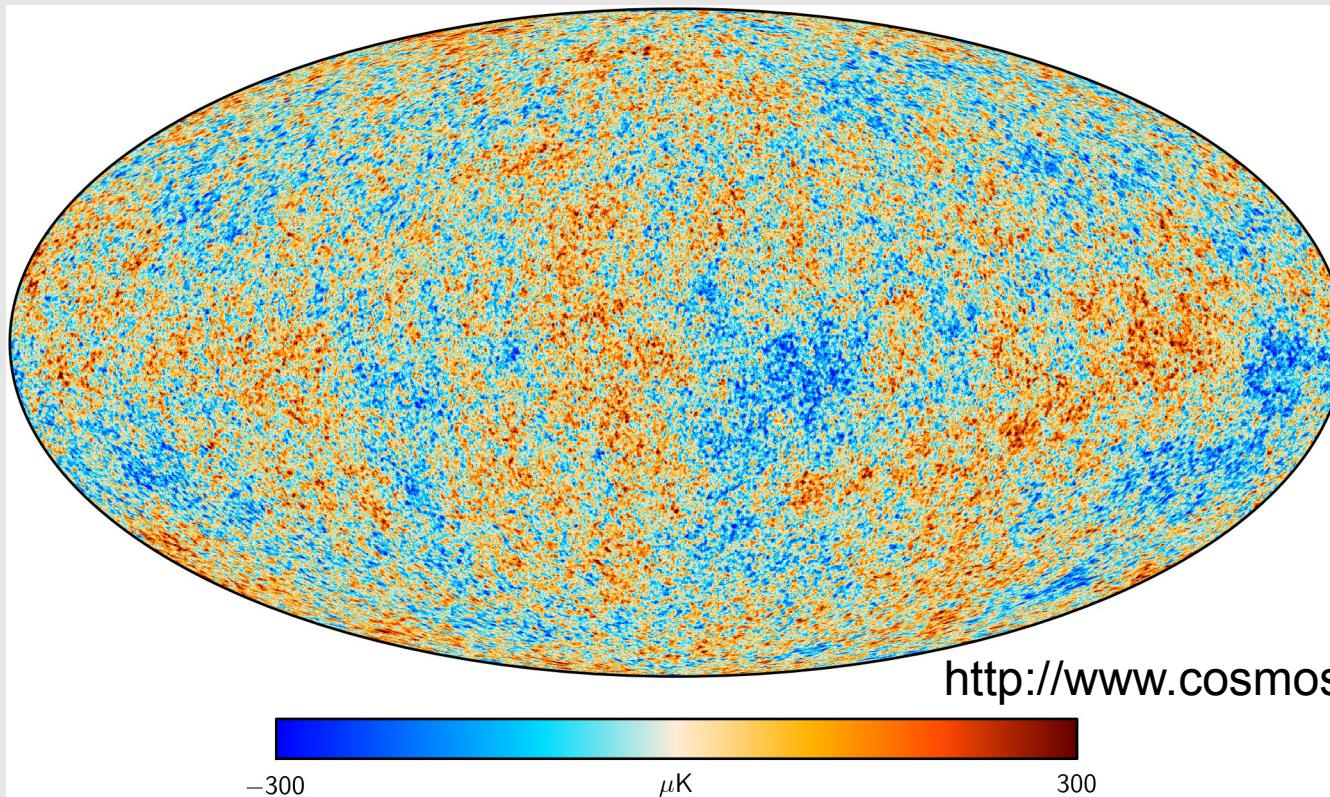


Cosmology: Theory

(PLANCK and all that)



- M. Bastero-Gil
- Dpto. Física Teórica y del Cosmos

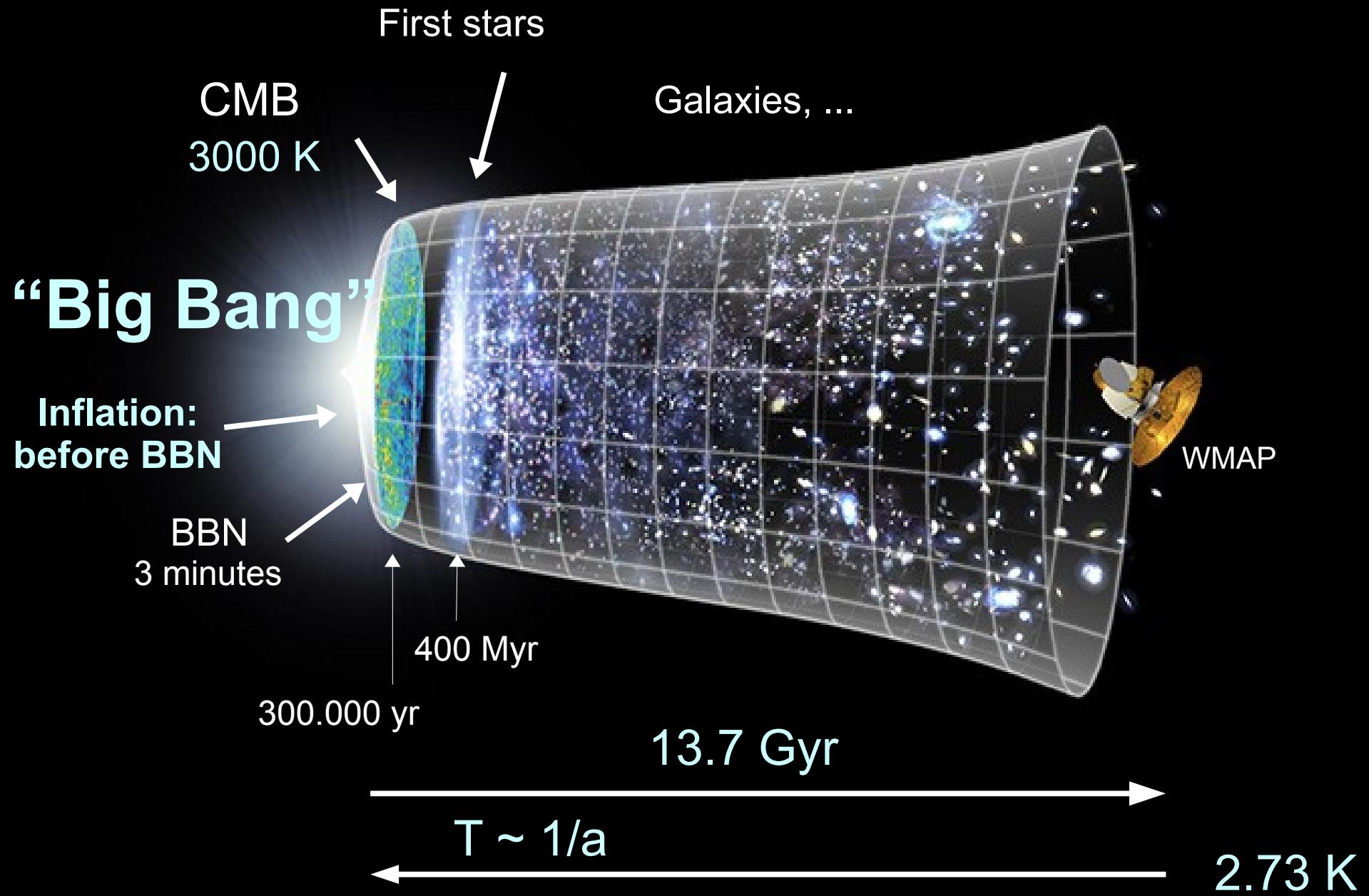


ugr

Universidad
de Granada

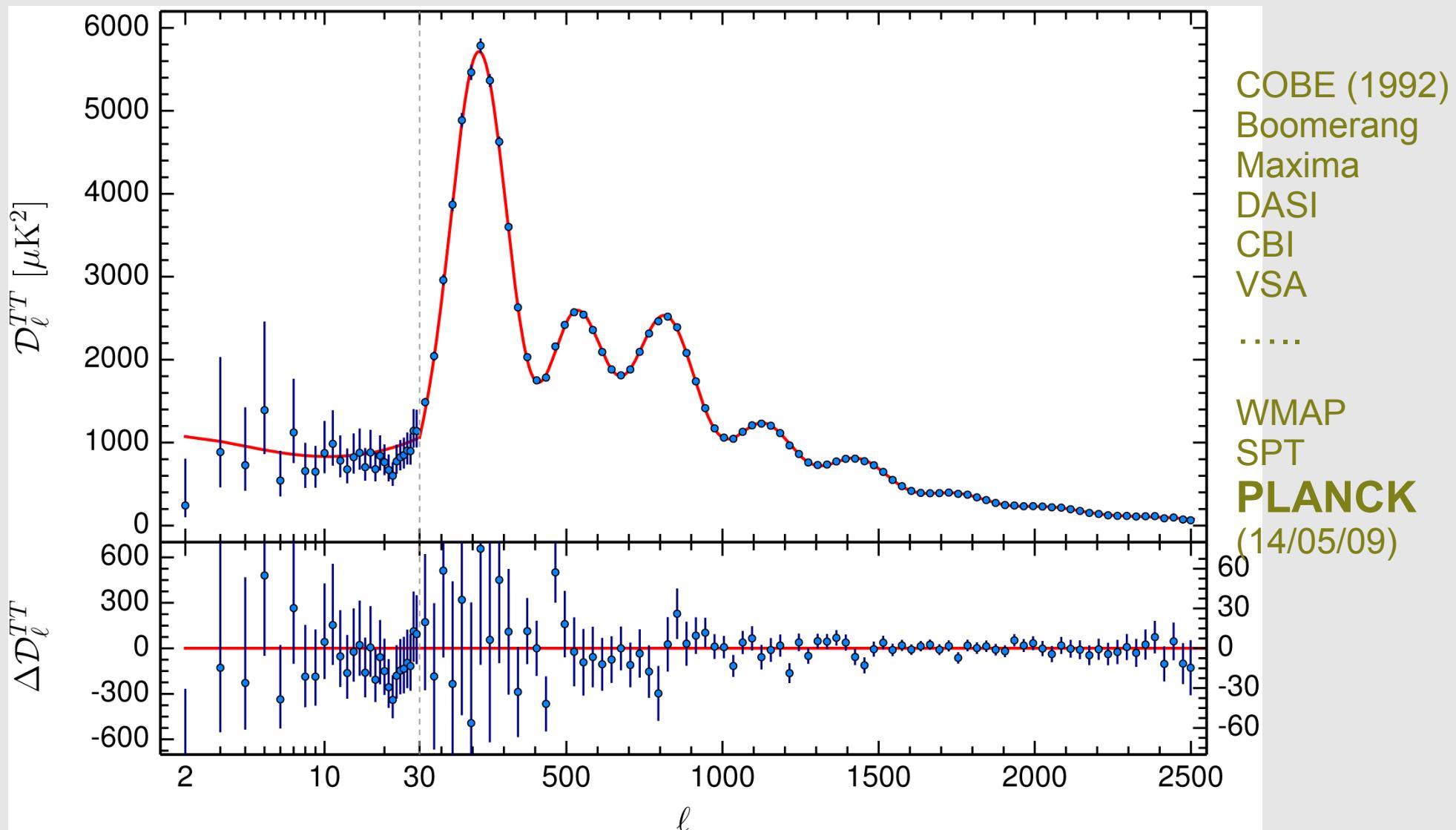


The expanding Universe



Cosmic microwave background radiation (CMB)

Spectrum of T fluctuations

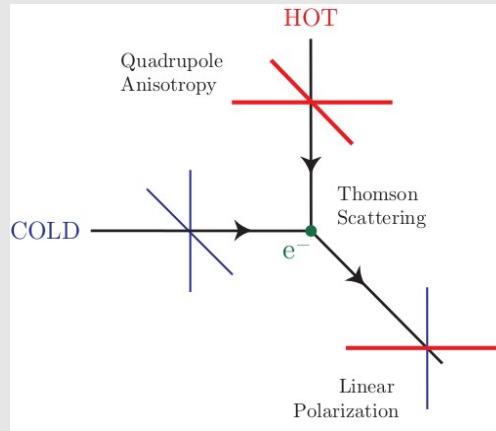


[Planck collab.: astro-ph/1502.01589]

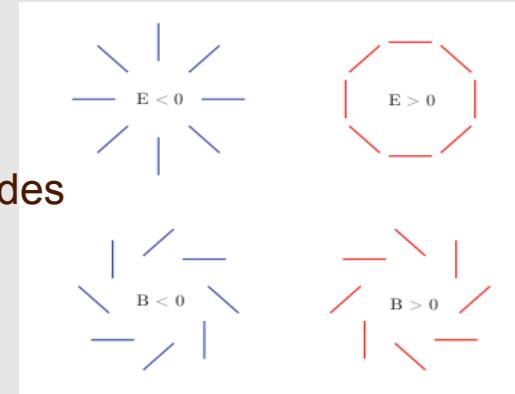
Polarization & Tensors (primordial GW)

Thompson scattering polarizes linearly the spectrum

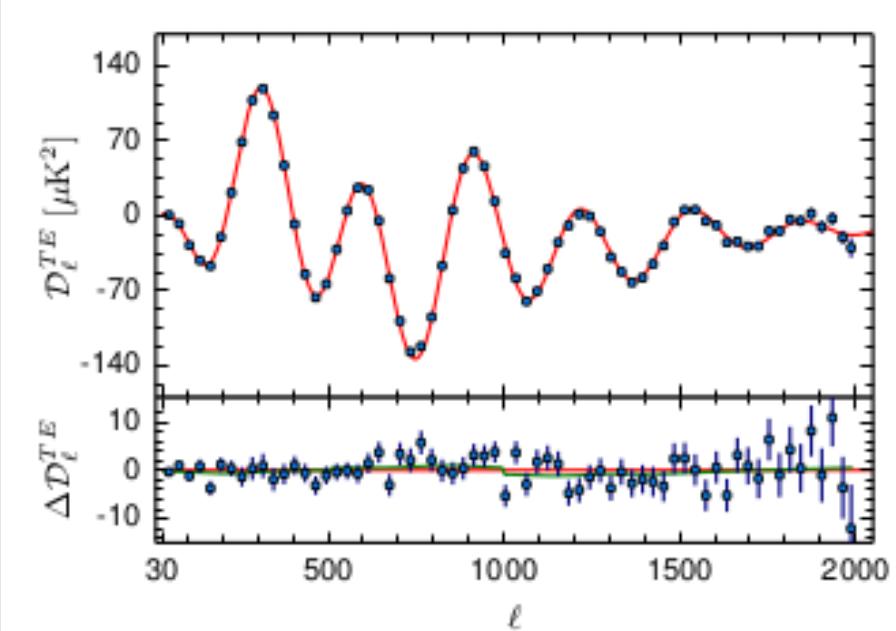
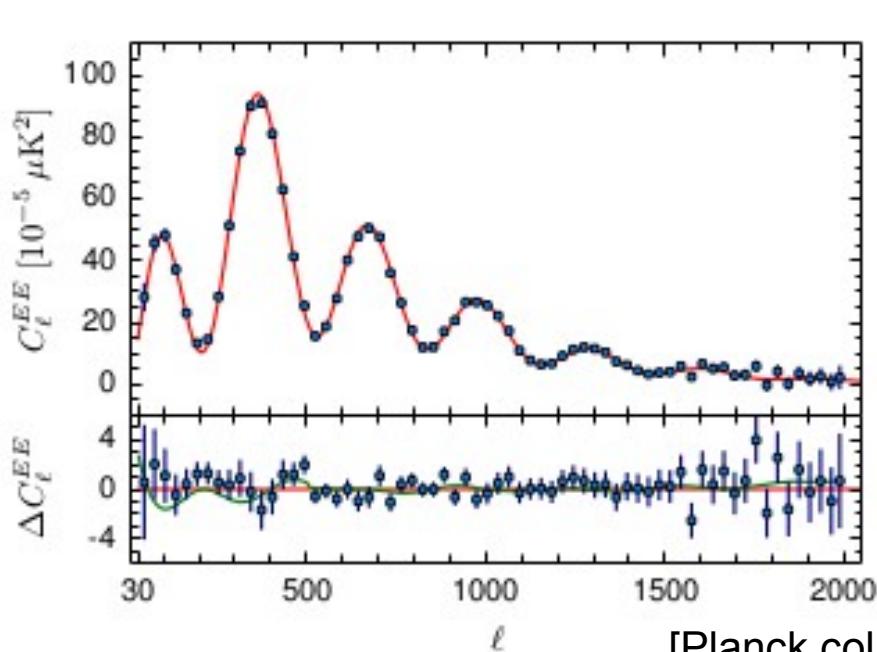
(DASI 2002)



Polarization: E & B modes



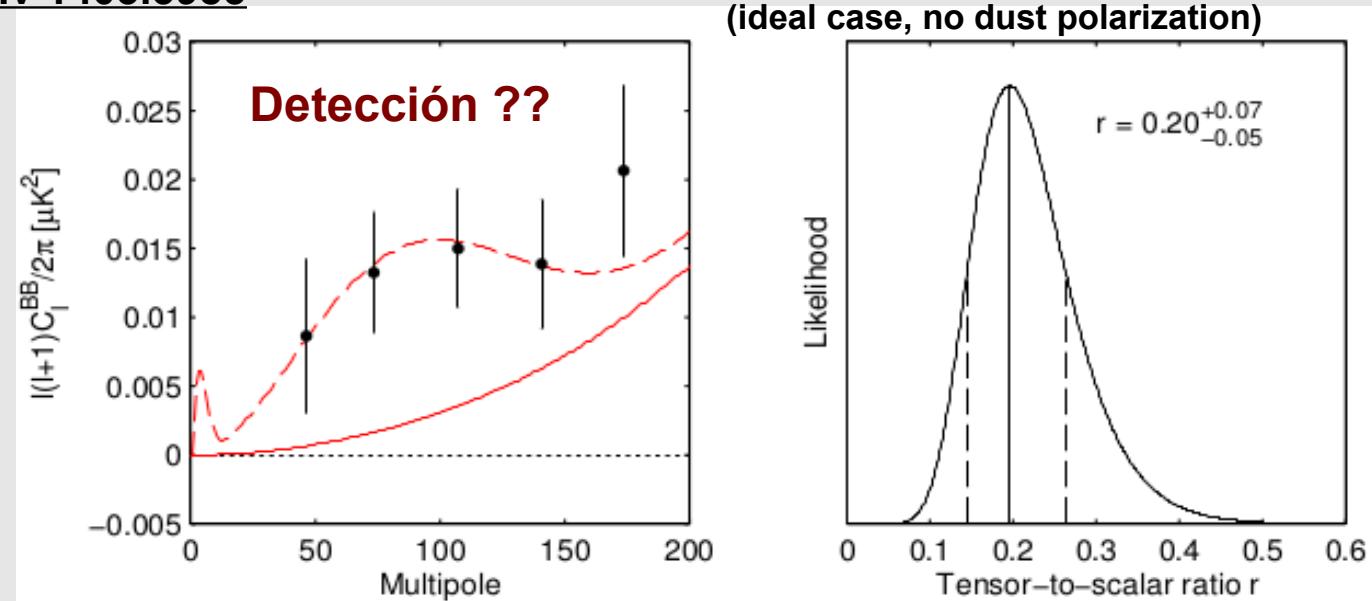
Spectra: C_{ℓ}^{TT} , C_{ℓ}^{EE} , C_{ℓ}^{TE} , C_{ℓ}^{BB} Tensors?



[Planck collab.: astro-ph/1502.01589]

BICEP2 & PLANCK

BICEP2: arXiv 1403.3985



So far, only galactic dust: $r_{0.05} < 0.12$ (95 % CL)

BICEP/KECK & PLANCK: arXiv 1502.00612

$r = 0.16^{+0.06}_{-0.05} ???$

Λ CDM Model: 6 parameters that fit the CMB

$$\left[P_R(k_0), n_s, \Omega_B, \Omega_M, \Omega_\Lambda, \tau_{\text{reio}} \right]$$

Primordial spectrum Matter content Optical depth to LSS due to reionization

$$P_R(k) = P_R(k_0) \left(\frac{k}{k_0} \right)^{n_s - 1}$$

Baryons+CDM+DE ($w_\Lambda = -1$)

Parameter	[1] <i>Planck</i> TT+lowP	[2] <i>Planck</i> TE+lowP	[3] <i>Planck</i> EE+lowP	[4] <i>Planck</i> TT,TE,EE+lowP	([1] - [4])/ $\sigma_{[1]}$
$\Omega_b h^2$	0.02222 ± 0.00023	0.02228 ± 0.00025	0.0240 ± 0.0013	0.02225 ± 0.00016	-0.1
$\Omega_c h^2$	0.1197 ± 0.0022	0.1187 ± 0.0021	$0.1150^{+0.0048}_{-0.0055}$	0.1198 ± 0.0015	0.0
$100\theta_{\text{MC}}$	1.04085 ± 0.00047	1.04094 ± 0.00051	1.03988 ± 0.00094	1.04077 ± 0.00032	0.2
τ	0.078 ± 0.019	0.053 ± 0.019	$0.059^{+0.022}_{-0.019}$	0.079 ± 0.017	-0.1
$\ln(10^{10} A_s)$	3.089 ± 0.036	3.031 ± 0.041	$3.066^{+0.046}_{-0.041}$	3.094 ± 0.034	-0.1
n_s	0.9655 ± 0.0062	0.965 ± 0.012	0.973 ± 0.016	0.9645 ± 0.0049	0.2
H_0	67.31 ± 0.96	67.73 ± 0.92	70.2 ± 3.0	67.27 ± 0.66	0.0
Ω_m	0.315 ± 0.013	0.300 ± 0.012	$0.286^{+0.027}_{-0.038}$	0.3156 ± 0.0091	0.0
σ_8	0.829 ± 0.014	0.802 ± 0.018	0.796 ± 0.024	0.831 ± 0.013	0.0
$10^9 A_s e^{-2\tau}$	1.880 ± 0.014	1.865 ± 0.019	1.907 ± 0.027	1.882 ± 0.012	-0.1

FLRW metric (homogeneous & isotropic): $a(t)$ =scale factor, $H=\frac{da}{dt}$ Hubble parameter

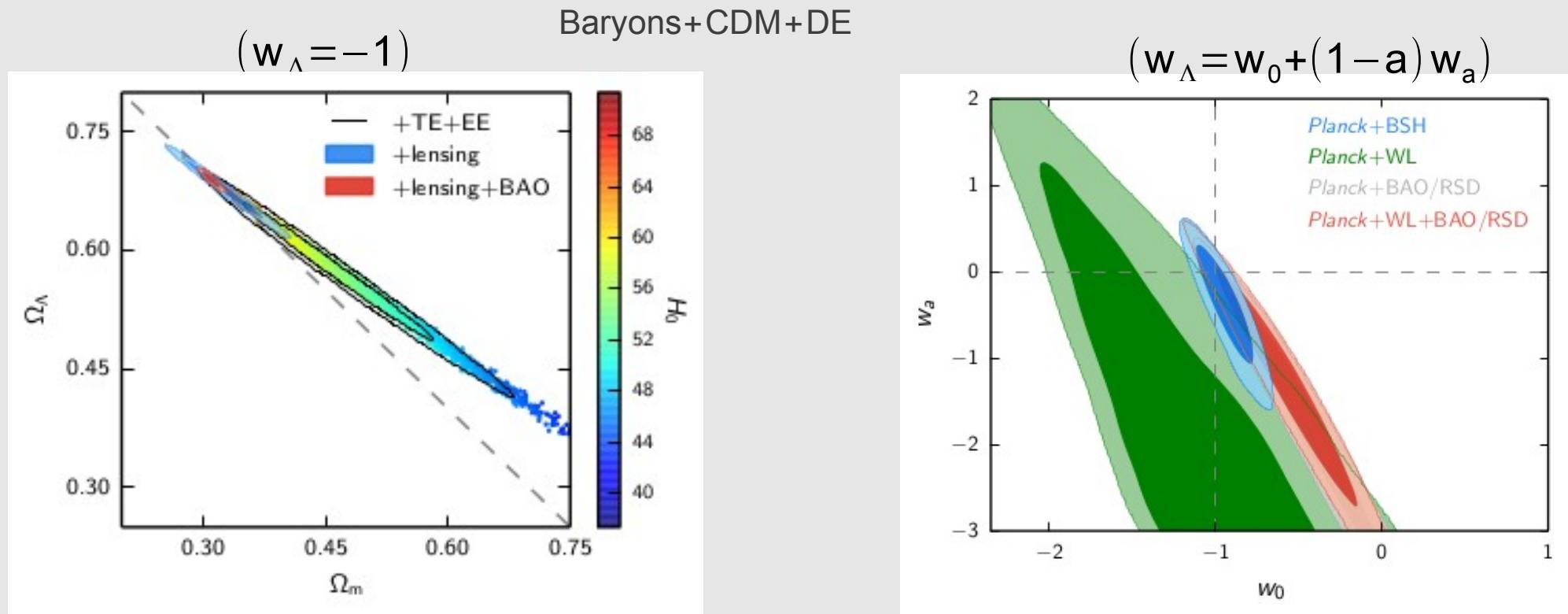
Flat Universe: $\Omega_M + \Omega_\Lambda = 1$

Primordial Fluctuations~ adiabatic, gaussian, scale-invarian spectrum

Λ CDM Model: 6 parameters that fit the CMB

$P_R(k_0), n_s,$ $\Omega_B, \Omega_M, \Omega_\Lambda,$ τ_{reio}

Primordial spectrum Matter content Optical depth to LSS due to reionization



Background \sim Flat Universe + “Cosmological constant”

FLRW metric (homogeneous & isotropic): $a(t)$ =scale factor), $H=d\ln a/dt$ Hubble parameter

Flat Universe: $\Omega_M + \Omega_\Lambda = 1$

[Planck collab.: astro-ph/1502.01589]

Cosmological Principle:

The Universe is homogeneous and isotropic at large scales

$$l_{\text{hom}} > 100 h^{-1} \text{hMpc}$$

Sarkar et al. 0906.3431 [SDSS DR6]; WiggleZ 1205.6812

Problems of the Big Bang Model

Homogeneous & isotropic  Horizon problem

The universe is spatially flat  Why?

$\Delta T/T \sim 10^{-5}$  Primordial fluctuations?

“Relics”: monopoles, topological defects, exotic particles... ?

Solution: Inflation

Period of accelerated expansion

$$\ddot{a} > 0 \iff \rho + 3p < 0$$

Energy & pressure

Flat Potential

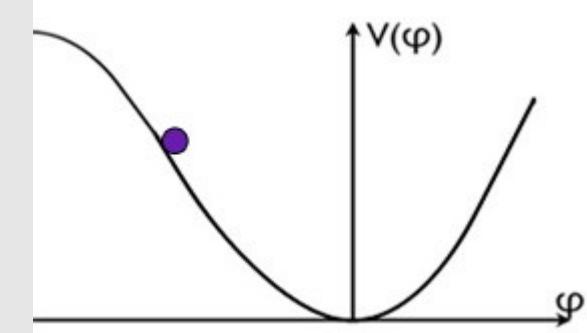
“Slow-roll” parameters

“slow-roll” evolution

Inflation: “slow-roll”

Kinetic energy << potential

$$\rho \approx V(\varphi) \quad p \approx -V(\varphi)$$

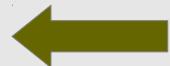


Negative pressure

The curvature and the slope of the potential less than the Hubble rate

$$H^2 \sim V/3m_P^2$$

$$|\eta_\varphi| = m_P^2 \left| \frac{d^2 V / d\varphi^2}{V} \right| < 1 \quad \epsilon_\varphi = \frac{m_P^2}{2} \left(\frac{dV / d\varphi}{V} \right)^2 < 1$$



curvature

slope

$$3H\dot{\varphi} \simeq -dV/d\varphi, \quad \ddot{\varphi} \ll 3H\dot{\varphi}, \quad \dot{\varphi}^2 \ll V(\varphi)$$

($\dot{\varphi} \equiv d\varphi/dt$)

CMB: T anisotropy spectrum

(background + linear fluctuations)

comoving scales

$$ds^2 = dt^2 - a(t)^2(\delta_{ij}(1+2R) + h_{ij})dx^i dx^j$$

curvature $R \approx \frac{H}{\dot{\varphi}} \delta\varphi$

sub-horizon

$$\langle R_k R_{k'} \rangle$$

\hat{R}_k
zero-point
fluctuations

horizon exit

$$\dot{R} \approx 0$$

super-horizon

$$|\delta\varphi|^2 \approx \left(\frac{H_{\text{ex}}}{2\pi}\right)^2$$

horizon re-entry

$$(aH)^{-1}$$

$$\Delta T \xrightarrow{\text{projection}} C_\ell$$

$$\Delta_{\text{TI}}(k)$$

CMB
recombination

today

time

$$\Delta T/T$$



$$C_l^{\text{TT}} = \frac{2}{\pi} \int k^2 dk P_R(k) \Delta_{\text{TI}}(k) \Delta_{\text{TI}}(k)$$

Primordial spectrum

$$(H_0, \Omega_i)$$

Recipe for inflation

- Choose a potential: $V(\phi)$ (model building: Particle physics)

- Get “slow-roll” parameters and no. of efolds

(Observable universe: $N_e \sim 50-60$)

- Amplitude of the primordial spectrum (input):

$$P_R \simeq (H/\dot{\phi})^2 P_{\delta\phi} \simeq (5 \times 10^{-5})^2 \quad (\text{CMB: COBE, WMAP, PLANCK})$$

(Fix one parameter of the potential)

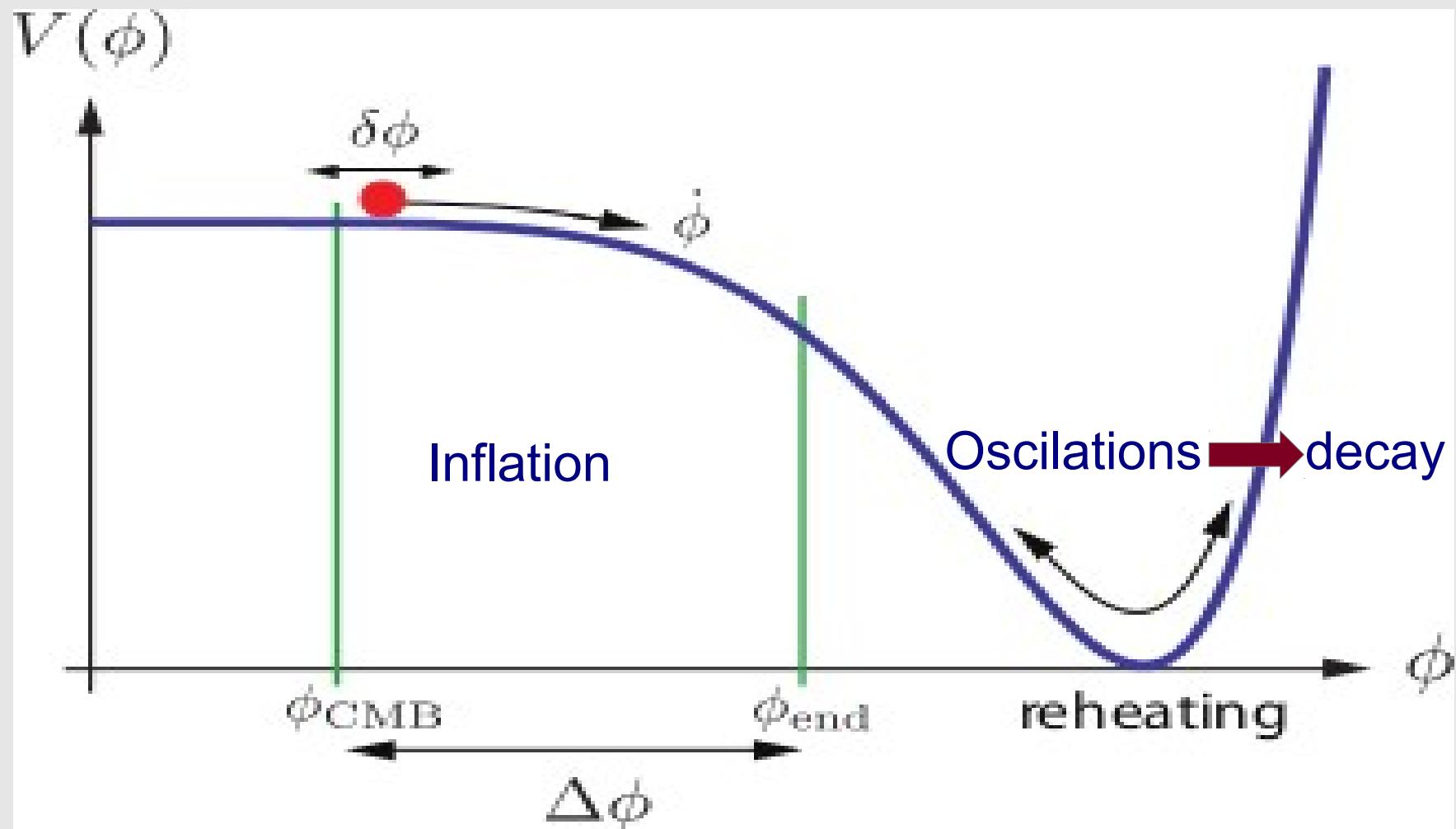
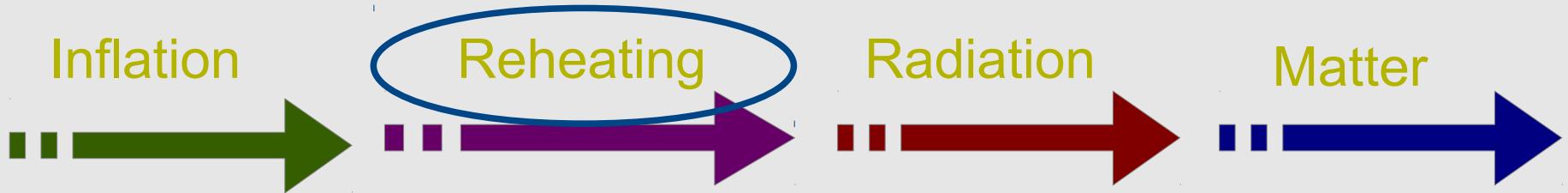
- Predictions (output): spectral index , tensor-to-scalar ratio,

$$n_s = \frac{d \ln P_R}{d \ln k} \simeq \frac{d \ln P_R}{d \ln a} \simeq 1 + 2 n_\phi - 6 \epsilon_\phi \quad (\text{spectral index}) \quad P_R = P_R(k_0) (k/k_0)^{n_s - 1}$$

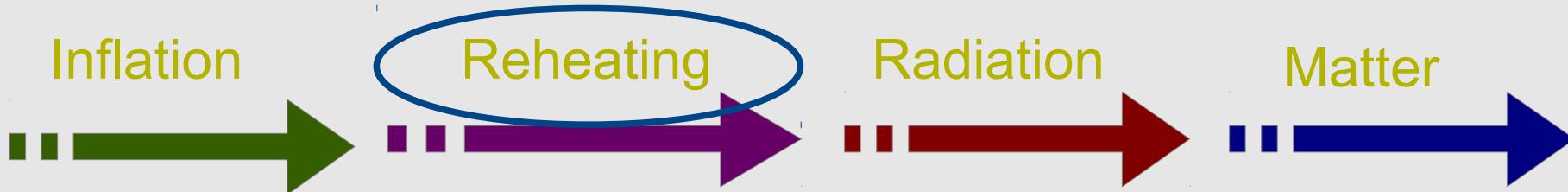
$$P_T = 8 \left(\frac{H}{2\pi m_P} \right)^2, \quad r = \frac{P_T}{P_R} < 1 \quad (\text{tensors : primordial gravitational waves})$$

- Compare with observations (CMB)

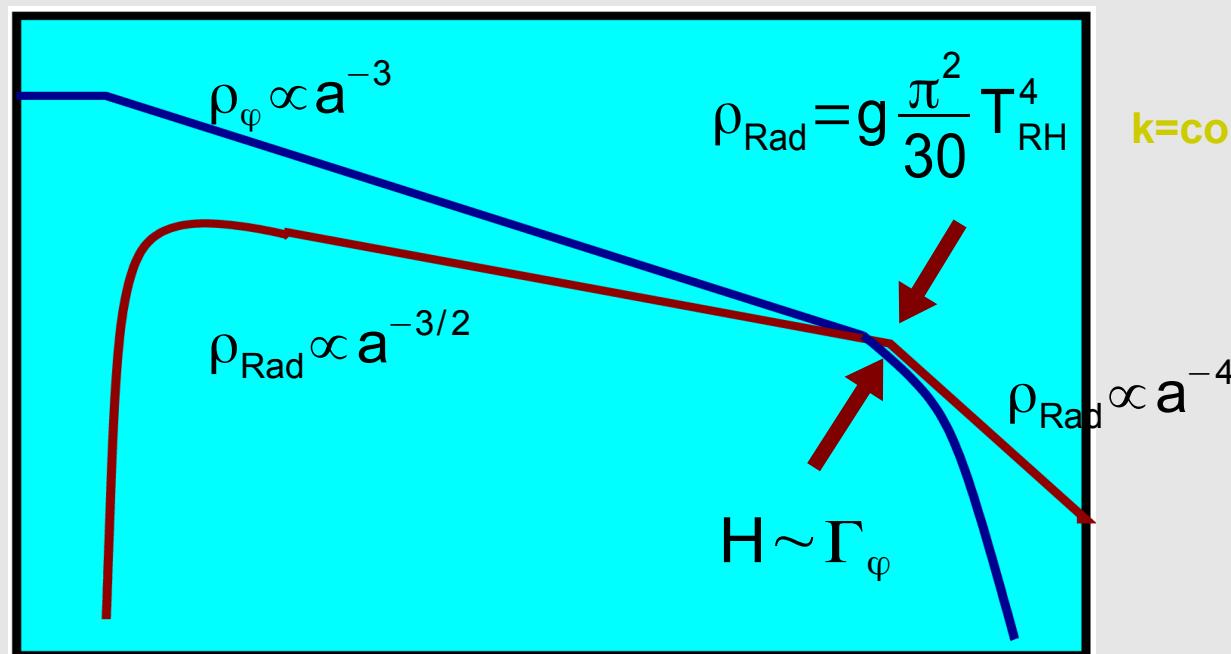
No. de e-folds: $N(k) = \ln \frac{a_{\text{end}}}{a_k}$



No. of e-folds: $N(k) = \ln \frac{a_{\text{end}}}{a_k}$



Horizon exit $k = \text{comoving wavenumber}$



$$N_e \simeq 56 + \frac{2}{3} \ln \frac{V_{\text{inf}}^{1/4}}{10^{15} \text{ GeV}} + \frac{1}{3} \ln \frac{T_{\text{RH}}}{10^9 \text{ GeV}} \sim 60 - 40$$

Primordial spectrum: $P_R = P_R(k_0) (k/k_0)^{n_s - 1}$

adiabatic, gaussian, scale-invariant spectrum

No evidence for:
non-gaussianity, isocurvature modes or running of the spectral index

- **Bispectrum:** $B_R(k_1, k_2, k_3) = \sum_{\text{cyc}} \langle R_1(k_1) R_1(k_2) R_2(k_3) \rangle = A_B(k) \bar{B}(k_1, k_2, k_3)$

$$f_{NL} = \frac{18}{5} \frac{A_B(k)}{P_R(k)^2} \quad \text{Non-linear parameter} \quad \text{shape}$$

$$f_{NL}^{\text{local}} = 2.5 \pm 5.7 \quad (k_3 \ll k_1 \sim k_2)$$

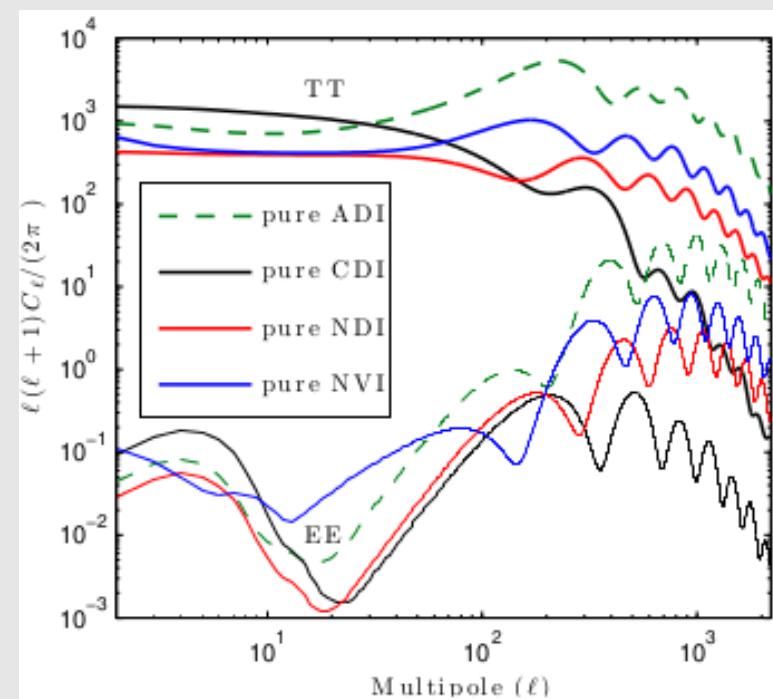
$$f_{NL}^{\text{equil}} = -16 \pm 70 \quad (k_3 \sim k_1 \sim k_2)$$

$$f_{NL}^{\text{orth}} = -34 \pm 33$$

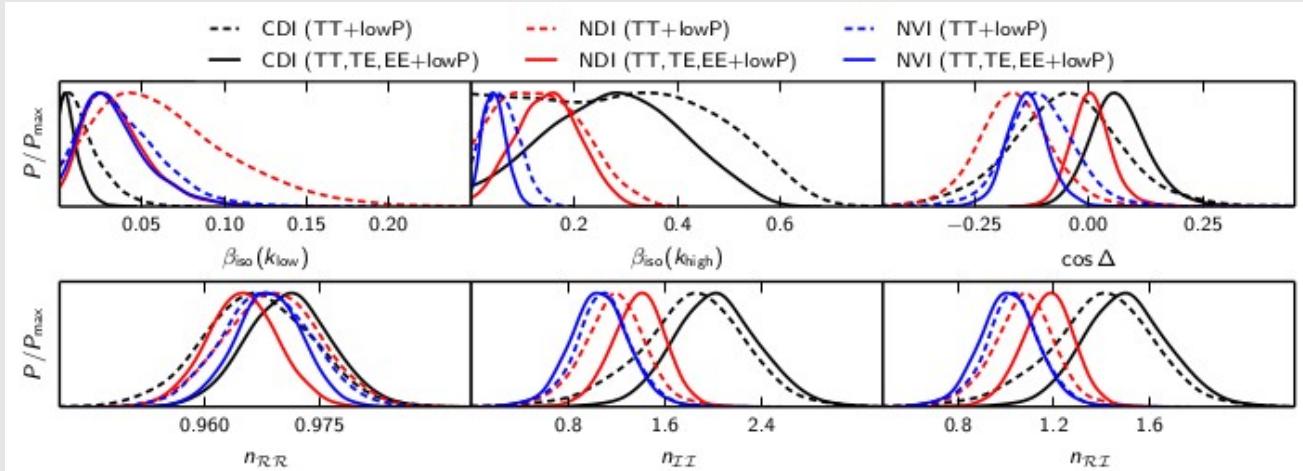
Primordial spectrum: $P_R = P_R(k_0)(k/k_0)^{n_s-1}$

adiabatic, gaussian, scale-invariant spectrum

No evidence for:
non-gaussianity, isocurvature modes or running of the spectral index



$$\beta_{\text{iso}} = \frac{P_{\parallel}(k)}{P_{RR}(k) + P_{\parallel}(k)}, \quad \cos \Delta = \frac{P_{ab}}{(P_{aa}P_{bb})^{1/2}}$$



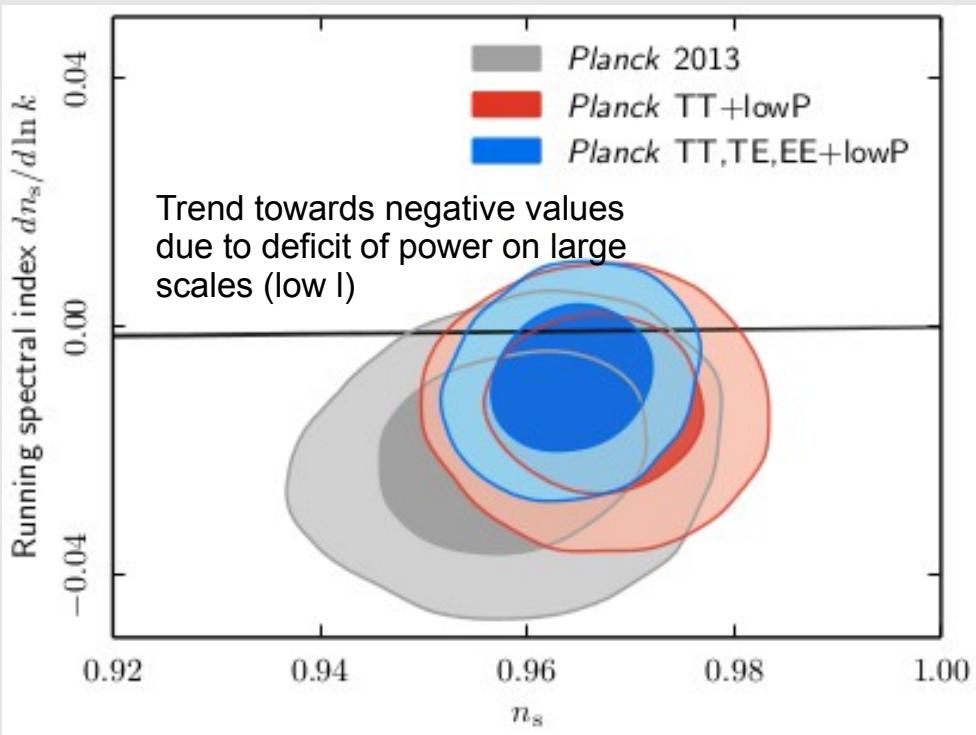
$$k_{\min} = 0.002 \text{ Mpc}^{-1}, \quad k_{\max} = 0.1 \text{ Mpc}^{-1}$$

Pure isocurvature mode ruled-out as only source for the spectrum

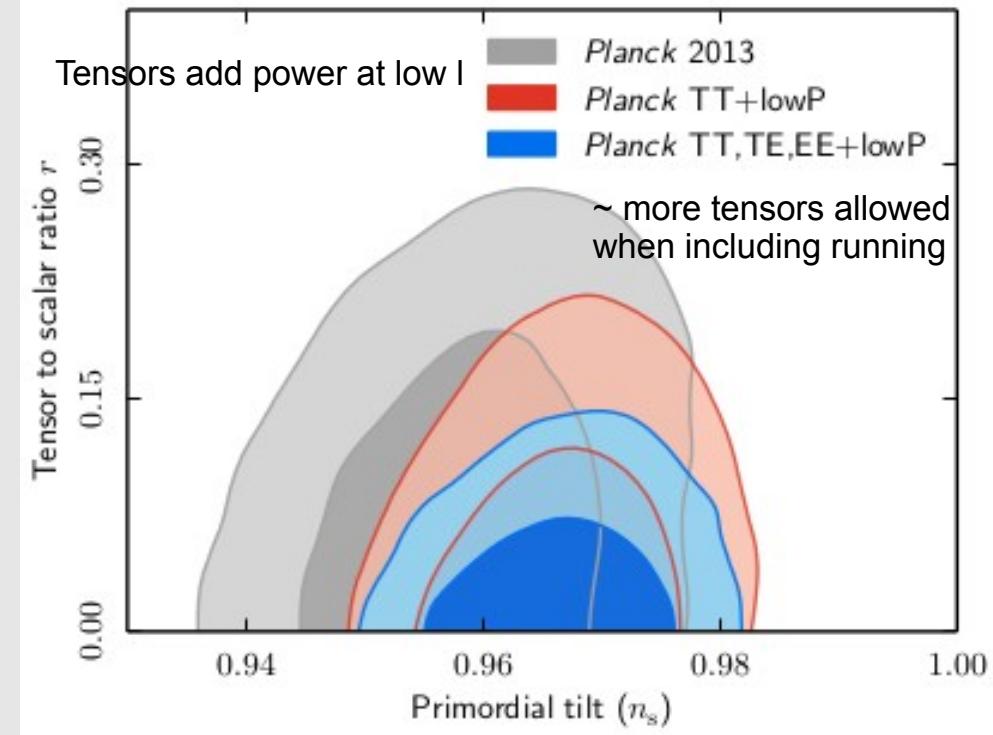
Primordial spectrum: $P_R = P_R(k_0)(k/k_0)^{n_s - 1 + \frac{1}{2}\alpha_s \ln k/k_0 + \dots}$ $k_0 = 0.05 \text{ Mpc}^{-1}$

adiabatic, gaussian, scale-invariant spectrum

No evidence for:
non-gaussianity, isocurvature modes or running of the spectral index



$$\alpha_s = \frac{dn_s}{d \ln k} = -0.0057 \pm 0.0071$$



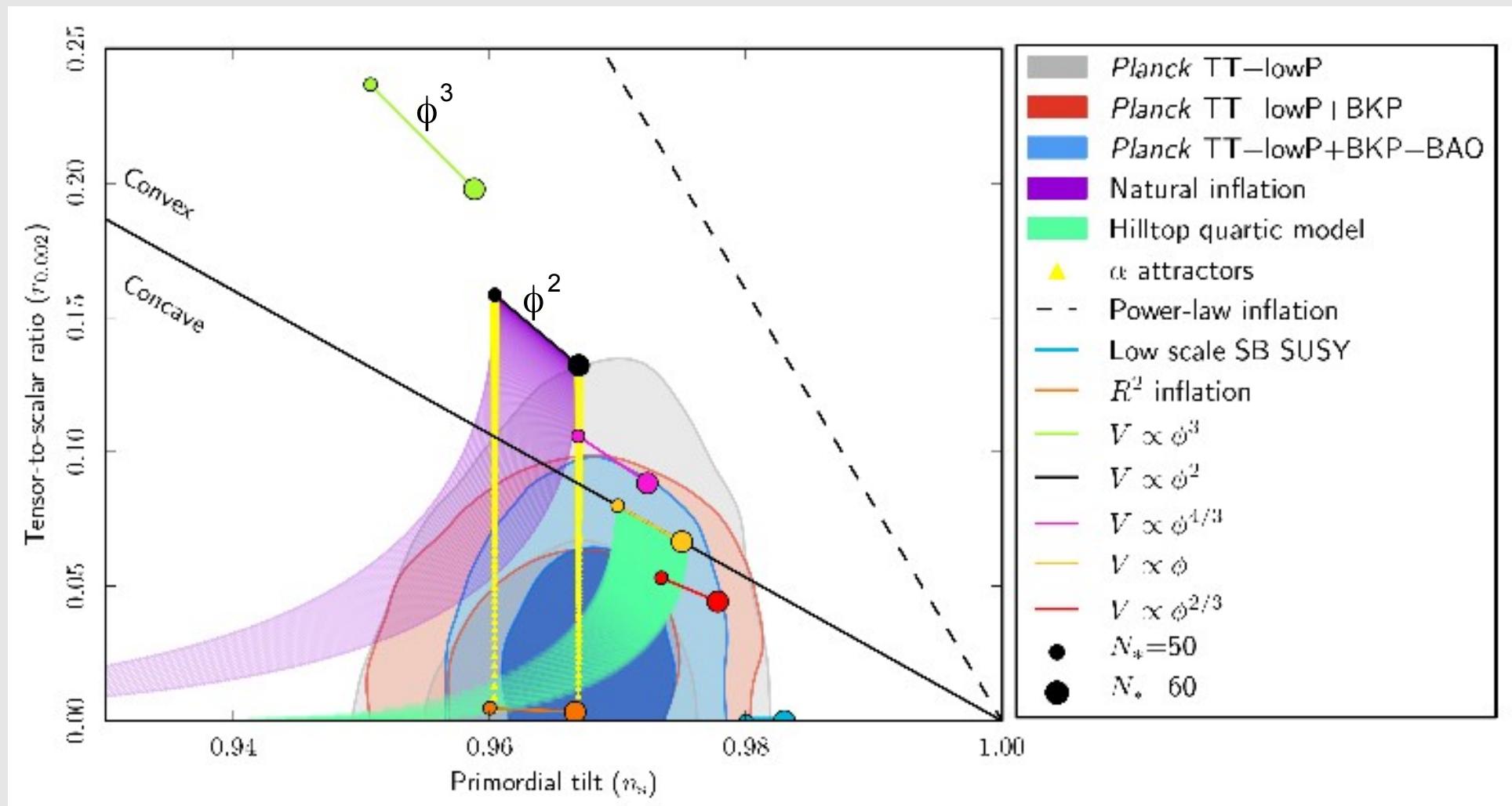
$$\alpha_s = -0.013 \pm 0.01, \quad r_{0.002} < 0.18 \\ [n_t = -r/8 < 0]$$

Primordial spectrum: Λ CDM extensions

Extended model, Λ CDM+r+	Parameter	<i>Planck</i> TT+lowP +lensing	<i>Planck</i> TT+lowP +BAO	<i>Planck</i> TT,TE,EE +lowP
+general reionization	r	$r < 0.11$	$r < 0.10$	< 0.10
	n_s	0.975 ± 0.006	0.971 ± 0.005	0.968 ± 0.005
$+N_{\text{eff}}$	r	< 0.14	< 0.12	< 0.11
	n_s	$0.977^{+0.016}_{-0.017}$	0.972 ± 0.009	0.964 ± 0.010
	N_{eff}	$3.24^{+0.30}_{-0.38}$	3.19 ± 0.24	$3.02^{+0.20}_{-0.21}$
$+Y_{\text{He}}$	r	< 0.14	< 0.12	< 0.12
	n_s	0.975 ± 0.007	0.973 ± 0.009	0.969 ± 0.008
	Y_{He}	0.258 ± 0.022	0.257 ± 0.022	0.252 ± 0.014
$+m_\nu$	r	< 0.11	< 0.11	< 0.11
	n_s	0.963 ± 0.007	0.967 ± 0.005	0.962 ± 0.005
	m_ν	< 0.67	< 0.21	< 0.58
$+{\Omega}_K$	r	< 0.15	$r < 0.11$	< 0.15
	n_s	0.971 ± 0.007	0.971 ± 0.007	0.969 ± 0.005
	Ω_K	$-0.008^{+0.010}_{-0.008}$	-0.001 ± 0.003	$-0.045^{+0.016}_{-0.020}$
$+w$	r	< 0.14	< 0.11	< 0.12
	n_s	0.969 ± 0.006	0.967 ± 0.006	0.966 ± 0.005
	w	$-1.46^{+0.20}_{-0.40}$	$-1.02^{+0.08}_{-0.07}$	$-1.57^{+0.17}_{-0.37}$
$+{\Omega}_K + dn_s/d \ln k$	r	< 0.20	< 0.18	< 0.19
	n_s	0.971 ± 0.007	0.969 ± 0.007	0.969 ± 0.005
	$dn_s/d \ln k$	-0.006 ± 0.009	-0.013 ± 0.009	-0.004 ± 0.008
$+N_{\text{eff}}+m_{\text{eff}}$	r	$r < 0.14$	$r < 0.13$	< 0.12
	n_s	$0.980^{+0.010}_{-0.014}$	$0.978^{+0.008}_{-0.011}$	$0.968^{+0.006}_{-0.008}$
	m_{eff}	< 0.27	< 0.21	< 0.83
	N_{eff}	< 3.45	< 3.73	< 3.47

Primordial spectrum: $P_R = P_R(k_0)(k/k_0)^{n_s-1}$ $k_0 = 0.002 \text{ Mpc}^{-1}$

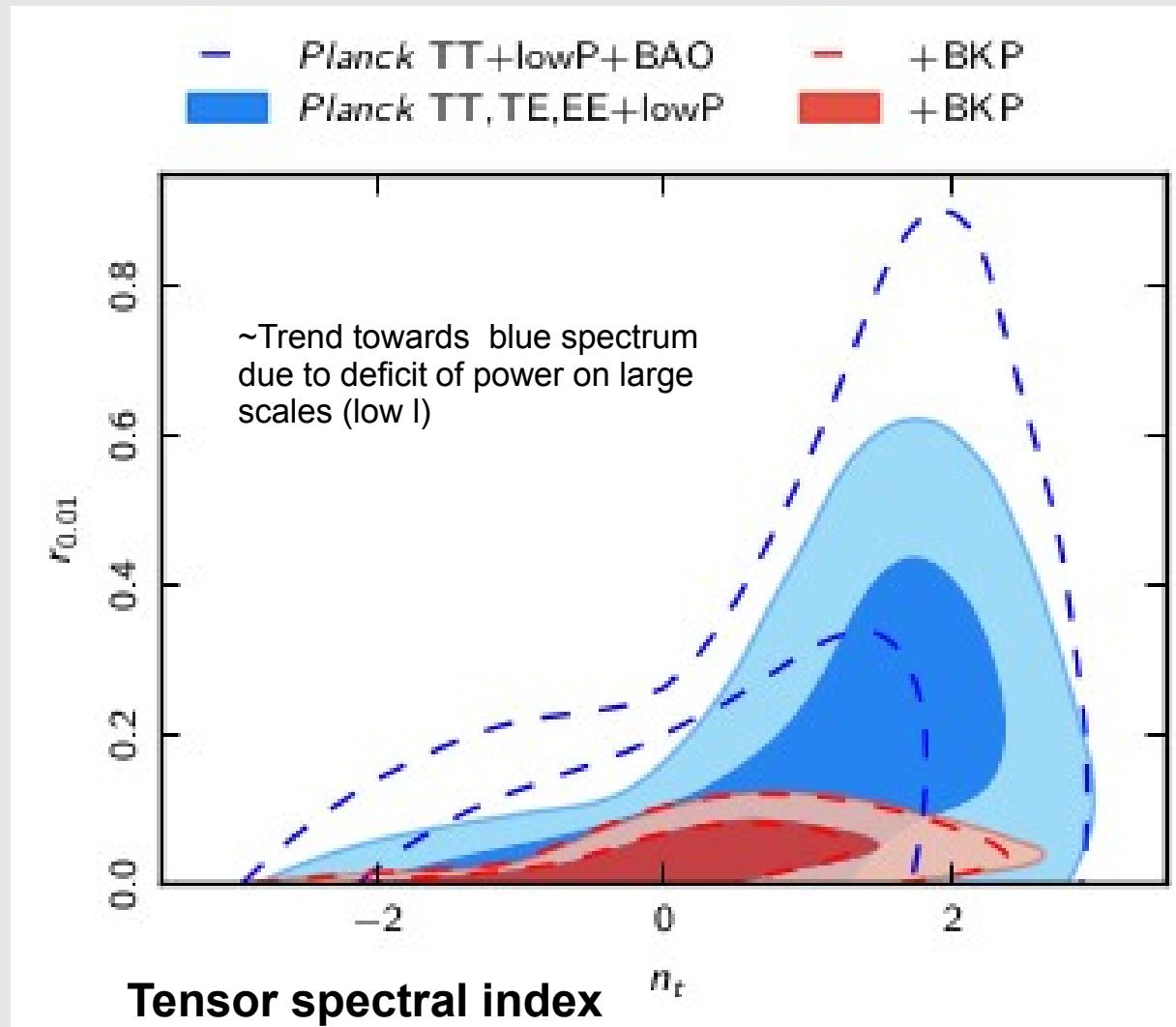
Tensor-to-scalar ratio : $r = P_T/P_R$ $P_R = 2.2 \times 10^{-9}$ ($n_T = -r/8$)



[Planck collab.: astro-ph/1502.02114]

Primordial spectrum: $P_R = P_R(k_0)(k/k_0)^{n_s-1}$ $k_0 = 0.01 \text{ Mpc}^{-1}$

Tensor-to-scalar ratio : $r = P_T/P_R$ $P_T = P_T(k_0)\left(\frac{k}{k_0}\right)^{n_t}$

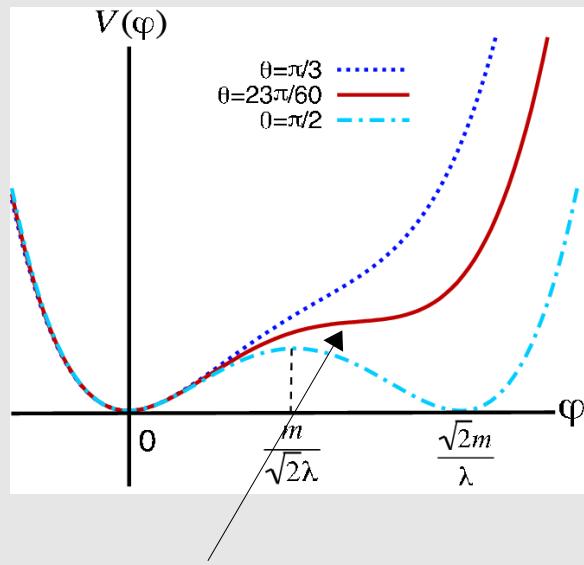


- Simple monomials models $V(\varphi) = V_0(\varphi/m_P)^p$ rule out for $p > 2$

$$n_s \approx 1 - \frac{2(p+2)}{4N_e + p}, \quad r \approx \frac{16p}{4N_e + p}$$

(minimal kinetic + minimal coupling to gravity)

- Adding more terms to the potential can save the model: “Polichaotic models”

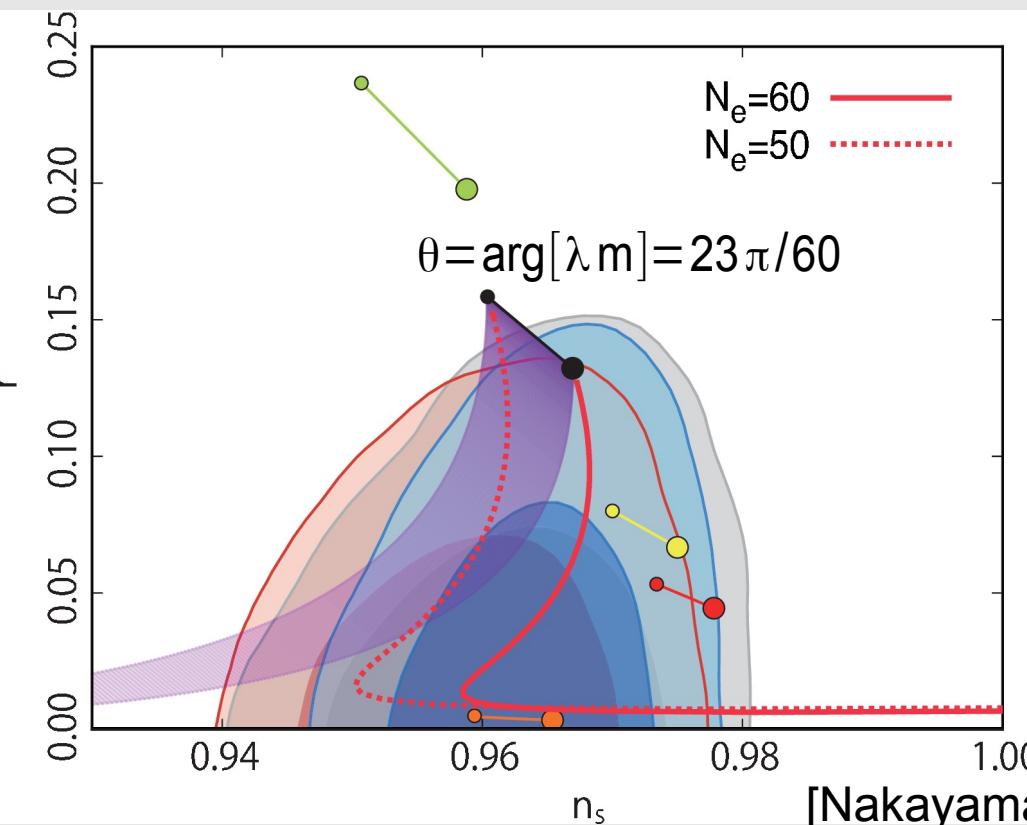


“Plateau”: lower H_*

$$r \sim \frac{H^2/m_P^2}{P_R} < 0.1$$

$$m \sim 10^{13} \text{ GeV}$$

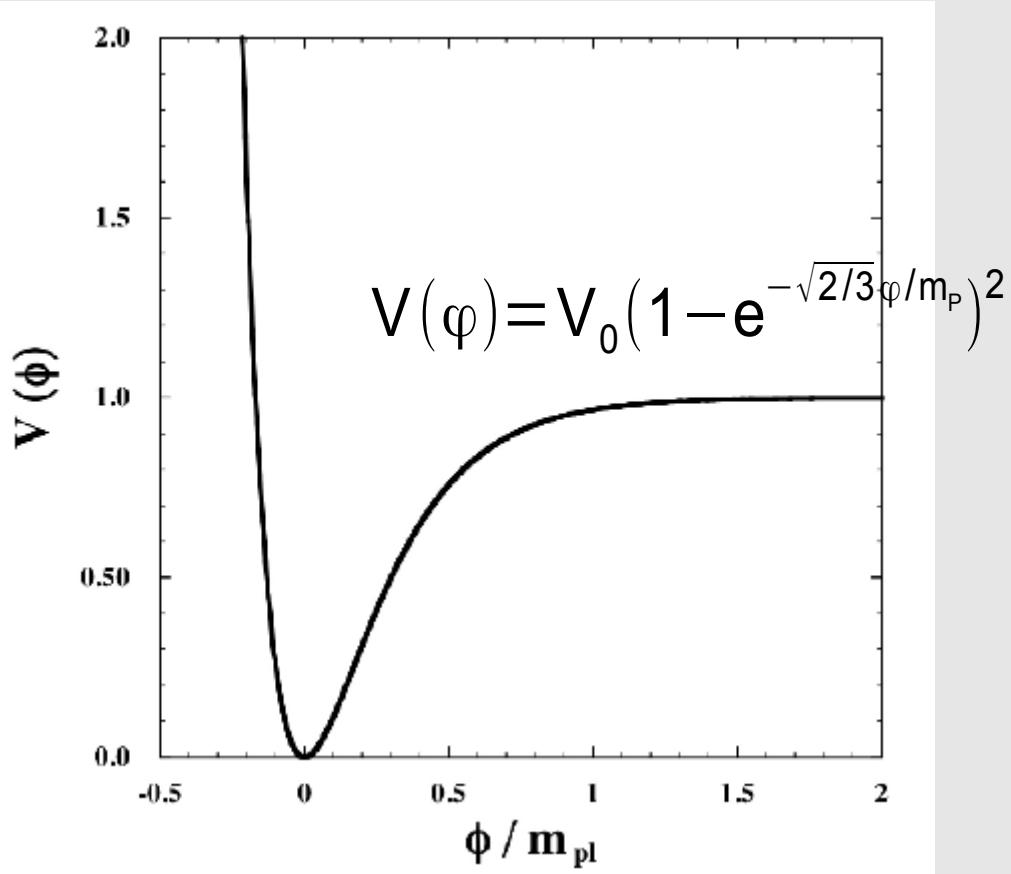
$$\lambda^2 \sim 10^{-14}$$



[Nakayama et al. 1303.7315]

R^2 inflation: gravitational action + radiative corrections

[Starobinsky '80]



Einstein grav. + Radiative corrections

$$S_{\text{gravity}} = \int \sqrt{-g} d^4x \left(\frac{m_p^2}{2} R + \frac{R^2}{12} + \dots \right)$$

R= Ricci scalar

Metric conformal transformation

$$S_{\text{gravity}} = \int \sqrt{-g} d^4x \left(\frac{m_p^2}{2} R + \frac{1}{2} (\partial_\mu \varphi)^2 - V(\varphi) \right)$$

$$n_s \simeq 1 - \frac{2}{N_e} \quad r \simeq \frac{12}{N_e^2}$$

R^2 inflation from supergravity:

Buchmuller, Domcke, Kamada, 1306.3471

Farakos, Kehagias, Riotto, 1307.1137

Ellis, Nanopoulos, Olive, 1307.3537

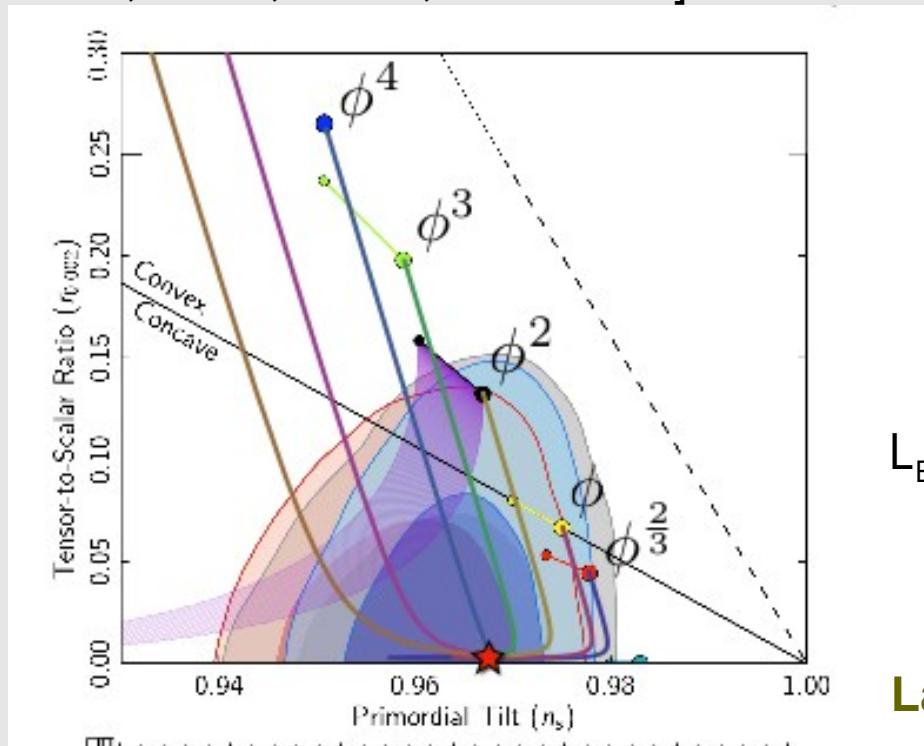
Ferrara, Kallosh, Van Proeyen, 1309.4052

“Starobinsky model”



Non-minimal coupling to gravity

[Kallosh, Linde, Roest, 1310.3950]



$$n_s \approx 1 - \frac{2}{N_e}, \quad r \approx \frac{12}{N_e^2}$$

Universal attractor point for models
with spontaneously broken (super)conformal invariance

Sugra embedding: $K = -3 \log(F[\Omega(\phi), \phi, S])$, $W = \lambda S f(\Phi)$

[Non-minimal kinetic term: $K_{\phi\bar{\phi}} \partial_\mu \phi \partial^\mu \bar{\phi}$]

Non-minimal coupling: Jordan frame

$$L_J = \sqrt{-g} \left(\frac{\Omega(\phi)}{2} R - \frac{1}{2} (\partial \phi)^2 - V_J(\phi) \right)$$

Metric conformal transformation

$$\Omega(\phi) = 1 + \xi f(\phi)$$

$$L_E = \sqrt{-g} \left(\frac{m_P^2}{2} R - \frac{1}{2} \left(\Omega(\phi)^{-1} + \frac{3}{2} (\partial \Omega(\phi))^2 \right) (\partial_\mu \phi)^2 - \frac{V_J(\phi)}{\Omega(\phi)^2} \right)$$

Non-minimal kinetic: Einstein frame

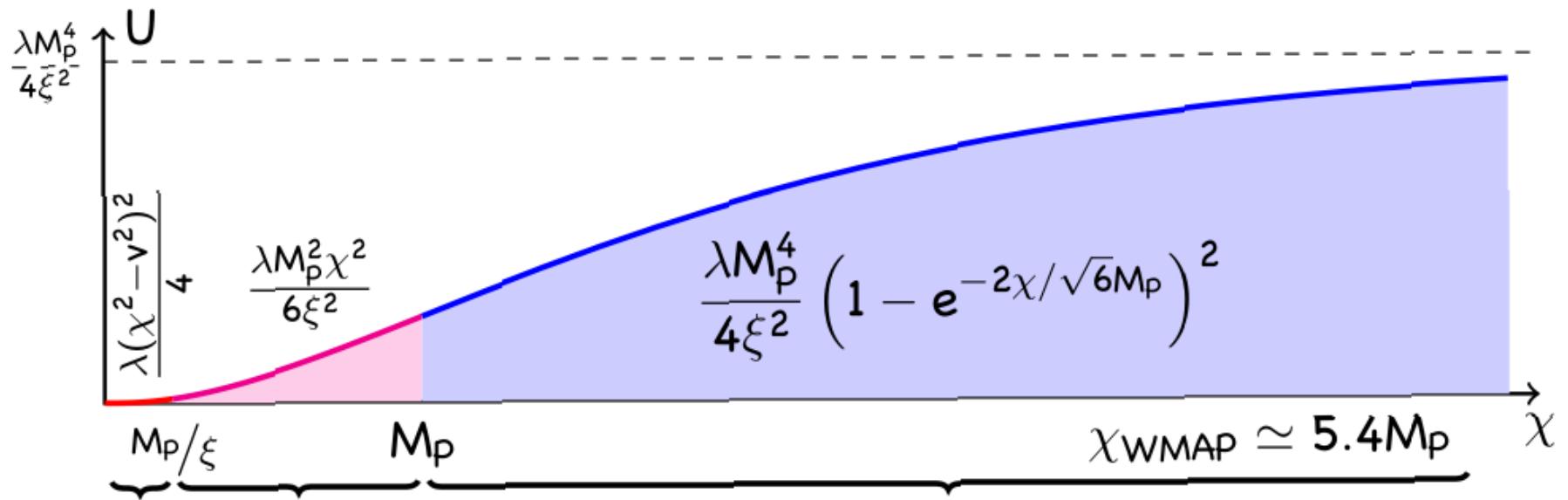
Large ξ limit + canonical kinetic term:

$$L_E = \sqrt{-g} \left(\frac{m_P^2}{2} R - \frac{1}{2} (\partial_\mu \phi)^2 - V_0 (1 - e^{-\sqrt{2}\phi/3})^2 \right)$$

[p=4 large ξ limit : $\xi \geq 0.1$]

Inflaton = Higgs

[Bezrukov & Shaposhnikov '08]



$$S = \int \sqrt{-g} d^4x \left(\frac{m_P^2}{2} R + \xi R \varphi^2 + \frac{1}{2} (\partial_\mu \varphi)^2 - V(\varphi) \right)$$

$$V(\varphi) = \frac{\lambda}{4} (\varphi^2 - v^2)^2 \xrightarrow{\text{Conformal transformation}} V(\chi) = \frac{\lambda_{\text{eff}}}{4} m_P^4 (1 - e^{-\sqrt{2/3} \chi / m_P})$$

$\varphi \simeq \chi \ll m_P/\xi$

$$\lambda \simeq O(0.1)$$

$n_s \simeq 0.97, \quad r \simeq 0.003$

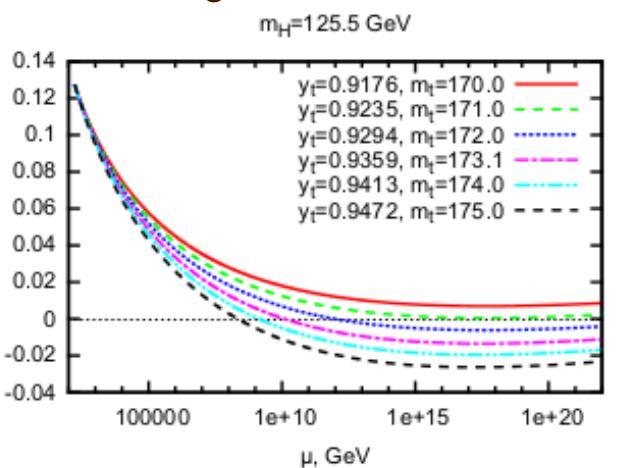
$$\lambda_{\text{eff}} \simeq 10^{-13}$$

[COBE]

Reheating: decay into SM degrees of freedom !!

Inflaton = Higgs: Stability of the Higgs potential?

RGE running:



[Bezrukov, Rubio & Shaposhnikov '14]

[Bezrukov & Shaposhnikov, '14]

SM: $\lambda(\mu) < 0$ at $\mu \simeq 10^{10}$ GeV for $m_t \geq 172$ GeV

$$V_{\text{SM}} = \frac{1}{4} \lambda \phi^4 - \frac{1}{2} m^2 \phi^2 \longrightarrow V_{\text{inf}} \simeq \frac{\lambda_{\text{eff}}}{4} m_P^4 (1 - e^{-\sqrt{2/3} \chi / m_P} (1 + \dots))$$

“Low energy”

$$\mu \leq m_P / \xi$$

“High Energy”

$$m_P / \xi \leq \mu \leq m_P / \sqrt{\xi}$$

Matching conditions at $\mu \sim m_P / \xi$: depend on higher order operators (UV completion)

$$\lambda(\mu_1) = \lambda^{\text{SM}}(\mu_1) + \delta\lambda, \quad h_t(\mu_1) = h_t^{\text{SM}}(\mu_1) + \delta h_t$$

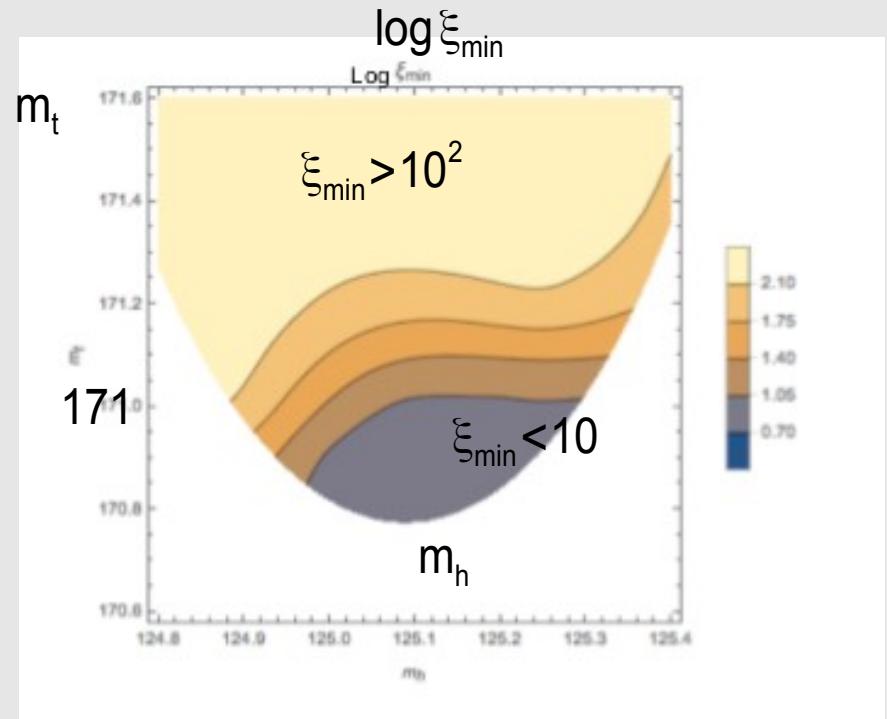
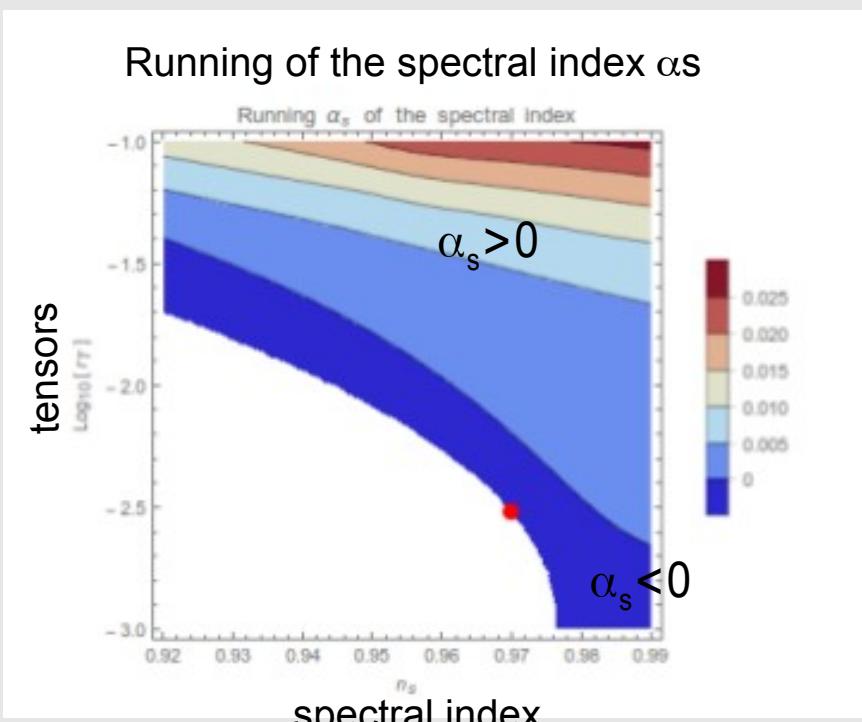
The negative minimum is lifted by thermal corrections after inflation:

$$\Delta V(T_{\text{reh}}, \chi_{\min}) > -V(\chi_{\min})$$

[Enqvist, Enckell, Nurmi '16]

[Fumagalli & Postma '16]

Inflaton = Higgs: Stability of the Higgs potential?



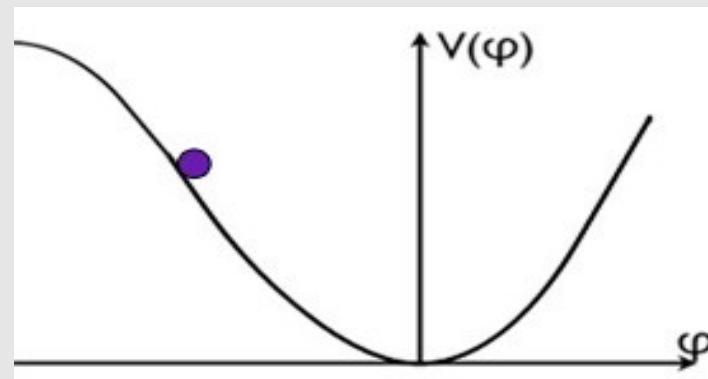
[Enqvist, Enckell, Nurmi '16]

Higgs inflation can be realised for $\xi < 10$ for low top masses

It can be falsified: a detection of a negative running of the spectral index + $m_t < 171.8$ GeV will rule out the model

Which inflationary potential?

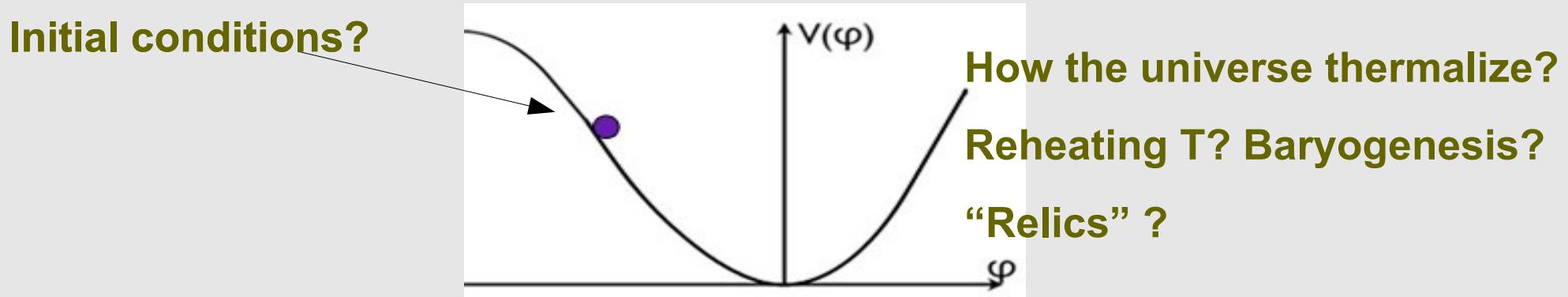
$$V(\varphi_i, \dots) = \underbrace{V_{\Delta N_e = 10}(\varphi_i)}_{\text{Primordial spectrum}} + \underbrace{V_{\text{end}}(\varphi_i, \dots)}_{\text{End of inflation}} + \underbrace{V_{\text{reh}}(\varphi_i, \dots)}_{\text{Inflation} \rightarrow \text{Radiation}}$$



J. Martin, C. Ringeval, & Vennin, arXiv: 1303.3787: “Encyclopedia Inflationaris”
(~ 100 models)

Which inflationary potential? :

$$V(\varphi_i, \dots) = \underbrace{V_{\Delta N_e=10}(\varphi_i)}_{\text{Primordial spectrum}} + \underbrace{V_{\text{end}}(\varphi_i, \dots)}_{\text{End of inflation}} + \underbrace{V_{\text{reh}}(\varphi_i, \dots)}_{\text{Inflation} \rightarrow \text{Radiation}}$$



J. Martin, C. Ringeval, & Vennin, arXiv: 1303.3787: “Encyclopedia Inflationaris”
(~ 100 models)

Initial conditions for inflation

(canonically normalized, minimally coupled field to gravity)

Homogeneous & isotropic background

- Small field models: $\Delta\phi < m_P$, $H_{\text{inf}}/m_P < 10^{13}$ GeV

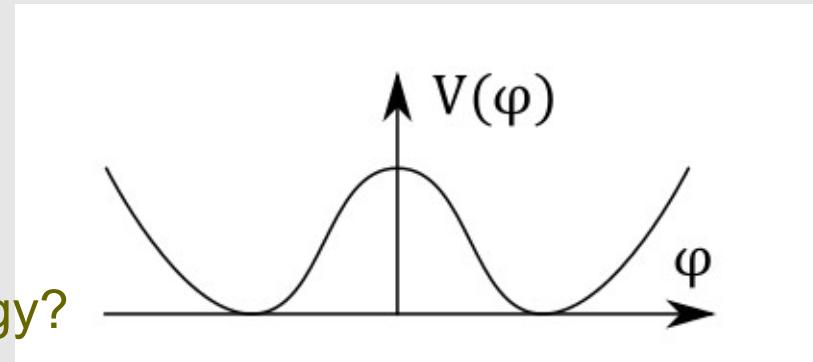
[Goldwirth & Piran, Phys. Rept. 214 '92]

The initial field velocity has to be tuned: $\dot{\phi} < H^2$

Tunneling from a false vacuum? Quantum cosmology?

[Garriga, Montes, Sasaki, Tanaka '99]

[Hartle&Hawking '83; Linde '84; Vilenkin '82]



- Large field models: $\Delta\phi \sim m_P$, $H_{\text{inf}}/m_P \sim 10^{13}$ GeV

The inflationary slow-roll trajectory is a local attractor in the initial conditions space

There is enough time for the kinetic energy to get redshifted and become subdominant

[Albrecht & Brandenberger '85; Brandenberger & Kung '90; Mukhanov '14]

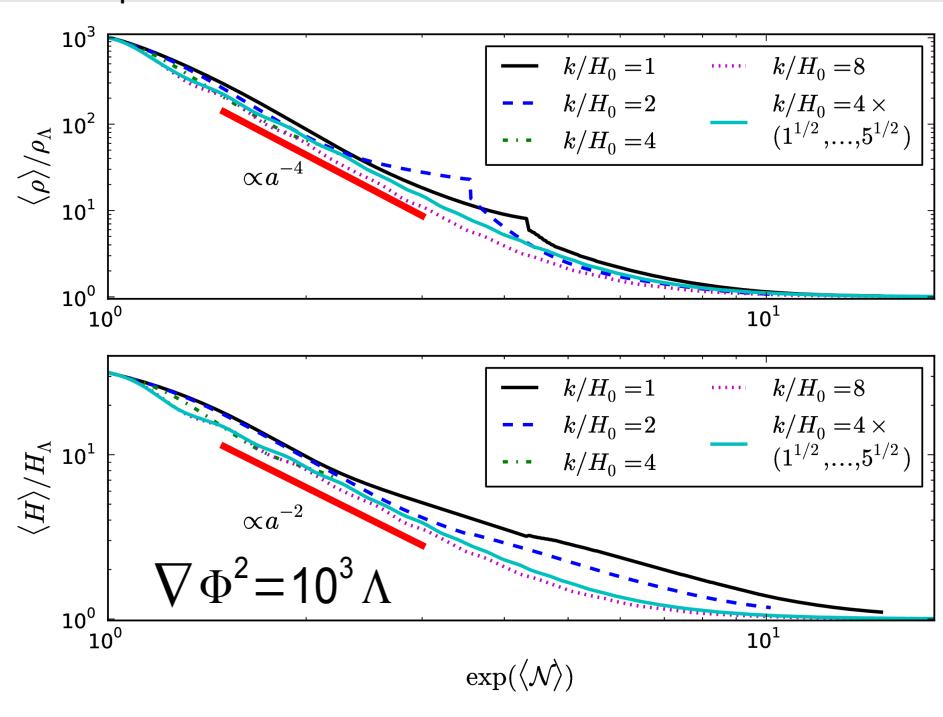
Initial conditions for inflation

(canonically normalized, minimally coupled field to gravity)

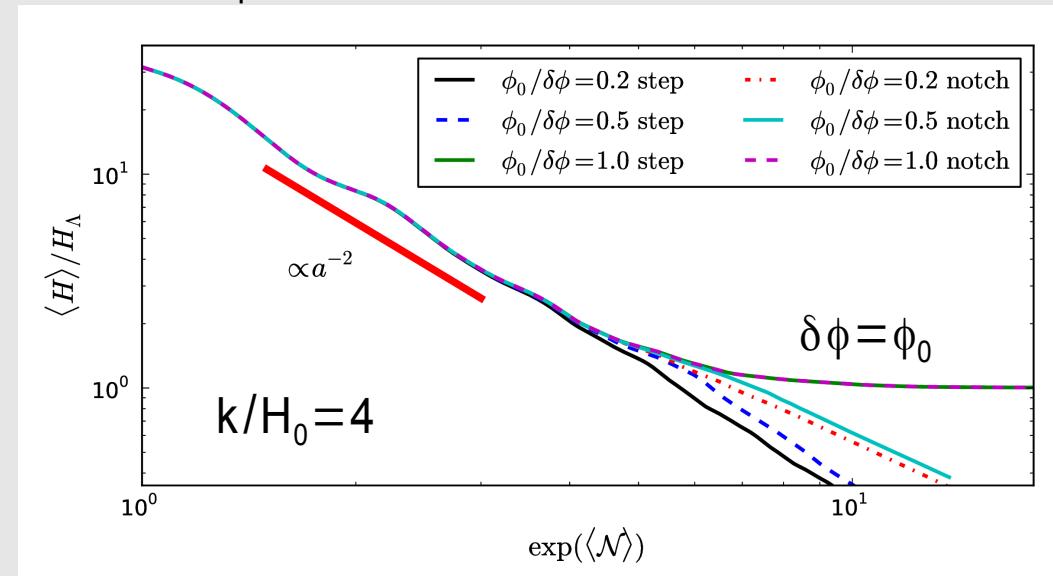
Inhomogeneous background: Field + metric (Einstein Eqs.)

$$\nabla \Phi^2 \geq V(\Phi) \quad \Phi(x) = \phi_0 + \delta\phi \sum \cos(kx + \theta_k)$$

$\delta\phi$ within the range of inflationary values



$\delta\phi$ outside the inflationary plateau



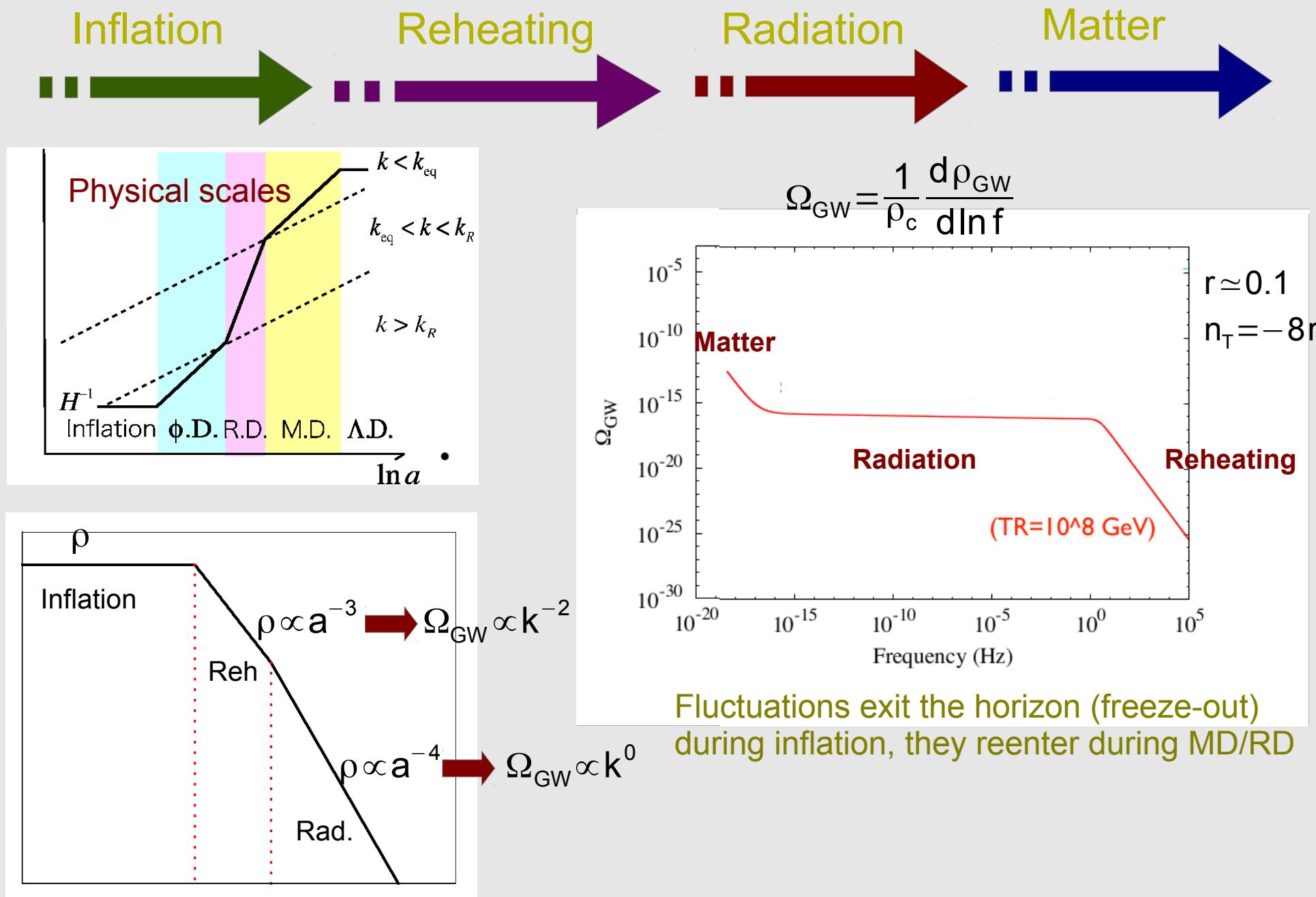
[East, Kleban, Linde, Senatore 1511.05143]

Exponential expansion occurs even in the presence of large gradients

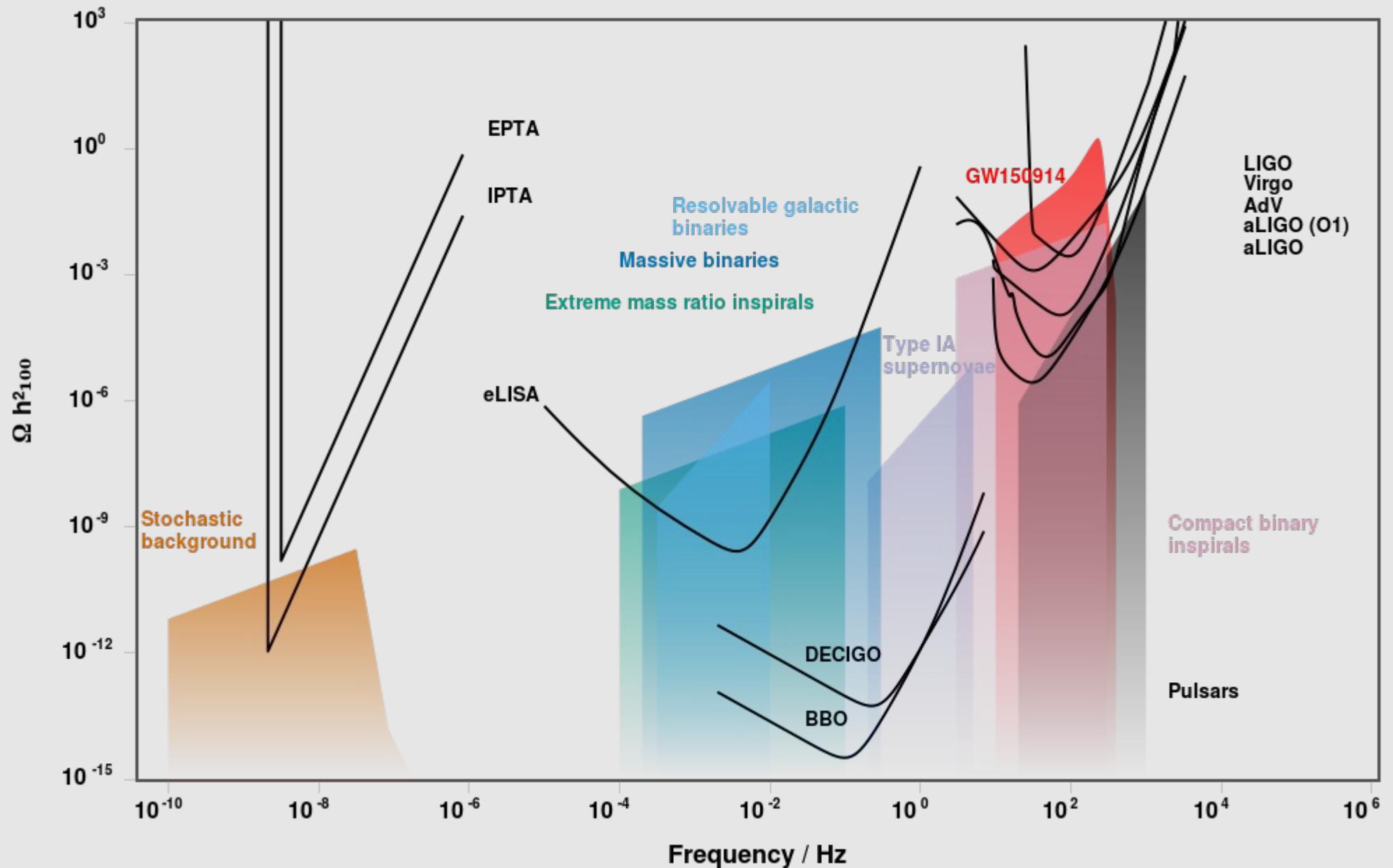
Large initial values of the inhomogeneous field (outside the “plateau”) may prevent inflation

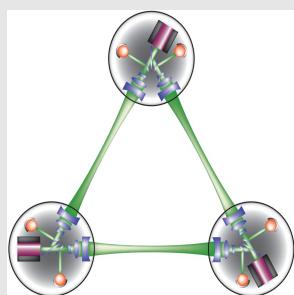
There is no need to assume a Hubble-sized homogeneous initial patch for inflation to occur

Primordial gravitational waves: direct detection?

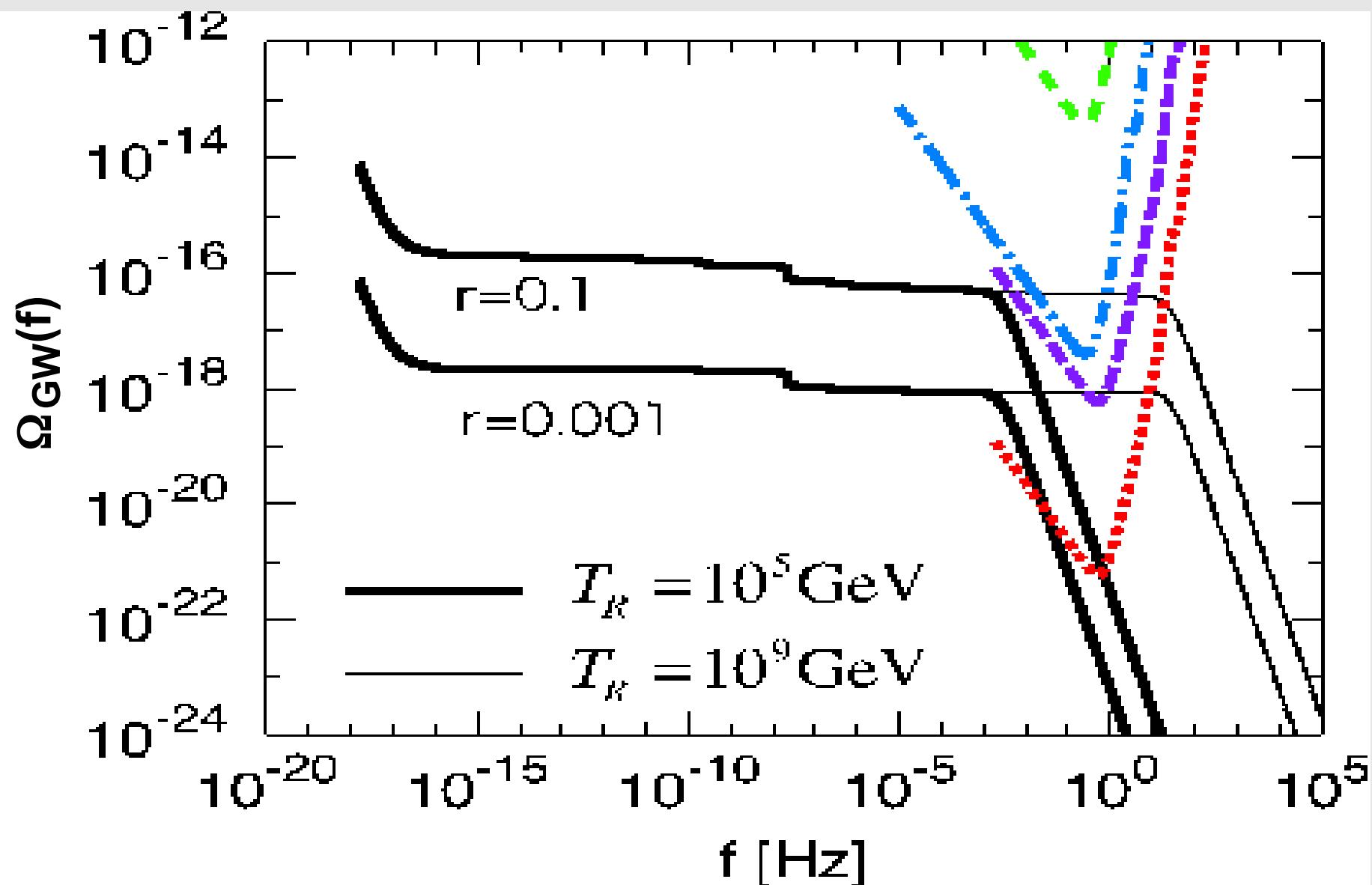


Gravitational Waves: Direct detection





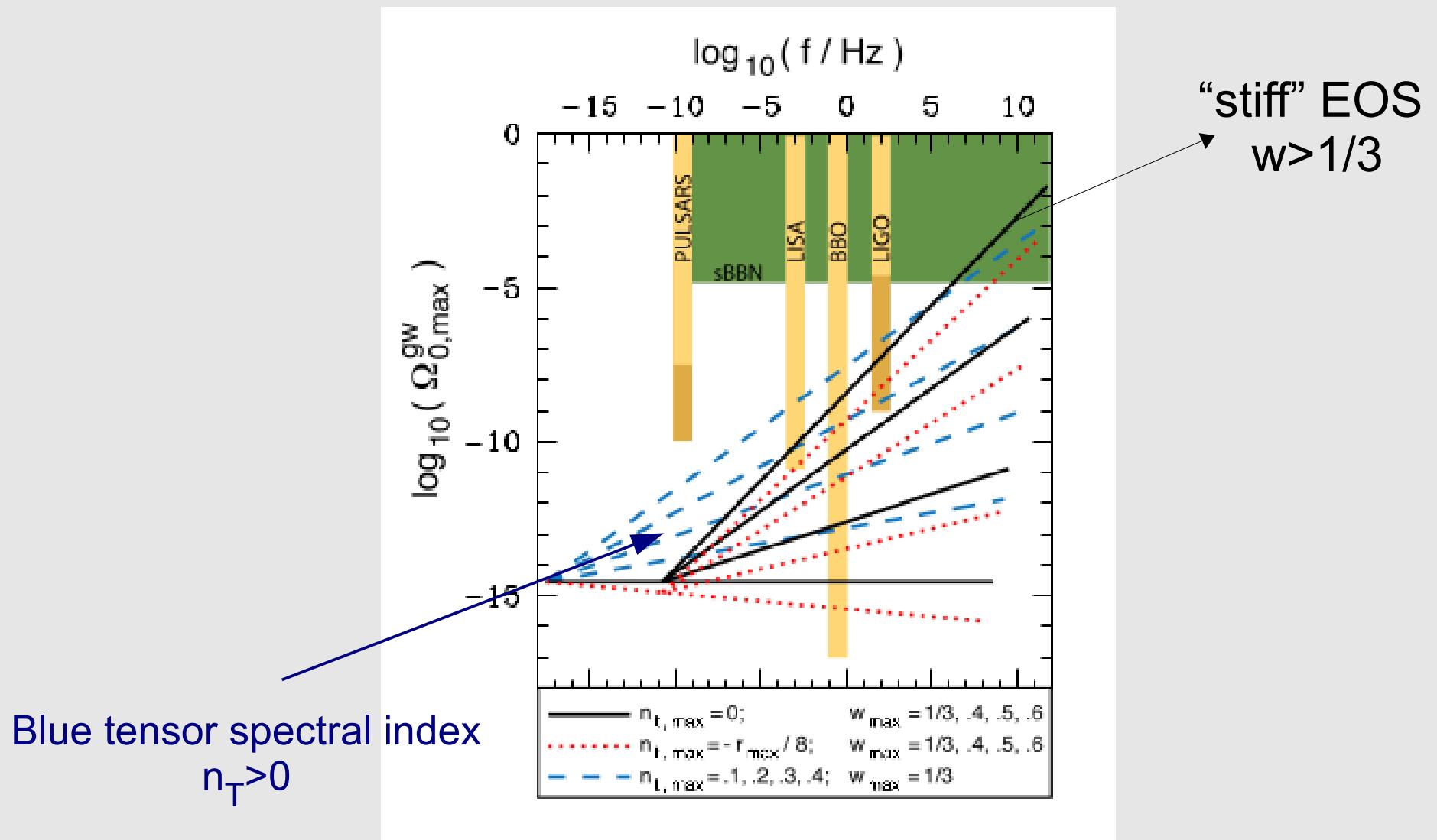
DECIGO



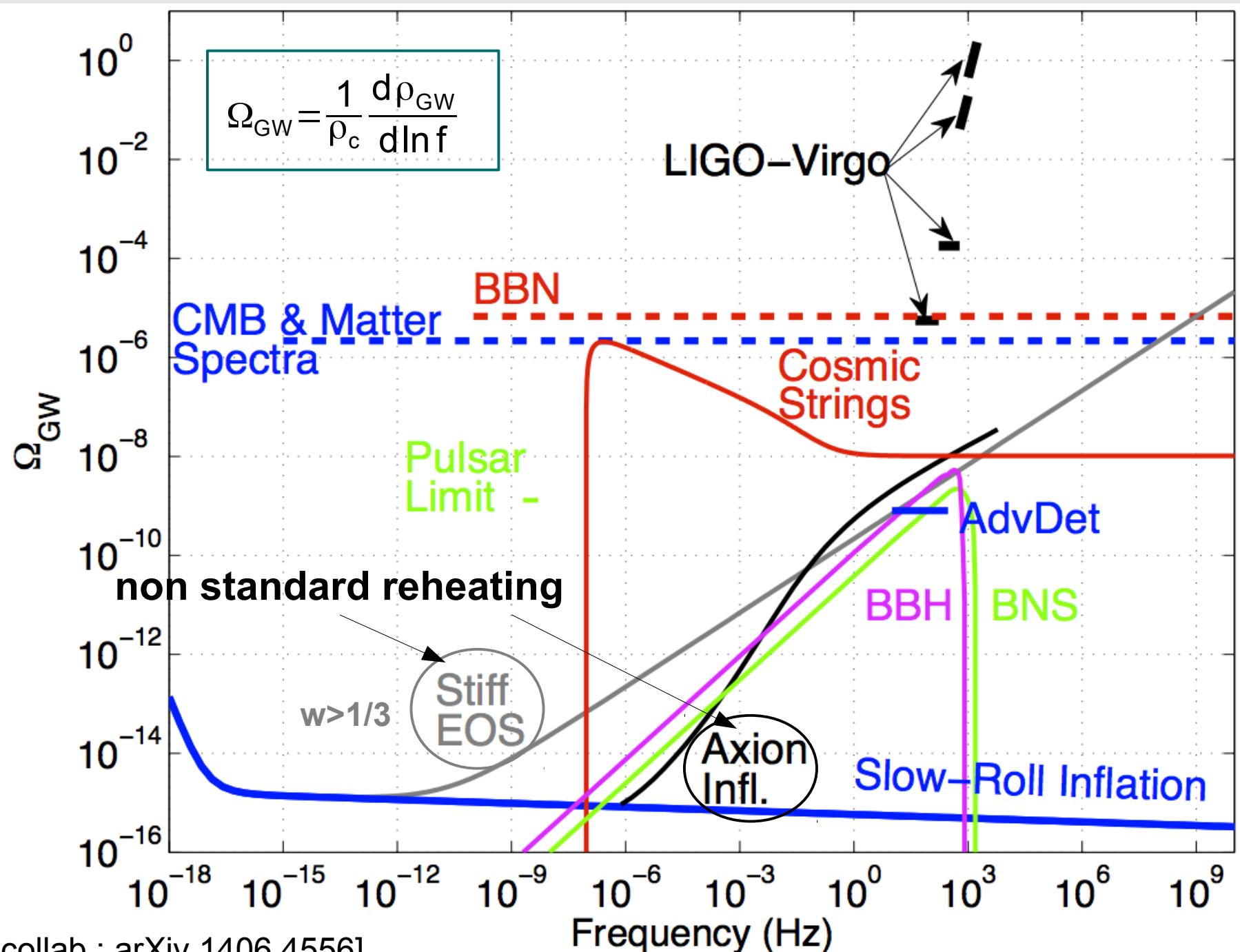
Primordial gravitational waves: direct detection?

$$\Omega_{\text{GW}}(k) = 4.2 \times 10^{-2} P_R(k_0)^2 \left(\frac{k}{k_0}\right)^{n_T} r_0 \left(\frac{a_{\text{eq}}}{a}\right) \left(\frac{k_{\text{eq}}}{k}\right)^\alpha$$

$\alpha = 2$, MD $w=0$
 $\alpha = 0$, RD $w=1/3$
 $\alpha < 0$, $w > 1/3$



Primordial gravitational waves: direct detection?



Preheating & Gravity Waves

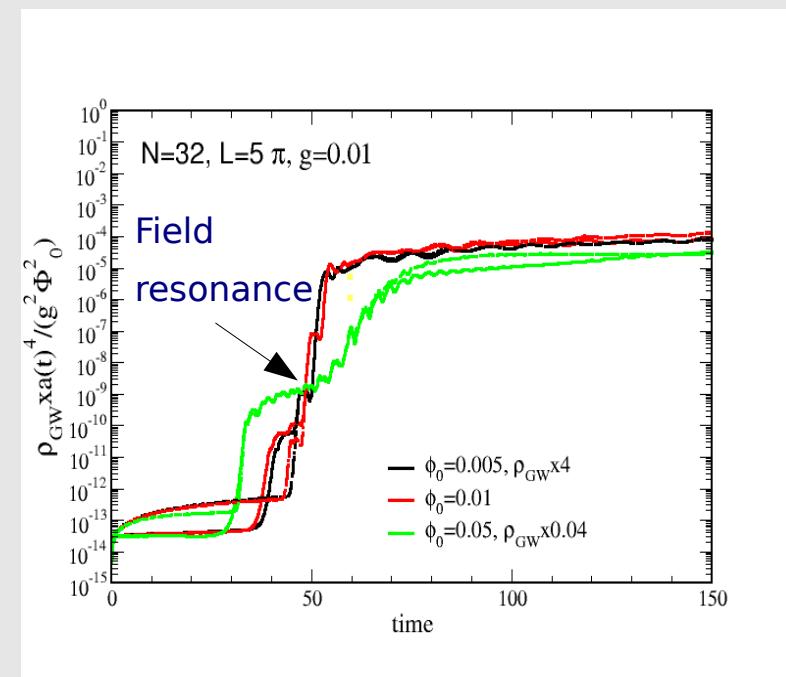
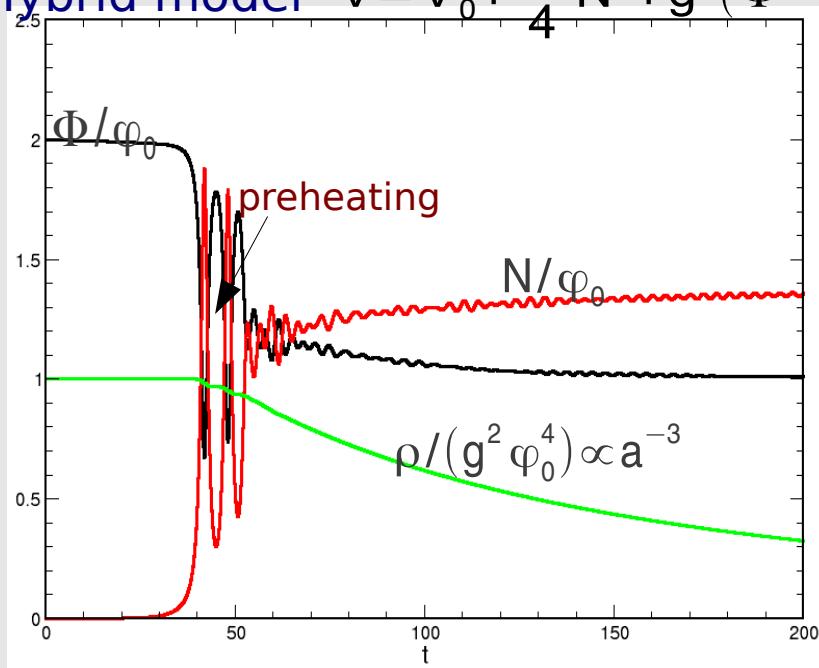
Non-adiabatic change in the time dependent effective masses leads to parametric amplification of the field fluctuations \rightarrow Source of GW

$$\ddot{h}_{ij} + 3H\dot{h}_{ij} - \nabla^2 h_{ij}/a^2 = S_{ij}^{TT}/a^2$$

Scalors
Fermions
Gauge bosons

Scalar matter $S_{ij}^M = 16\pi G T_{ij} \sim 16\pi G \nabla_i \chi \nabla_j \chi$

Hybrid model $V = V_0 + \frac{g^2}{4} N^4 + g^2 (\Phi^2 - \varphi_c^2) N^2 + \frac{1}{2} m_\varphi^2 \Phi^2$



Easther, Giblin & Lim '06 ; García-Bellido & Figueroa '07; Dufaux et al. '08
Dufaux, Figueroa & García-Bellido '10; Enqvist, Figueroa & Meriniemi '12

Frequency (scalars):

$$f_0 = \frac{k_0}{2\pi} \simeq \frac{k_{\text{res}}}{2\pi} \frac{a_{\text{res}}}{a_0} \simeq 4 \times 10^{10} \frac{k_{\text{res}}}{\rho_{\text{res}}^{1/4}} \text{Hz} \sim 10^{10} g^{1/2} \text{Hz}$$

($k_{\text{res}} \simeq g \varphi_0$, $\rho_{\text{res}} \simeq g^2 \varphi_0^4$)

(No entropy production: $g_s(T_R) a^3(t_R) T_R^3 = g_s(T_0) a^3(t_0) T_0^3$)

Power spectrum (scalars):

$$\Omega_{\text{GW}}^{(0)} h^2 = \Omega_{\text{GW}}^{(\text{res})} h^2 \left(\frac{a_{\text{res}}}{a_0} \right)^4 \simeq 10^{-6} \frac{1}{\rho_{\text{res}}} \frac{d\rho_{\text{GW}}^{\text{res}}}{d \ln k} \propto 10^{-6} \left(\frac{\Phi_0}{m_P} \right)^2$$

Redshift: radiation

Low scale hybrid inflation (very weak coupling $g < 10^{-12}$) could lead to a stochastic background of GW within the reach of GW detectors....

Frequency (fermions):

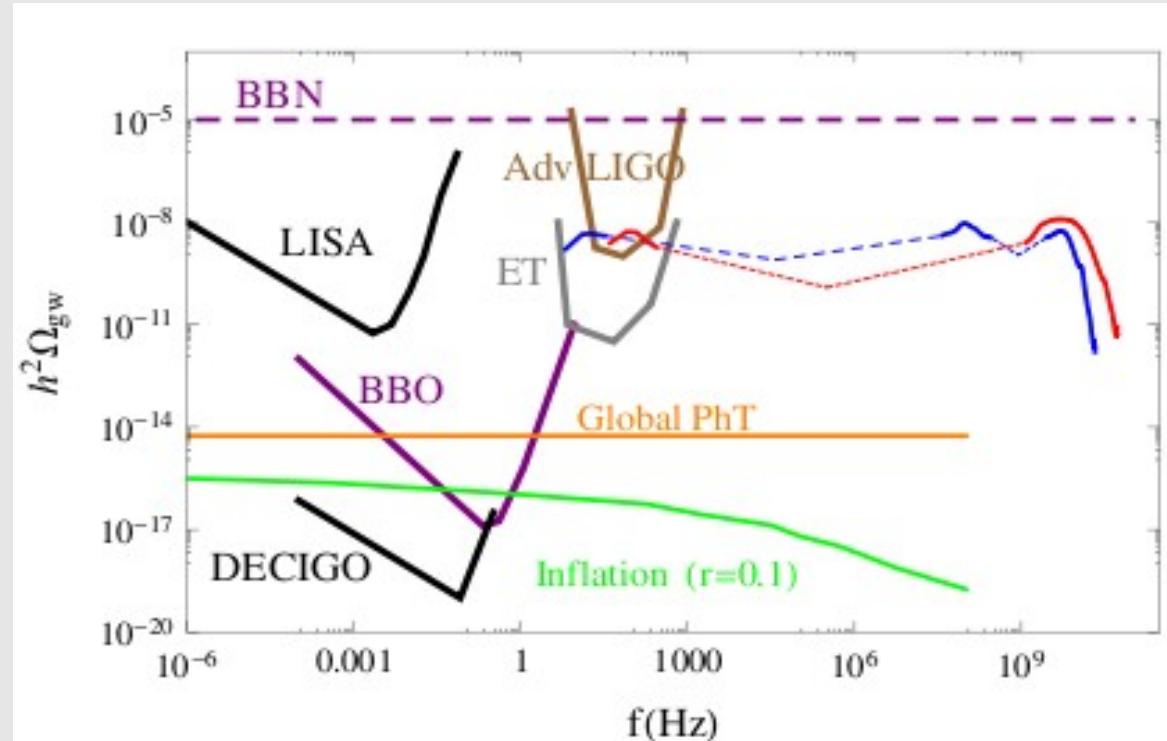
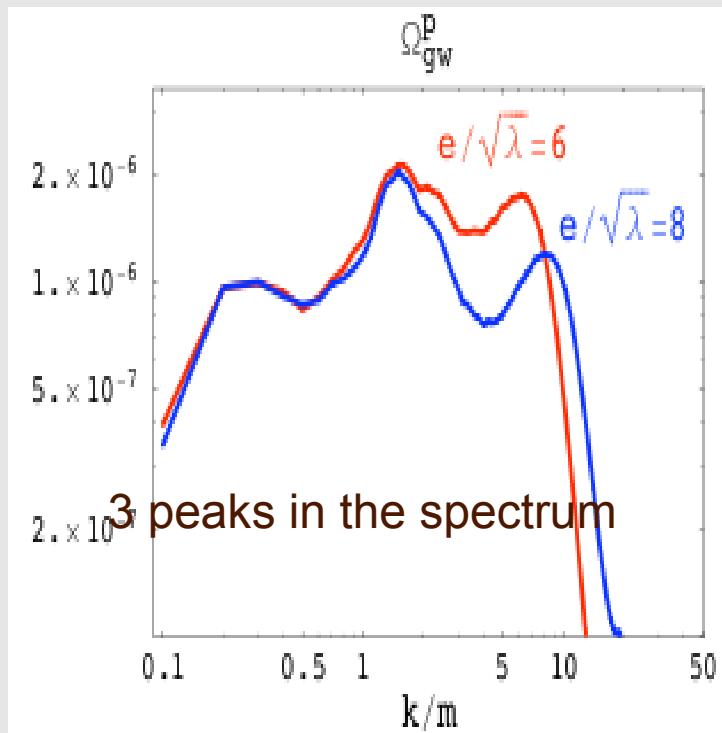
$$f_0 = q^{1/4} \omega_0 \quad (\omega_0 = \text{frequency of inflaton osc.})$$

Frequency too large (outside observable range) for $h_0^2 \Omega_{\text{GW}} > 10^{-14}$

GW from preheating: Abelian Higgs model

$$L = D_\mu N D^\mu N + \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - V(\varphi, N), \quad D_\mu = \partial_\mu - ie A_\mu$$

Scalar couplings: g, λ ; gauge coupling: e



$$f_{IR} \simeq \begin{cases} \lambda^{1/4} g 10^{11} \text{Hz}, & g^2 > \lambda \\ \lambda^{-3/4} g 10^{10} \text{Hz}, & g^2 < \lambda \end{cases}$$

$$f_{mid} \simeq \lambda^{1/4} 10^{11} \text{Hz}$$

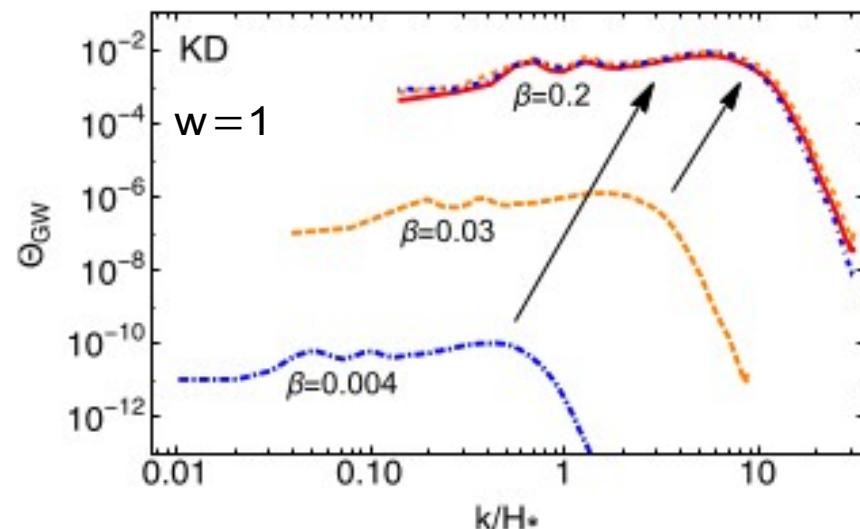
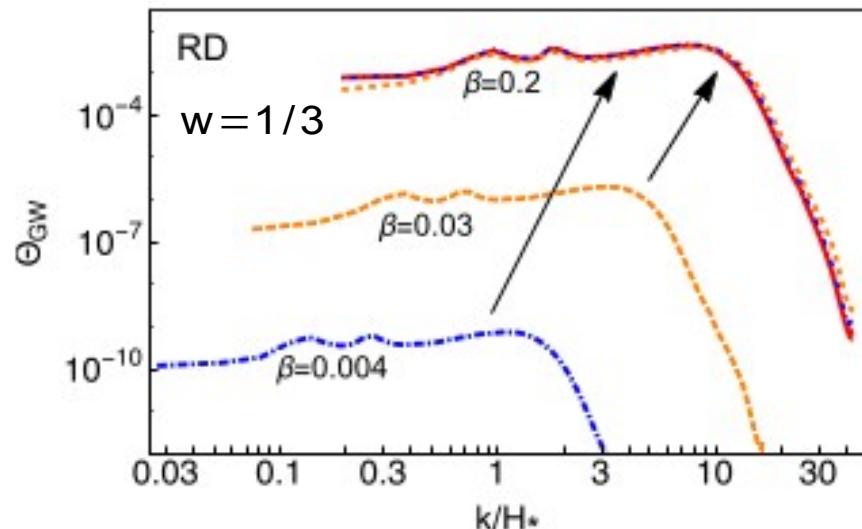
$$f_{UV} \simeq \frac{e}{\sqrt{\lambda}} \lambda^{1/4} 10^{11} \text{Hz}$$

GW from preheating: Higgs (condensate) decay into gauge bosons after inflation

$$\Omega_{\text{GW}}(k, z) = \left(\frac{H_{\text{end}}}{m_P}\right)^4 \left(\frac{a}{a_{\text{end}}}\right)^{3w-1} \Theta_{\text{GW}}$$

$$\beta = \lambda^{1/2} \frac{\phi_{\text{end}}}{H_{\text{end}}}$$

w= eos after inflation



$$h^2 \Omega_{\text{GW}}(f_p) \leq \begin{cases} 10^{-29} & \text{RD} \quad f_p \leq 3 \times 10^8 \text{ Hz} \quad (\text{too small}) \\ 10^{-16} & \text{KD} \quad f_p \leq 3 \times 10^{11} \text{ Hz} \end{cases}$$

Summary

- Planck 2015: the most precise cosmological data up today

6 parameters to fit the Universe: Λ CDM

$$\left(P_R(k_0), n_s, \Omega_B, \Omega_M, \Omega_\Lambda, \tau_{\text{reio}} \right)$$

[Some tension with SNIa (H_0) & CFHTLenS (σ_8).....]

- Open questions: Dark Matter? Dark Energy? WDE=-1?
- Primordial spectrum: gaussian? Adiabatic? Tensors?

Inflation provides a solution to the standard cosmological problems, and a causal mechanism for the primordial spectrum

Many models still consistent with observations

Reheating? Thermalization of the Universe after inflation?

Alternatives: loop quantum cosmology? String gas cosmology ($n_T > 0$)?...

Detection of primordial tensors may help with model selection

- GW may offer a unique window also into the physics of the very early universe (before BBN) !