

Axionic signals at colliders with a dynamical Higgs

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In collaboration : Brivio, Gavela, Mimasu, Merlo, No, Sanz. Paper to appear.

Why an extra Goldstone?

$$\frac{a(x)}{f_a}$$

GBs are included in many BSM proposals.

Eg. The Peccei-Quinn $U(1)_A$ axion that solves the Strong CP problem.

... but also Majoron, ALPS, etc.

We will analyse at colliders for scales lower than the QCD axion scale.
=> Look for derivative couplings

Why a composite Higgs?

It is a solution to the Hierarchy Problem.

Is the Higgs...

Elementary?

Composite?

As in certain solutions of the Hierarchy problem. e.g. SUSY...

e.g. a pseudo-Goldstone Boson (GB) in $SO(5)$

LINEAR EWSB

NON-LINEAR EWSB

The composite Higgs appears as the pseudo-GB of a spontaneously broken symmetry ... this explains its "lightness".

Linear

The Higgs is a **doublet**.

The doublet structure determines much of its phenomenology.

- The Higgs appears in the Lagrangian as $(v+h(x))^n$.

$$\Phi(x) = \frac{v+h(x)}{\sqrt{2}} e^{i\vec{\pi}\vec{\sigma}/v} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

The "pions" are parametrised by a dimensionless matrix.
(As in the QCD Non-linear Sigma model)

$$\mathbf{U}(x) \iff \begin{matrix} W^+ & \pi_1, \pi_2 \\ & Z \\ & & \pi_3 \end{matrix}$$

Non-linear

The Higgs is **not necessarily a doublet**, and can be inserted as a gauge singlet.

- $(v+h(x))^n$ is replaced by a more general expansion

$$\mathcal{F}_i(h) = 1 + 2a_i \frac{h}{v} + b_i \left(\frac{h}{v}\right)^2 + \dots$$

- $\mathcal{F}(h)$ and $\mathbf{U}(x)$ are separate building blocks in the EFT
→ Generically, one obtains a larger NLO basis.

Linear

[Georgi, Randall, Kaplan -1984]

$$\mathcal{A}_\Phi = \Phi^\dagger i \overleftrightarrow{D}_\mu \Phi \frac{\partial^\mu a}{f_a}$$

is NLO in linear (d=5)

Main differences in scalar couplings:

Non-Linear

[Us - to appear]

→ **New LO couplings**
for example $aZh, aZh\gamma$...

$$\mathcal{A}_{2D} = \frac{iv^2}{4} \text{Tr}(\mathbf{T}\mathbf{V}_\mu) \frac{\partial^\mu a}{f_a} \mathcal{F}_{2D}(h)$$

becomes LO (2 derivatives)

$$i(D^\mu \Phi^\dagger \square \Phi - \square \Phi^\dagger D^\mu \Phi) \frac{\partial_\mu a}{f_a}$$

doesn't appear in linear until d=7

→ **New NLO couplings**
and $aAh^n, WWah$...

$$\mathcal{A}_9 = i \text{Tr}(\mathbf{T}\mathbf{V}_\mu)^2 \text{Tr}(\mathbf{T}\mathbf{V}_\nu) \frac{\partial^\nu a}{f_a} \mathcal{F}_9(h)$$

becomes NLO (4 derivatives)

$$\epsilon^{\mu\nu\rho\alpha} p_{\alpha\sigma}$$

→ **New Lorentz structures at NLO**

And also for $aWW, aZWW$...

$$(g^{\mu\rho} p_a^\nu + g^{\nu\rho} p_a^\mu) g^{\mu\nu} p_a^\rho$$

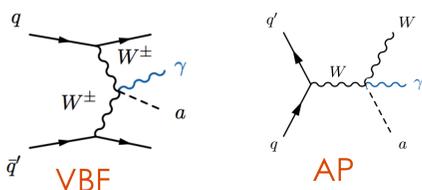
$$\mathbf{T} = \mathbf{U}\sigma^3\mathbf{U}^\dagger$$

$$\mathbf{V}_\mu = (D_\mu \mathbf{U})\mathbf{U}^\dagger$$

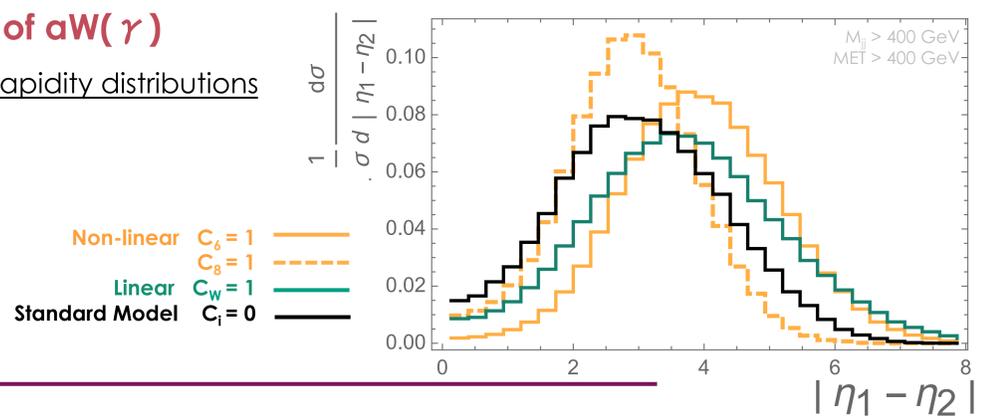
Distinctive signatures of non-linearity at colliders

1 Vector Boson Fusion – $a_{jj}(\gamma)$ – & Associated Production of $aW(\gamma)$

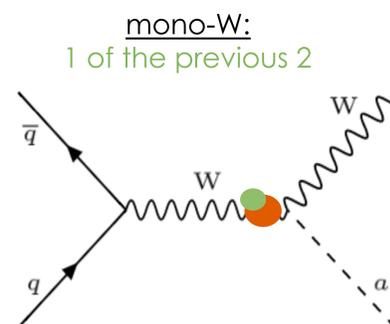
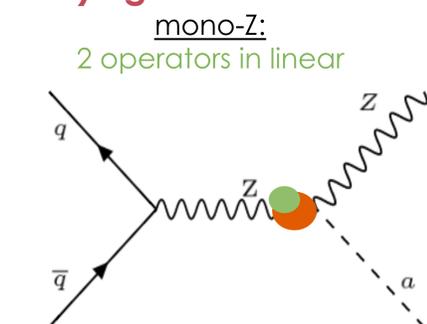
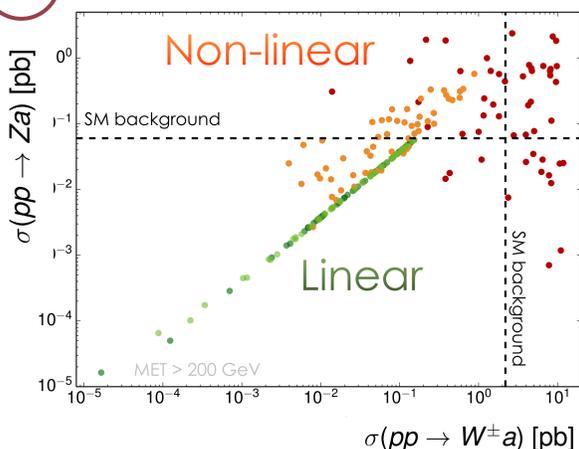
for example: VBF rapidity distributions



4 signals depending on 4 Wilson coefficients offer a possibility to **disentangle!**



2 Mono-Z and mono-W (and mono-h) signatures



+3 new ones in non-linear...

+2 more!

More operators: can "escape" axion-photon constraint and break correlations.

Decorrelation of signals which are linearly correlated are **smoking guns** for non-linearity!

We have derived the **effective chiral Lagrangian** for **ALPs + bosonic sector** with a **dynamical Higgs**.

300 fb⁻¹ of LHC data can test for **new GBs** with **f_a ≥ TeV**.