SHEDDING LIGHT ON DARK MATTER USING BOSE-EINSTEIN CONDENSATES

19TH MULTIDARK CONSOLIDER WORKSHOP

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Most part of this work was developed, simultaneously, at the Institute for Theoretical Physics (theory) and the Kirchhoff Institute for Physics (experiment) of the University of Heidelberg, Germany.





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Fig. 1. Left: theory team. Right: experimental team.

MOTIVATION

- QFT in curved spacetimes \rightarrow (DM) Particle production
- Direct detection \rightarrow Very hard
- Analog model \rightarrow Simulate with a dynamically equivalent system
- Weakly interacting BEC is described by

 $\Phi = \text{background} + \text{fluctuations} = \sqrt{n_0}e^{iS_0} + \delta\Phi$ (1)

Real world	BEC experiment
Spacetime	n_0, S_0
FLRW metric	Acoustic metric
Massless scalar particle	Phonons (sound)

THE ACOUSTIC METRIC

• We parametrise the field as

$$\Phi = \phi_0 + \frac{1}{\sqrt{2}} \left(\phi_1 + i \phi \right), \tag{2}$$

and expand

$$S[\Phi] = S[\phi_0] + \text{linear in } \phi_1, \phi + S_2[\phi_1, \phi]$$
(3)

• EOM for the fluctuations is given by

$$\delta S[\Phi] = \delta S_2[\phi_1, \phi] \stackrel{!}{=} 0 \tag{4}$$

- 1. Assume excitations have low momentum
- 2. Integrate out $\phi_1(\varphi)$
- 3. Adjust the background to be static

The action for the fluctuating field $\boldsymbol{\varphi}$ can be written as

$$S_{2}[\phi] = \frac{1}{2} \int_{x} \sqrt{g} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi, \qquad (5)$$

with the line element

$$ds^{2} = -dt^{2} + \frac{1}{c^{2}(t,r)} \left(dr^{2} + r^{2}d\phi^{2} \right) \equiv -dt^{2} + \frac{a^{2}(t)}{n_{0}(r)} \left(dr^{2} + r^{2}d\phi^{2} \right),$$
(6)

where the speed of sound is

$$c^{2}(t,r) = \frac{\lambda(t) n_{0}(r)}{m}.$$
(7)

We make $\mathbf{r} \rightarrow \mathbf{u}$ so that

$$\frac{dr^2}{n_0(r)} = \frac{du^2}{1 - \kappa u^2}, \quad \frac{r^2}{n_0(r)} d\phi^2 = u^2 d\phi^2, \tag{8}$$

and the line element reads

$$ds^{2} = -dt^{2} + a^{2}(t) \left(\frac{du^{2}}{1 - \kappa u^{2}} + u^{2}d\phi^{2} \right).$$
 (9)



Fig. 2. Density distribution for hyperbolic (a) and spherical (b) geometry and wave packet propagation. Viermann et al. [2202.10399]



Fig. 3. Analogy between the expanding universe and the condensate. The relation between the interaction strength λ and the scale factor is sketched.

PARTICLE PRODUCTION

- We make predictions for the spectrum of fluctuations S_k
- We measure correlations of density contrast $\langle \delta_c(t,u) \delta_c(t,u') \rangle$

We can relate $\langle \delta_c(t, u) \delta_c(t, u') \rangle \sim \langle \dot{\varphi}(t, u) \dot{\varphi}(t, u') \rangle \sim S_k(t)$.



Fig. 4. Time evolution of the scattering length $\lambda(t)$ and the scale factor a(t) for a polynomial expansion during the experiment. Tolosa-Simeón et al [2202.10441].



Fig. 5. Phase and amplitude of the oscillating spectum $S_k(t)$. The phase allows for a qualitative distinction between expansion with different scale factors $a(t) \sim t^{\gamma}$. Viermann et al. [2202.10399].

OUTLOOK

The abundance of DM can be explained by gravitational production during inflation alone.



Fig. 6. Abundance for a scalar field ϕ in logarithmic scale, considering $T_{rh} = 10^{15}$ GeV. White contours represent the observed value. Here, ξ is the coupling to the Ricci scalar and m is the mass of the DM particle. Cembranos et al, JHEP 2020, 84.

- We achieved a 1-1 mapping from BECs to curved FLRW universes
- We measured particle production in an actual experiment and extracted information of the type of expansion
- Tools from analogue gravity (curvature and phase of S_k) can be used in cosmological analyses.
- Particle production is very sensible to the inflaton scheme
- Spin-1 BECs allow to extend the analogy to massive complex scalar fields

Any questions?