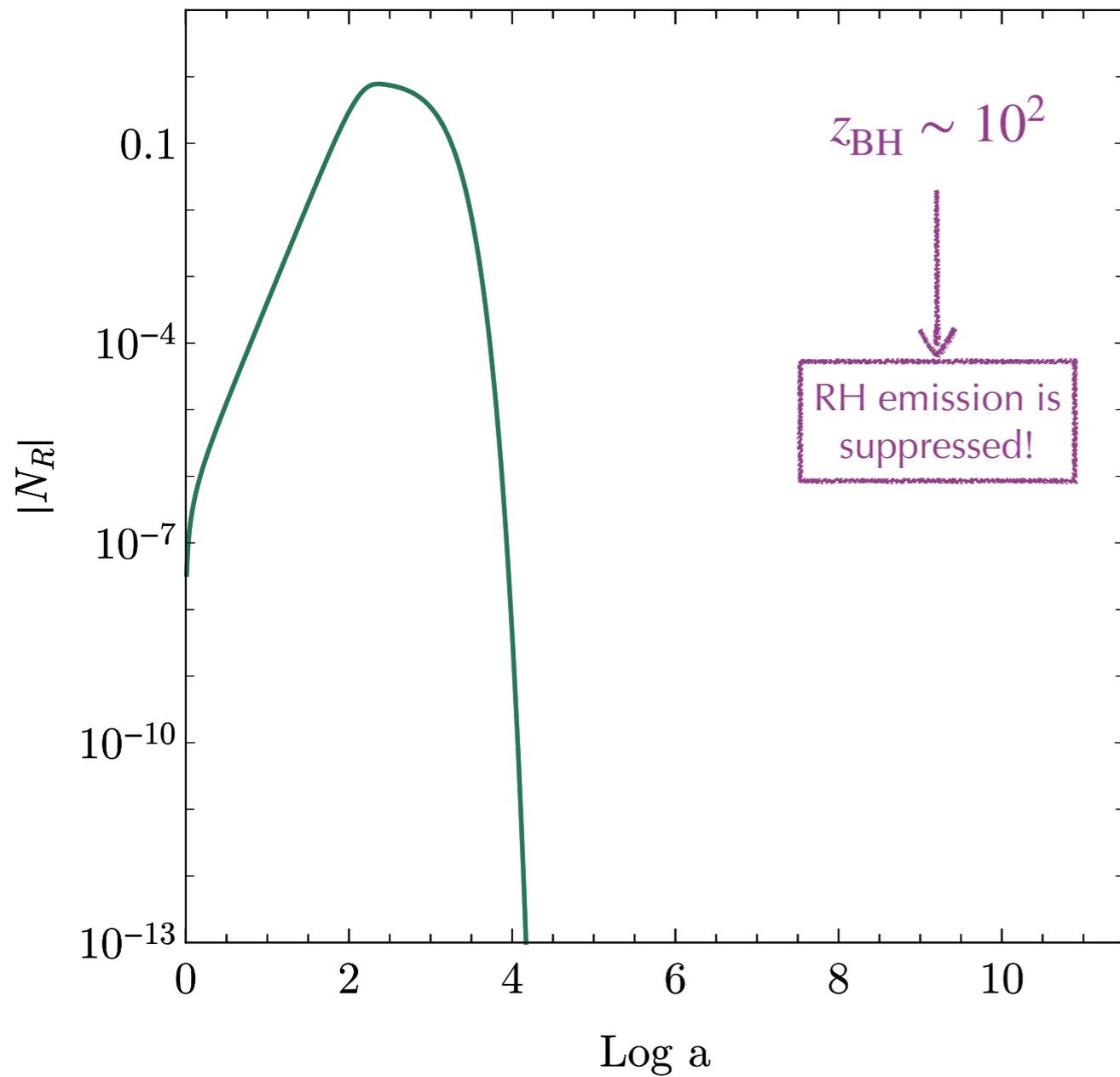
Final baryon asymmetry



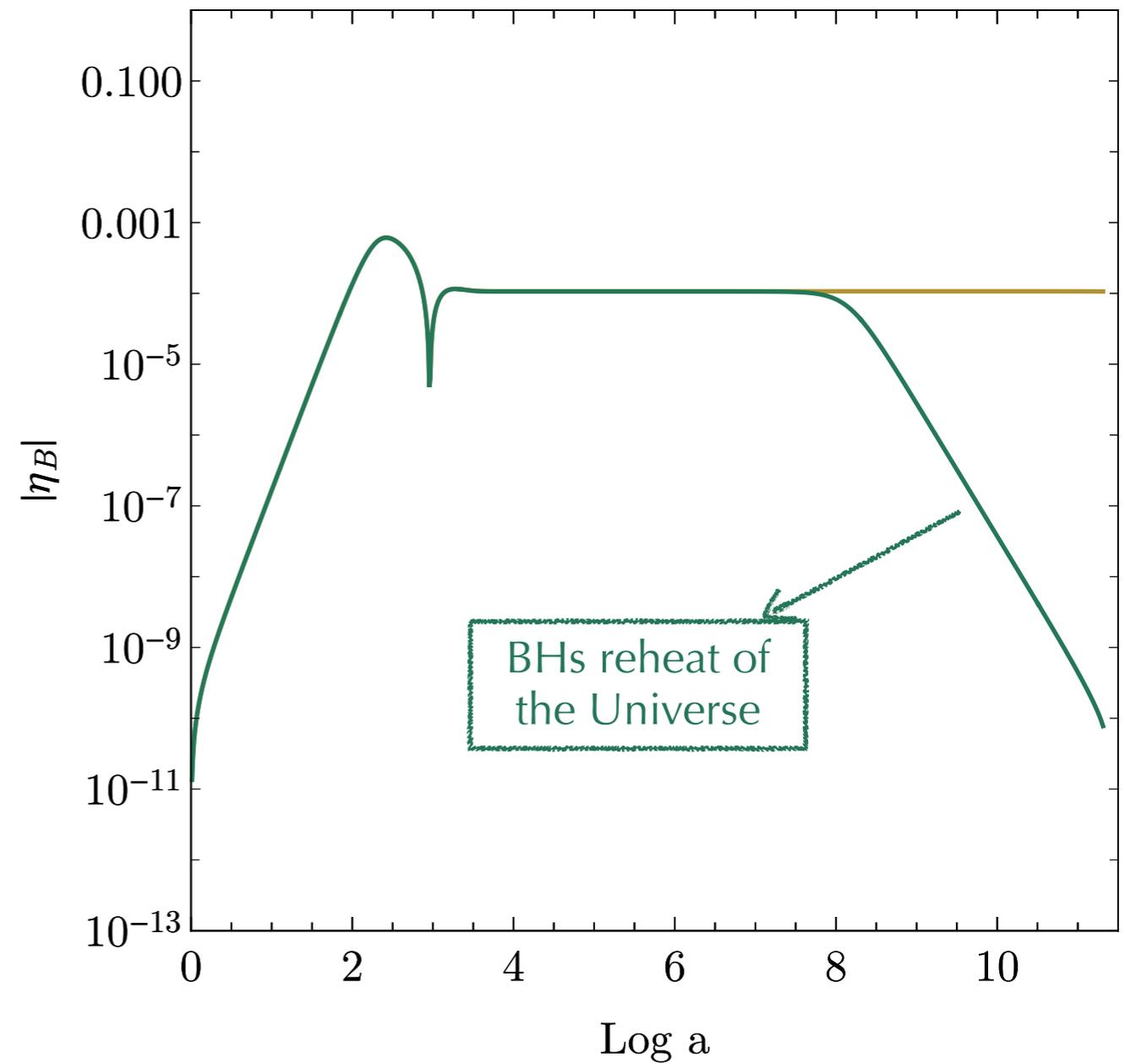
C. PBH create a RH density after the thermal case

$$M_i = 10^4 \text{ g}$$
$$\beta'_i = 10^{-3}$$

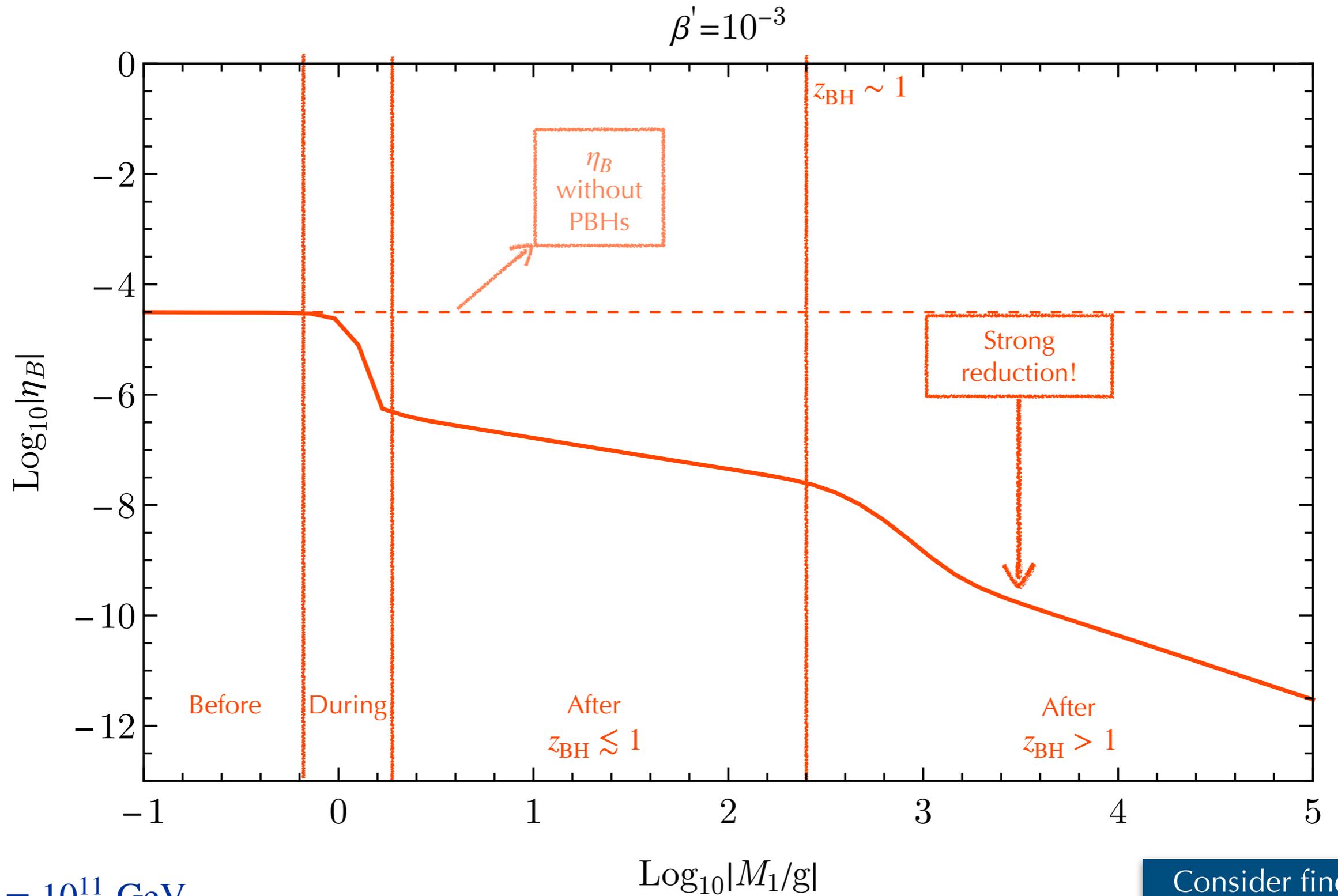
RH neutrino abundance



Final baryon asymmetry



$$Y_\nu = \frac{1}{v} U_{\text{PMNS}} \sqrt{m_\nu} R^T M_N^T$$



$$M_{N_1} = 10^{11} \text{ GeV}$$

Consider fine-tuned Yukawas

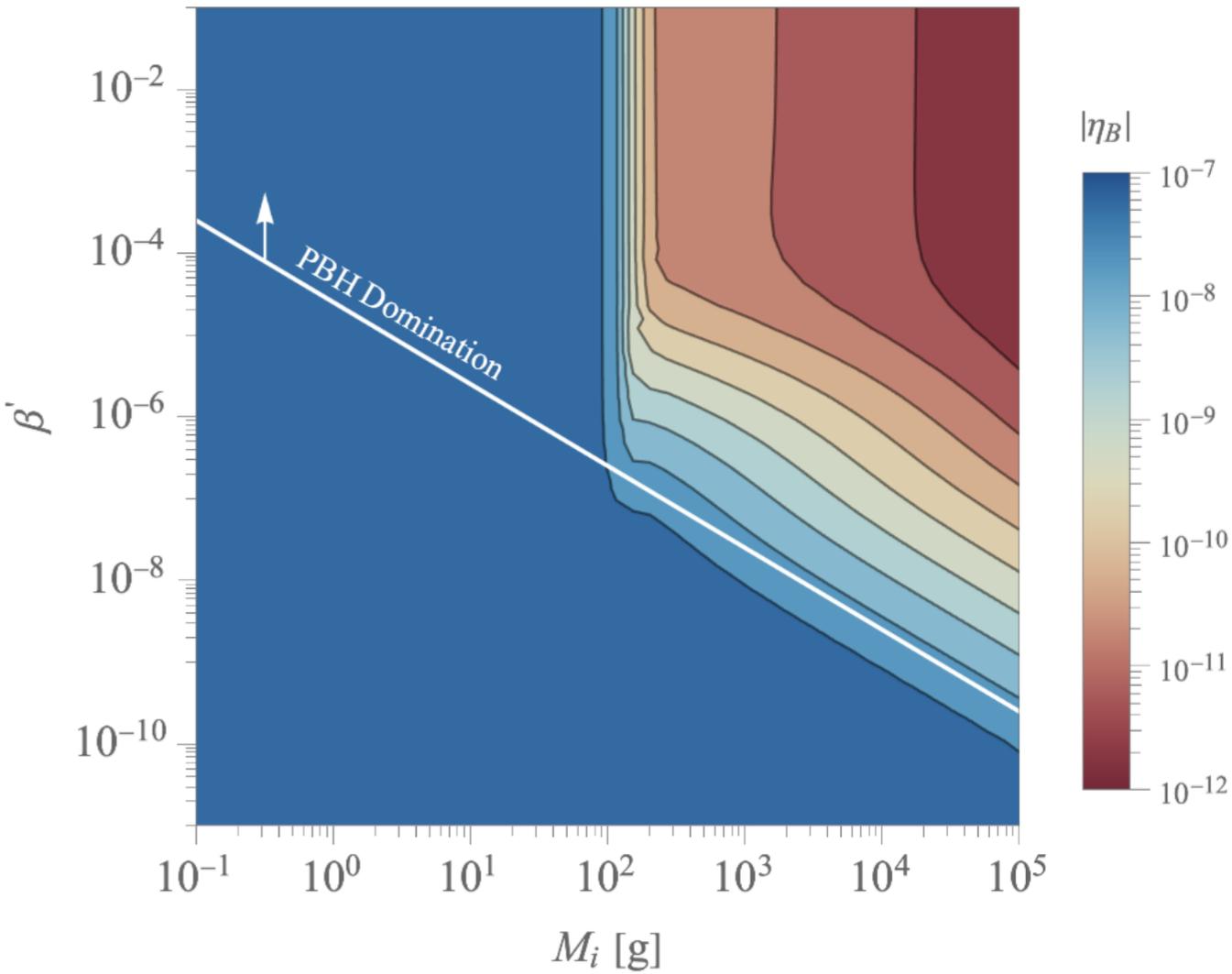
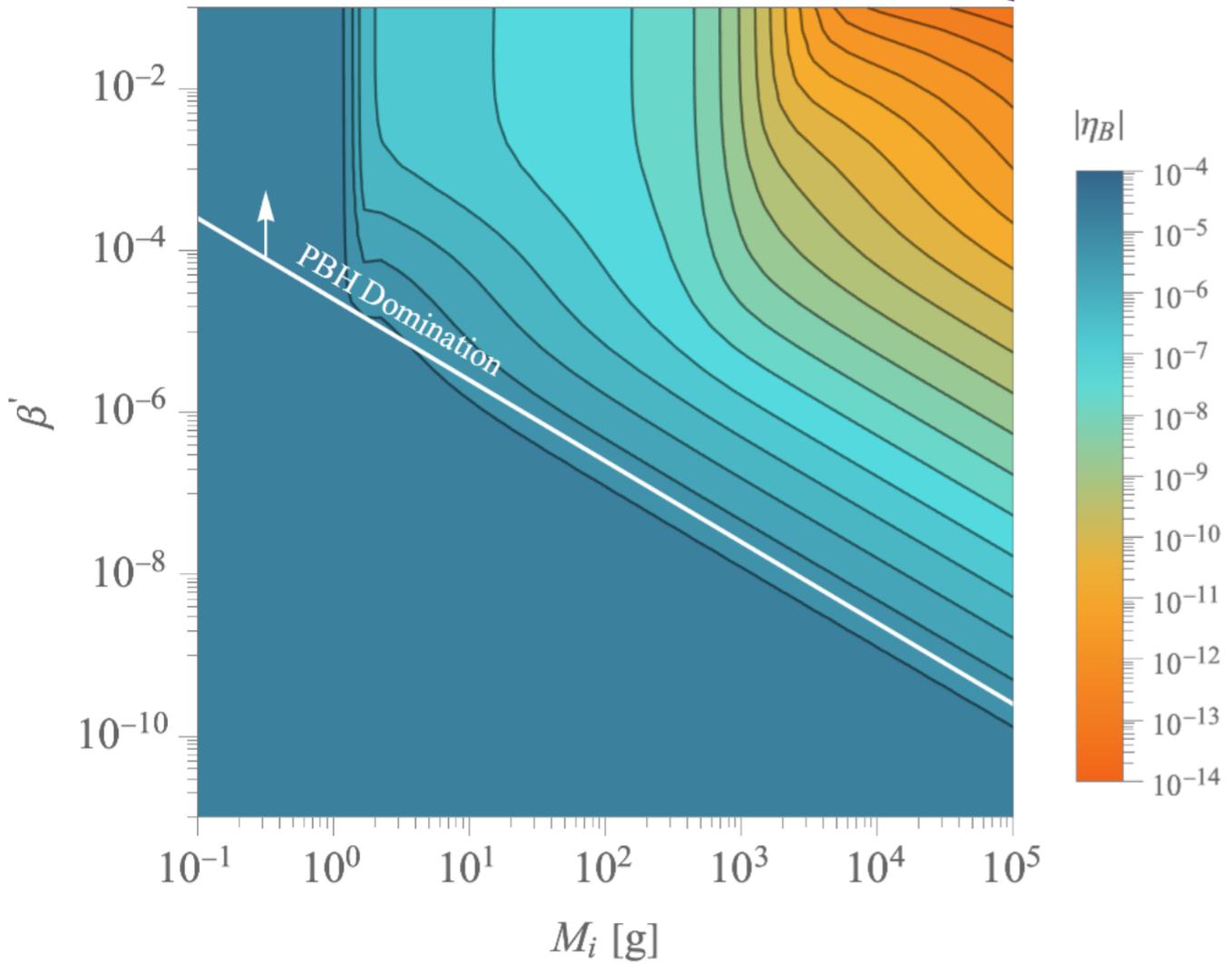
# Dependence on $\beta'$

How the depletion depends on the initial fraction?

$$T_{ev} \gtrsim 150 \text{ GeV}^*$$

$M_{N_1} = 10^{11} \text{ GeV}$

$M_{N_1} = 10^8 \text{ GeV}$

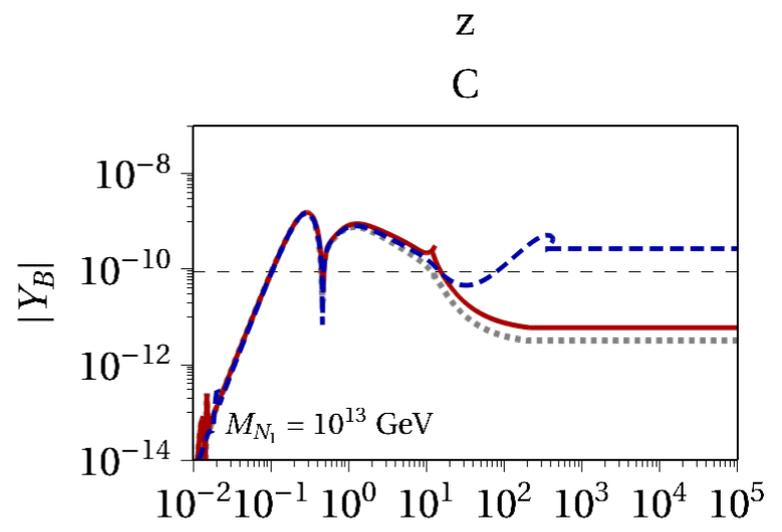
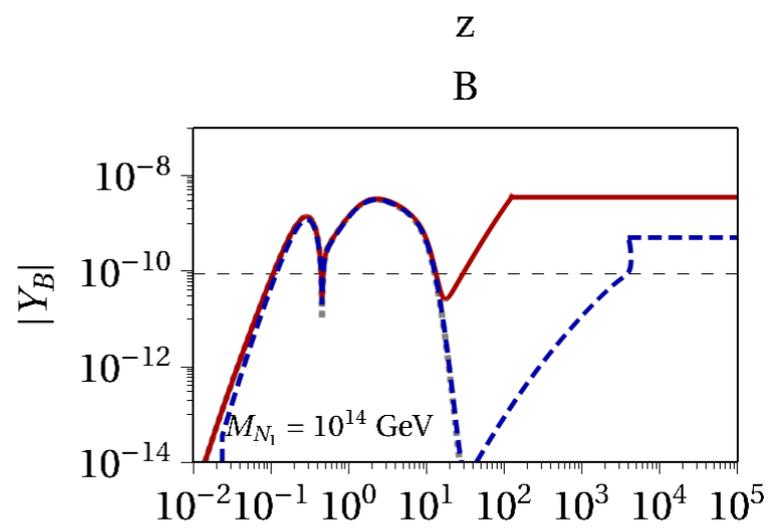
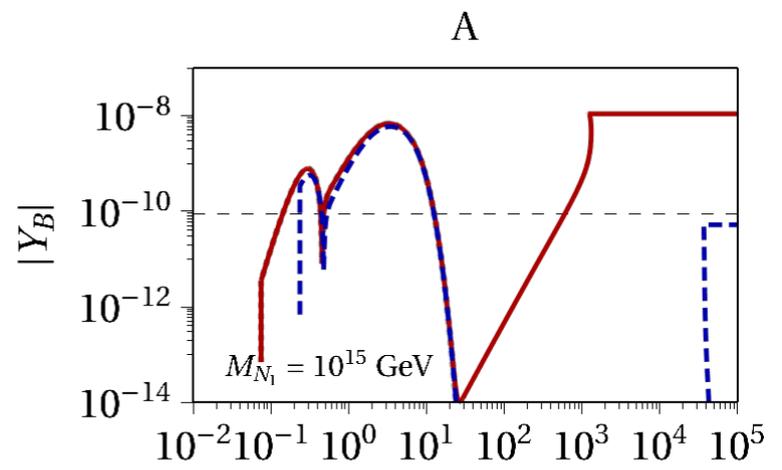
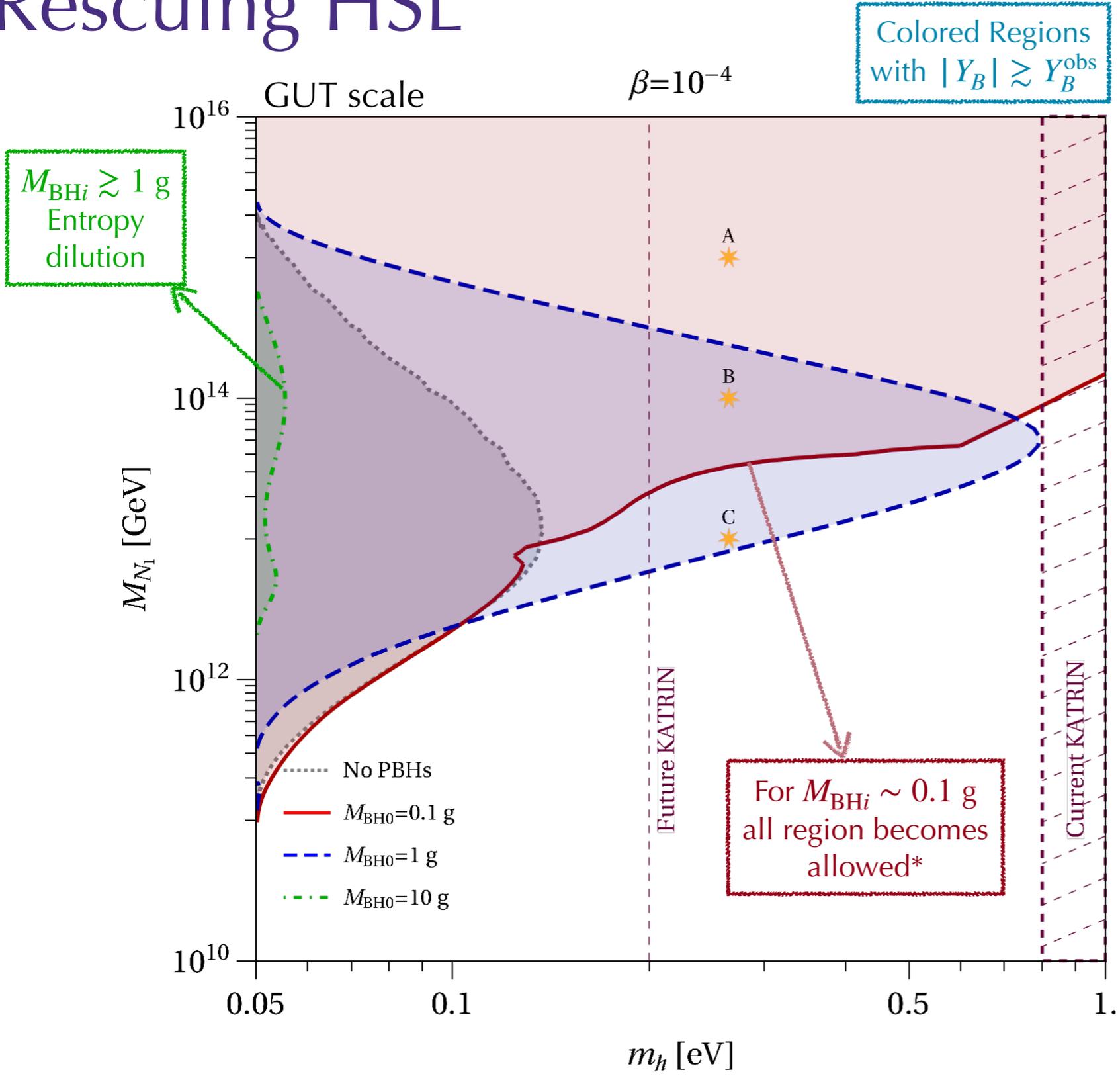


Entropy dilution is a generic and important feature to be taken into account

\*Sphalerons in the SM plasma

Sphalerons active around PBHs  
 García-Bellido et al, 1904.11482  
 De Luca et.al., 2102.07408

# Rescuing HSL



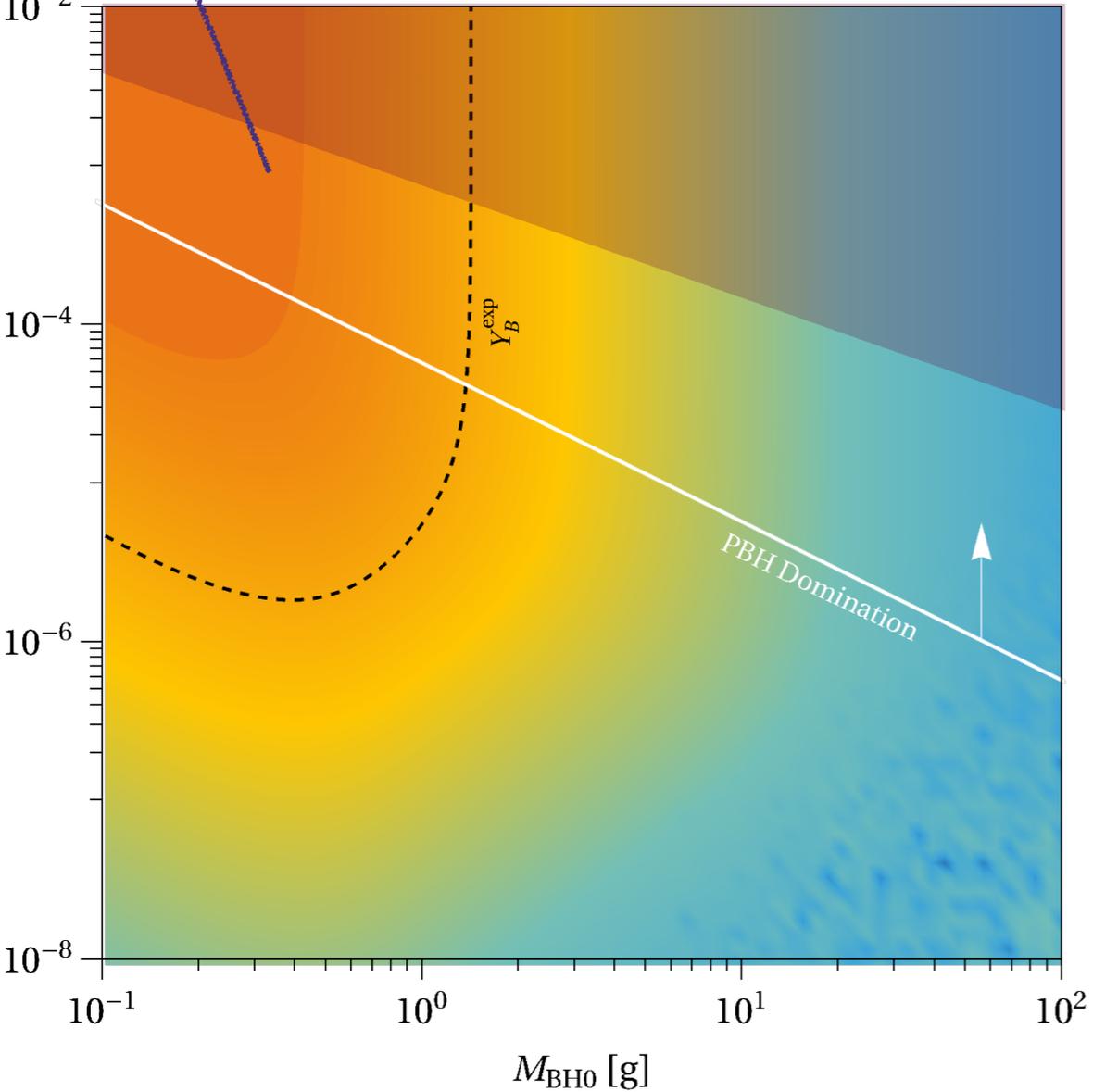
Maximizing over Yukawa parameters

\*Up to perturbativity

# Rescuing HSL

$\mathcal{O}(10^6)$   
enhancement

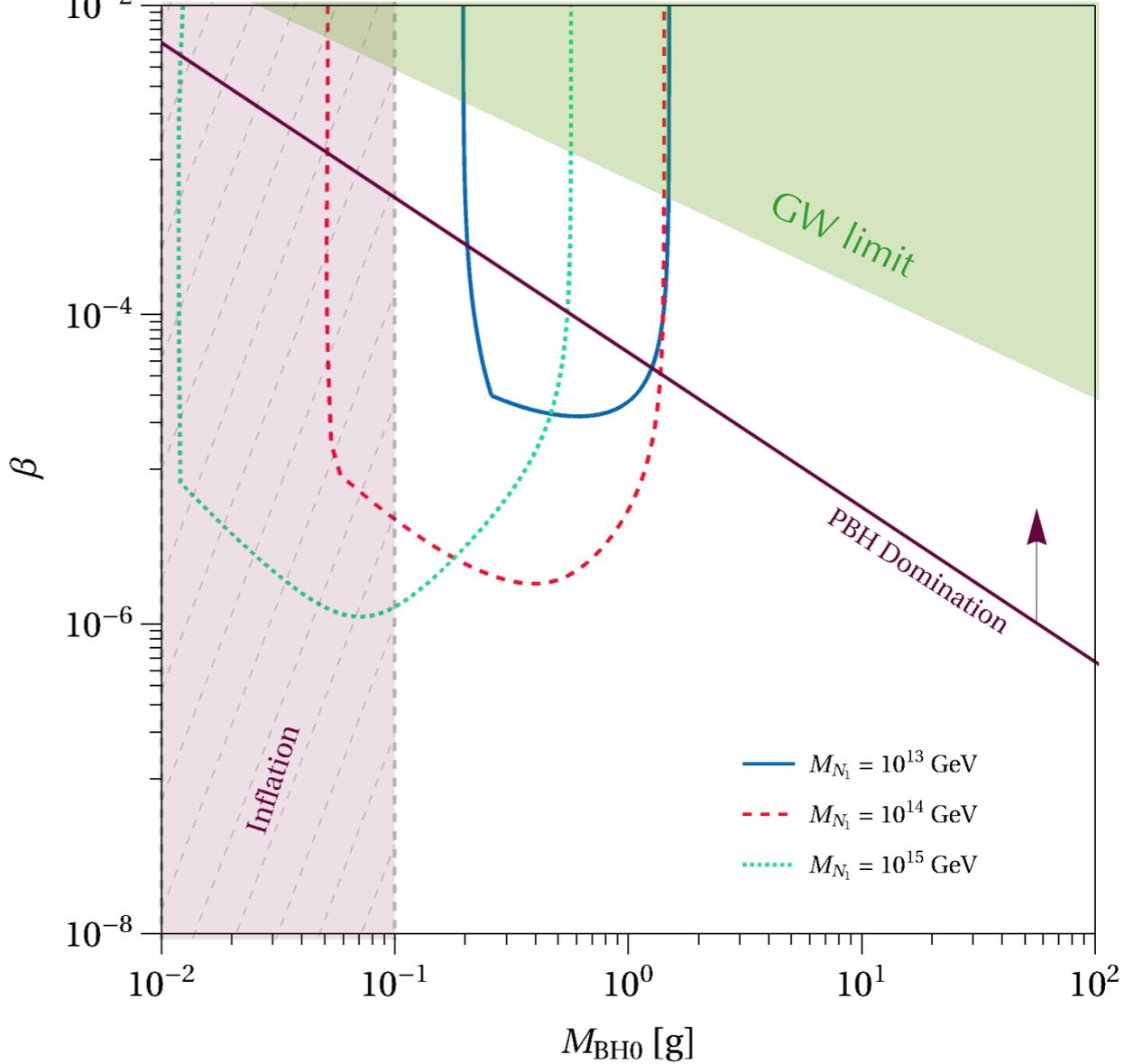
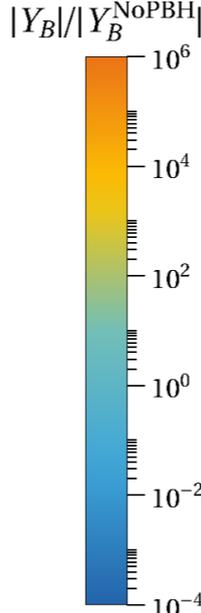
$M_{N_1} = 10^{14}$  GeV,  $m_h = 0.27$  eV



Washout process still active during evaporation

Efficient Production

PBH+Lepto,  $m_h = 0.27$  eV



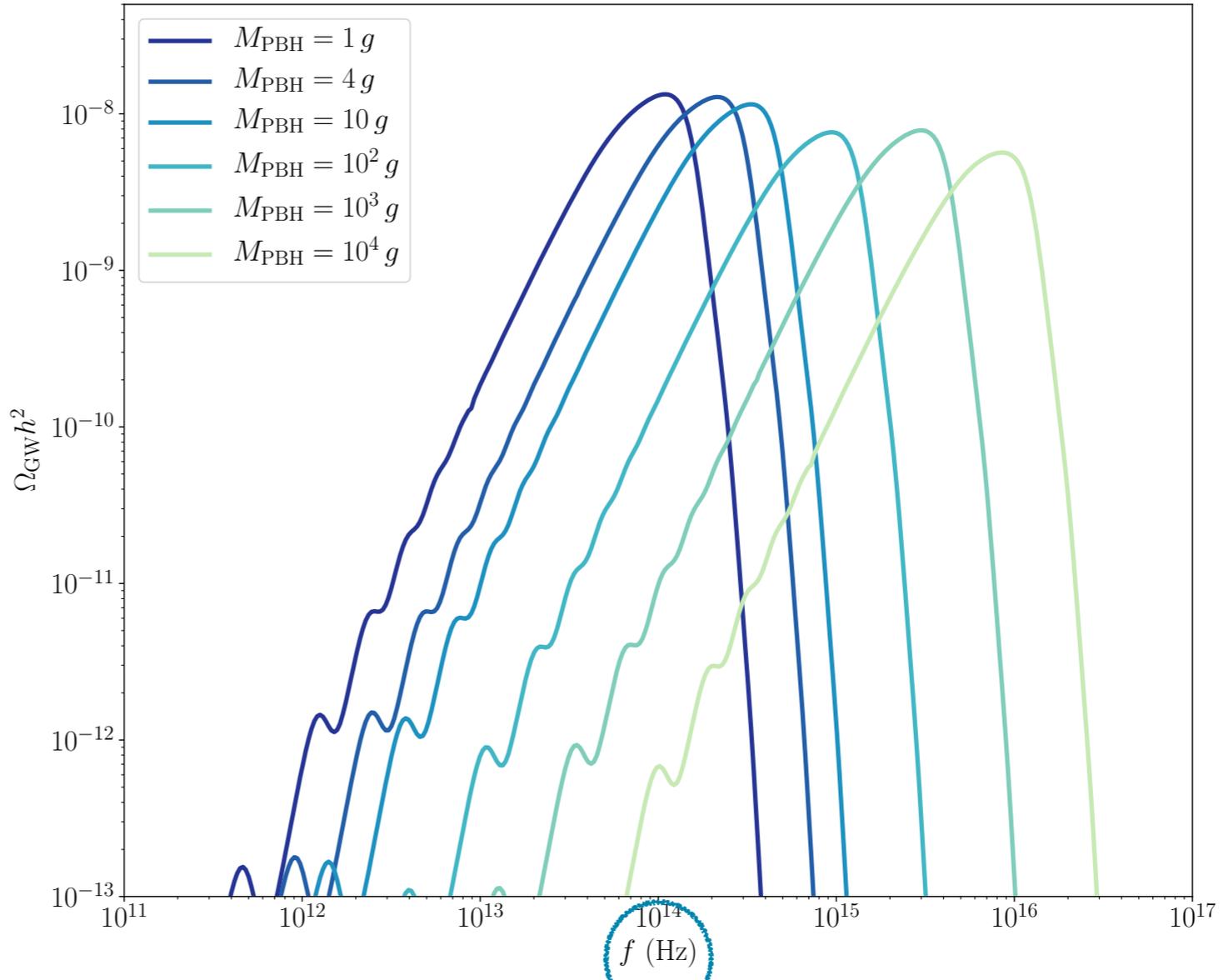
Testing these scenarios?

# Gravitational Waves?

How to prove that the early Universe had a PBH-dominated era?

Example: GW produced by the evaporation

- Arvanitaki, Geraci, 2013
- Ito et al, 2020
- Chen et al, 2020

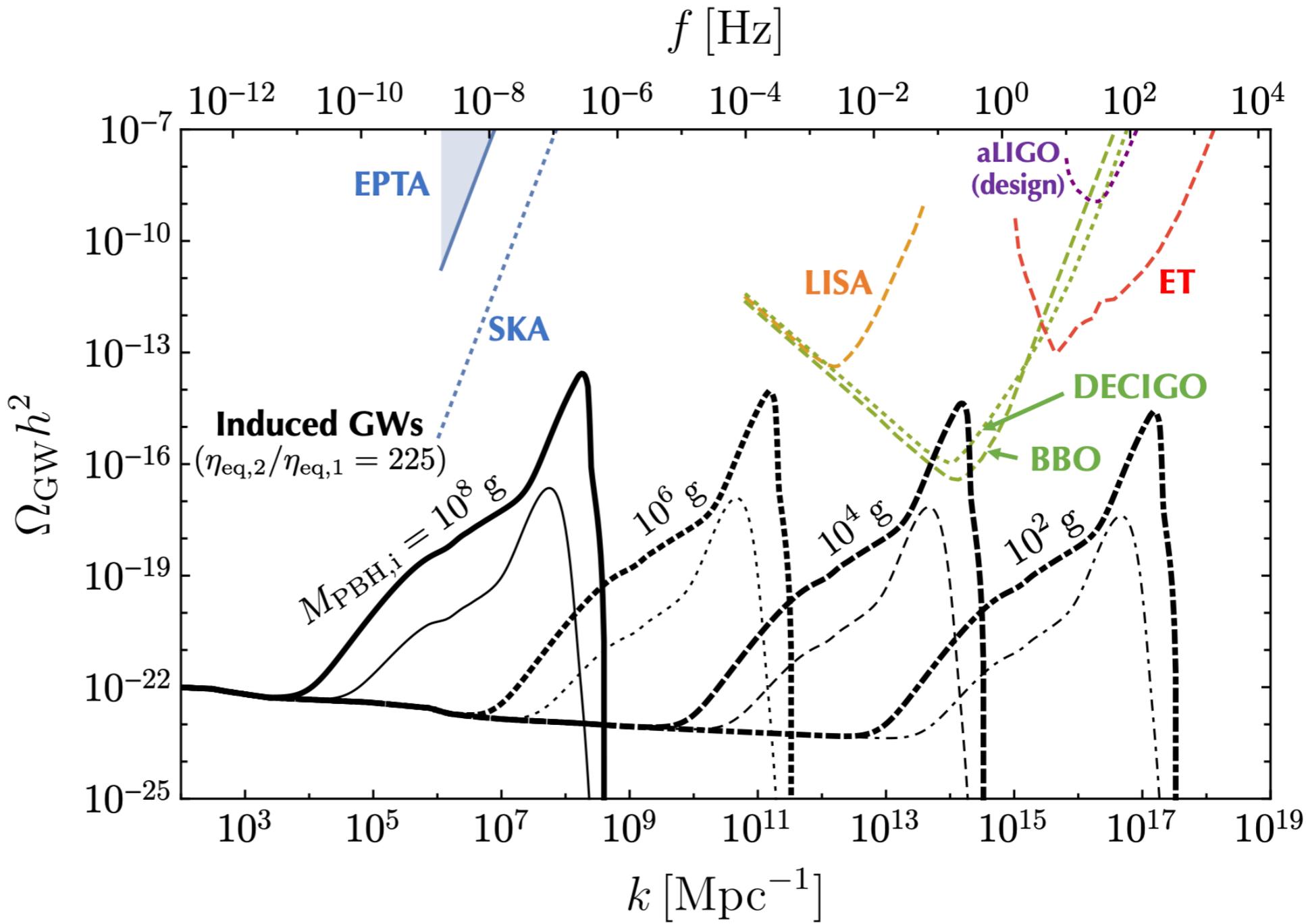


Huge frequency!

# Gravitational Waves?

How to prove that the early Universe had a PBH-dominated era?

Poltergeist Mechanism



Sudden transition from a matter to radiation dominated U enhances the GWs

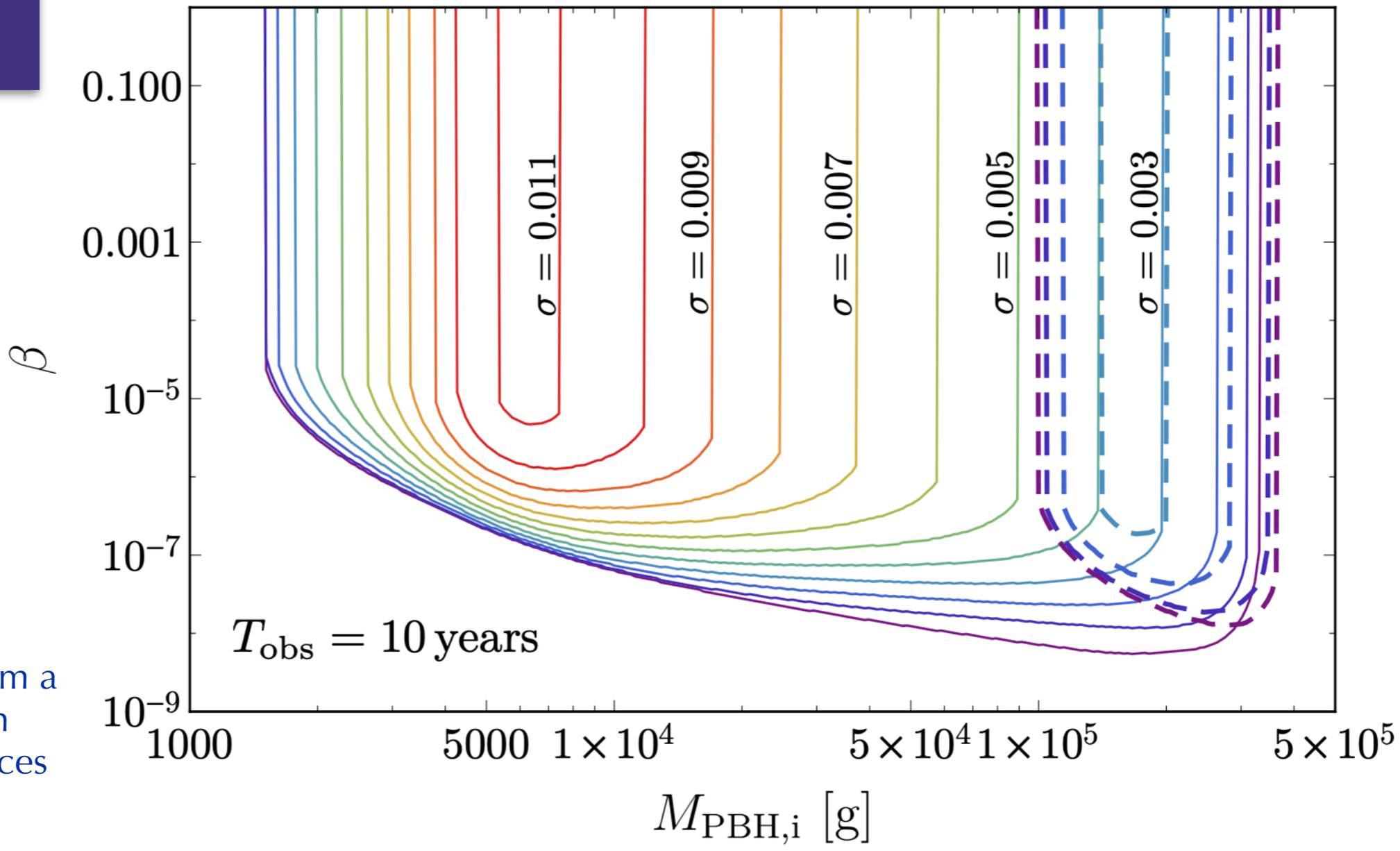
Inomata et al, 2003.10455

See also: 2205.06260

# Gravitational Waves?

How to prove that the early Universe had a PBH-dominated era?

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Inomata et al, 2003.10455

See also: 2205.06260

Falsifying Intermediate scale Leptogenesis?

DECILIGO = Full  
LISA = Dashed

# In progress

Based on:  
Cheek, Heurtier, YFPG, Turner. To appear

# Mass Distributions

Monochromatic approximation *too* approximated?

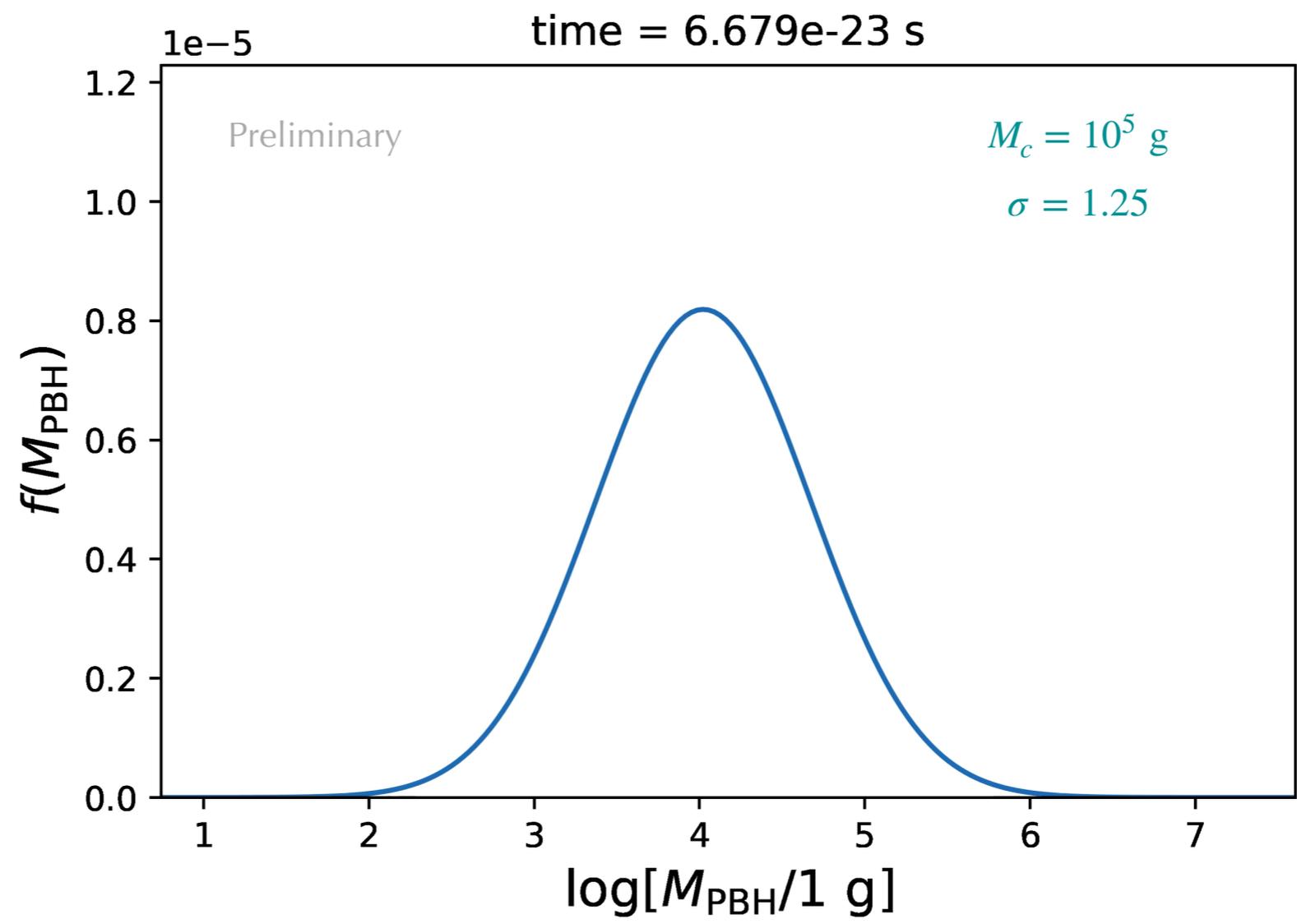
Connection with different formation mechanisms?

Having PBHs with different masses could have a distinct impact on the previous results

$$n_{\text{PBH}} = \int dM f(M) \quad f(M) = \frac{n_{\text{BH}}}{\sqrt{2\pi}\sigma M} \exp\left(-\frac{\log^2(M/M_c)}{2\sigma^2}\right)$$

Log-normal distribution

Dolgov, 93  
Green, 2016  
Kannike, 2017



We aim to make our code public!

# Mass Distributions

Connection with different formation mechanisms?

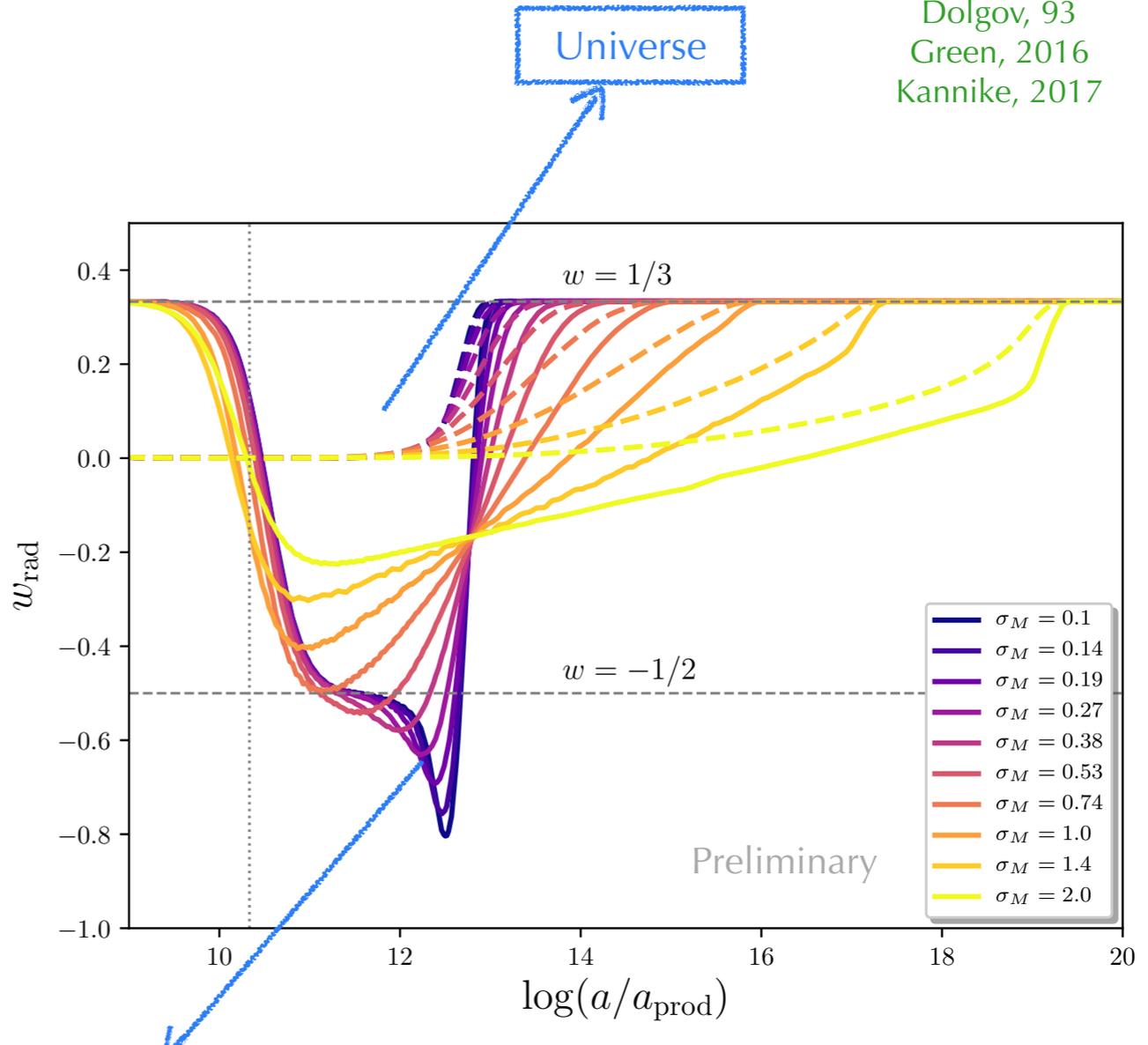
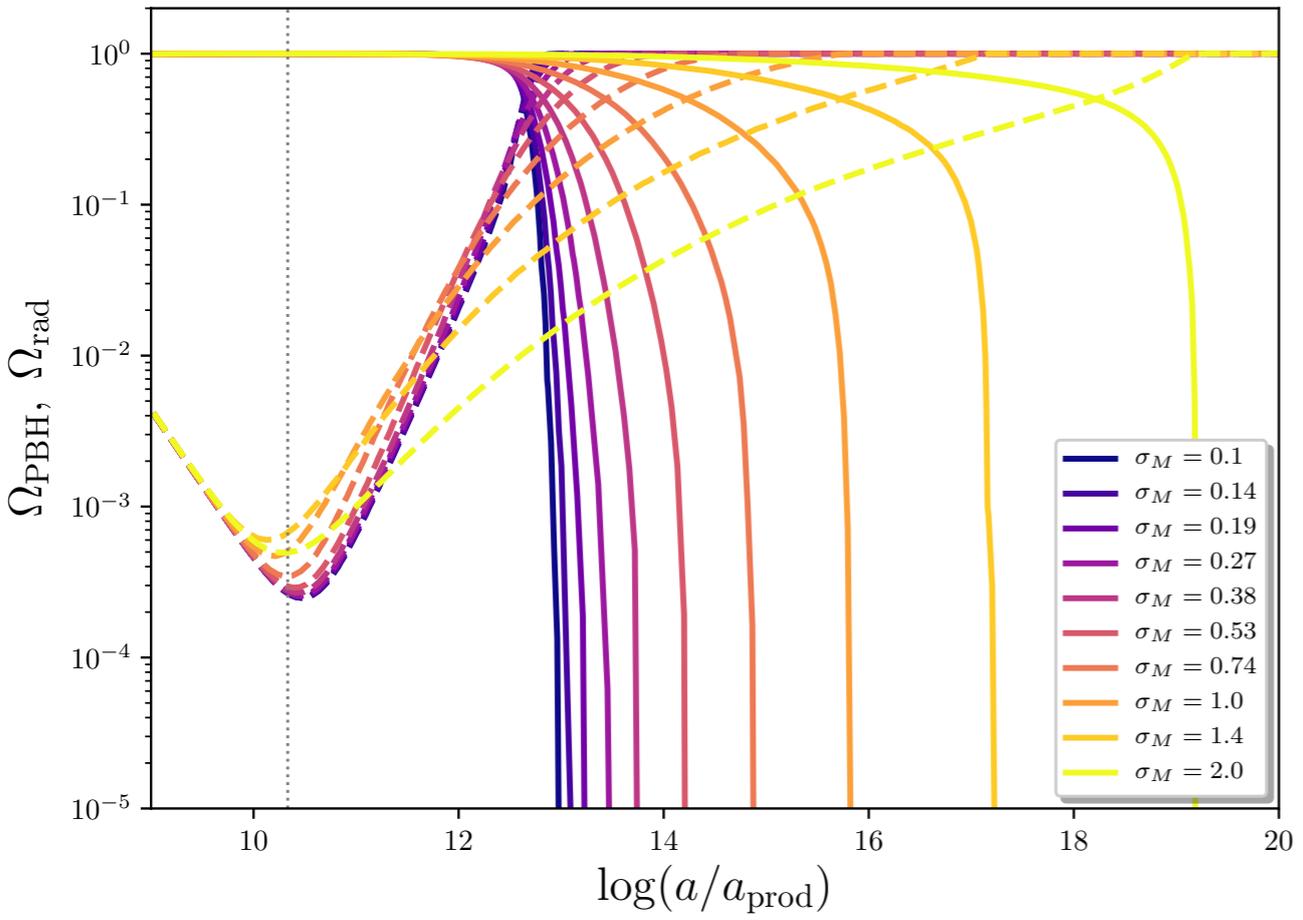
$$n_{\text{PBH}} = \int dM f(M)$$

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Log-normal distribution

Dolgov, 93  
Green, 2016  
Kannike, 2017

$m_{\text{DM}} = 10 \text{ GeV}, M_{\text{PBH}}^c = 10^4 \text{ g}, a_c = 0.1, \sigma_a = 0.015$



Universe

SM rad

Equation of state changes significantly

$$p = \omega\rho$$

$$\rho \propto a^{-3(1+\omega)}$$

Preliminary

# Mass Distributions

Monochromatic approximation *too* approximated?

Connection with different formation mechanisms?

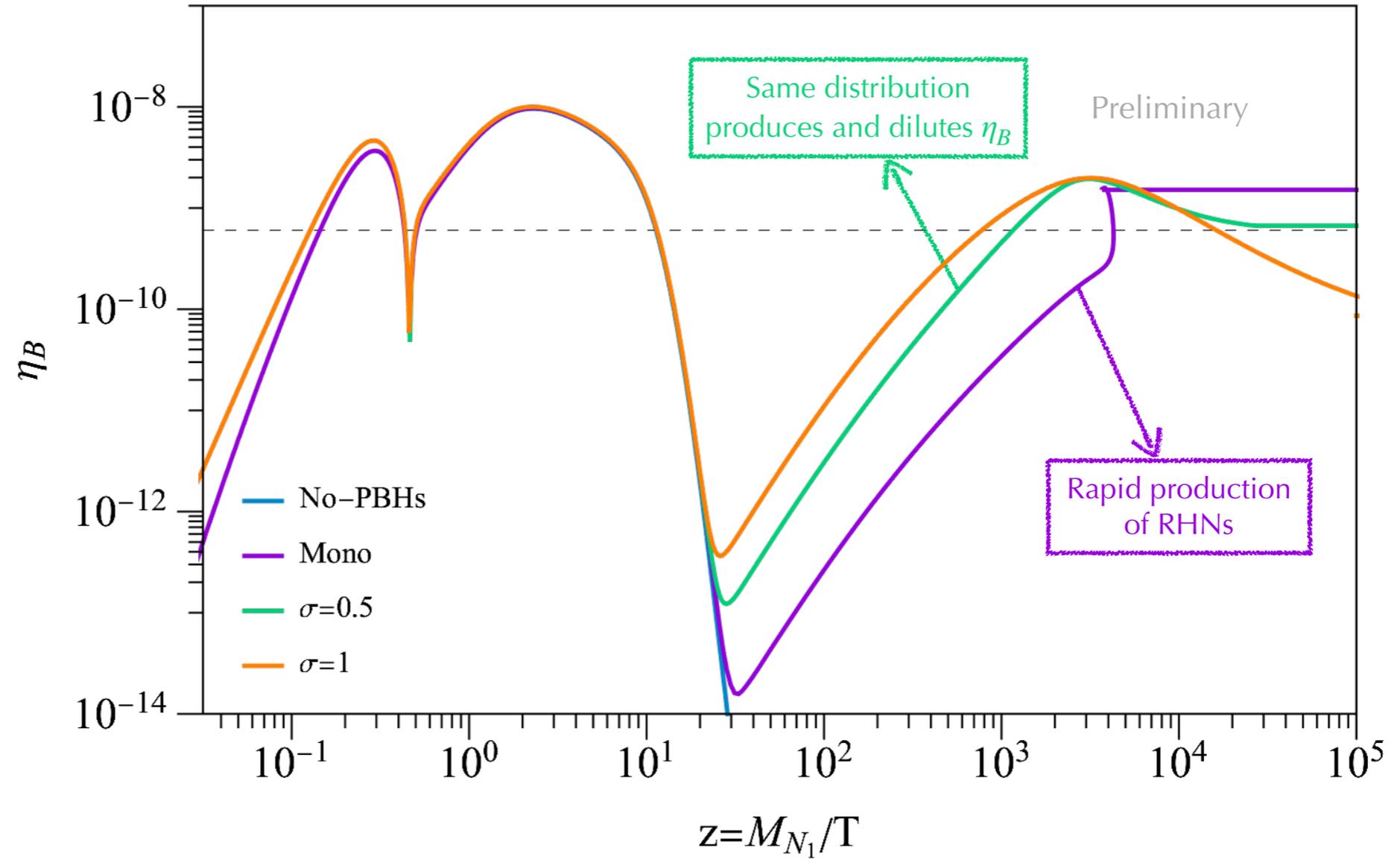
Having PBHs with different masses could have a distinct impact on the previous results

$$n_{\text{PBH}} = \int dM f(M) \quad f(M) = \frac{n_{\text{BH}}}{\sqrt{2\pi\sigma M}} \exp\left(-\frac{\log^2(M/M_c)}{2\sigma^2}\right)$$

$$M_c = 1 \text{ g}$$

Log-normal distribution

Dolgov, 93  
Green, 2016  
Kannike, 2017



We aim to make our code public!

# Conclusions

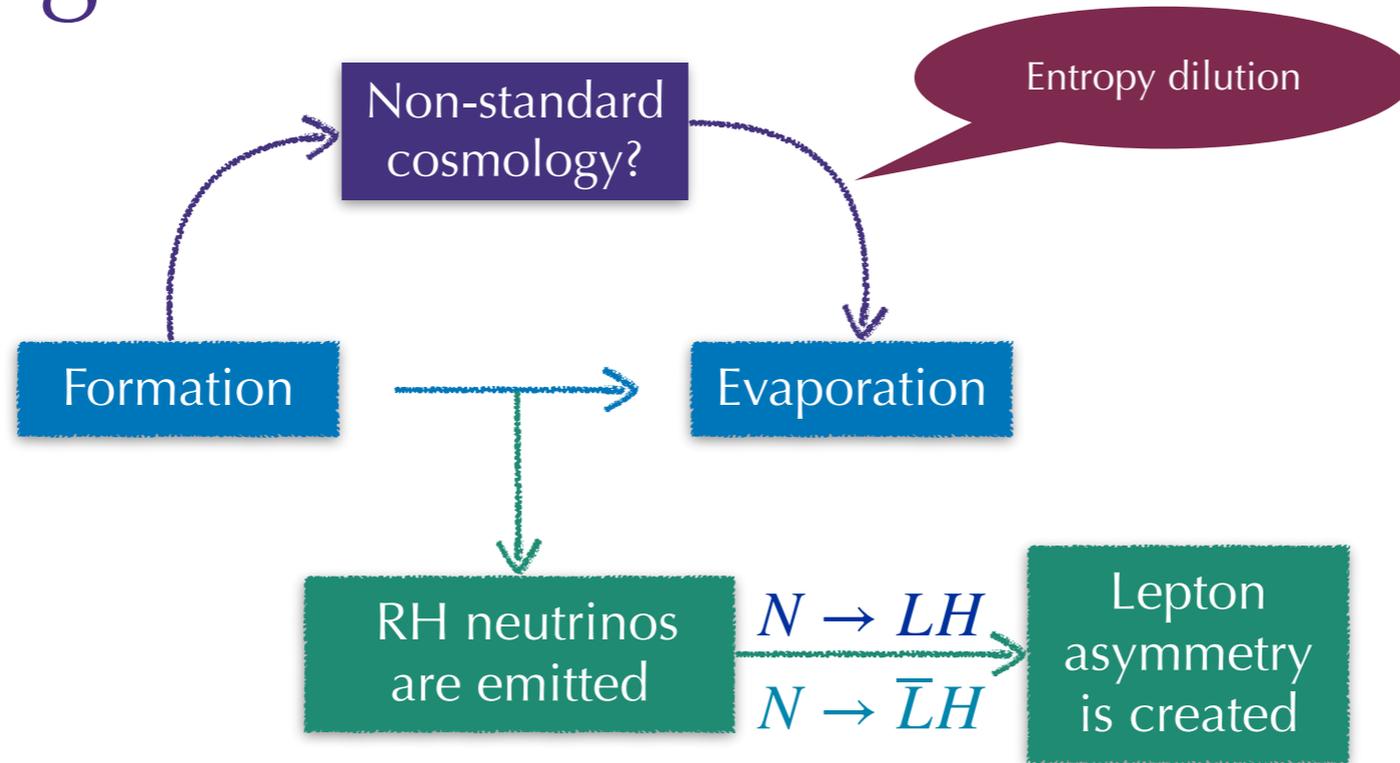
- ❖ PBH evaporation offers a unique mechanism to produce particles in the Early Universe
- ❖ The effects are threefold:
  - Universal particle emission
  - Modifying the Cosmological Background
  - Entropy dilution
- ❖ We explored the effects on leptogenesis assuming the existence of a PBH population
- ❖ Future directions:
  - Relating to “more realistic” PBH formation mechanisms (connected to PBH-DM?)
  - Low scale leptogenesis? Sphalerons around PBHs after EWPT?
  - Kerr PBH  $\longrightarrow$  Additional interesting properties!

Thank you for your attention!

# Backup

$$z_{\text{BH}} \equiv \frac{M_{N_1}}{T_{\text{BH}}}$$

# PBH + Leptogenesis



## Boltzmann Equations

$$aH \frac{dn_{N_1}^{\text{TH}}}{da} = - (n_{N_1}^{\text{TH}} - n_{N_1}^{\text{eq}}) \Gamma_{N_1}^T$$

Thermal contribution

$$aH \frac{dn_{N_1}^{\text{BH}}}{da} = - n_{N_1}^{\text{BH}} \Gamma_{N_1}^{\text{BH}} + n_{\text{BH}} \Gamma_{\text{BH} \rightarrow N_1}$$

From the evaporation

$$aH \frac{dn_{\alpha\beta}^{\text{B-L}}}{da} = \epsilon_{\alpha\beta}^{(1)} \left[ (n_{N_1} - n_{N_1}^{\text{eq}}) \Gamma_{N_1}^T + n_{N_1}^{\text{BH}} \Gamma_{N_1}^{\text{BH}} \right] + \mathcal{W}_{\alpha\beta}$$

Thermally averaged RH neutrino decay

$\alpha, \beta$  → Flavor indexes

# Evaporation function

$$\begin{aligned}
 \varepsilon_N(M) = & 2f_1 + 4f_{1/2}^1 \left\{ \sum_{\ell=e,\mu,\tau} \exp \left[ -\frac{M}{\beta_{1/2} M_\ell} \right] + 3 \sum_q \exp \left[ -\frac{M}{\beta_{1/2} M_q} \right] \right\} \\
 & + 2 \eta_\nu^N f_{1/2}^0 \sum_{a=1,2,3} \exp \left[ -\frac{M}{\beta_{1/2} M_a} \right] \\
 & + 16f_1 \exp \left[ -\frac{M}{\beta_1 M_g} \right] \\
 & + 3f_1 \left\{ 2 \exp \left[ -\frac{M}{\beta_1 M_W} \right] + \exp \left[ -\frac{M}{\beta_1 M_Z} \right] \right\} + f_0 \exp \left[ -\frac{M}{\beta_0 M_H} \right]
 \end{aligned}$$

$$\beta_s = \begin{cases} 2.66 & \text{for } s = 0 \\ 4.53 & \text{for } s = \frac{1}{2} \\ 6.04 & \text{for } s = 1 \end{cases}$$

$$f_s = \begin{cases} 0.267 & \text{for } s = 0 \\ 0.060 & \text{for } s = 1 \\ 0.007 & \text{for } s = 2 \end{cases}$$

$$f_{1/2}^q = \begin{cases} 0.147 & \text{for } q = 0 \text{ (neutral)} \\ 0.142 & \text{for } q = 1 \text{ (charged)} \end{cases}$$

# Dirac vs Majorana

## Dirac neutrinos

$$\sigma_{\text{abs}}^{\nu}(+1/2) = \sigma_{\text{abs}}^{\nu}(-1/2)$$

Unruh, 1976

No helicity suppression

Production of light RH neutrinos!

Cecilia Lunardini, YFPG  
JCAP08(2020)014

## Majorana neutrinos

Heavy RH neutrinos



PBH-induced Leptogenesis

Yamada and Iso, 1610.02586  
Morrison et al, 1812.10606  
Baldes et al, 2004.14773  
Hooper and Krnjaic, 2010.01134

YFPG and Jessica Turner,  
PRD 104(2021) 103021  
Bernal, Fong, YFPG, Turner  
2203.08823

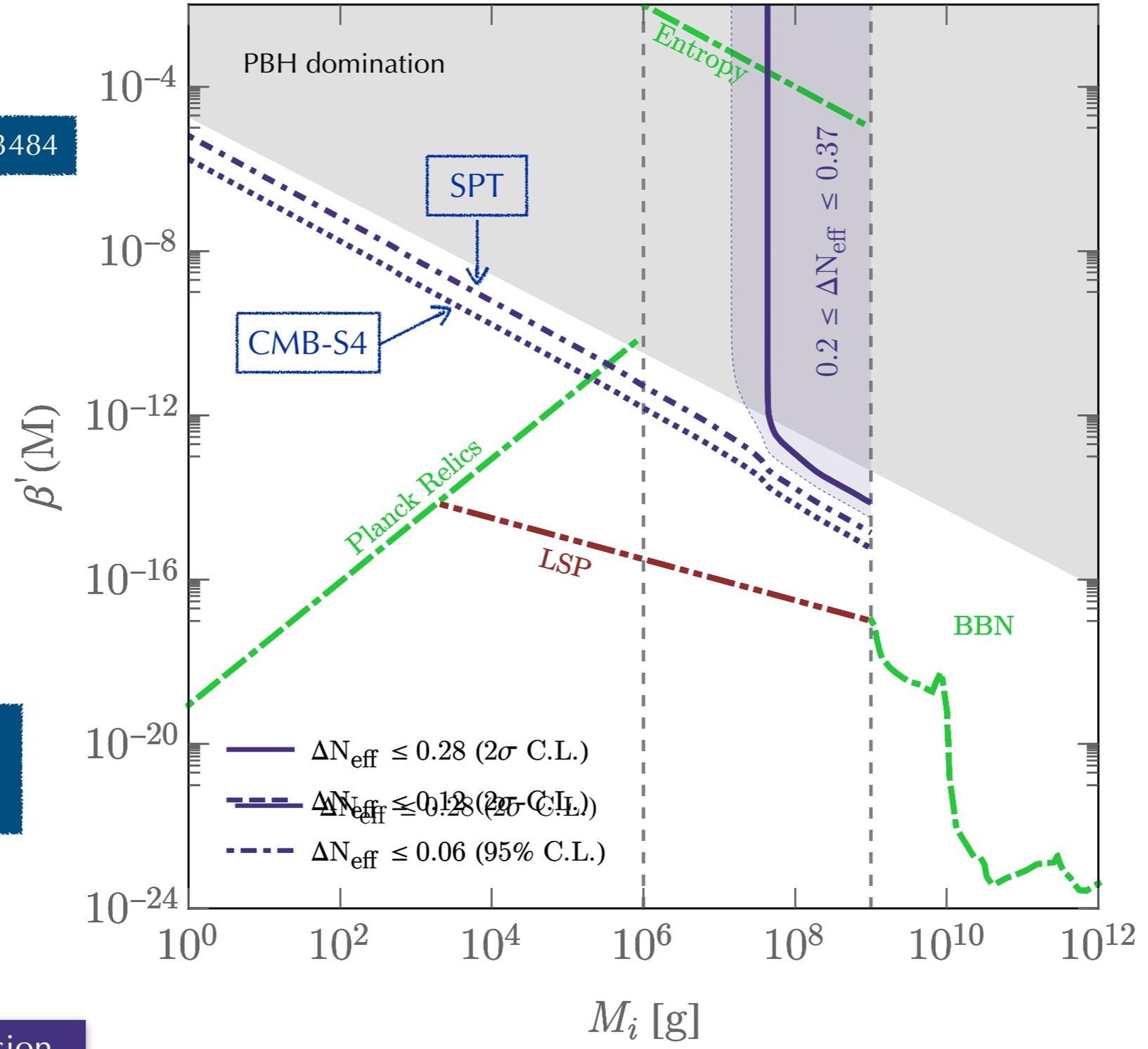
# Dirac Neutrinos

Green and Liddle, 9903484

Zel'dovich et al, 1977

MacGibbon, 1987  
Barrow et al, 1992  
Carr et al, 1994

Carr et al,  
0912.5297



Alleviate the  $H_0$  tension

What happens when  $M \rightarrow M_{Pl}$  ?

# Boltzmann Equations

$$aH \frac{dn_{N_1}^{\text{TH}}}{da} = - (n_{N_1}^{\text{TH}} - n_{N_1}^{\text{eq}}) \Gamma_{N_1}^T$$

Thermal contribution

$$aH \frac{dn_{N_1}^{\text{BH}}}{da} = - n_{N_1}^{\text{BH}} \Gamma_{N_1}^{\text{BH}} + n_{\text{BH}} \Gamma_{\text{BH} \rightarrow N_1}$$

RH neutrinos emitted from the evaporation

$$aH \frac{dn_{\alpha\beta}^{\text{B-L}}}{da} = \epsilon_{\alpha\beta}^{(1)} \left[ (n_{N_1} - n_{N_1}^{\text{eq}}) \Gamma_{N_1}^T + n_{N_1}^{\text{BH}} \Gamma_{N_1}^{\text{BH}} \right] + \mathcal{W}_{\alpha\beta}$$

Thermally averaged RH neutrino decay

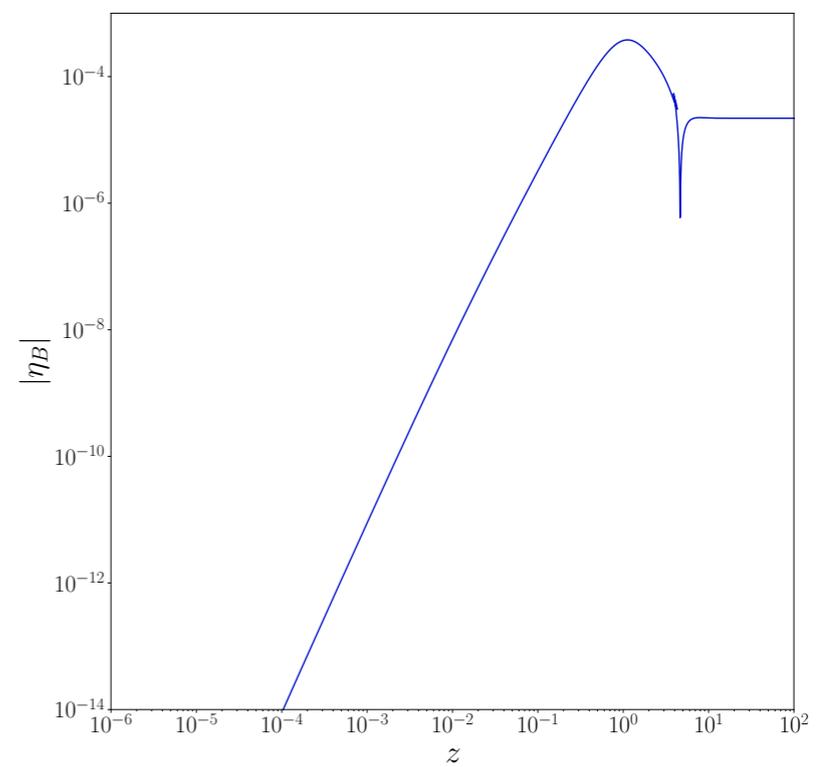
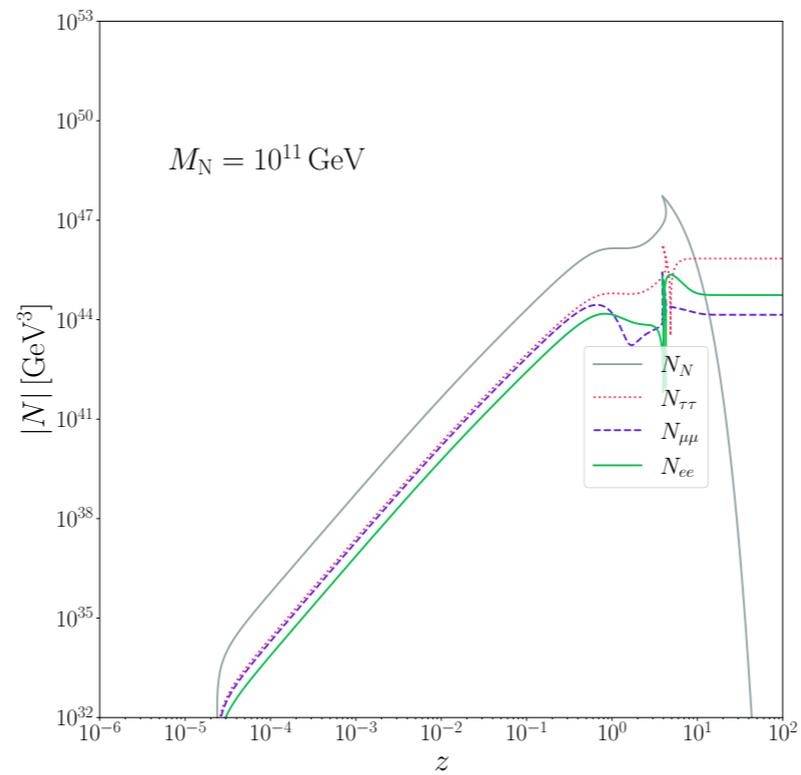
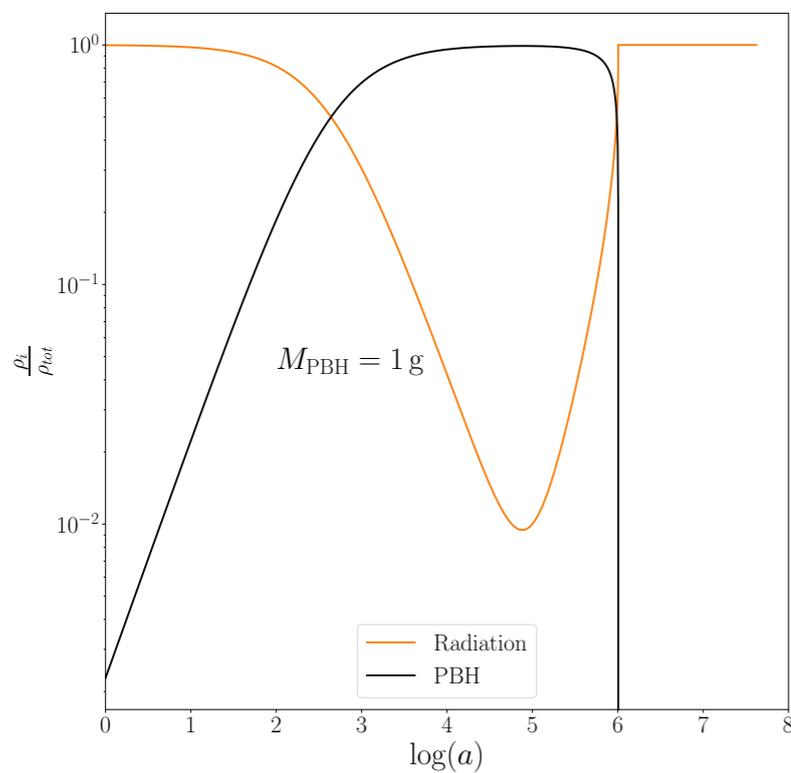
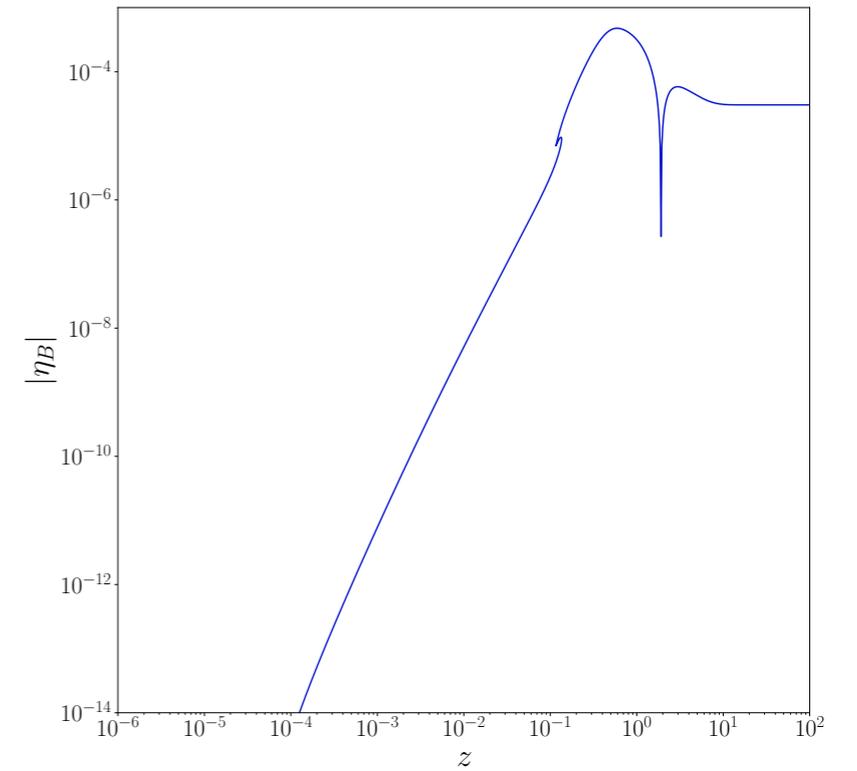
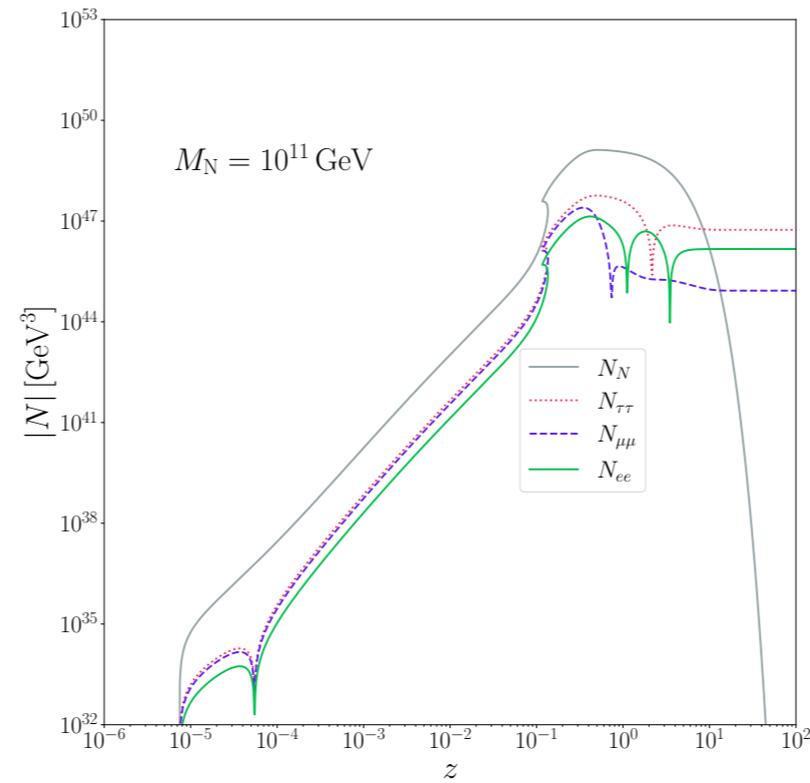
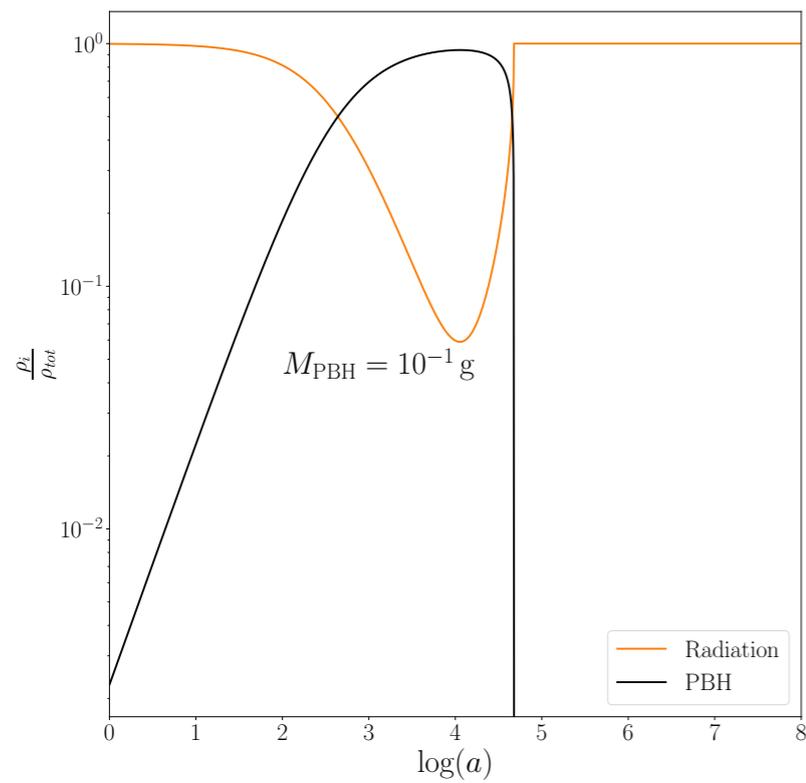
$\alpha, \beta$  → Flavor indexes

$$z_{\text{BH}} \equiv \frac{M_{N_1}}{T_{\text{BH}}}$$

$$\Gamma_{\text{BH} \rightarrow N_1} \equiv \int_0^\infty \frac{d^2 \mathcal{N}_{N_1}}{dp dt} dp$$

$$\approx \frac{27 T_{\text{BH}}}{32 \pi^2} \left( -z_{\text{BH}} \text{Li}_2(-e^{-z_{\text{BH}}}) - \text{Li}_3(-e^{-z_{\text{BH}}}) \right)$$

$$\Gamma_{N_1}^{\text{BH}} \equiv \left\langle \frac{M_{N_1}}{E_{N_1}} \right\rangle_{\text{BH}} \Gamma_{N_1}^0 \approx \frac{\mathcal{K}_1(z_{\text{BH}})}{\mathcal{K}_2(z_{\text{BH}})} \Gamma_{N_1}^0$$



# Gravitational Waves?

How to prove that the early Universe had a PBH-dominated era?

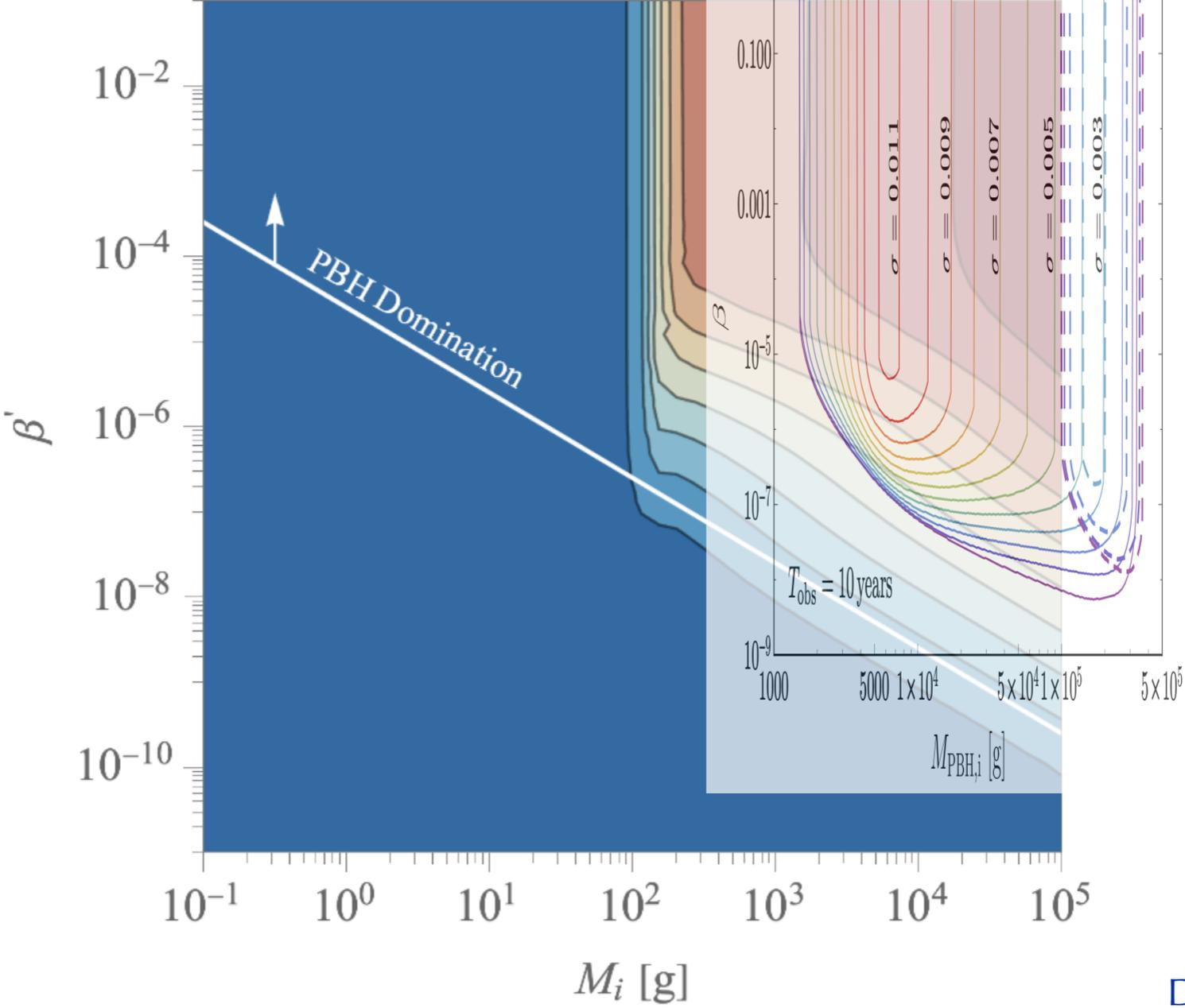
Poltergeist Mechanism

$$M_{N_1} = 10^8 \text{ GeV}$$

Sudden transition from a matter to radiation dominated U enhances the GWs

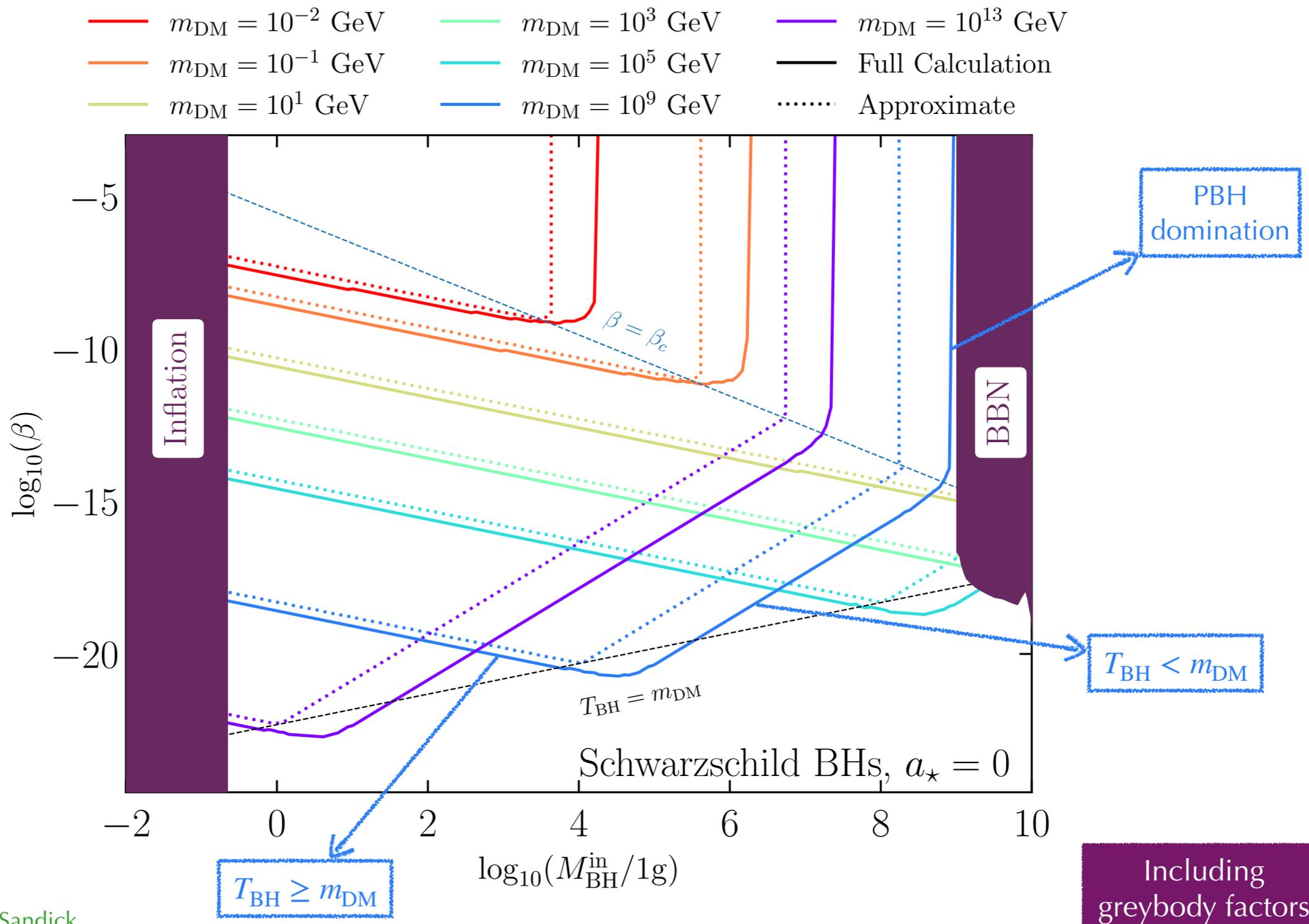
Inomata et al, 2003.10455

See also: 2205.06260



DECILIGO = Full  
LISA = Dashed

# Only Gravitational Interacting DM



Dashed from Gondolo, Sandick, Shams Es Haghi 2009.02424

# Freeze-Out

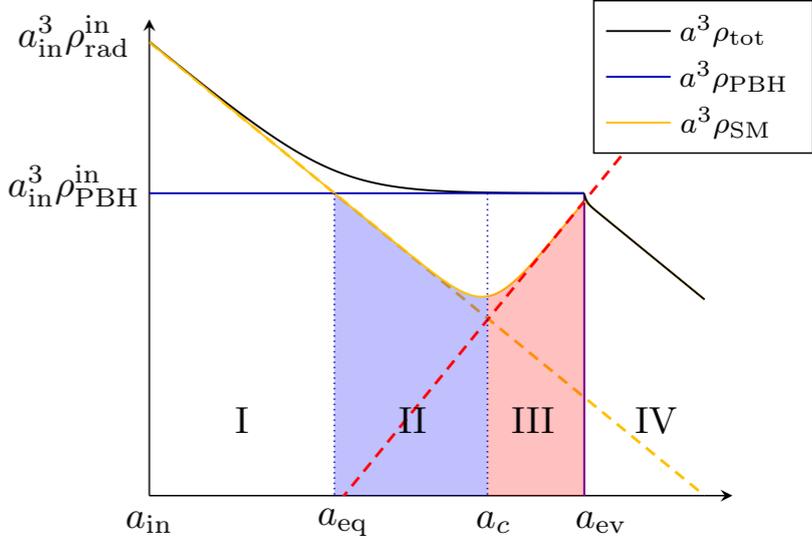
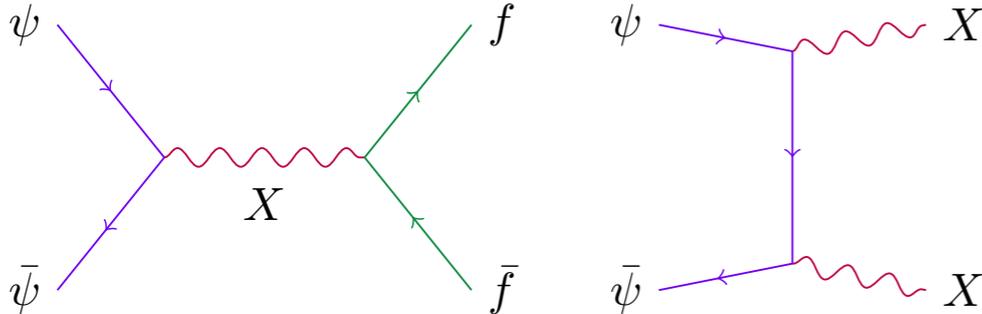
What would be the interplay with other DM production mechanisms?

Assume a dark  $U(1)_X$

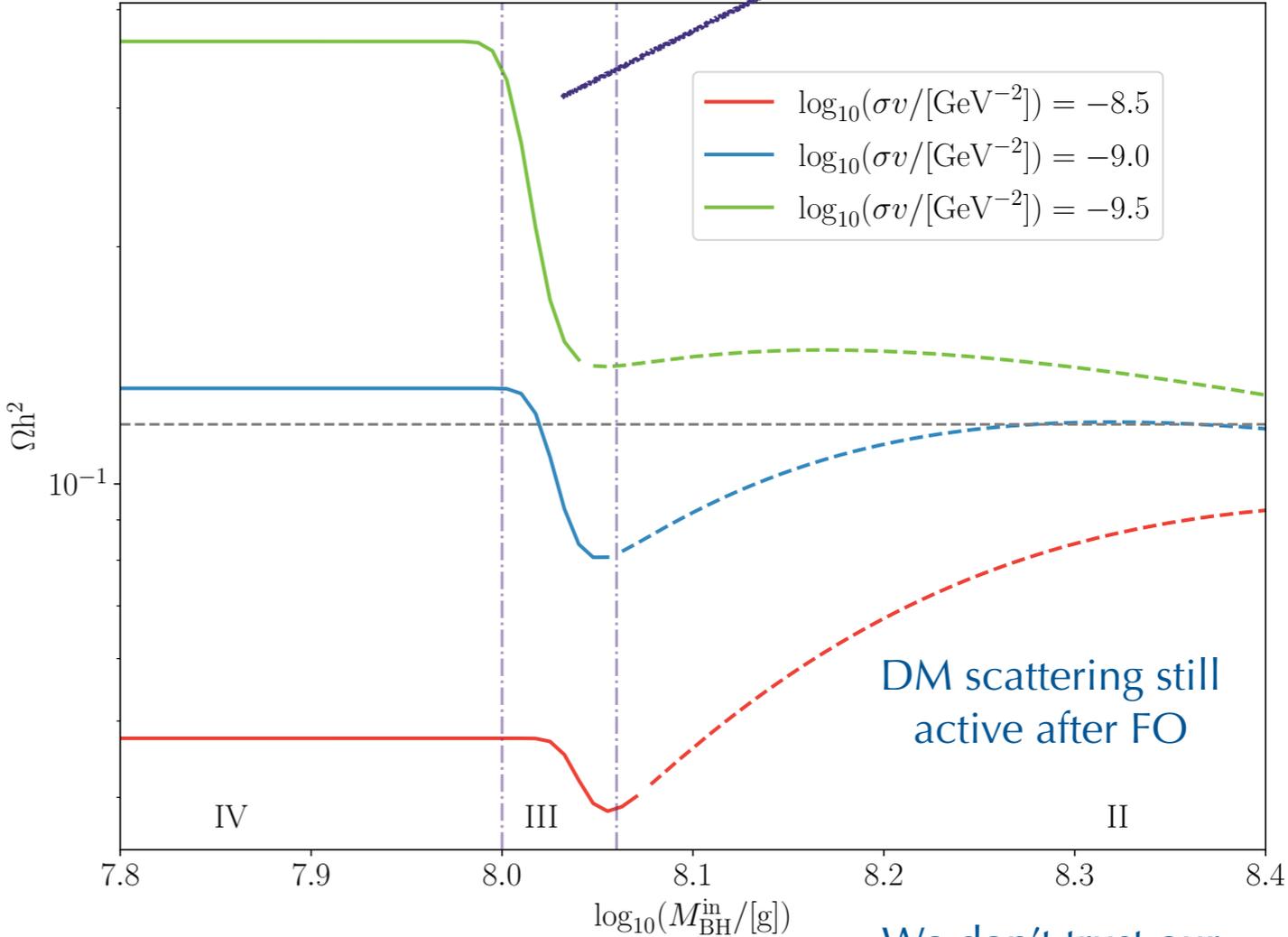
$\psi \longrightarrow$  DM

$X \longrightarrow$  Mediator

$f \longrightarrow$  SM fermions



Entropy dilution

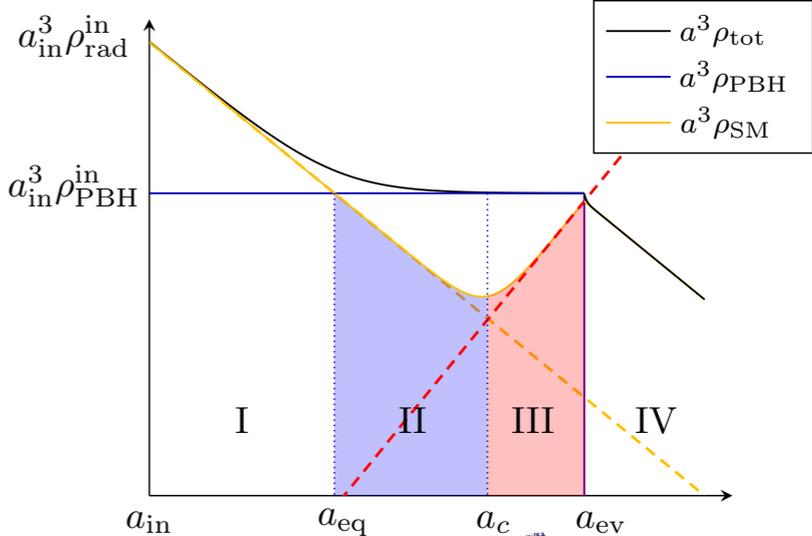


DM scattering still active after FO

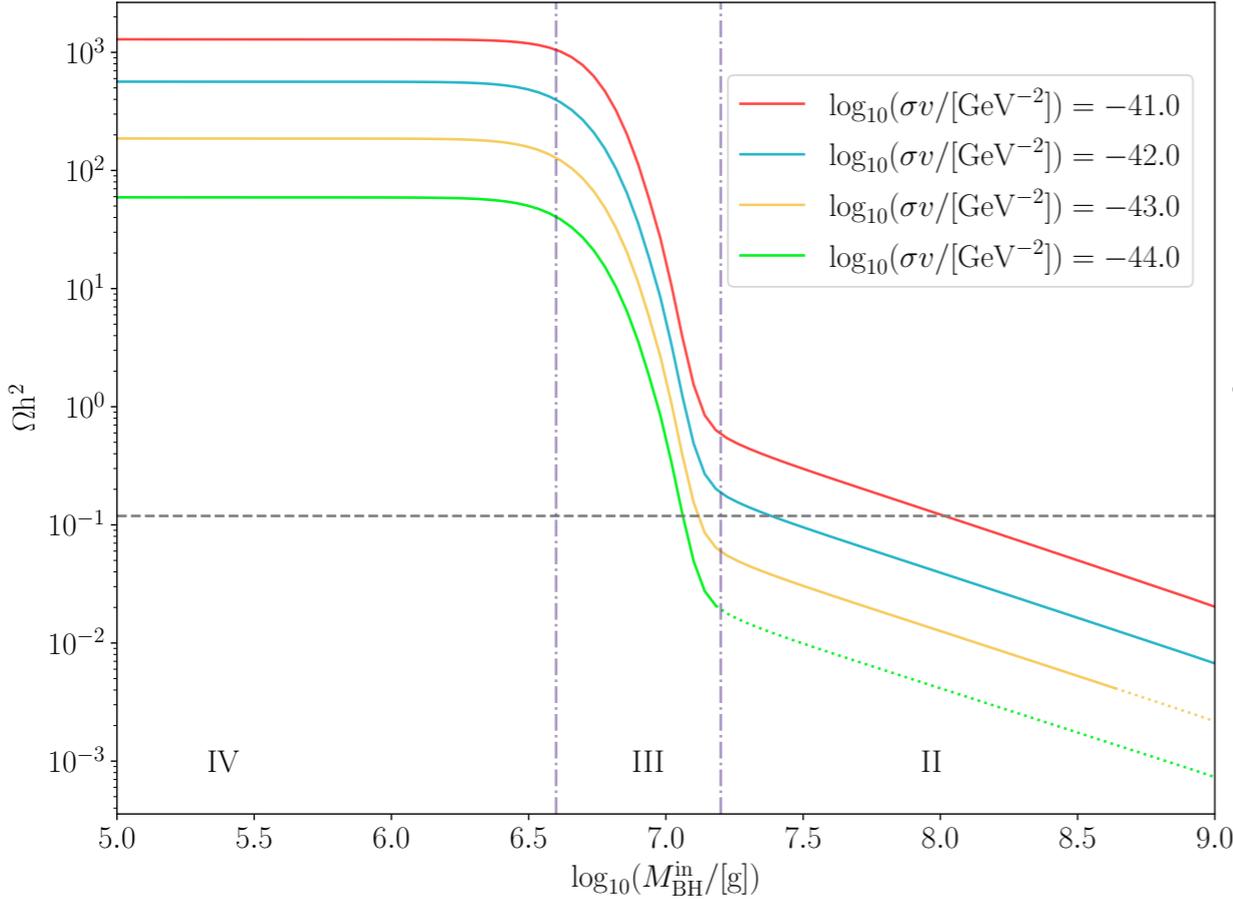
We don't trust our code in this region...

# Freeze-In

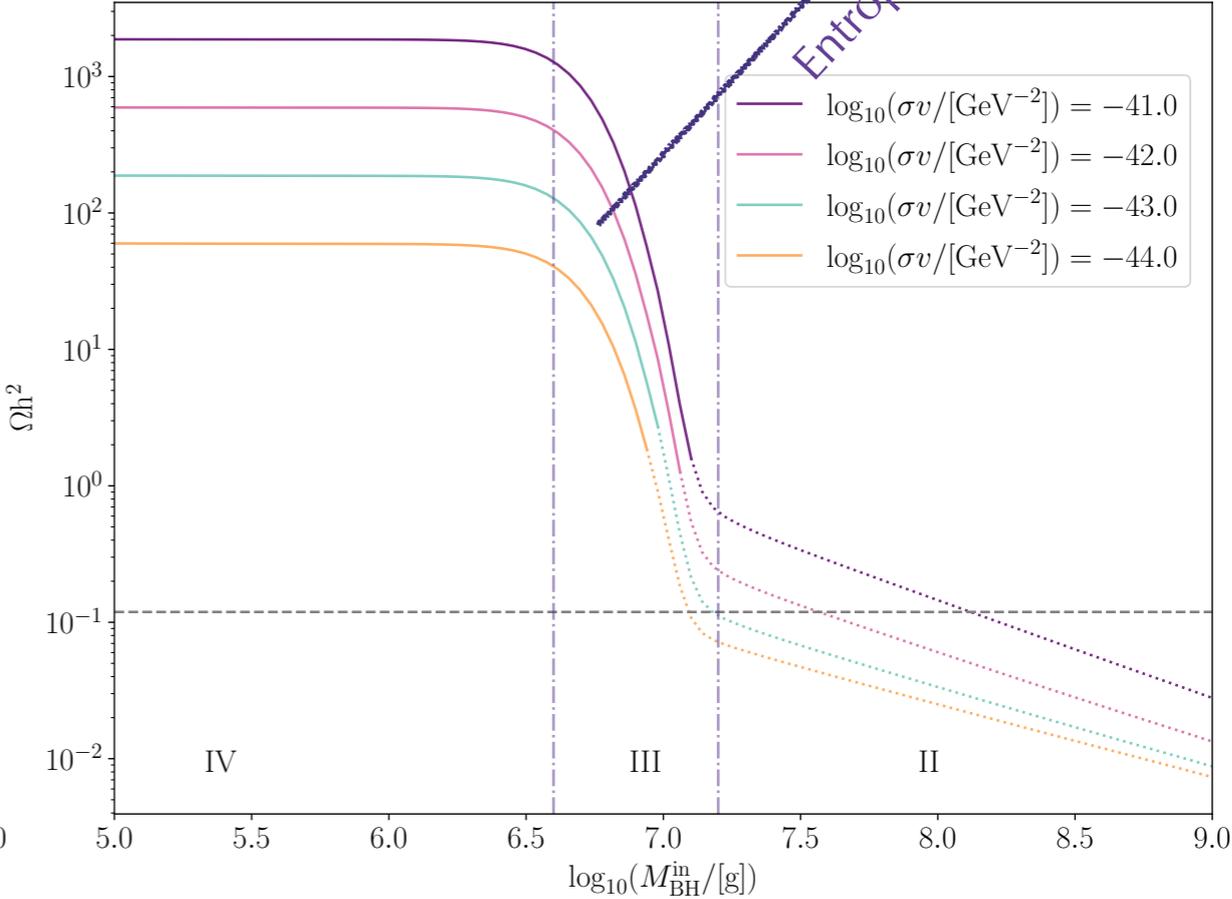
Overproduction is "cured" with the dilution



$m_{\text{DM}} = 10^{-3} \text{ GeV}$



$m_{\text{DM}} = 10^{-1} \text{ GeV}$



$m_X = 10 \text{ GeV}$

Dotted lines: Warm DM constraints