Are there late-time solutions to the  $H_0$  and  $\sigma_8$  tensions?

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# $H_0$ and $\sigma_8$



[Millennium Simulation]



Clustering amplitude

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Both can be inferred from the CMB (after choosing a cosmoloy!!)

#### Cosmological tensions



Toy model of late-time 'solution': phantom dark energy with w = const. < -1











Why?



Overall position of the peaks: acoustic scale

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$$\theta_* = \frac{r_{\rm s}(z_*)}{d_A(z_*)}$$

$$\begin{split} r_{\rm s}(z_*) &= \int_{z_*}^{\infty} \frac{\mathrm{d}z}{H} c_{\rm s} &\to \quad \text{Early } H(z) \\ d_A(z_*) &= \int_0^{z_*} \frac{\mathrm{d}z}{H} &\to \quad \text{Late } H(z) \end{split}$$



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To infer a larger  $H_0$  from the CMB, we need models such that  $\theta_* \downarrow$  (fixed  $H_0$ )

- Early-time solutions:  $r_{\rm s}(z_*)\downarrow$
- Late-time solutions:  $d_A(z_*) \uparrow$

Both usually worsen the  $\sigma_8$  tension!

 $\Lambda \text{CDM}$  cosmology:

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$$H(z) = H_{\Lambda \text{CDM}}(z) + \frac{\delta H(z)}{\delta H(z)}, \qquad \frac{\delta H(z)}{H(z)} \ll 1 \qquad \delta H(z > 100) = 0$$

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- Model-independent approach: we **do not** parameterize  $\delta H(z)$
- Fully analytical results
- Generalization: we will also include (some) additional effects on the perturbations ( $G_{\rm eff}=G_N+\delta G(z)$ )

## $\Lambda \text{CDM}$

$$\left\{ \begin{array}{l} (H_0,\Omega_m) \\ \\ H=H_{\Lambda {\rm CDM}}(z) \\ \\ G_{\rm eff}=G={\rm const.} \end{array} \right.$$

# $\begin{cases} (H_0, \Omega_m) \\ H = H_{\Lambda CDM}(z) \\ G_{\text{eff}} = G = \text{const.} \end{cases}$

ΛCDM

#### Alternative model

 $\left\{ \begin{array}{l} (H_0,\Omega_m,+{\rm new \ params.})\\ H=H_{\Lambda{\rm CDM}}(z)+\delta H(z)\\ G_{\rm eff}=G+\delta G(z) \end{array} \right.$ 

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After we compare with observations, the preferred values will be different from  $\Lambda \text{CDM}$ 

 $(H_0 + \Delta H_0, \ \Omega_m + \Delta \Omega_m, \ \dots \ | \ \sigma_8 + \Delta \sigma_8, \ \theta_* + \Delta \theta_*, \ \dots)$ 

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These changes are connected to the properties of the new model

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(we will fix  $\omega_m = \Omega_m h^2$  here, more details in the paper)

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 → Effect of, e.g., dark energy, alternative dark matter

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$$\frac{\Delta\chi(z)}{\chi(z)} = I_{\chi}(z)\frac{\delta h}{h} + \int_{0}^{\infty} \frac{\mathrm{d}x_{z}}{1+x_{z}}R_{\chi}(x_{z},z)\frac{\delta H(x_{z})}{H(x_{z})}$$

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We can relate  $\delta h$  and  $\delta H(z)$  by fixing  $\Delta \theta_* = 0$  (CMB angular scale)

Fixing the angular scale of the CMB ( $\Delta \theta_* = 0$ ), we can relate  $\delta h$  and  $\delta H(z)$ 

$$\frac{\delta h}{h} = \int_0^\infty \frac{\mathrm{d}z}{1+z} \mathcal{R}_h(z) \frac{\delta H(z)}{H(z)}$$

$$\frac{\Delta\sigma_8}{\sigma_8} = I_{\sigma_8}\frac{\delta h}{h} + \int_0^\infty \frac{\mathrm{d}z}{1+z} R_{\sigma_8}(z) \frac{\delta H(z)}{H(z)}$$

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$$= \int_0^\infty \frac{\mathrm{d}z}{1+z} \mathcal{R}_{\sigma_8}(z) \frac{\delta H(z)}{H(z)}$$

$$\stackrel{-0.25}{-0.50}$$

$$\stackrel{-0.75}{-0.75}$$

$$\stackrel{-0.$$

8

 $10^{2}$ 



(Both  $\mathcal{R}_h$  and  $\mathcal{R}_{\sigma_8}$  are fully analytical!)











**O** Solving  $H_0$  tension  $(\delta h > 0) \Rightarrow \exists z \mid \delta H(z) < 0 \Rightarrow \exists z \mid w(z) < -1$ 



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(1) Solving both tensions ( $\delta h > 0$ ,  $\Delta \sigma_8 < 0$ ):

a) If  $G_{\text{eff}} = G \implies \delta H(z)$  changes sign  $\implies w(z)$  crosses -1

b) If  $G_{\text{eff}} \neq G \implies ?$ 

# Modifying the perturbations

The matter density perturbation on small scales evolves as

$$\frac{\mathrm{d}^2 \delta_m}{\mathrm{d}a^2} + \frac{\mathrm{d}\log\left(a^3 H\right)}{\mathrm{d}a} \frac{\mathrm{d}\delta_m}{\mathrm{d}a} - \frac{3\Omega_m H_0^2}{2a^5 H^2} \frac{G_{\mathrm{eff}}}{G} \delta_m = 0$$

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We will assume that the effects of the new model can be described with a scale-independent, small deviation  $\delta G(z)$ 

$$G_{\text{eff}} = G + \delta G(z)$$

Then, only the growth factor D is modified

$$\delta_m(z,k) \propto \left( D(z) + (\Delta D)|_{\delta G} \right) T(k)$$

Following the same steps, we compute the response to  $\delta G(z)$  so now the full variation is

$$\frac{\Delta\sigma_8}{\sigma_8} = \int_0^\infty \frac{\mathrm{d}z}{1+z} \mathcal{R}_{\sigma_8} \frac{\delta H}{H} + \int_0^\infty \frac{\mathrm{d}z}{1+z} \mathcal{G}_{\sigma_8} \frac{\delta G}{G}$$

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**(**) If we want to solve the  $H_0$  and  $\sigma_8$  tensions ( $\delta h > 0$  and  $\Delta \sigma_8 < 0$ ) and if  $\delta H(z)$  does not change sign ( $\delta H(z) < 0$ )

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$$\implies \qquad \qquad \frac{\delta G(z)}{G} < \alpha(z) \frac{\delta H(z)}{H(z)} < 0 \quad \text{for some } z$$

Again,  $\alpha(z)$  is a known, analytical function

General conditions for any late dark energy model to simultaneously solve the  $H_0$  tension (i.e.  $H_0 \uparrow$ ) and the  $\sigma_8$  tension (i.e.  $\sigma_8 \downarrow$ ):

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 tension  $\implies \delta H(z) < 0$  for some  $z$   
 $\implies w(z) < -1$  for some  $z$ 

**()** Without modifying the perturbations  $(G_{\text{eff}} = G)$ :

- **a** Both tensions cannot be solved if  $\delta H(z)$  does not change sign
- **(b)** Solving the  $H_0$  and  $\sigma_8$  tensions  $\implies w(z)$  crosses the phantom divide

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# Example 1: wCDM



	$100 \times \delta h/h$		$100 \times$	$\Delta \sigma_8 / \sigma_8$
w	class	Analytical	class	Analytical
-0.8	-8.57	-10.97	-8.98	-7.13
-0.85	-6.47	-7.76	-6.29	-5.32
-0.9	-4.35	-4.89	-3.93	-3.53
-0.95	-2.19	-2.32	-1.85	-1.76
-1.05	2.22	2.10	1.65	1.75
-1.1	4.47	4.01	3.12	3.48
-1.15	6.75	5.74	4.45	5.21
-1.2	9.07	7.33	5.66	6.93

# Example 2: $w_0 w_a CDM$



$w_0 = -1.05$	$100 \times \delta h/h$		$100 \times$	$\Delta \sigma_8 / \sigma_8$
$w_a$	class	Analytical	class	Analytical
-0.05	2.81	2.61	2.24	2.08
-0.01	2.35	2.22	1.86	1.75
0.03	1.89	1.80	1.47	1.39
0.07	1.42	1.37	1.07	1.03
0.11	0.946	0.93	0.66	0.64
0.14	0.465	0.46	0.24	0.23
0.174	0.095	0.093	-0.089	-0.092
0.18	-0.022	-0.025	-0.19	-0.20
0.22	-0.52	-0.53	-0.64	-0.65
0.26	-1.02	-1.07	-1.09	-1.13
0.3	-1.54	-1.62	-1.56	-1.64

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# Additional response functions

