



# SYMMETRIC PRODUCT ORBIFOLDS

Alejandra Castro  
DAMTP



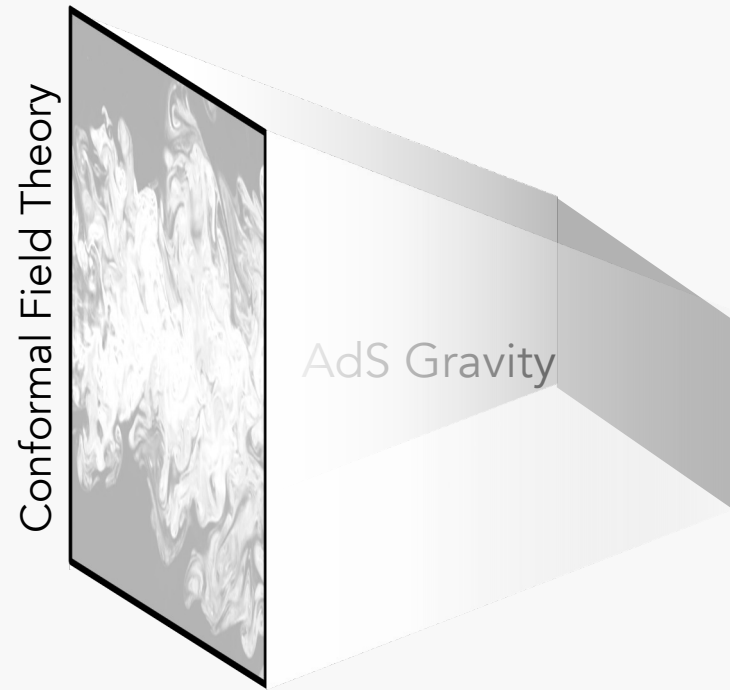
29th IFT Xmas Workshop  
Madrid, December, 2023

# HOW TO ENGINEER QUANTUM GRAVITY IN $ADS_3/CFT_2$

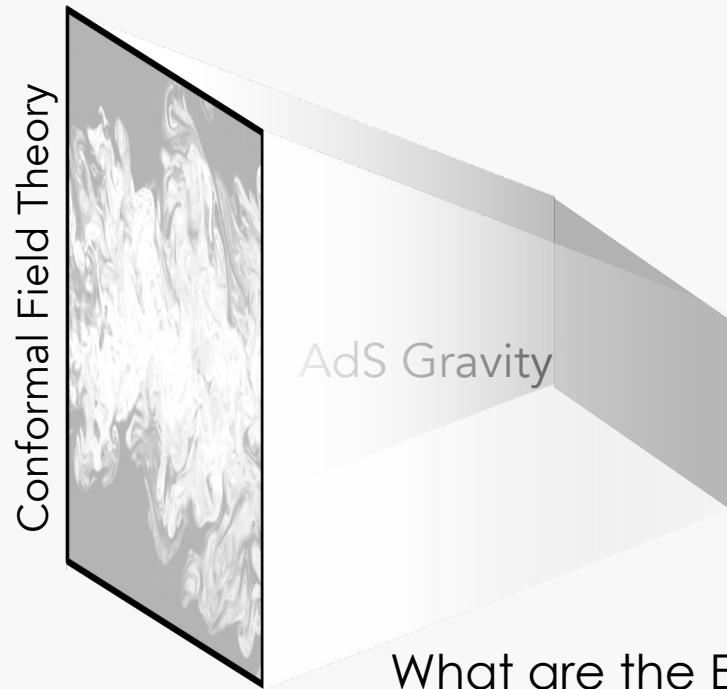
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$\text{CFT}_D \rightarrow \text{AdS}_{D+1} \text{ Gravity}$



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What are the EFT that admit a UV completion in QG?

Landscape

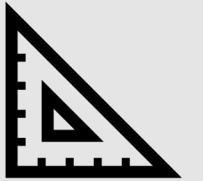
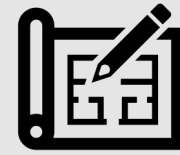
What are the consequences of this connection?

Principles & Laws in QG

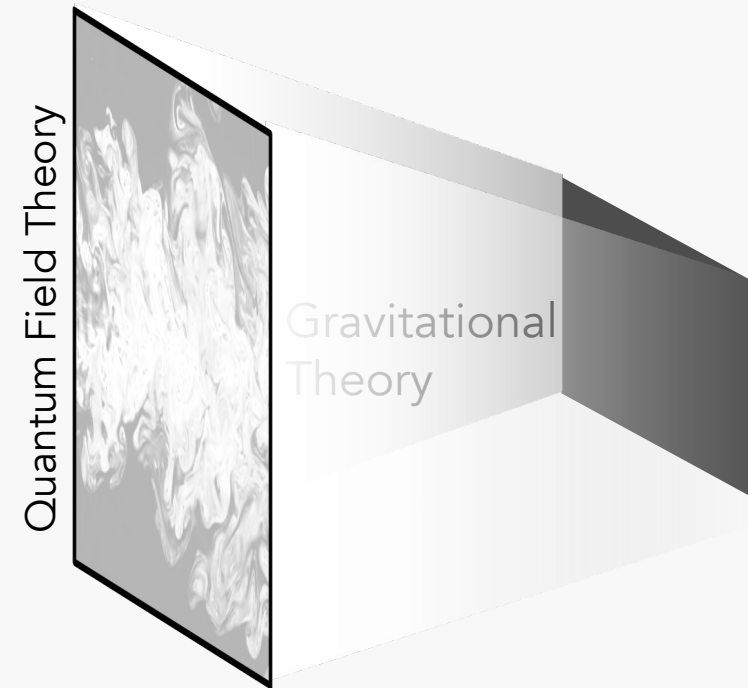
Engineering  $AdS_3$  Quantum Gravity

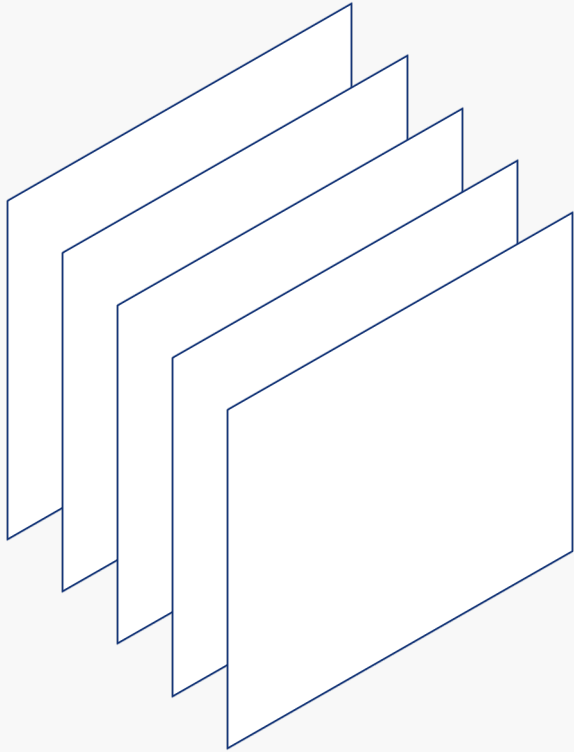
Symmetric Product Orbifolds

# Engineering AdS<sub>3</sub> Quantum Gravity



- Define gravity via the dual CFT
- Identify necessary conditions
- Possible designs we can achieve





- Implement conditions
- Precise outcomes (with surprises)
- New features in the design of AdS/CFT

A. Belin, J. Gomes, C. Keller, [AC](#), 2016, 2018

A. Belin, C. Keller, B. Mühlmann, [AC](#), 2019 (x2)

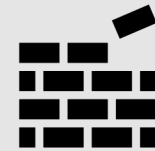
A. Belin, N. Benjamin, C. Keller and S. Harrison, [AC](#), 2020

L. Apolo, A. Belin, S. Bintanja, C. Keller, [AC](#) 2022

N. Benjamin, S. Bintanja, J. Hollander, [AC](#) 2022

L. Apolo, A. Belin, S. Bintanja, C. Keller, [AC](#) 2023

## Symmetric Product Orbifolds

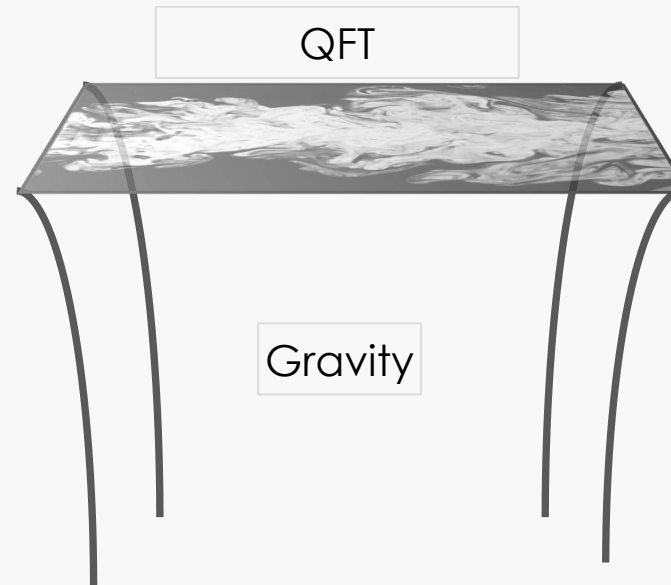
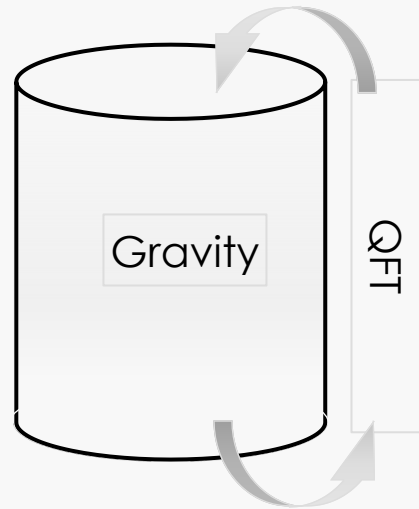


# Engineering AdS<sub>3</sub> Quantum Gravity



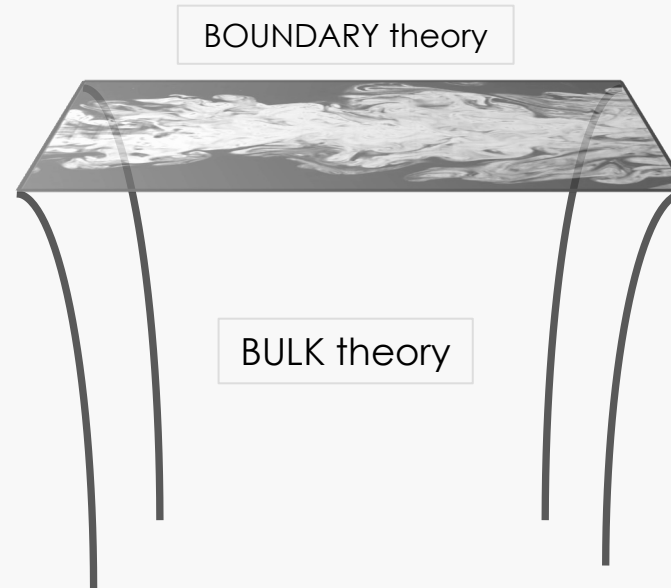
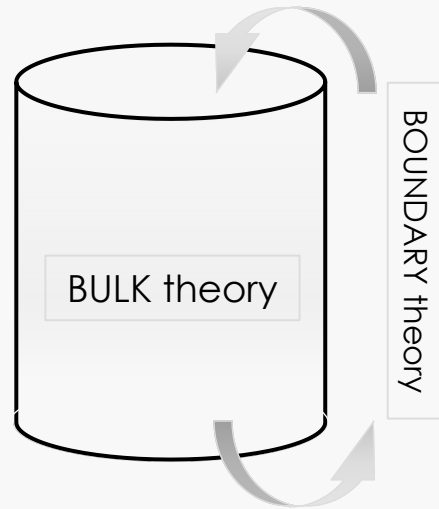
# Holographic Principle

Gravitational theory in  $D+1$  dimensions is equivalent to quantum field theory in  $D$  dimensions



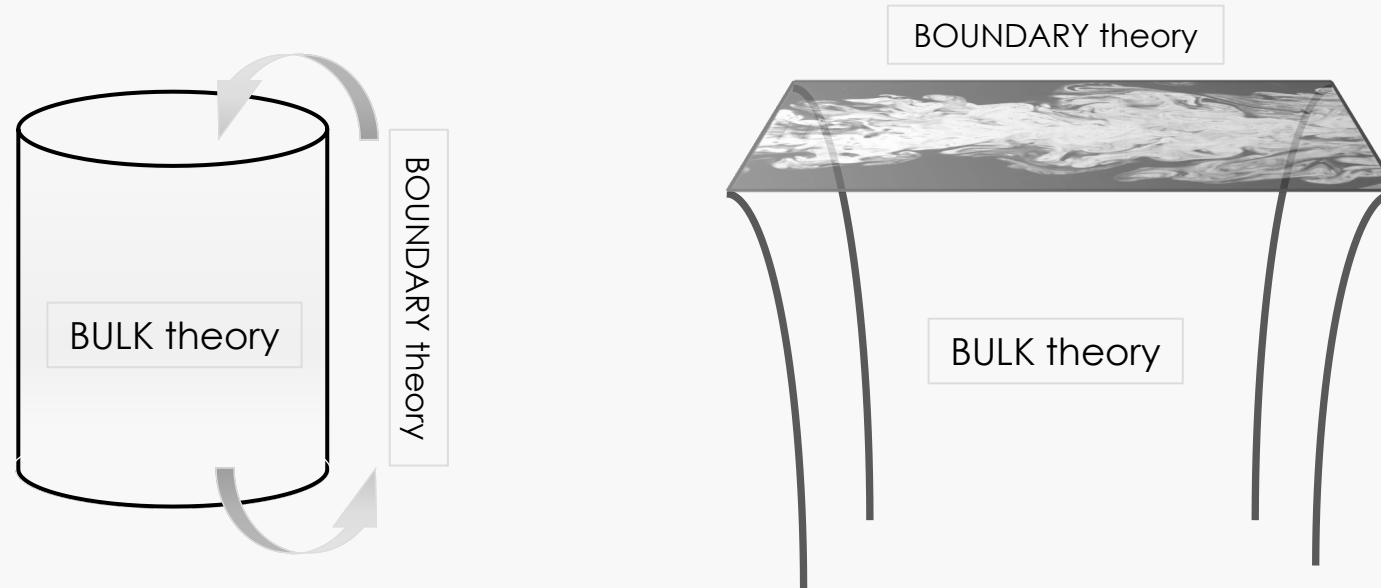
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Control and precision:  
 $AdS_{D+1}/CFT_D$

# AdS<sub>D+1</sub>/CFT<sub>D</sub>

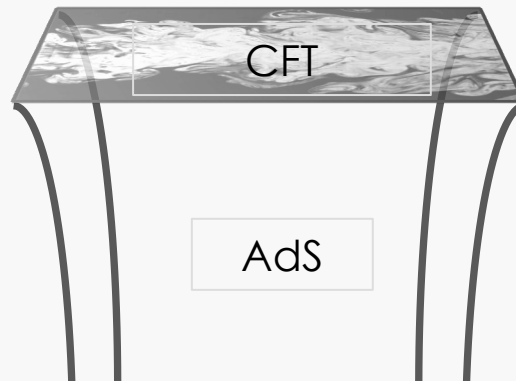
Born and refined in **String Theory**, the most precise incarnation states:

AdS Gravity

AdS<sub>D+1</sub> = gravity with negative cosmological constant in D+1 spacetime dimensions.

Conformal Field Theory

CFT<sub>D</sub> = quantum system invariant under dilations in D spacetime dimensions.



# AdS<sub>D+1</sub>/CFT<sub>D</sub>

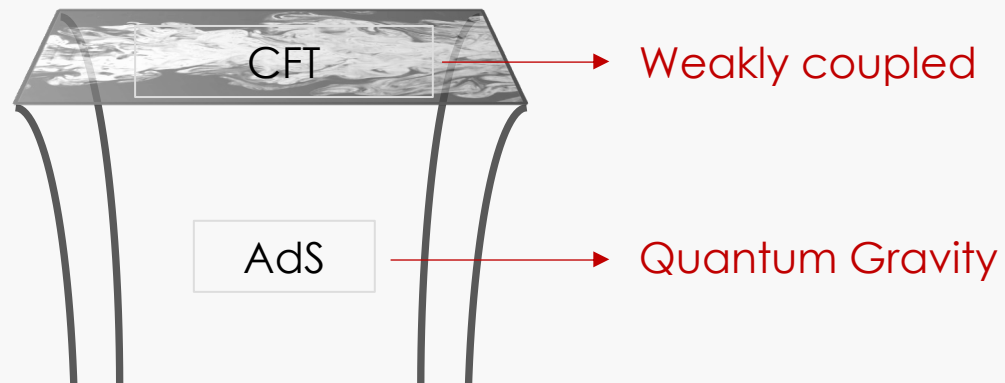
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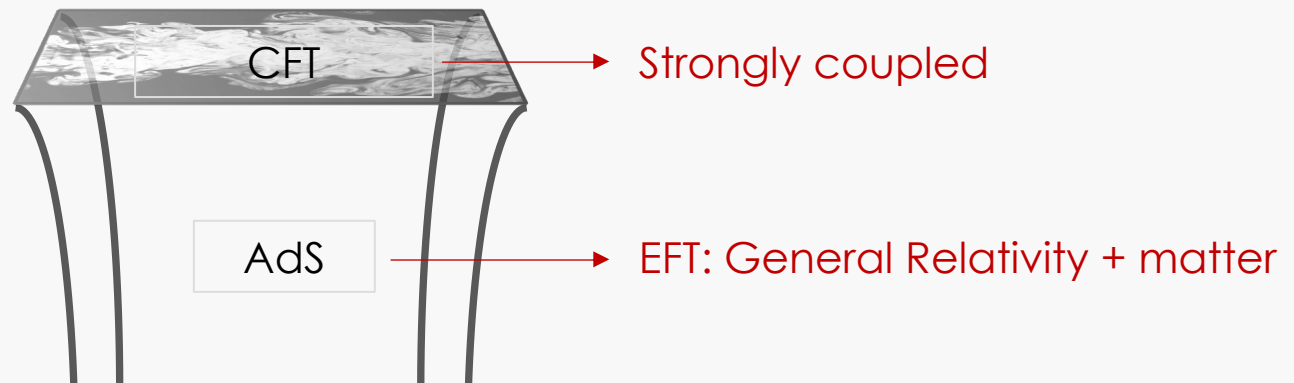
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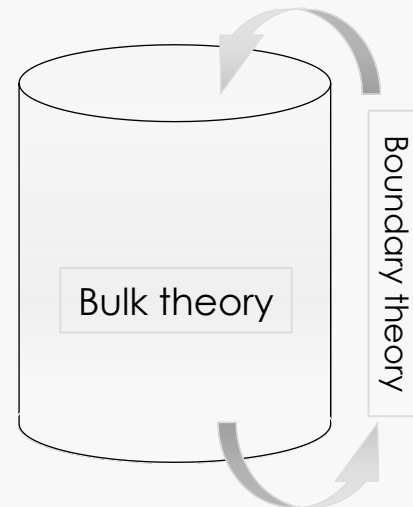
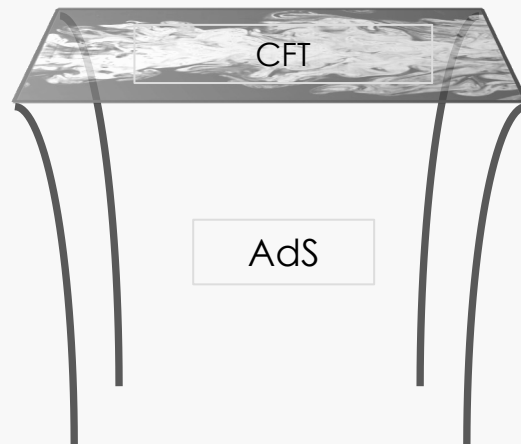
# AdS<sub>D+1</sub>/CFT<sub>D</sub>

It is a **duality**:

It maps a gravitational theory to a CFT and viceversa.

It is a **framework**:

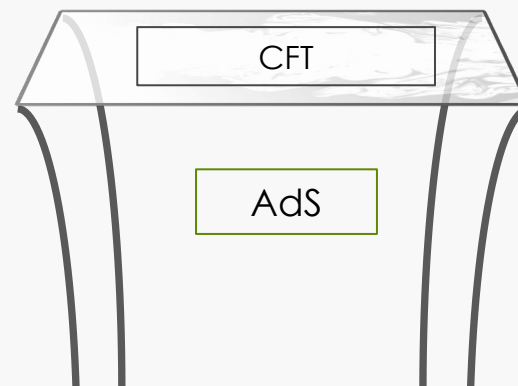
The dual CFT predicts the adequate observables in QG.



If we want to capture the fundamental mechanism of **AdS/CFT**, we need more than a few examples.

Possible Directions:

- Given a CFT, how to organize it such that quantum gravity is manifest?
- Which CFTs capture classical (geometric) properties of gravity?
- Select a theory of gravity, what are properties of the dual?

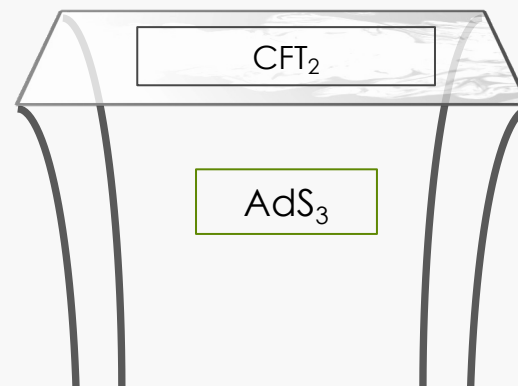




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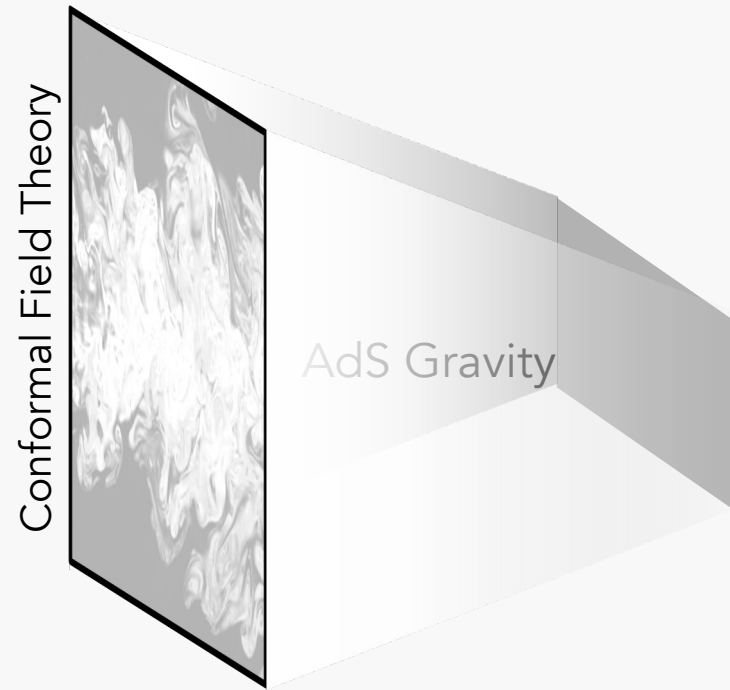
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- Given a CFT, how to organize it such that quantum gravity is manifest?
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We will focus on the difficulties you encounter in **AdS<sub>3</sub>/CFT<sub>2</sub>**. Not universal, but it illustrates the challenges.

$\text{CFT}_2 \rightarrow \text{AdS}_3 \text{ Gravity}$



# AdS<sub>3</sub> Gravity

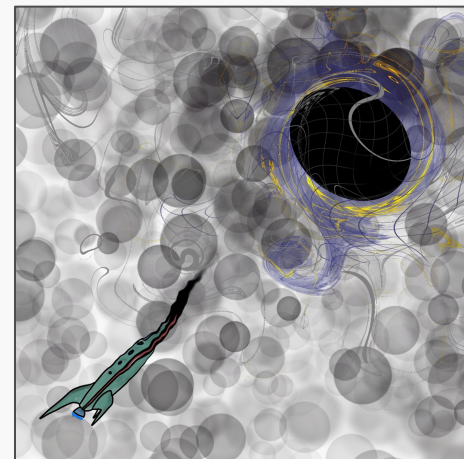
The theory:

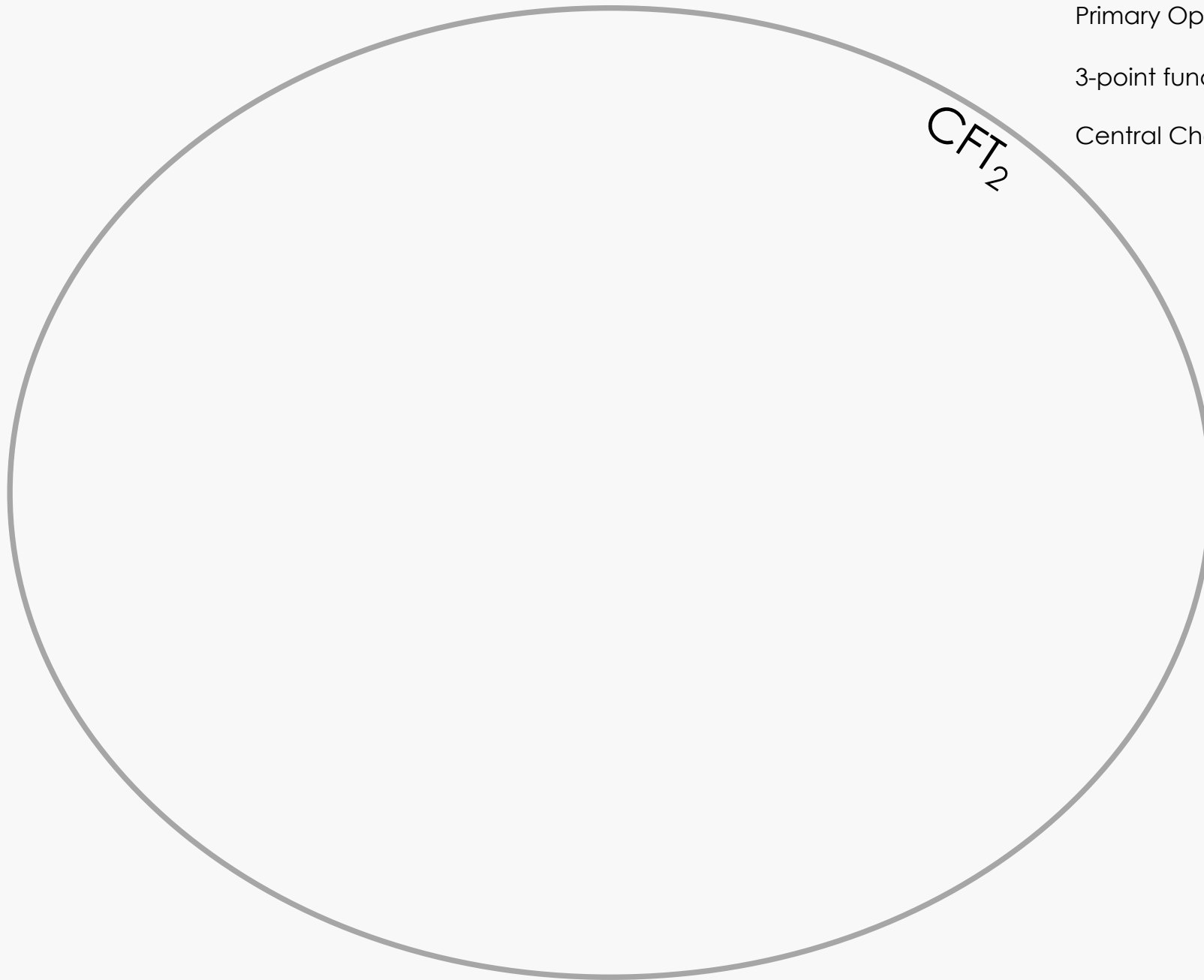
$$I_{3D} = \frac{1}{16\pi G_N} \int d^3x \sqrt{-g} \left( R + \frac{2}{\ell^2} \right) + \text{matter content: scalars, gauge fields, fermions}$$

Basic Requirements:

Gravitational theories that have the property that their low-energy description is given by a local EFT.

Universal entry in AdS<sub>3</sub>/CFT<sub>2</sub>:  $c = \frac{3\ell}{2G_N} \gg 1$



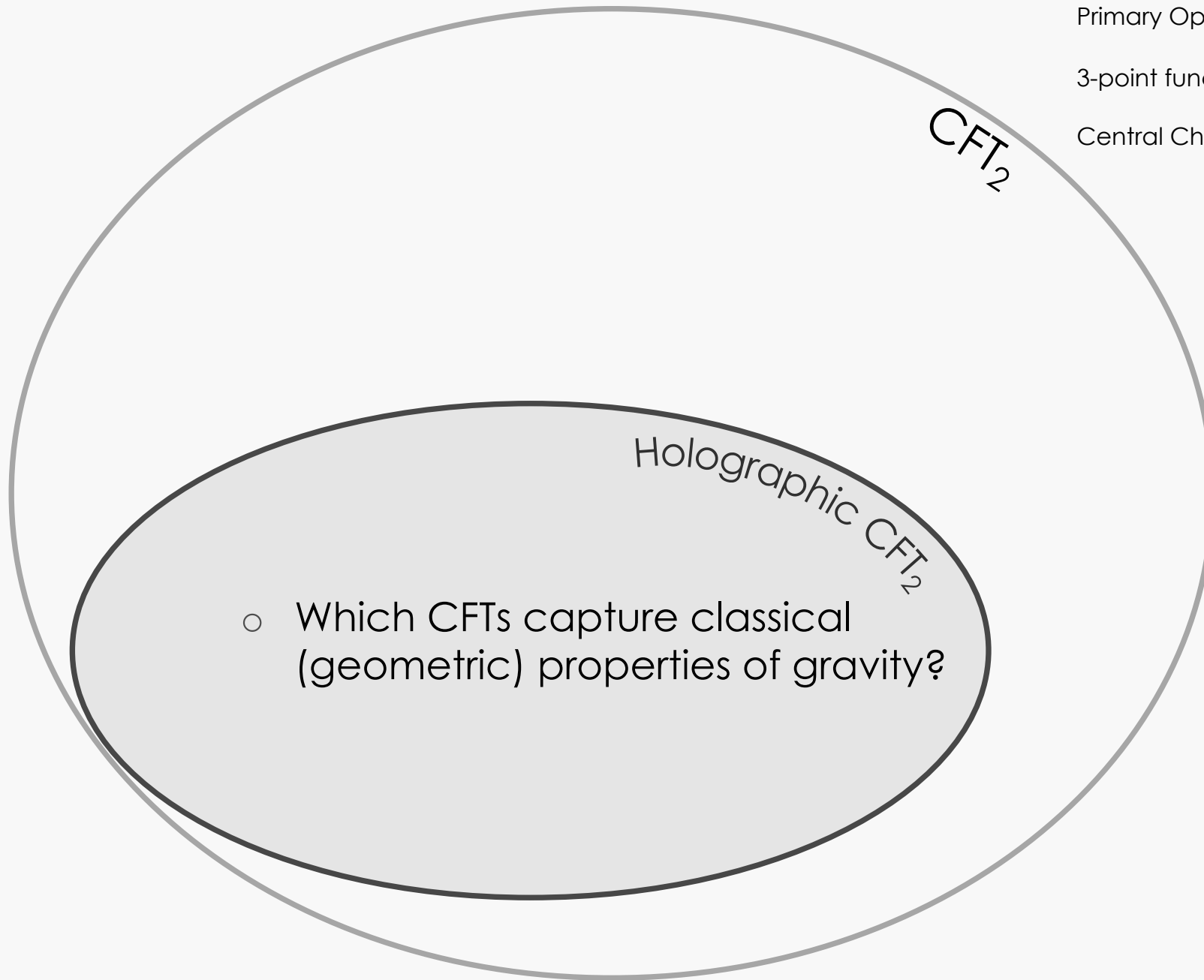


CFT<sub>2</sub>

Primary Operators:  $\Delta_J$  (Conformal dimensions)

3-point functions:  $C_{IJK}$  (OPE coefficients)

Central Charge:  $c$

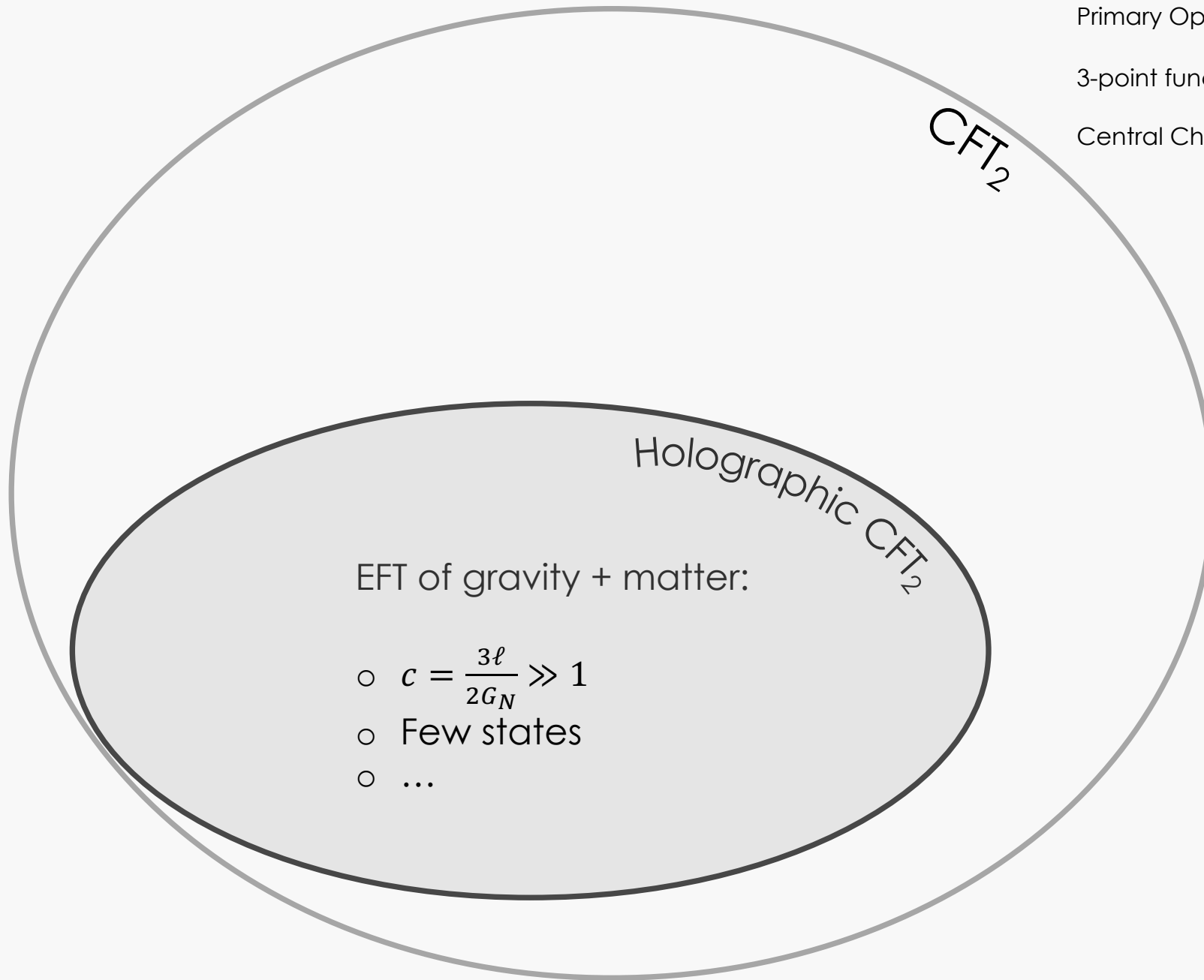


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*warning:* Sizes are not meaningful.

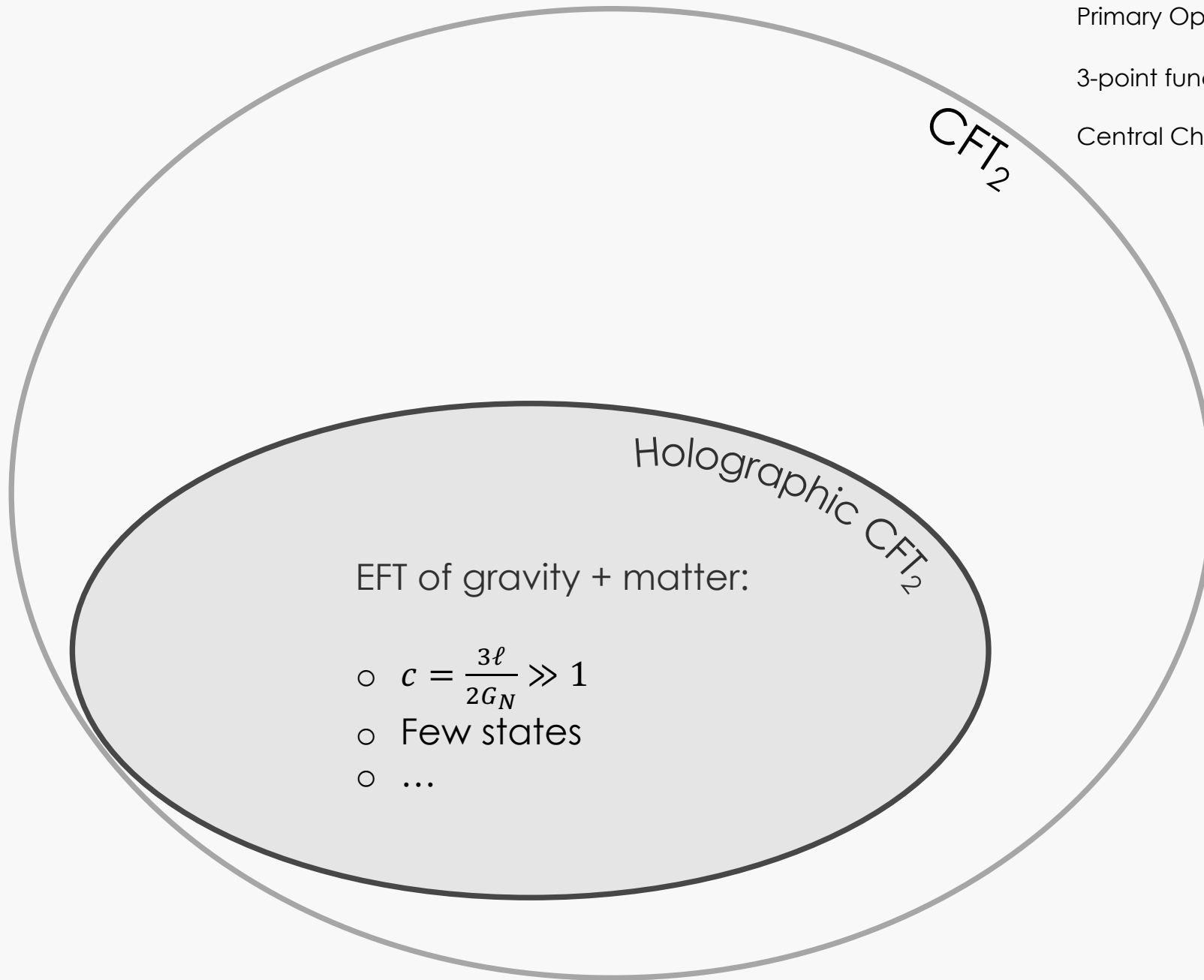


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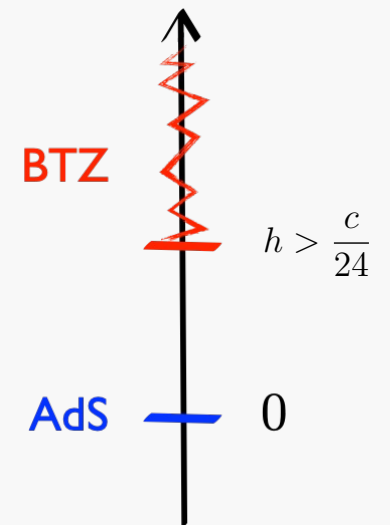
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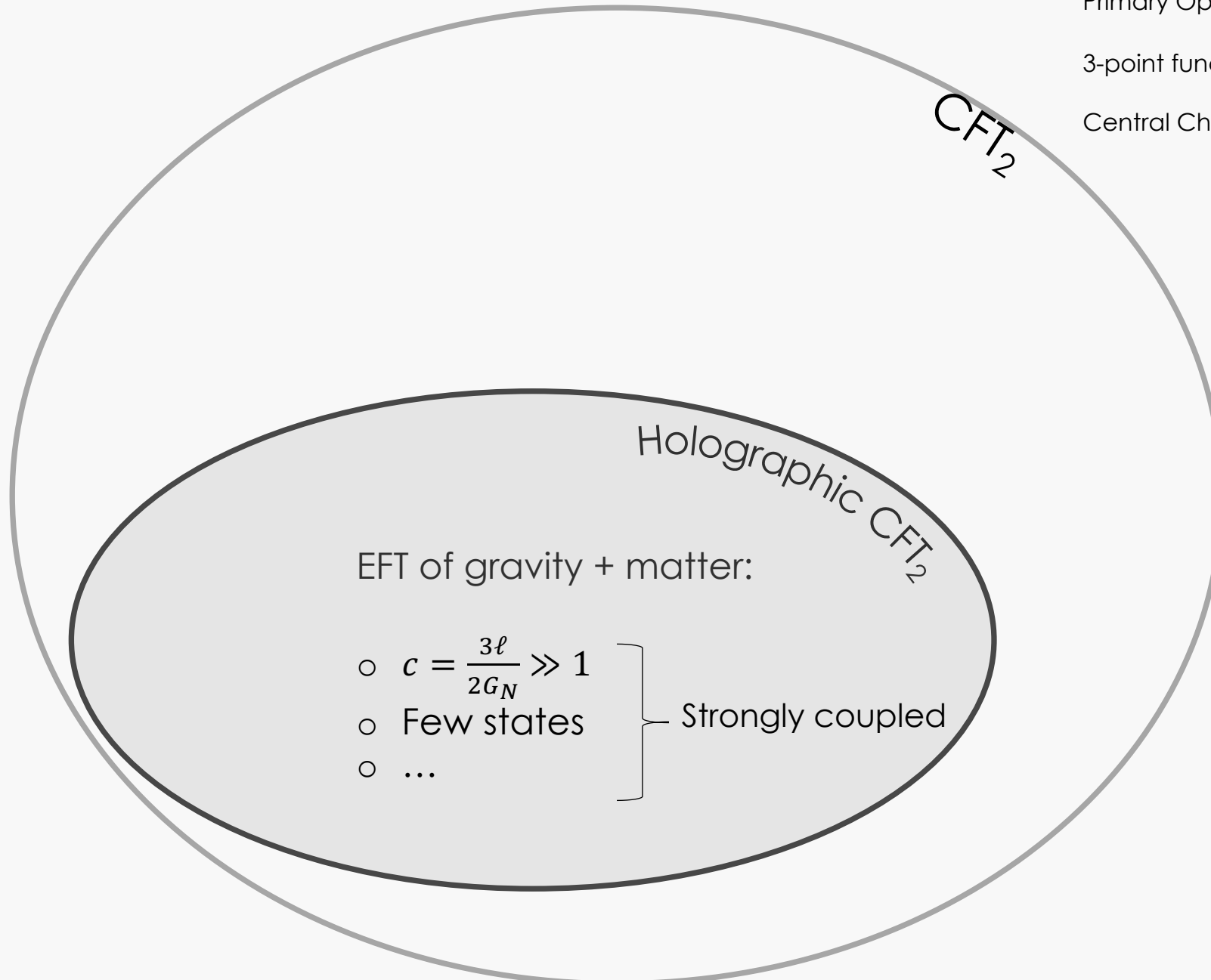


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## Holographic CFT<sub>2</sub>

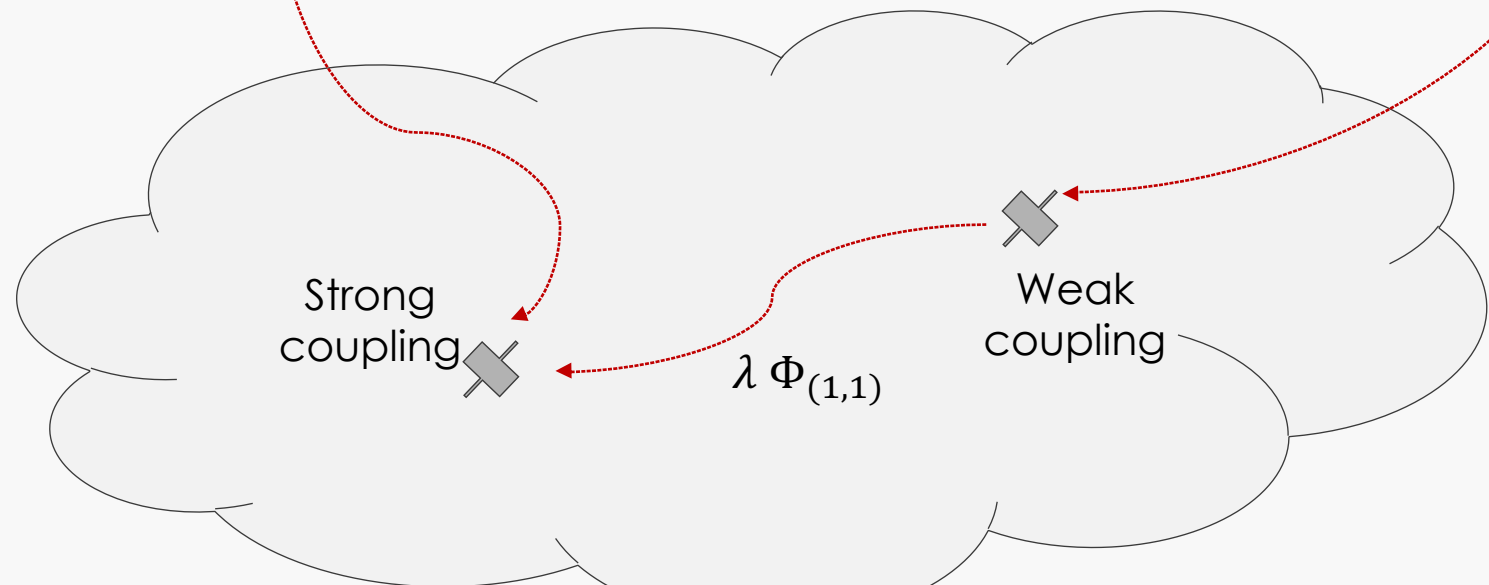
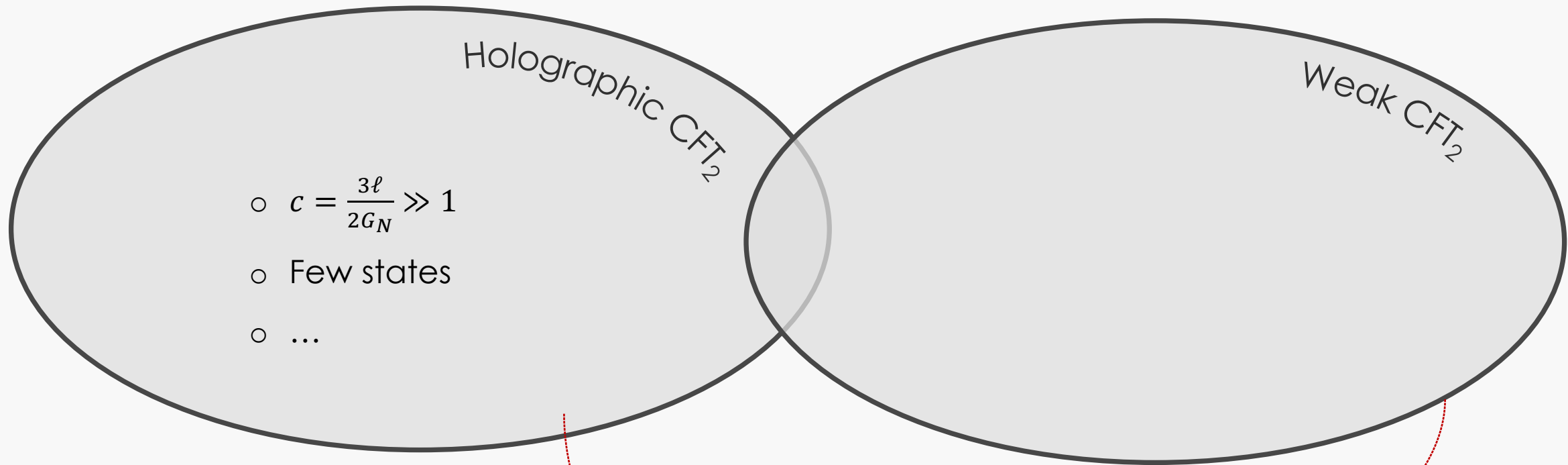
- $c = \frac{3\ell}{2G_N} \gg 1$
- Few states
- ...

Strong  
coupling

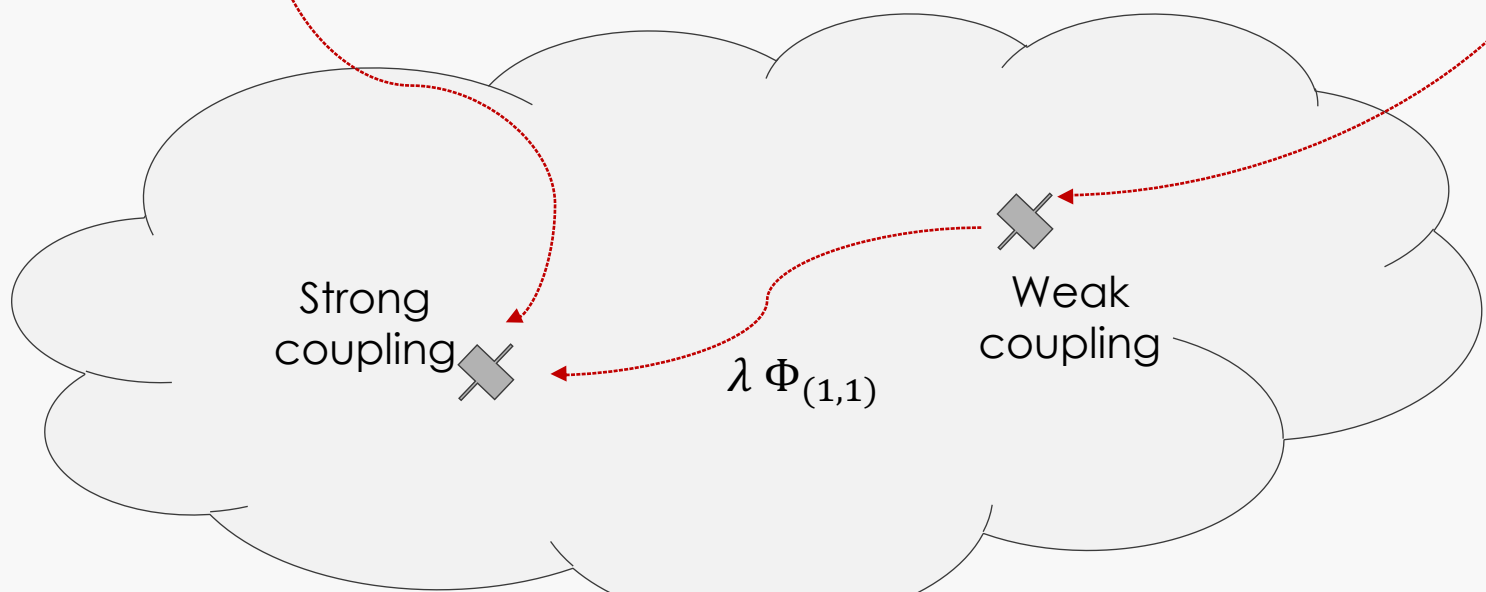
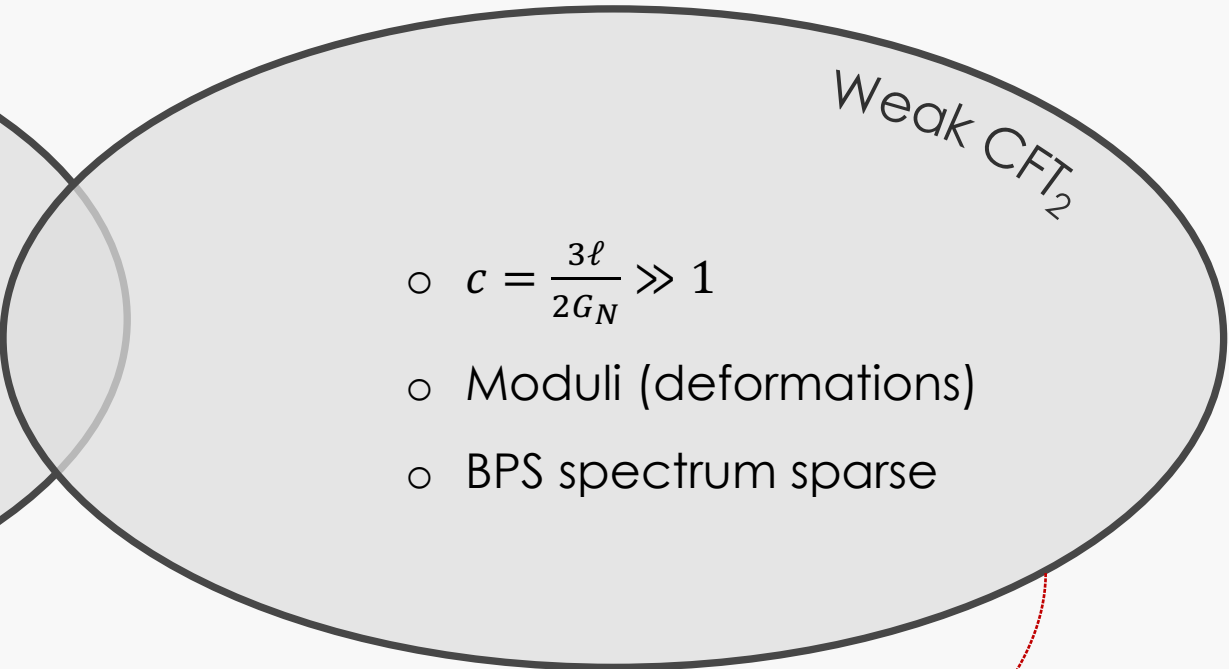
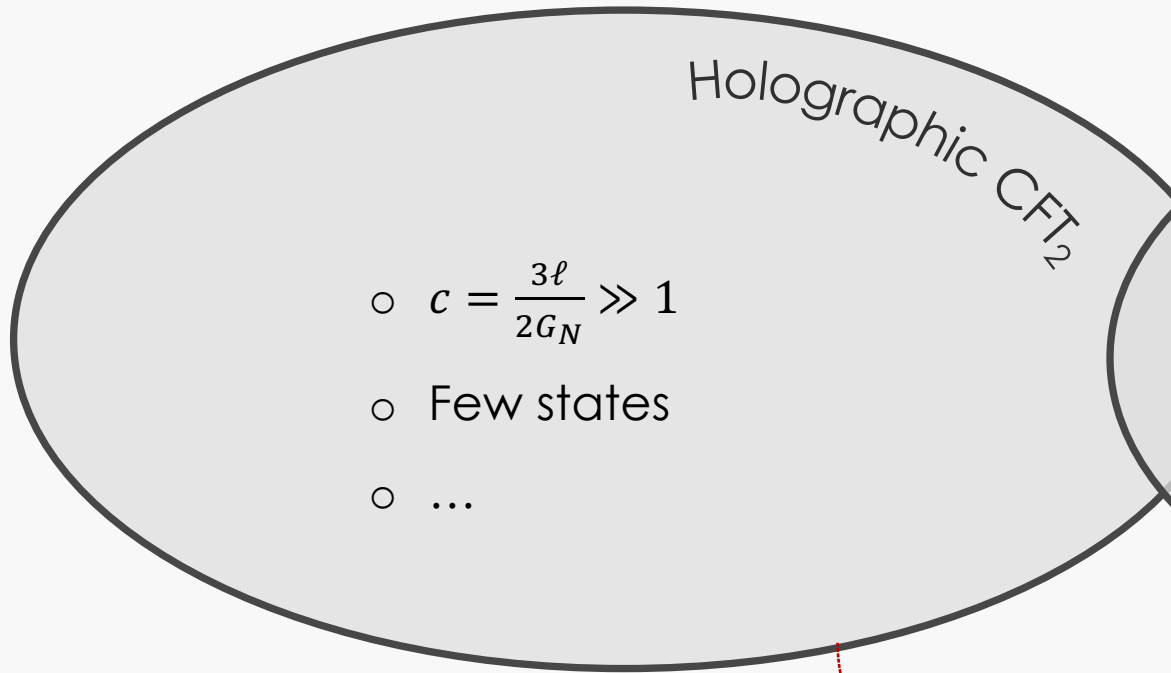


$\lambda \Phi_{(1,1)}$

Moduli space: set of exactly marginal deformations



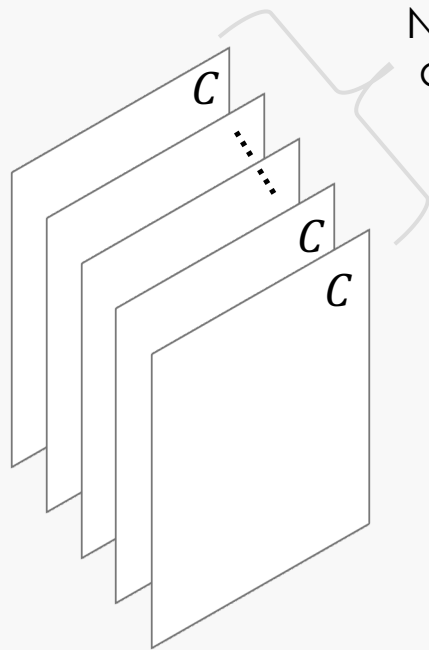
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# Symmetric Product Orbifolds

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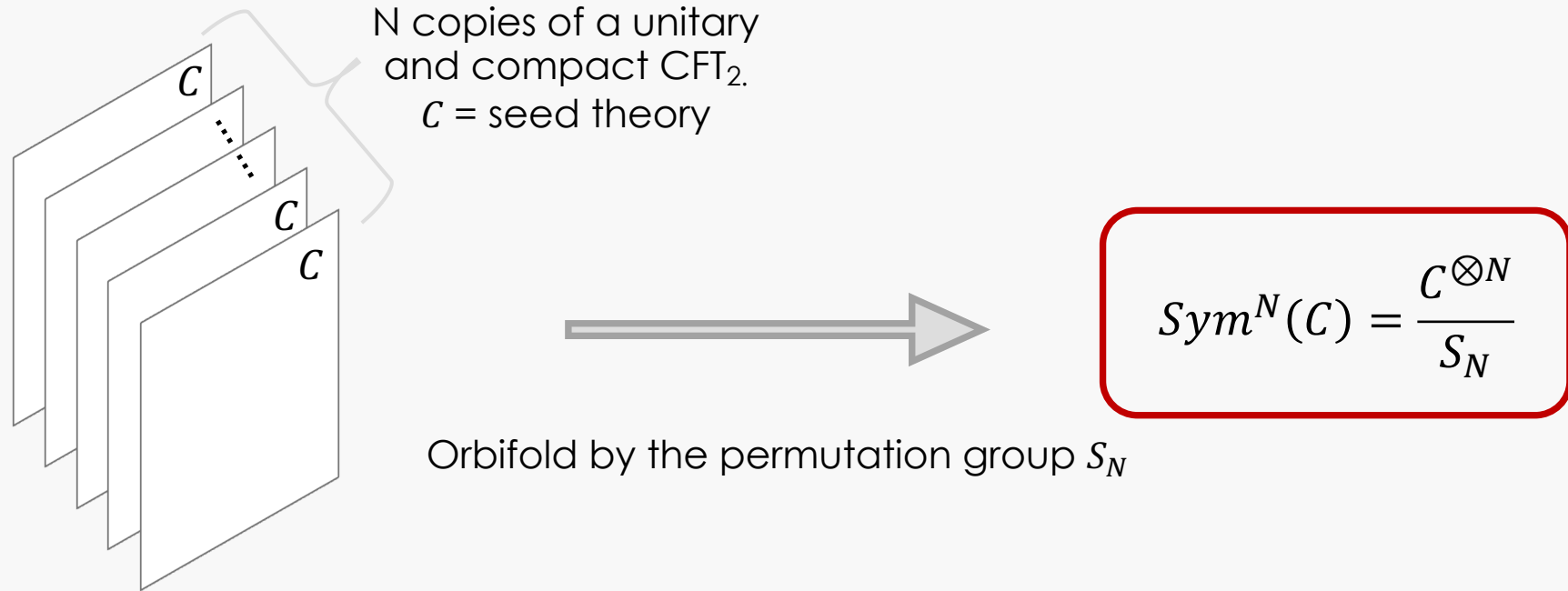
$N$  copies of a unitary  
and compact  $\text{CFT}_2$ .  
 $\mathcal{C}$  = seed theory



Orbifold by the permutation group  $S_N$

$$\text{Sym}^N(\mathcal{C}) = \frac{\mathcal{C}^{\otimes N}}{S_N}$$

# Symmetric Product Orbifolds

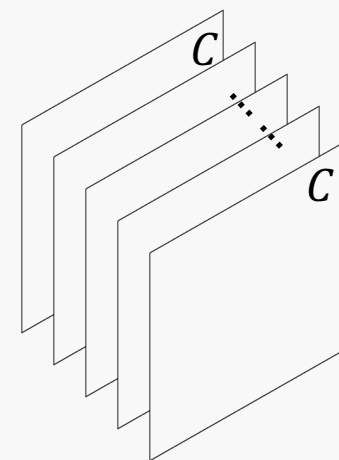


The orbifold introduces two class of states:

- **untwisted sector:** it keeps states that are invariant under  $S_N$ .
- **twisted sectors:** new states labelled by conjugacy classes of  $S_N$ .

# Symmetric Product Orbifolds

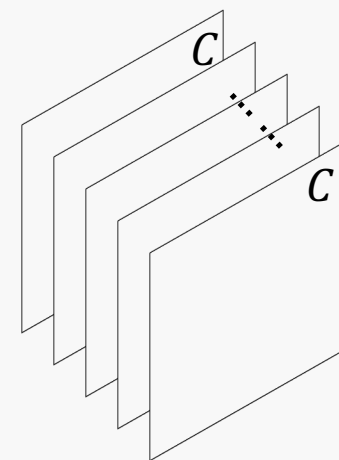
- **Appeal:** Mathematical and analytic control, e.g., DMVV formula.
- **Familiarity:** D1D5 CFT.
- **Universality:** large-N behavior is robust.
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Today: non-universal properties.  
Demonstrate that there are different  
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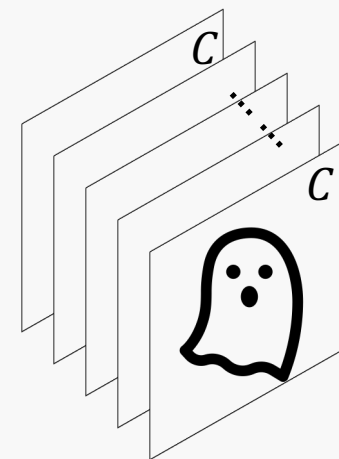




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# Universal Aspects

All symmetric product orbifolds satisfy:

- Correlation functions comply with large-N factorization.

[Pakman et.al., Mathur et.al., Belin et.al., Hael et.al., ...]

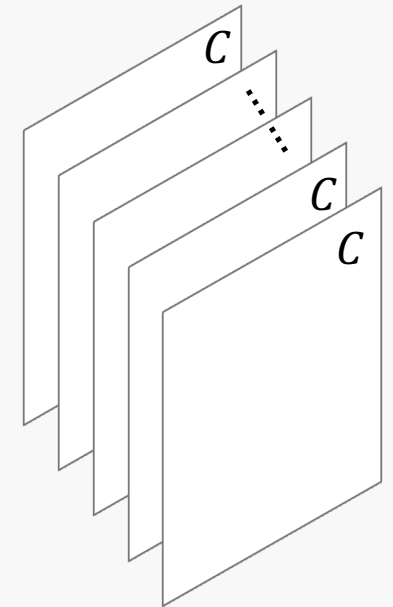
- Hawking-Page transition at large-N.

[Keller 2011; Hartman, Keller, Stoica 2014; Benjamin et.al. 2015]

- Higher spin currents due to orbifold structure.

- Universal Hagedorn growth of light states.


[Keller 2011]

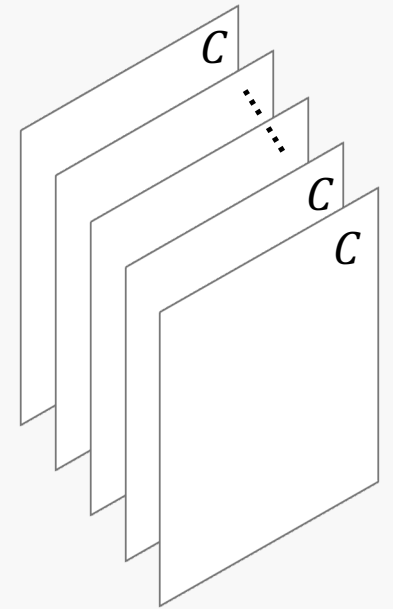


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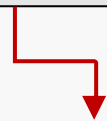


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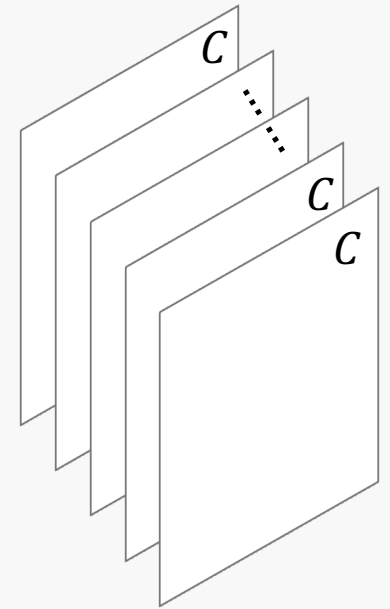
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AdS/CFT interpretation: Dual of  $Sym^N(C)$  looks like a tensionless string theory (or higher spin gravity).




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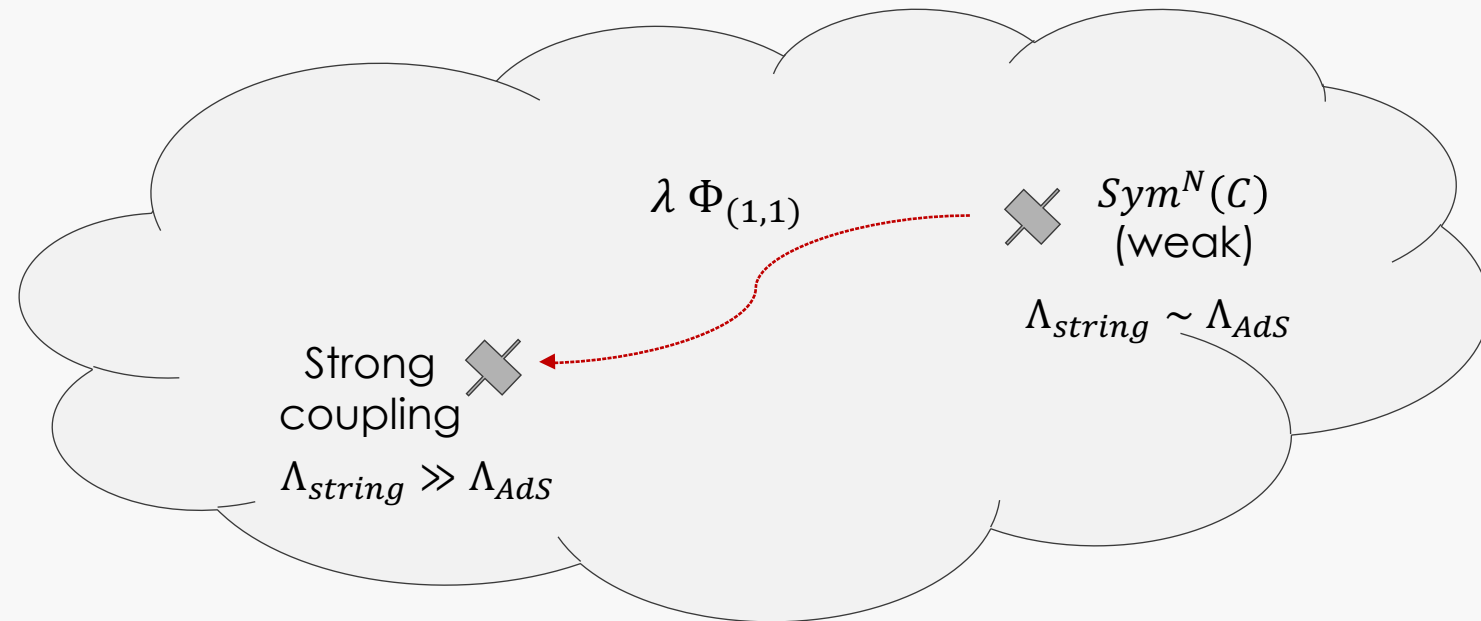


AdS/CFT interpretation: Dual of  $Sym^N(\mathcal{C})$  looks like a tensionless string theory (or higher spin gravity).

- Higher spin currents due to orbifold structure.
- Universal Hagedorn growth of light states. 🖱️

**Question:** Which  $Sym^N(\mathcal{C})$  could admit in their moduli space a dual supergravity point?

**Strategy:** Impose necessary conditions. Identify which  $Sym^N(\mathcal{C})$  comply with them.



Moduli space: set of exactly marginal deformations

## Holographic CFT<sub>2</sub>

Some requirements:

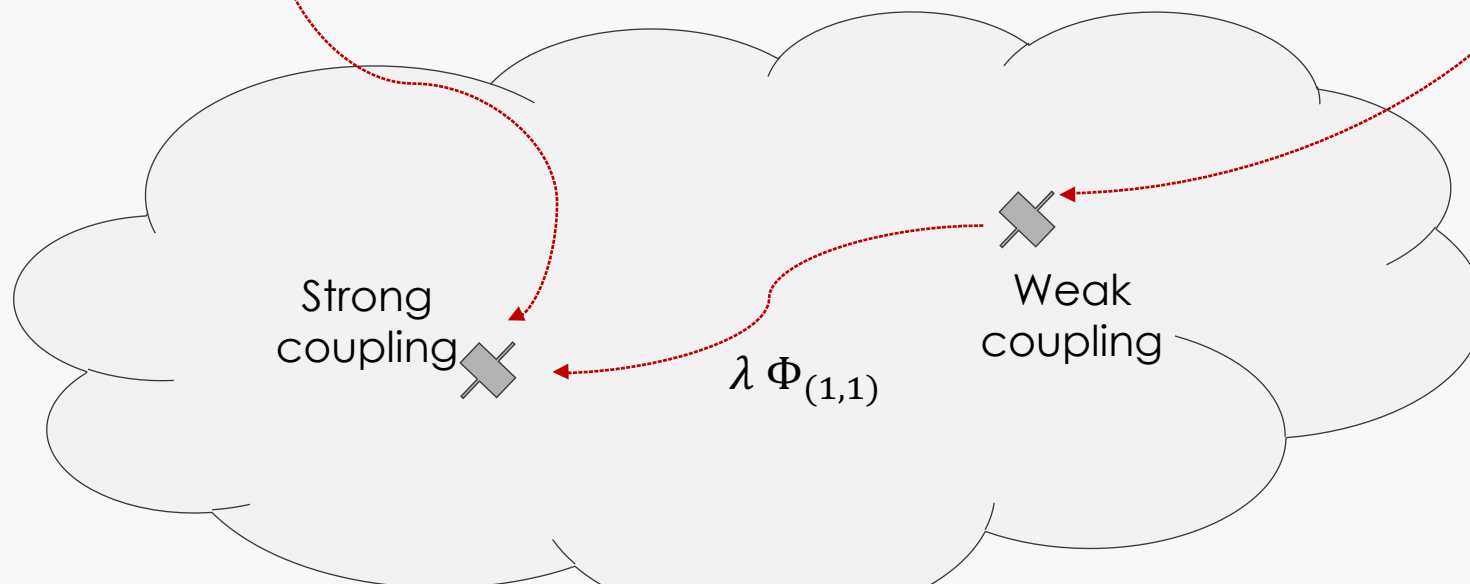
- Large-N:  $c = \frac{3\ell}{2G_N} \gg 1$
- Sparse spectrum
- Large gap spectrum
- ...

## Symmetric Product orbifolds

Focus on theories with at least N=(2,2).

At large-N, classify them according to:

- Moduli (deformation)
- BPS spectrum



Moduli space: set of exactly marginal deformations

# Necessary conditions

- **Criterion 1:** Existence of suitable moduli (single trace, twisted, BPS).
- **Criterion 2:** Sparseness condition on the elliptic genera (index that captures BPS states).

Based on these two criteria, we will classify  $Sym^N(\mathcal{C})$  theories, and label them as

Type I:  
Both criteria

Type II:  
Only criterion 1

Type III:  
Neither criteria

Type IV:  
Only criterion 2

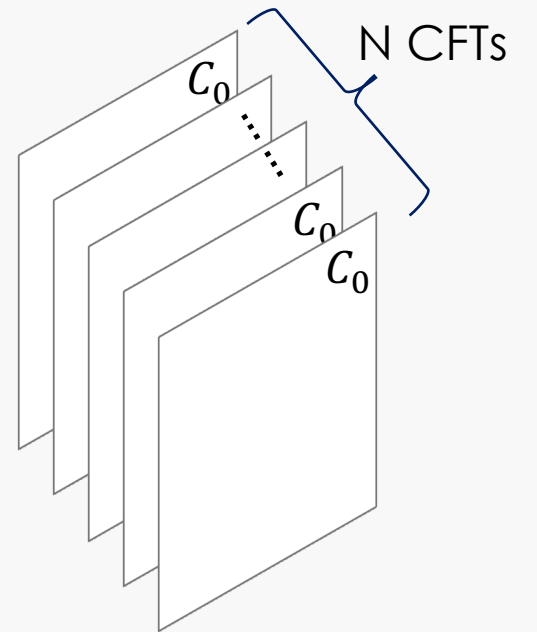


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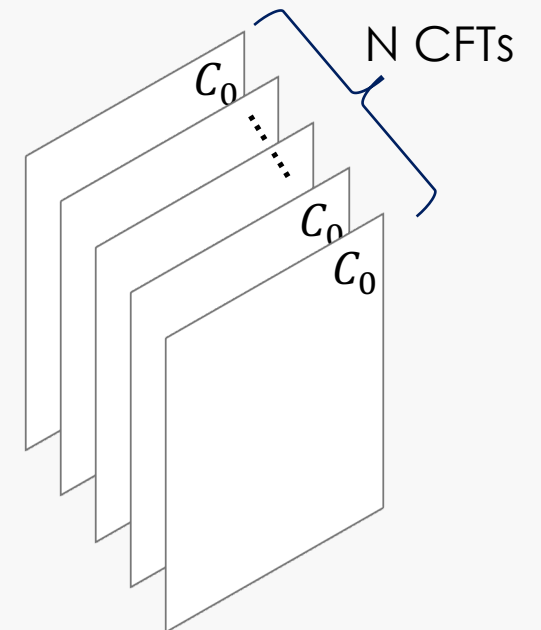
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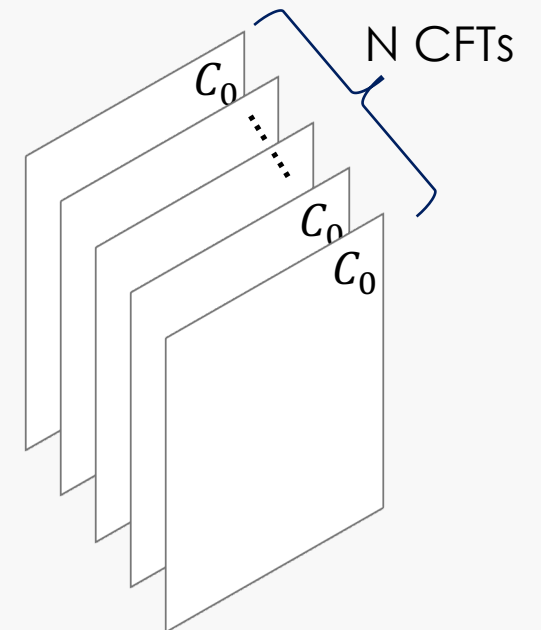
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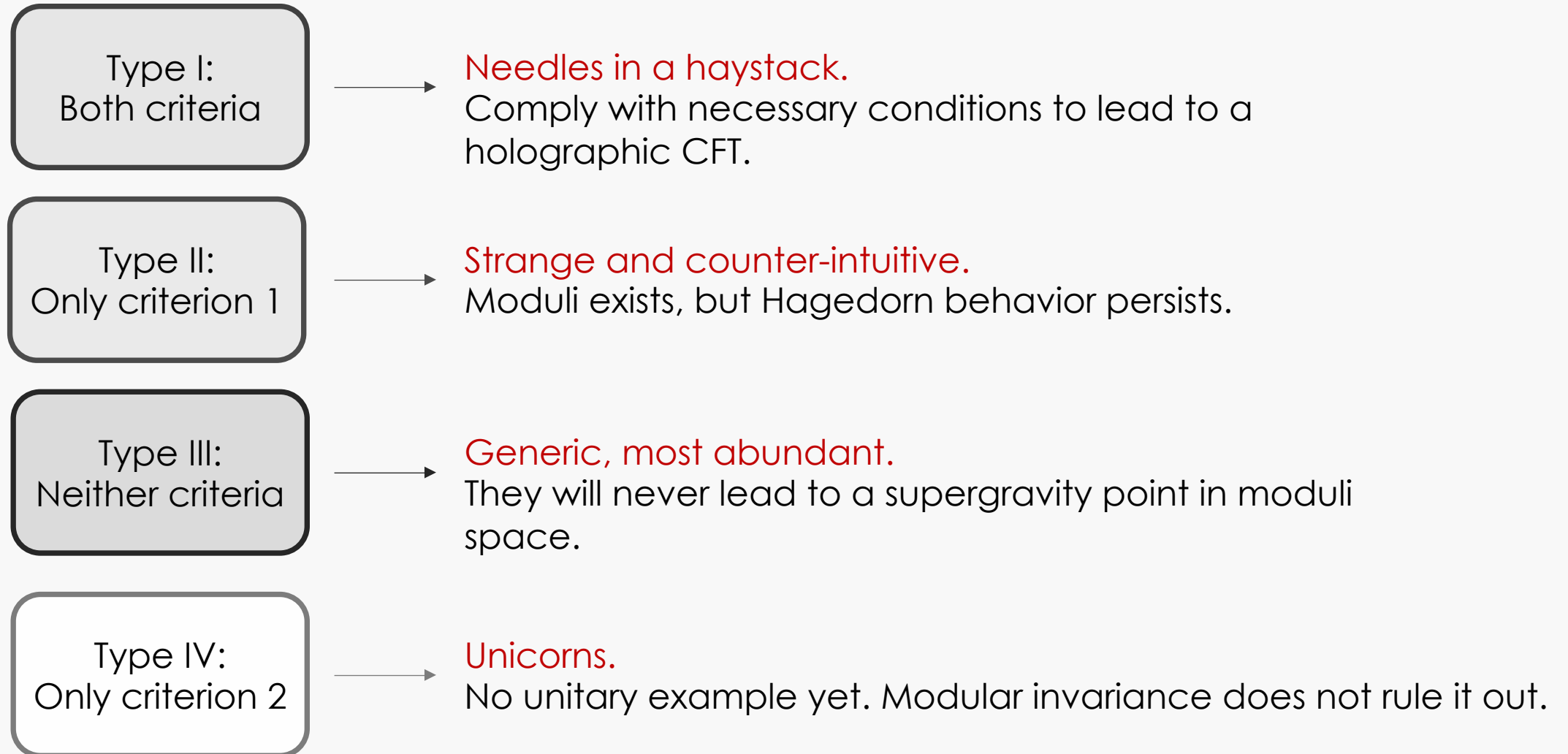
2. Criterion 2 can be done systematically and is exhaustive.

3. If Criterion 2 is satisfied, we proved that one always gets

$$d(\Delta) \sim e^{\sqrt{\Delta}} \quad \text{where} \quad \Delta \gg 1, \quad \Delta \sim O(N^0)$$



# Classification



# Classification

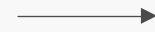
Type I:  
Both criteria



**Needles in a haystack.**

Comply with necessary conditions to lead to a holographic CFT.

Type II:  
Only criterion 1



**Strange and counter-intuitive.**

Moduli exists, but Hagedorn behavior persists



Type III:  
Neither criteria



**Generic, most abundant.**

They will never lead to a supergravity point in moduli space.

Type IV:  
Only criterion 2



**Unicorns.**

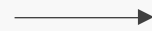
No unitary example yet. Modular invariance does not rule it out.

	Theory	Sparse?	Moduli?	Composition
	$A_6 \otimes A_{41}$	✓	✓	(11,88), (22,22)
	$A_7 \otimes A_{23}$	✓	✓	(11,55),(22,22)
	$A_8 \otimes A_{17}$	✓	✓	(11,44),(22,22)
	$A_9 \otimes A_{14}$	✓	✓	(22,22)
	$A_{11} \otimes A_{11}$	✓	✓	(11,33),(33,11),(22,22)
	$A_6 \otimes D_{22}$	✗	✗	
	$A_7 \otimes D_{13}$	✗	✓	(11,55)
	$A_{23} \otimes D_5$	✗	✓	(55,11)
	$A_8 \otimes D_{10}$	✗	✗	
	$A_{14} \otimes D_6$	✗	✗	
	$A_{11} \otimes D_7$	✓	✓	(11,33),(33,11)
	$A_8 \otimes E_7$	✗	✗	
Type II	$A_{11} \otimes E_6$	✗	✓	(33,11)
	$D_5 \otimes D_{13}$	✗	✓	(11,55)
	$D_7 \otimes D_7$	✓	✓	(11,33),(33,11)
Type I	$D_7 \otimes E_6$	✓	✓	(33,11)
	$E_6 \otimes E_6$	✗	✗	
	$A_2 \otimes A_5 \otimes A_5$	✓	✓	(11,11,22),(11,22,11)
	$A_2 \otimes A_5 \otimes D_4$	✓	✓	(11,22,11)
	$A_2 \otimes D_4 \otimes D_4$	✗	✗	
	$A_3 \otimes A_3 \otimes A_5$	✓	✓	(11,11,22)
	$A_3 \otimes A_3 \otimes D_4$	✗	✗	

Examples of theories where the seed has  $c_0 = 5$

# Comparisson

Type I:  
Both criteria



**Needles in a haystack.**

Comply with necessary conditions to lead to a holographic CFT.

Type II:  
Only criterion 1



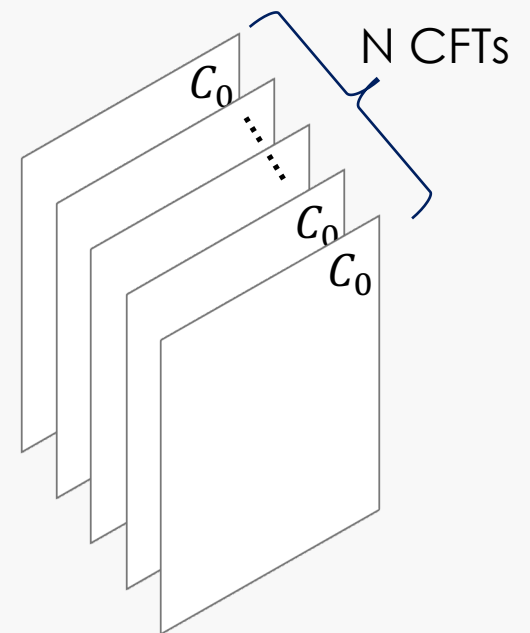
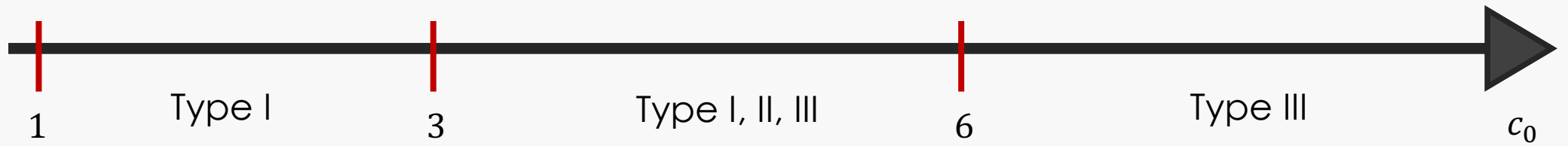
**Strange and counter-intuitive.**

Moduli exists, but Hagedorn behavior persists



- Evaluated anomalous dimension of several holomorphic operators (currents).
- Type I and II theories exhibit no difference at leading order in perturbation theory. 🤔
- What is the key feature that guarantees a supergravity point in moduli space?

# Summary



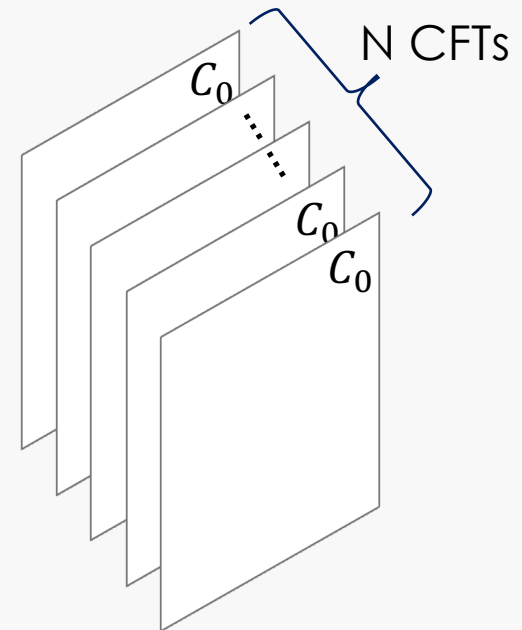


# Summary



Comments:

- Only consider CFTs that are unitary and compact.
- Assume that the elliptic genus does not vanish.
- D1D5 on K3 sits at  $c_0 = 6$ .
- Search between  $1 \leq c_0 < 3$  is exhaustive: N=2 Minimal Models.
- Search between  $3 \leq c_0 \leq 6$  is not exhaustive (but systematic).



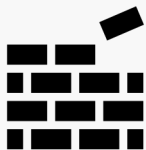
# Type I: Examples

Series	$k$	untwisted moduli	twisted moduli	single trace twisted
$A_2$	1	1	28	1 twist 5, 1 twist 7
$A_3$	2	3	26	1 twist 3, 1 twist 4, 1 twist 5
$A_5$	4	9	24	1 twist 2, 1 twist 3, 1 twist 4
$A_{k+1}$	odd, $\geq 3$	$P(k+2) - 2$	9	1 twist 3
$A_{k+1}$	even, $\geq 6$	$P(k+2) - 2$	$10 + \sum_{r=1}^{\frac{k}{2}+2} P(r)$	1 twist 2, 1 twist 3
$D_4$	4	6	20	1 twist 2, 2 twist 3, 1 twist 4
$D_{\frac{k}{2}+2}$	$0 \pmod 4, \geq 8$	$P(\frac{k}{2}+1) + P(\frac{k}{4}+1)$	$8 + \sum_{r=1}^{\frac{k}{4}+1} P(r)$	1 twist 2, 1 twist 3
$D_{\frac{k}{2}+2}$	$2 \pmod 4, \geq 6$	$P(\frac{k}{2}+1)$	7	1 twist 3
$E_6$	10	4	5	1 twist 2
$E_7$	16	6	5	1 twist 2
$E_8$	28	6	5	1 twist 2

## N=2 Virasoro Minimal Models

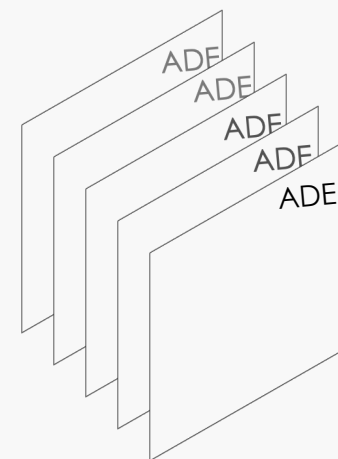
$$c_0 = \frac{3k}{k+2} < 3$$

where  $k = 1, 2, \dots$



Necessary conditions:

- Criterion 1: Exactly marginal operator
- Criterion 2: Sparse spectrum for elliptic genera



# Type I: Examples

Series	$k$	untwisted moduli	twisted moduli	single trace twisted
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$D_4$	4	6	20	1 twist 2, 2 twist 3, 1 twist 4
$D_{\frac{k}{2}+2}$	$0 \pmod 4, \geq 8$	$P(\frac{k}{2} + 1) + P(\frac{k}{4} + 1)$	$8 + \sum_{r=1}^{\frac{k}{4}+1} P(r)$	1 twist 2, 1 twist 3
$D_{\frac{k}{2}+2}$	$2 \pmod 4, \geq 6$	$P(\frac{k}{2} + 1)$	7	1 twist 3
$E_6$	10	4	5	1 twist 2
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$E_8$	28	6	5	1 twist 2

Responsible of lifting most states.  
Breaks higher spin symmetry

# Type I: Examples

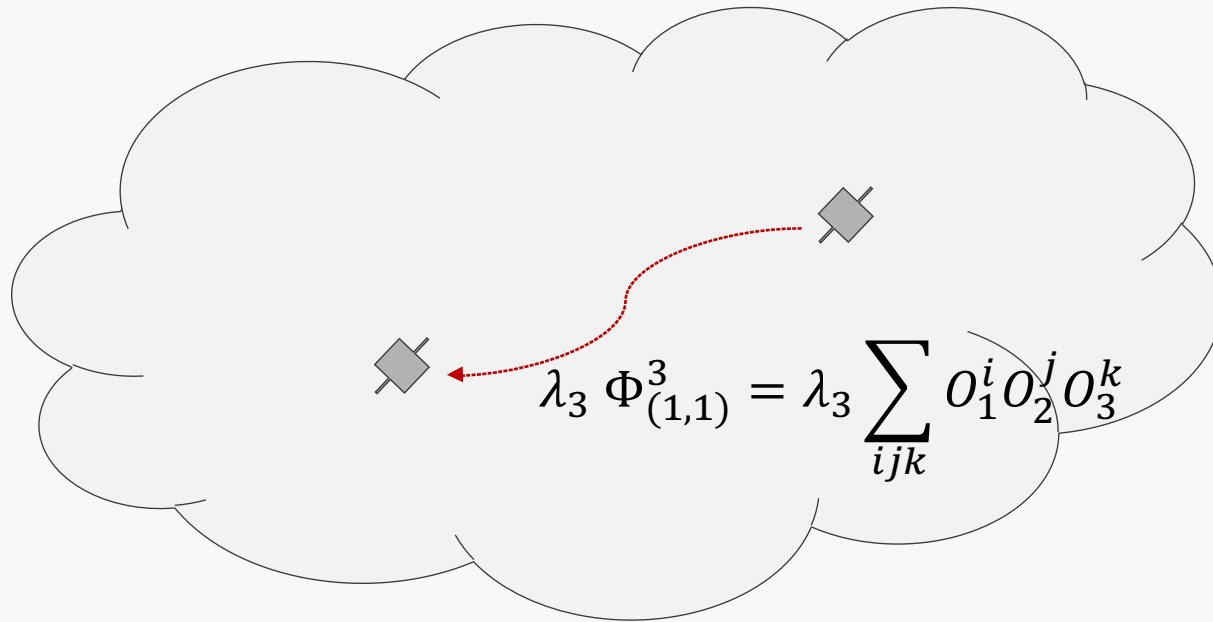
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Multi-trace deformations.

Formally known, first explicit example of CFT with these BPS deformations. They can turn on couplings, which hints to strongly coupled matter.

# Destroy Factorization

Consider any CFT that complies with **Large-N Factorization**



To make sure that the deformation is robust, we ask that  $\Phi_{(1,1)}^3$  is:

- Marginal
- $\frac{1}{2}$ -BPS

Note: it can be either twisted or untwisted.

The coupling  $\lambda$  is independent of N.

# Destroy Factorization

Consider any CFT that complies with **Large-N Factorization**.

Next deform the theory by  $\Phi_{(1,1)}^3 = \sum_{ijk} O_1^i O_2^j O_3^k$

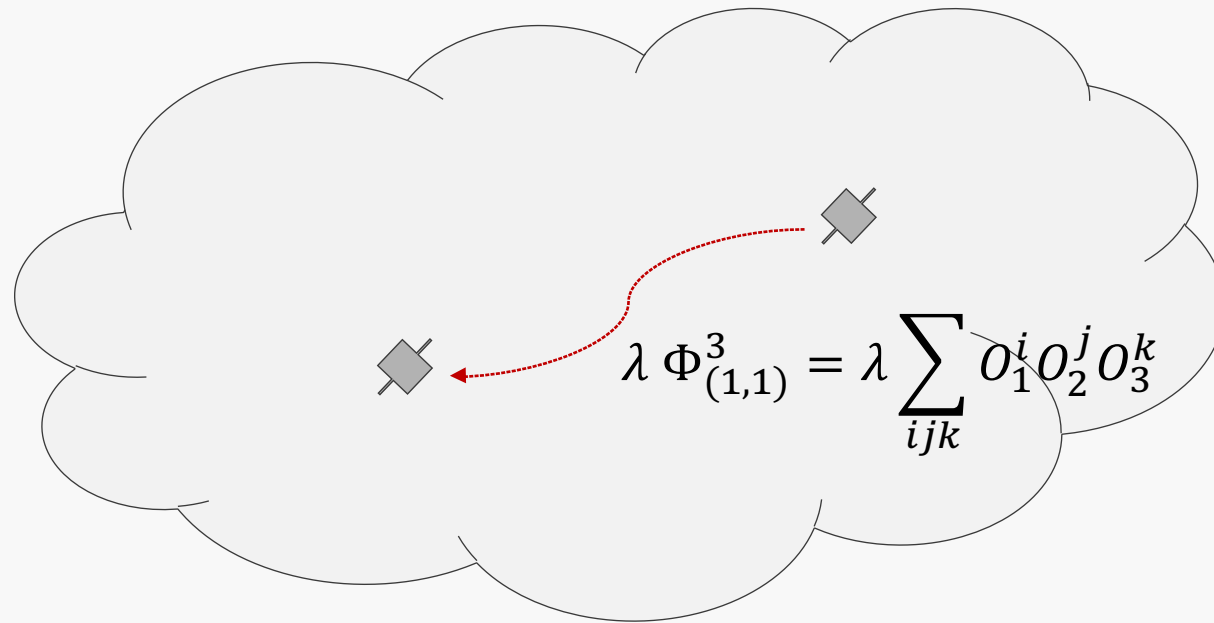
$$\langle O_1(z) O_2(w) O_3(w) \rangle_\lambda = \underbrace{\langle O_1(z) O_2(w) O_3(w) \rangle_0}_{\text{Suppressed by } O(N^{-\frac{1}{2}})} + \lambda_3 \int \underbrace{\langle O_1(z) O_2(w) O_3(w) \Phi_{1,1}^3(w) \rangle_0}_{\text{Leading order } O(N^0)} dw + \dots$$

Suppressed by  $O(N^{-\frac{1}{2}})$   
Large-N factorization

Leading order  $O(N^0)$   
Introduced a new coupling!

# Destroy Factorization

Consider any CFT that complies with **Large-N Factorization**



$$\langle O_1 O_2 O_3 \rangle_\lambda \sim \lambda N^0 + \dots$$

- Breaks large-N factorization
- Interactions that are not controlled by  $G_N$
- Type I theories have these deformations
- Argument is general: applies to  $\text{CFT}_D$

The coupling  $\lambda$  is independent of  $N$ .

This deformation does not affect the convergence of the large-N limit (observables finite).

# AdS<sub>3</sub>/CFT<sub>2</sub>

Typical theory: Couplings controlled by  $G_N$

$$I_{3D} = \frac{1}{16\pi G_N} \int d^3x \sqrt{-g} \left( R + \frac{2}{\ell^2} - (\partial\Phi)^2 - \frac{\hat{\lambda}_3}{\ell^2} \Phi^3 \right) + \dots$$

New flavors: multi-trace deformations could turn on independent couplings.

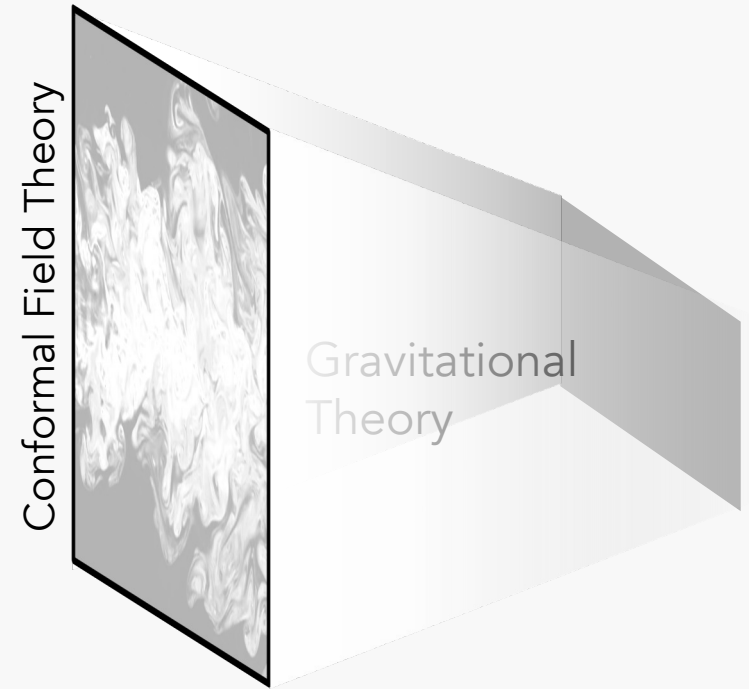
$$I_{3D} = \frac{1}{16\pi G_N} \int d^3x \sqrt{-g} \left( R + \frac{2}{\ell^2} \right) + \int d^3x \sqrt{-g} \left( -(\partial\Phi)^2 - \frac{\lambda_3}{\ell} \Phi^3 \right)$$



Outlook

Quantify the space of type I theories:

- Different from known examples
- Systematic and tractable
- Infinite family



Type I  $Sym^N(C)$

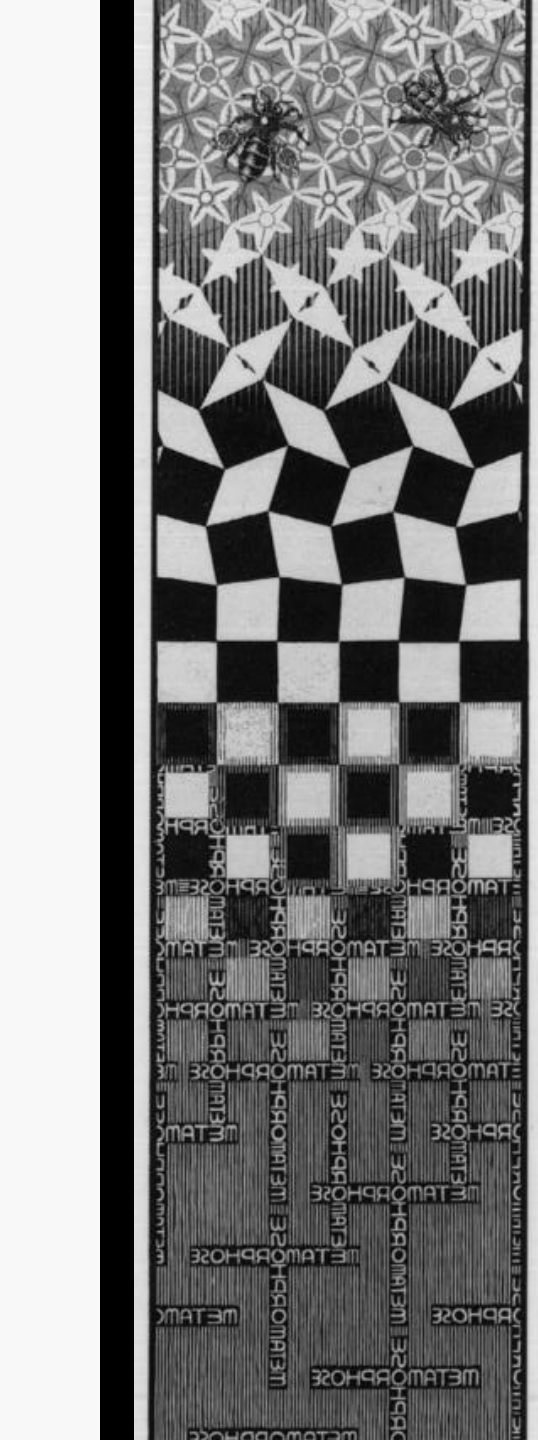
Conditions:

- Large-N
- Sparse elliptic genera
- Moduli

Holographic  $CFT_2$

Some requirements:

- Large-N
- Sparse spectrum
- Large gap spectrum

- 
- Which CFTs capture classical (geometric) properties of gravity?
  - What are possible theories of quantum gravity that can be designed?
  - What are the materials needed to assemble them?

Next steps:

- Multi-trace deformations (to appear by Apolo, Belin and Bintanja).
- String theory and supergravity description.
- Heavy states: contrast black holes among type I, II and III.
- Mock Modularity and asymptotic expansions.
- Type I vs II: lifting of generic operators.