

Standard Model Theory

**XLIV International Meeting
on Fundamental Physics**

Universidad Autonoma, MADRID

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MAX-PLANCK-INSTITUT FÜR PHYSIK, MÜNCHEN

Outline

- The Standard Model in the pre-Higgs era
- The Standard Model in the Higgs era
- The status of Standard Model
- Future prospects

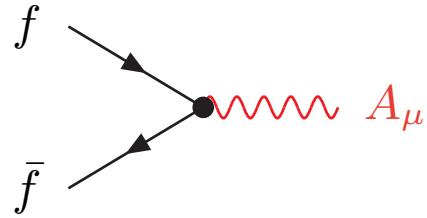
the rise of the EW Standard Model

- the symmetry group $SU(2) \times U(1)$
 - discovery of neutral currents 1973
- the W and Z bosons
 - discovery 1983 at SPS (CERN)
- the coupling structure from local gauge invariance
 - precise measurements at LEP/SLC 1989 - 2000
- the Higgs mechanism and Yukawa interactions
 - top discovery at Tevatron 1995
 - Higgs discovery at LHC 2012

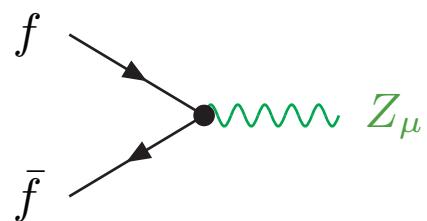
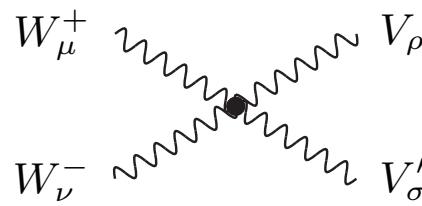
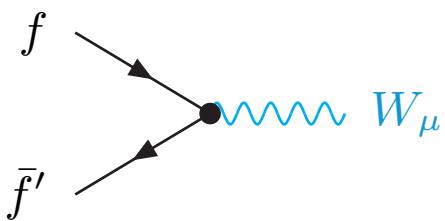
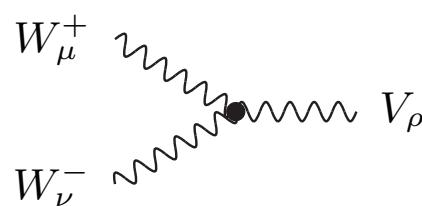
**precision physics: major role in establishing the SM
restricting BSM physics**

interactions from gauge-symmetric Lagrangian

fermion–vectorbosons



vectorboson self interactions



coupling constants:

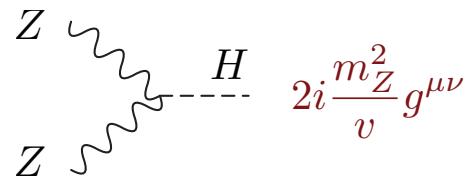
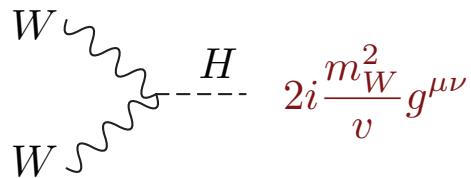
$$g_{\text{em}} = e, \quad g_{\text{weak}} = e / \sin \theta_w$$

group entries:

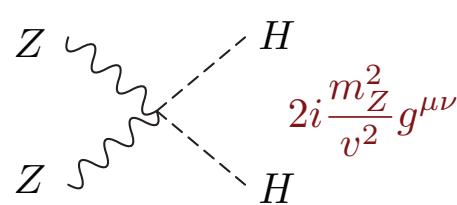
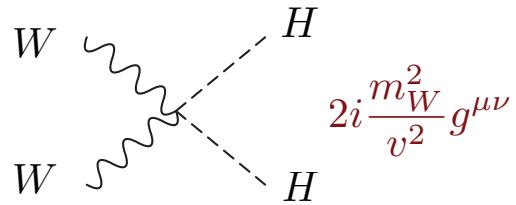
isospin I_3^f , charge Q_f

interactions from symmetry-breaking Lagrangian

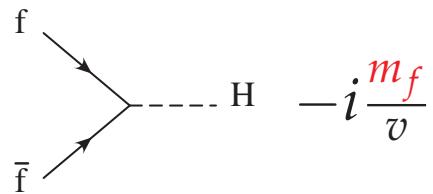
WWH and ZZH couplings



WWHH and ZZHH couplings

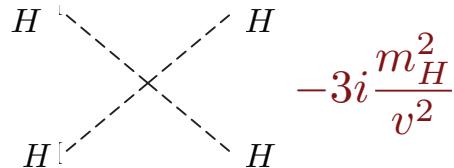
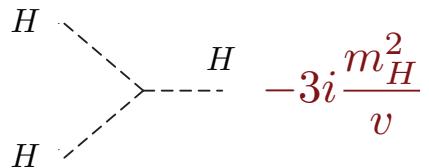


Fermion-Higgs coupling



couplings proportional to masses

Cubic and Quartic self-couplings



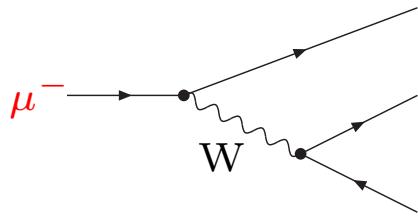
- Higgs boson probably found, all other particles confirmed
- consistent quantum field theory
 - in accordance with unitarity
 - renormalizable \Rightarrow predictions at higher orders
- formal parameters: $g_2, g_1, v, \lambda, g_f, V_{\text{CKM}}$
physical parameters: $\alpha, M_W, M_Z, M_H, m_f, V_{\text{CKM}}$

observables and experiments

- Muon decay:

$$\mu^- \rightarrow \nu_\mu e^- \bar{\nu}_e$$

determination of the Fermi constant



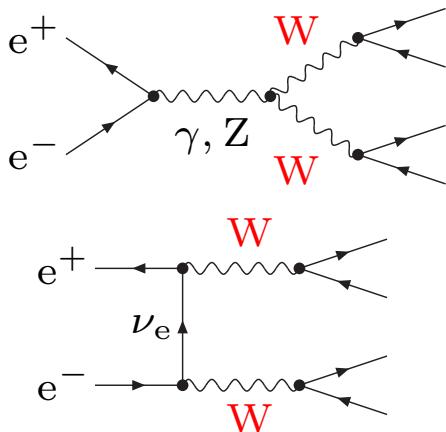
$$G_\mu = \frac{\pi \alpha M_Z^2}{\sqrt{2} M_W^2 (M_Z^2 - M_W^2)} + \dots$$

- Z production (LEP1/SLC):

$$e^+ e^- \rightarrow Z \rightarrow f \bar{f}$$

various precision measurements at the
Z resonance: M_Z , Γ_Z , σ_{had} , A_{FB} , A_{LR} , etc.
 \Rightarrow good knowledge of the $Z f \bar{f}$ sector

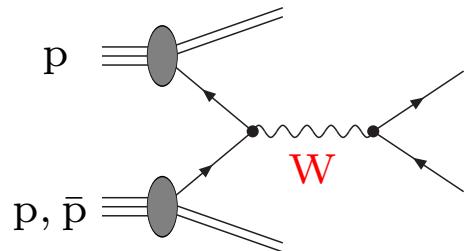
- W-pair production (LEP2/ILC): $e^+ e^- \rightarrow WW \rightarrow 4f (+\gamma)$



- measurement of M_W
- $\gamma WW/ZWW$ couplings
- quartic couplings: $\gamma\gamma WW$, γZWW

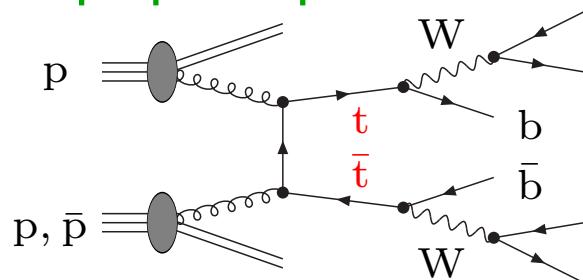
experiments at hadron colliders

- W production (Tevatron/LHC): $pp, p\bar{p} \rightarrow W \rightarrow l\nu_l (+\gamma)$



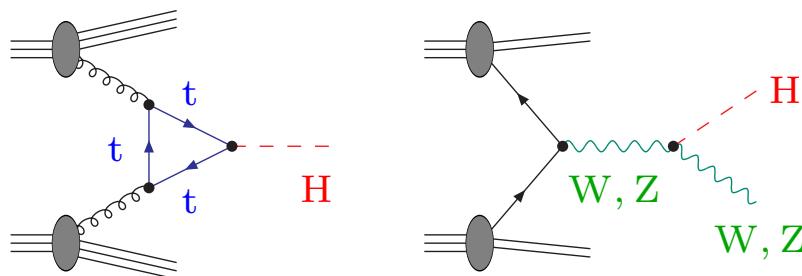
- measurement of M_W
- bounds on γWW coupling

- top-quark production (Tevatron/LHC): $pp, p\bar{p} \rightarrow t\bar{t} \rightarrow 6f$



- measurement of m_t

- Higgs production (LHC): $pp \rightarrow H \rightarrow \gamma\gamma, ZZ, \dots$



input quantities

$$G_F = 1.1663787(6) \cdot 10^{-5} \text{ GeV}^{-2}$$

$$M_Z = 91.1875 \pm 0.0021 \text{ GeV}$$

$$M_W = 80.385 \pm 0.015 \text{ GeV}$$

$$m_t = 173.2 \pm 0.9 \text{ GeV}$$

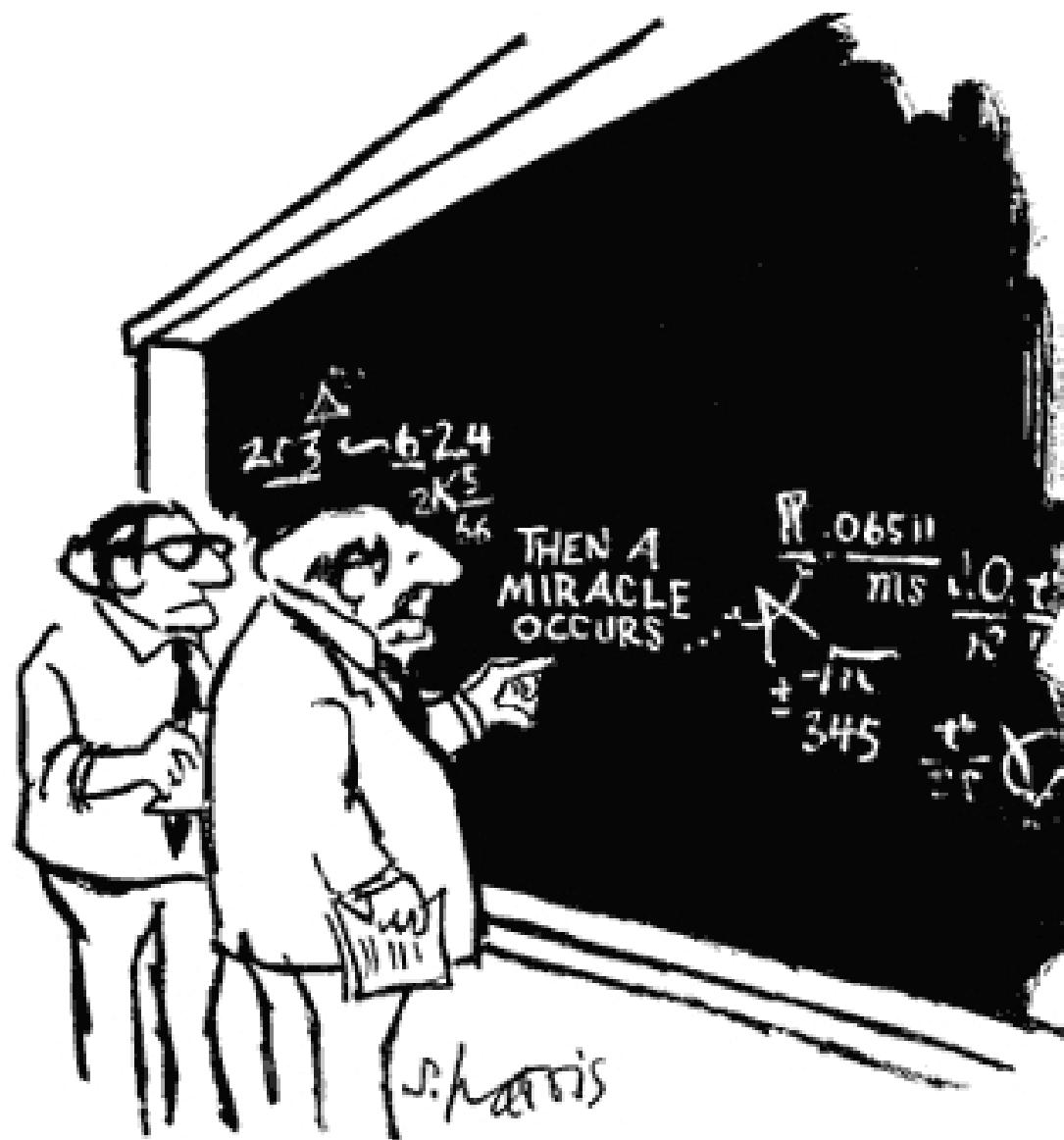
$$M_H = 125.09 \pm 0.24 \text{ GeV}$$

$$\alpha_s(M_Z) = 0.1185 \pm 0.0006$$

precise experiments ...

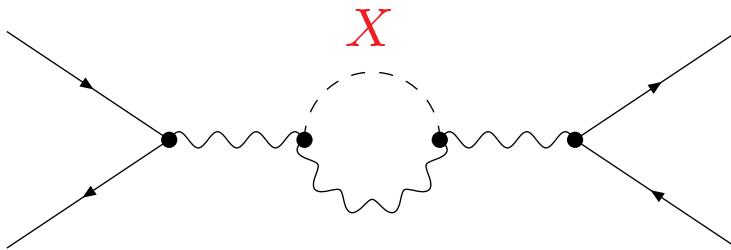


... need precise calculations



"I think you should be more explicit here in step two."

role of theory: exploiting the quantum structure



sensitivity to heavy internal particles (X)

Standard Model: $X = \text{Higgs, top}$

electroweak precision observables

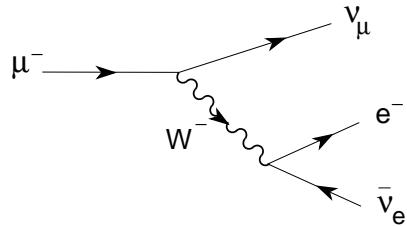
- μ lifetime: G_F
- Z observables: $M_Z, \Gamma_Z, \sin^2 \theta_{\text{eff}}, \dots$
- LEP 2, Tevatron, LHC: $M_W, m_t + M_H$
- low energy: $(g - 2)_\mu$

$M_W - M_Z$ correlation

Definition of Fermi constant G_F via muon lifetime:

$$\tau_\mu^{-1} = \frac{G_F^2 m_\mu^5}{192\pi^3} F \left(\frac{m_e^2}{m_\mu^2} \right) \left(1 + \frac{3}{5} \frac{m_\mu^2}{M_W^2} \right) (1 + \Delta q)$$

Δq : QED corrections in Fermi Model,



$$\frac{G_F}{\sqrt{2}} = \frac{\pi\alpha}{M_W^2 (1 - M_W^2/M_Z^2)} + [\text{higher orders}]$$

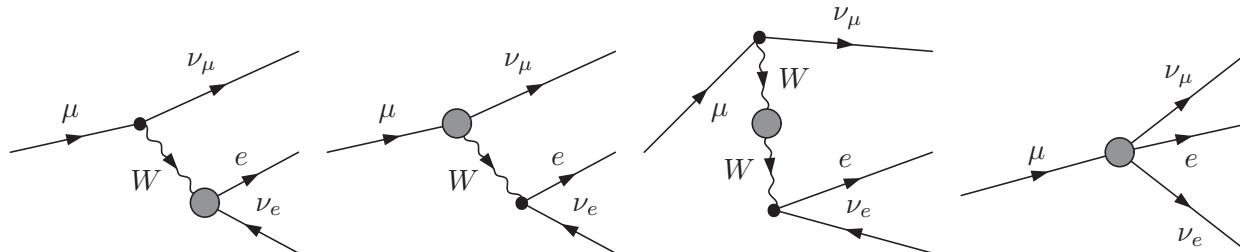
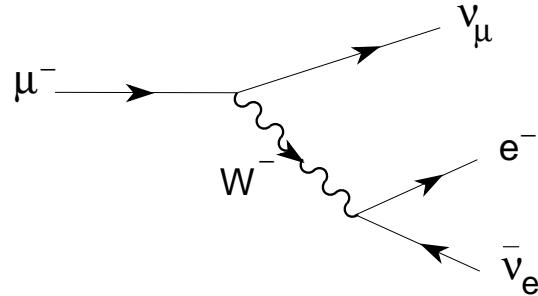
- ★ \Rightarrow prediction of M_W from other parameters
- ★ \Rightarrow each amplitude $\sim G_F$ requires these higher-order terms

$M_W - M_Z$ correlation

Definition of Fermi constant G_F via muon lifetime:

$$\tau_\mu^{-1} = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) \left(1 + \frac{3}{5} \frac{m_\mu^2}{M_W^2}\right) (1 + \Delta q)$$

Δq : QED corrections in Fermi Model,



vector-boson mass correlation

$$\frac{G_F}{\sqrt{2}} = \frac{\pi\alpha}{M_W^2 (1 - M_W^2/M_Z^2)} \cdot (1 + \Delta r)$$

Δr : quantum correction

$$\Delta r = \Delta r(M_Z, M_W, m_t, M_H)$$

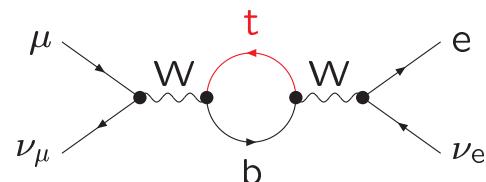
determines W mass

$$M_W = M_W(\alpha, G_F, M_Z, m_t, M_H)$$

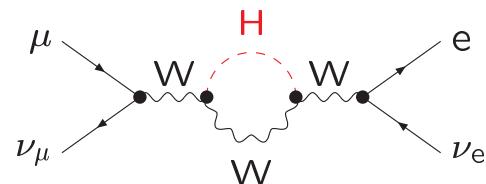
complete at 2-loop order α^2 and $\alpha\alpha_s$

1-loop entries

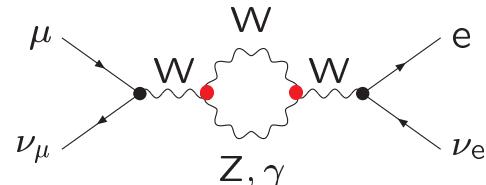
- top quark



- Higgs boson



- gauge-boson self-couplings



dominant contributions to Δr

$$\Delta r = \Delta\alpha - \frac{c_w^2}{s_w^2} \Delta\rho + \dots$$

Large universal terms:

$$\Delta\alpha = \Pi_{\text{ferm}}^\gamma(M_Z^2) - \Pi_{\text{ferm}}^\gamma(0) = 0.05907 \pm 0.0001$$

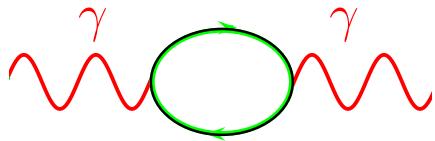
$$\Delta\rho = \frac{\Sigma_Z(0)}{M_Z^2} - \frac{\Sigma_W(0)}{M_W^2} = 3 \frac{G_F m_t^2}{8\pi^2 \sqrt{2}} = 0.0094 \quad [\text{one-loop}] \quad \sim \frac{m_t^2}{v^2} \sim \alpha_t$$

beyond 2-loop order: $\Delta\rho^{(3)} + \Delta\rho^{(4)} \sim \alpha_s^2 \alpha_t, \alpha_s \alpha_t^2, \alpha_t^3, \alpha_s^3 \alpha_t$

reducible higher order terms from $\Delta\alpha$ and $\Delta\rho$ via

$$1 + \Delta r \rightarrow \frac{1}{\left(1 - \Delta\alpha\right) \left(1 + \frac{c_w^2}{s_w^2} \Delta\rho\right) + \dots}$$

photon vacuum polarization



$$\Pi_{\text{ferm}}^\gamma(M_Z^2) - \Pi_{\text{ferm}}^\gamma(0) \equiv \Delta\alpha \quad \rightarrow \quad \alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha} \simeq \frac{1}{129}$$

$$\Delta\alpha = \Delta\alpha_{\text{lept}} + \Delta\alpha_{\text{had}},$$

$$\Delta\alpha_{\text{lept}} = 0.031498 \text{ (4-loop)} \quad \text{Steinhauser 1998; Sturm 2013}$$

$$\Delta\alpha_{\text{had}} = 0.02757 \pm 0.00010 \quad \text{Davier et al. 2010}$$

$$= 0.027626 \pm 0.000103 \quad \text{Hagiwara et al. 2011}$$

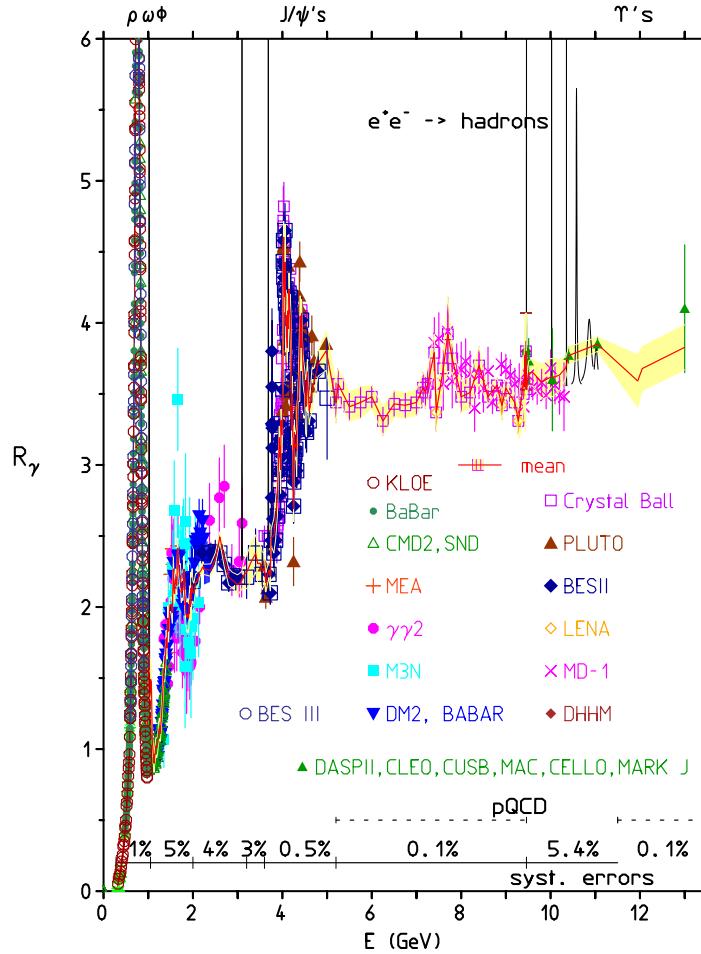
$$= 0.027504 \pm 0.000118 \quad \text{Jegerlehner 2015}$$

significant parametric uncertainty

$$\Delta\alpha_{\text{had}} = -\frac{\alpha}{3\pi} M_Z^2 \operatorname{Re} \int_{4m_\pi^2}^\infty ds' \frac{R_{\text{had}}(s')}{s'(s' - M_Z^2 - i\epsilon)}$$

$R_{\text{had}} =$

$$\frac{\sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \gamma^* \rightarrow \mu^+\mu^-)}$$



Jegerlehner 2015

EW 2-loop calculations for Δr

Freitas, Hollik, Walter, Weiglein

Awramik, Czakon

Onishchenko, Veretin

Degrassi, Gambino, Giardino

universal terms at 3- and 4-loops (EW and QCD)

van der Bij, Chetyrkin, Faisst, Jikia, Seidensticker

Faisst, Kühn Seidensticker, Veretin

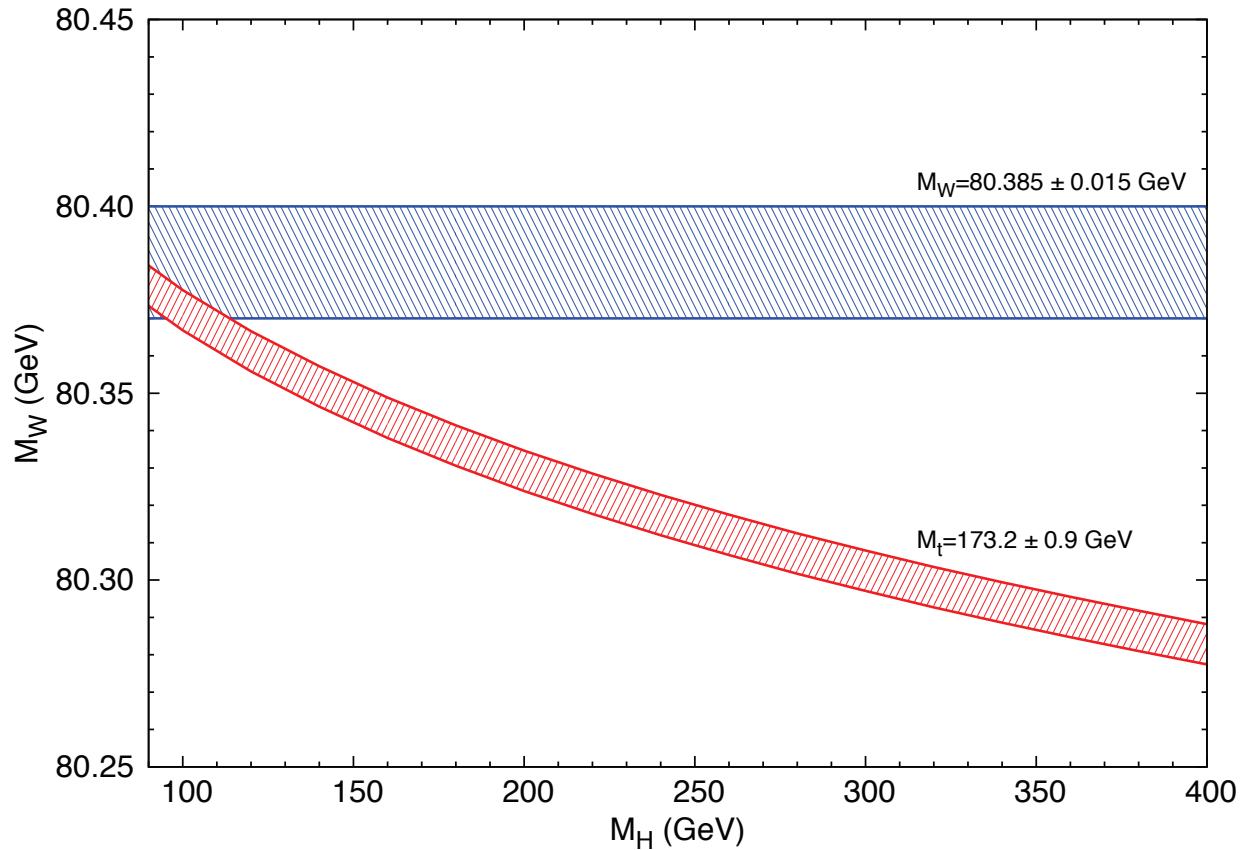
Boughezal, Tausk, van der Bij

Schröder, Steinhauser

Chetyrkin, Faisst, Kühn

Chetyrkin, Faisst, Kühn, Maierhofer, Sturm

Boughezal, Czakon

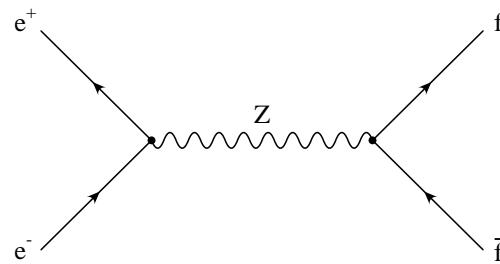
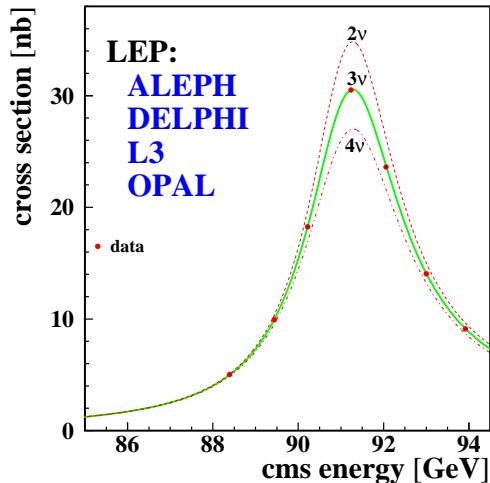


variation of Δr by 0.001 $\Rightarrow \delta M_W = 18 \text{ MeV}$

present exp. error: $\Delta M_W = 15 \text{ MeV}$ / theo: 4 MeV

$\delta(\Delta\alpha) = 10^{-4} :$ $\Delta M_W = 1.8 \text{ MeV}$

Z resonance



- effective Z boson couplings with higher-order $\Delta g_{V,A}$

$$g_V^f \rightarrow g_V^f + \Delta g_V^f, \quad g_A^f \rightarrow g_A^f + \Delta g_A^f$$

- effective ew mixing angle (for $f = e$):

$$\sin^2 \theta_{\text{eff}} = \frac{1}{4} \left(1 - \text{Re} \frac{g_V^e}{g_A^e} \right) = 1 - \frac{M_W^2}{M_Z^2} + \frac{M_W^2}{M_Z^2} \Delta \rho + \dots$$

two classes of observables;

- line-shape observables: M_Z , Γ_Z , *partial widths*, σ_{peak} , . . .

$$R_\ell = \frac{\Gamma_{\text{had}}}{\Gamma_{\text{lept}}} , \quad R_b = \frac{\Gamma_b}{\Gamma_{\text{had}}}$$

EW 2-loop fermionic corrections [Freitas; Freitas, Y.-C. Huang]

hadronic width:

$$\Gamma_{\text{had}} = \Gamma_{\text{had}}^{\text{EW}} \cdot R_{\text{QCD}}$$

$$R_{\text{QCD}} = 1 + \frac{\alpha_s}{\pi} + A_2 \left(\frac{\alpha_s}{\pi} \right)^2 + A_3 \left(\frac{\alpha_s}{\pi} \right)^3 + A_4 \left(\frac{\alpha_s}{\pi} \right)^4$$

[Baikov, Chetyrkin, Kühn, Rittinger]

two classes of observables;

- **asymmetries:** A_{FB} , A_{LR}, \dots

$$\Rightarrow \text{forward-backward asymmetry} \quad A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} A_e A_f$$

$$A_f = \frac{2 g_V^f / g_A^f}{1 + (g_V^f / g_A^f)^2} \quad \Rightarrow \quad \sin^2 \theta_{\text{eff}}^f$$

polarized cross section for $e_{L,R}^-$:

$$\Rightarrow \text{left-right asymmetry} \quad A_{LR} = \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = A_e \quad \Rightarrow \quad \sin^2 \theta_{\text{eff}}$$

asymmetries determine $\sin^2 \theta_{\text{eff}}$

EW 2-loop calculations for $\sin^2 \theta_{\text{eff}}$

Awramik, Czakon, Freitas, Weiglein

Awramik, Czakon, Freitas

Hollik, Meier, Uccirati

universal terms at 3- and 4-loops (EW and QCD)

van der Bij, Chetyrkin, Faisst, Jikia, Seidensticker

Faisst, Kühn Seidensticker, Veretin

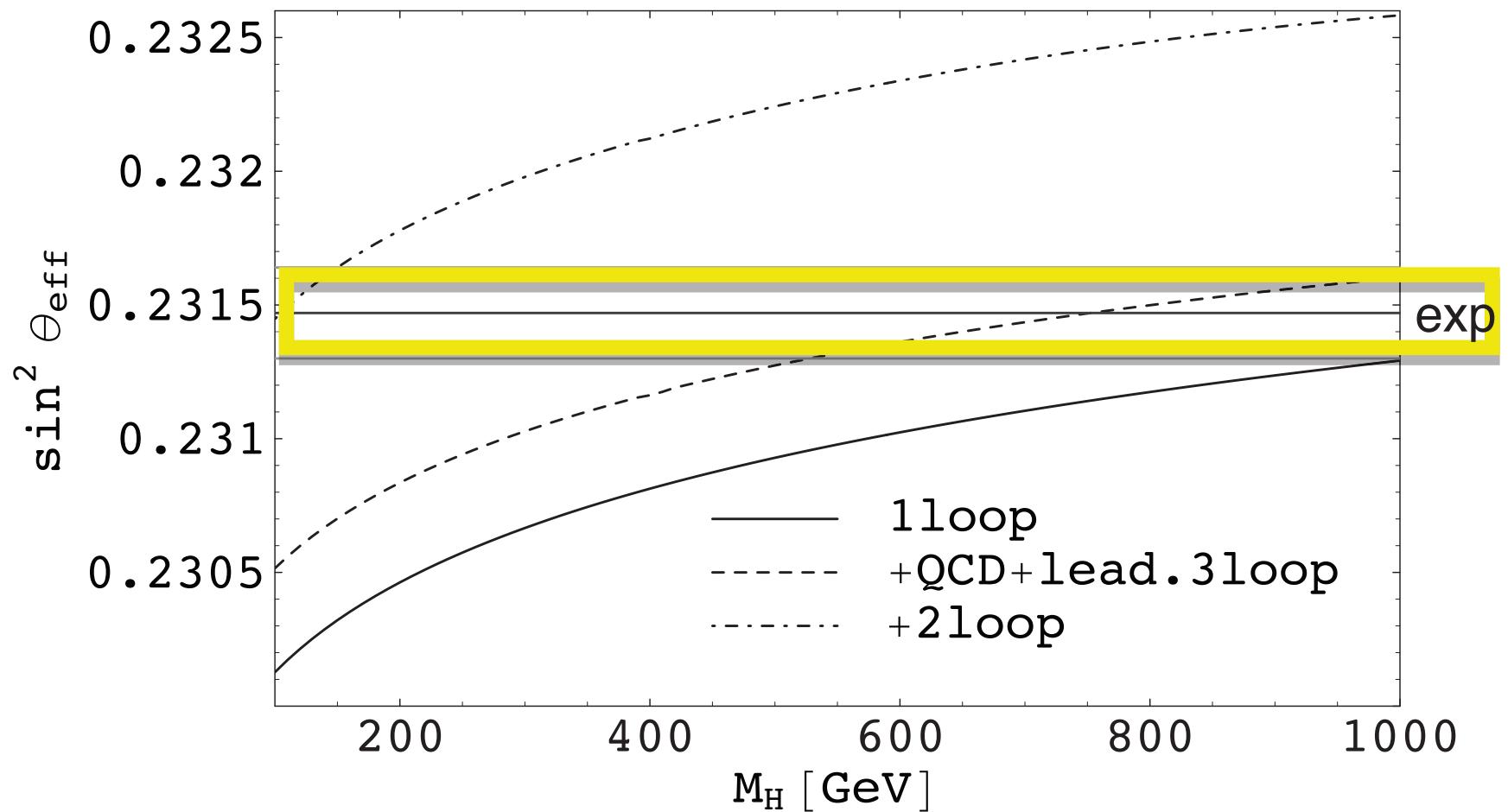
Boughezal, Tausk, van der Bij

Schröder, Steinhauser

Chetyrkin, Faisst, Kühn

Chetyrkin, Faisst, Kühn, Maierhofer, Sturm

Boughezal, Czakon



theory uncertainty: $\delta \sin^2 \theta_{\text{eff}} = 5 \cdot 10^{-5}$

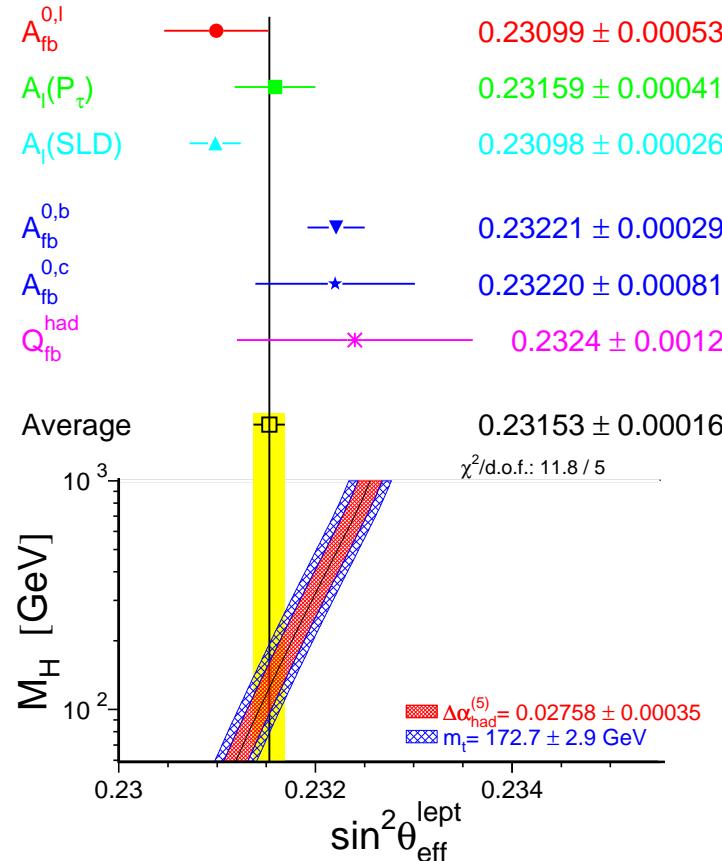
$\delta(\Delta\alpha) = 10^{-4}$: $\delta \sin^2 \theta_{\text{eff}} = 3.5 \cdot 10^{-5}$

some observables with not-so-good agreement

- in general, SM is in overall agreement with data
 - yet a few quantities stand a bit apart ($\sim 3\sigma$)
 - the forward-backward asymmetry for b quarks,
 $A_{\text{FB}}^{b\bar{b}}$ at the Z peak
 - the anomalous magnetic moment of the muon
 - the forward-backward asymmetry for top quarks
at the Tevatron, $p\bar{p} \rightarrow t\bar{t}$
- new NNLO QCD calculation: *Czakon, Fiedler, Mitov 2015*
data now compatible with SM prediction (QCD + EW)

forward-backward asymmetry for b quarks

- is around since the days of LEP/SLC
- indicates presence of some extra right-handed contributions to the Zbb coupling
- no convincing explanation in current new physics models



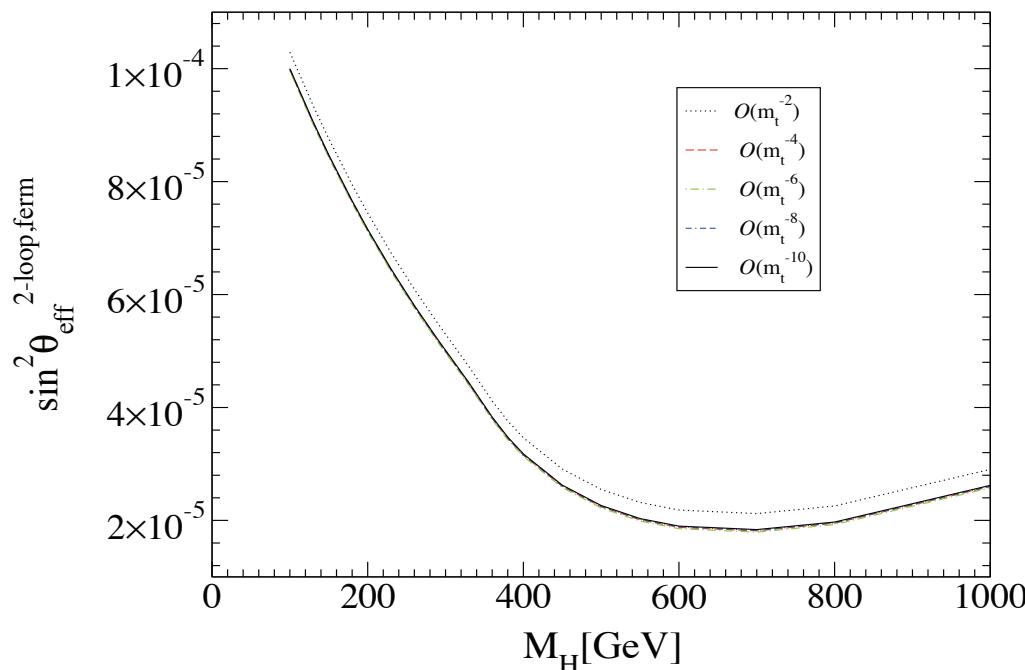
A_{FB}^b determined by $\sin^2 \theta_e$ and $\sin^2 \theta_b$:

$$A_{\text{FB}}^b = \frac{3}{4} A_e A_b$$

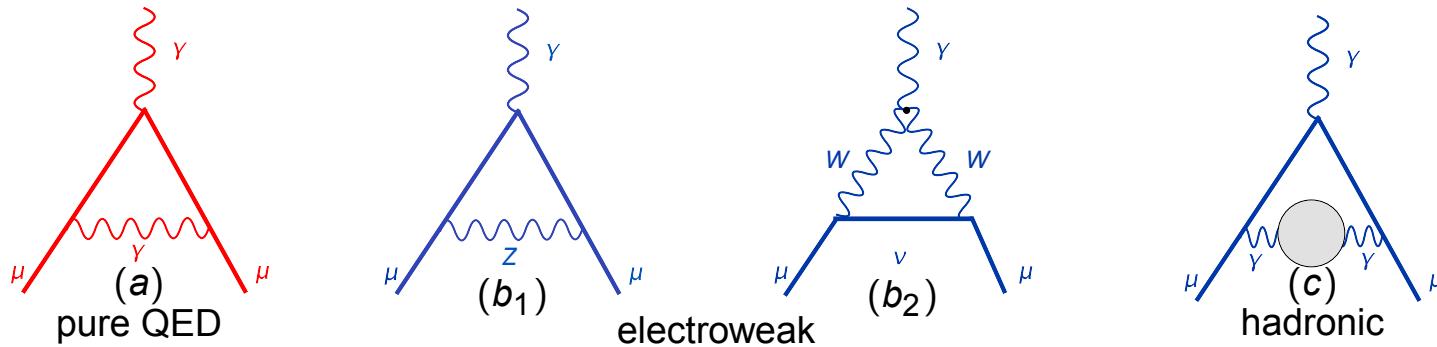
$$A_e = \frac{1-4 \sin^2 \theta_e}{1+(1-4 \sin^2 \theta_e)^2}, \quad A_b = \frac{1-4/3 \sin^2 \theta_b}{1+(1-4/3 \sin^2 \theta_b)^2}$$

recent electroweak 2-loop calculation – too small

Awramik, Czakon, Freitas, Kniehl 08

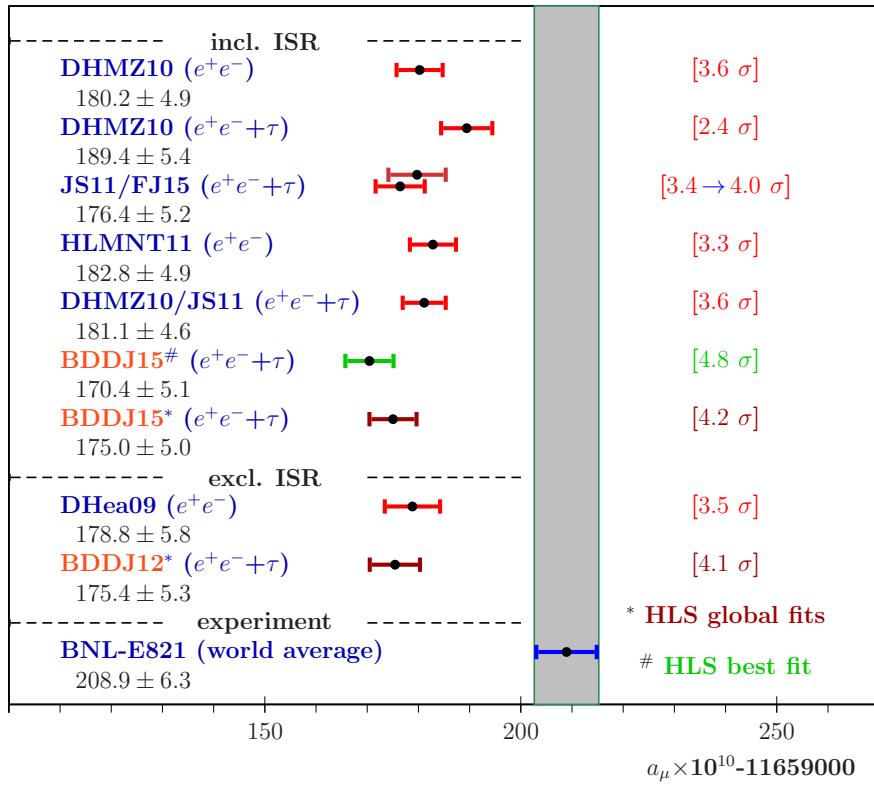


anomalous magnetic moment of the muon



- The QED part is known to 5 loops
Kinoshita et al.
- The EW part is known to 2 loops
Czarnecki, Krause, Marciano;
Gnendiger, Stöckinger, Stöckinger-Kim
- The hadronic part is known with limited accuracy,
hadronic vacuum polarization from data at low energies
many authors, most recent: Jegerlehner 2015

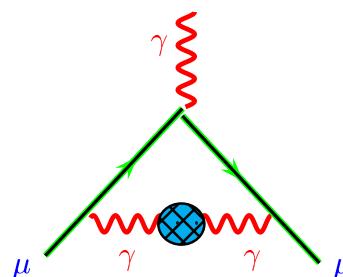
Anomalous g-factor of the muon



Jegerlehner 2015

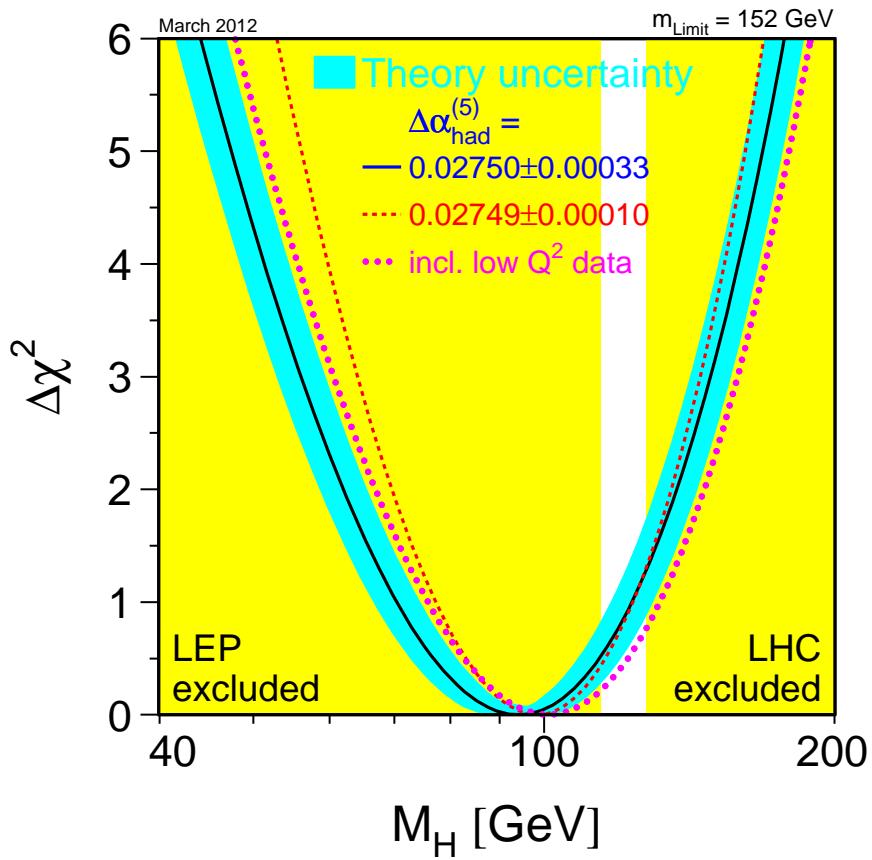
SM prediction: 4σ below exp. value

theory uncertainty from hadronic vacuum polarization



Global fit to the Higgs boson mass

before the discovery in 2012



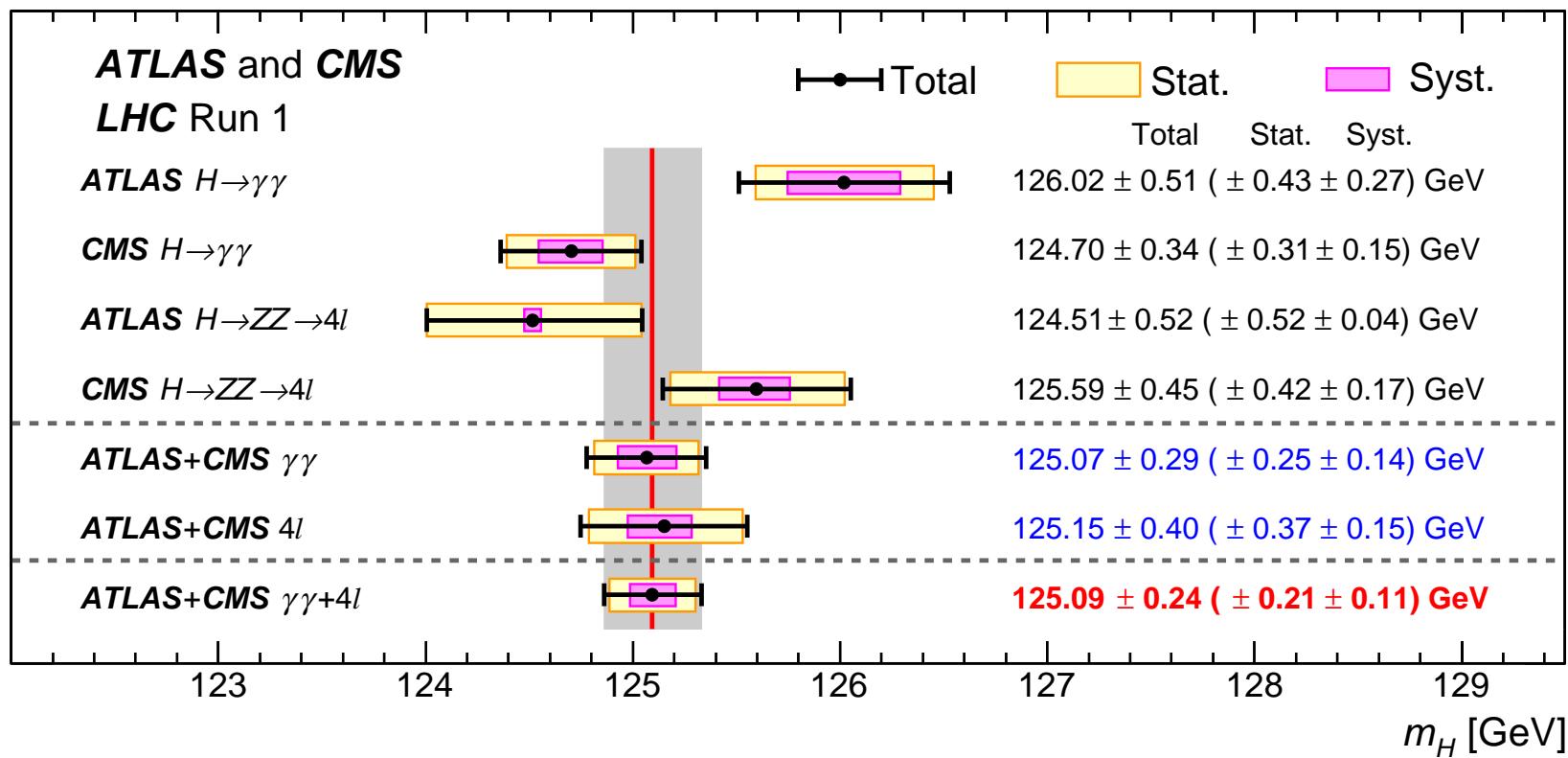
blueband: Theory uncertainty

$M_H < 152 \text{ GeV} \text{ (95\%C.L.)}$

$M_H = 94^{+29}_{-24} \text{ GeV}$

The Higgs era

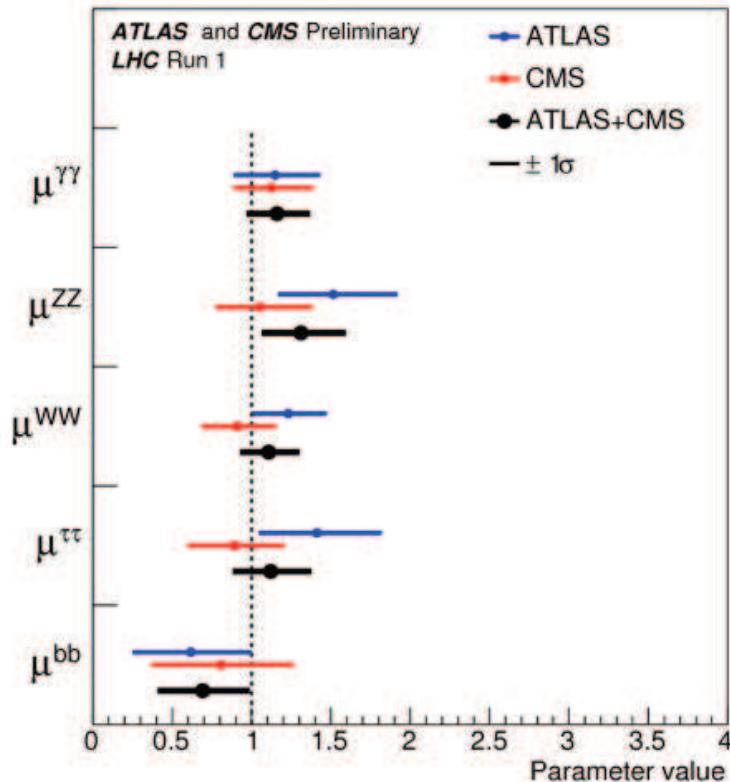
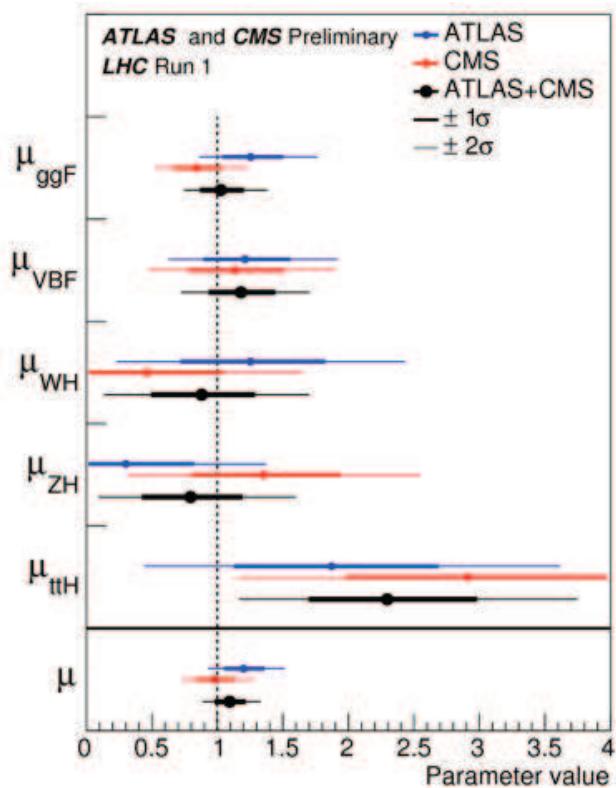
2015: first combined ATLAS/CMS Higgs-boson mass determination





A Standard Model Higgs boson at the LHC?

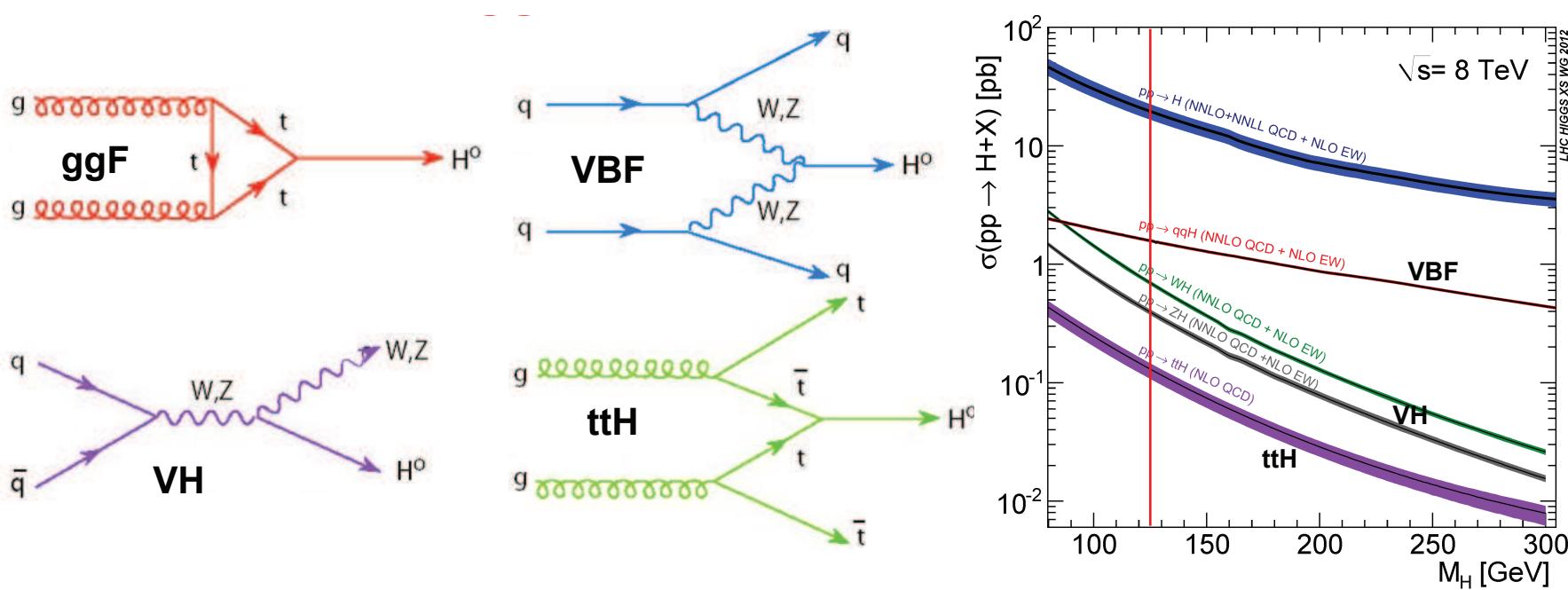
Moriond 2016



$$\mu = 1.09^{+0.11}_{-0.10}$$

$$\mu = \sigma(pp \rightarrow H) \cdot BR(H \rightarrow X) / [\text{SM}_{\text{theory}}]$$

Higgs production at the LHC



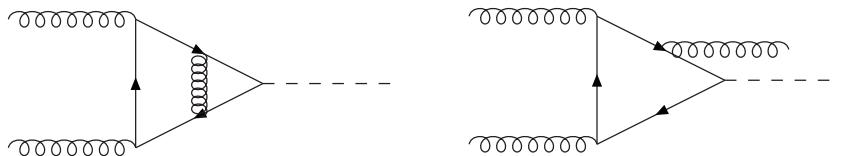
Handbook of Higgs Cross Sections, Vol. 1, 2, 3

CERN-2011-002, CERN-2012-002, CERN-2013-004

NLO exact

Dawson (1991); Djouadi, Spira, Zerwas (1991)

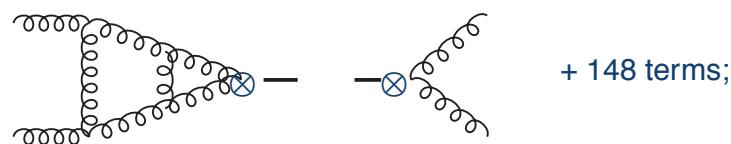
Graudenz, Spira, Zerwas (1993)



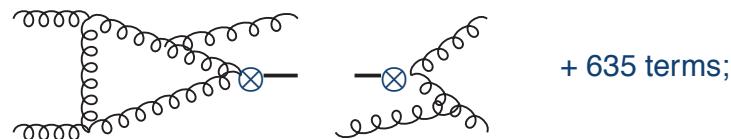
NNLO computed using effective vertex

Harlander, Kilgore (2002)
Anastasiou, Melnikov (2002)
Ravindran, Smith, van Neerven (2003)

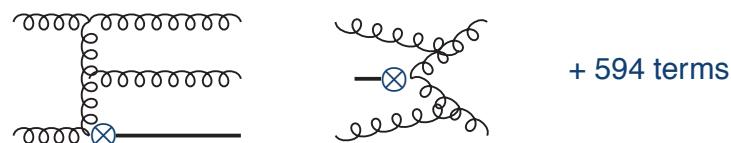
Double Virtual:



Real Virtual:



Double Real:

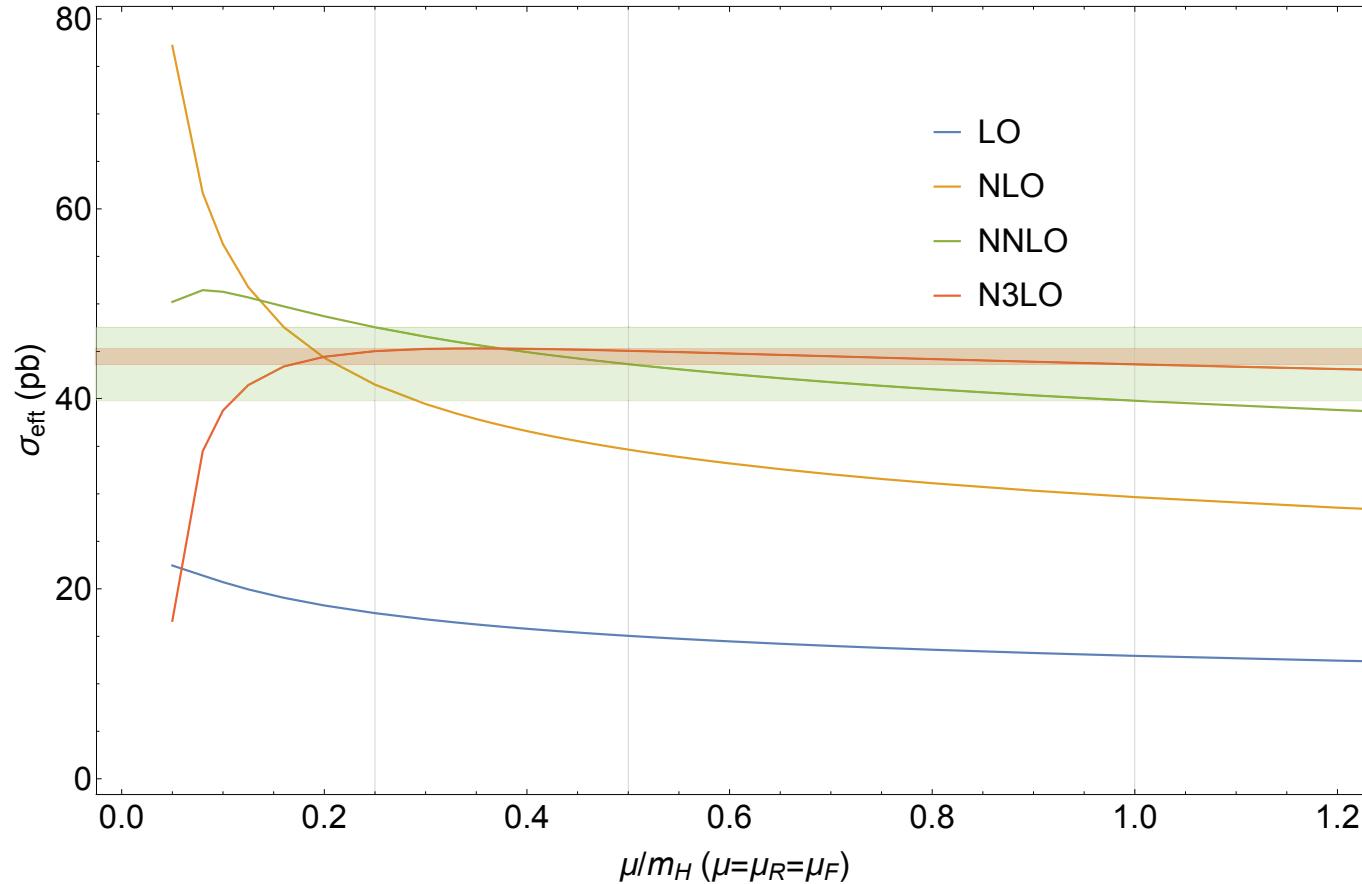


In addition:

$$q + g \rightarrow h + X(q, \bar{q}, g) \quad q_i + q_j(\bar{q}_j) \rightarrow h + X(q, \bar{q}, g)$$

NNNLO for $gg \rightarrow H$

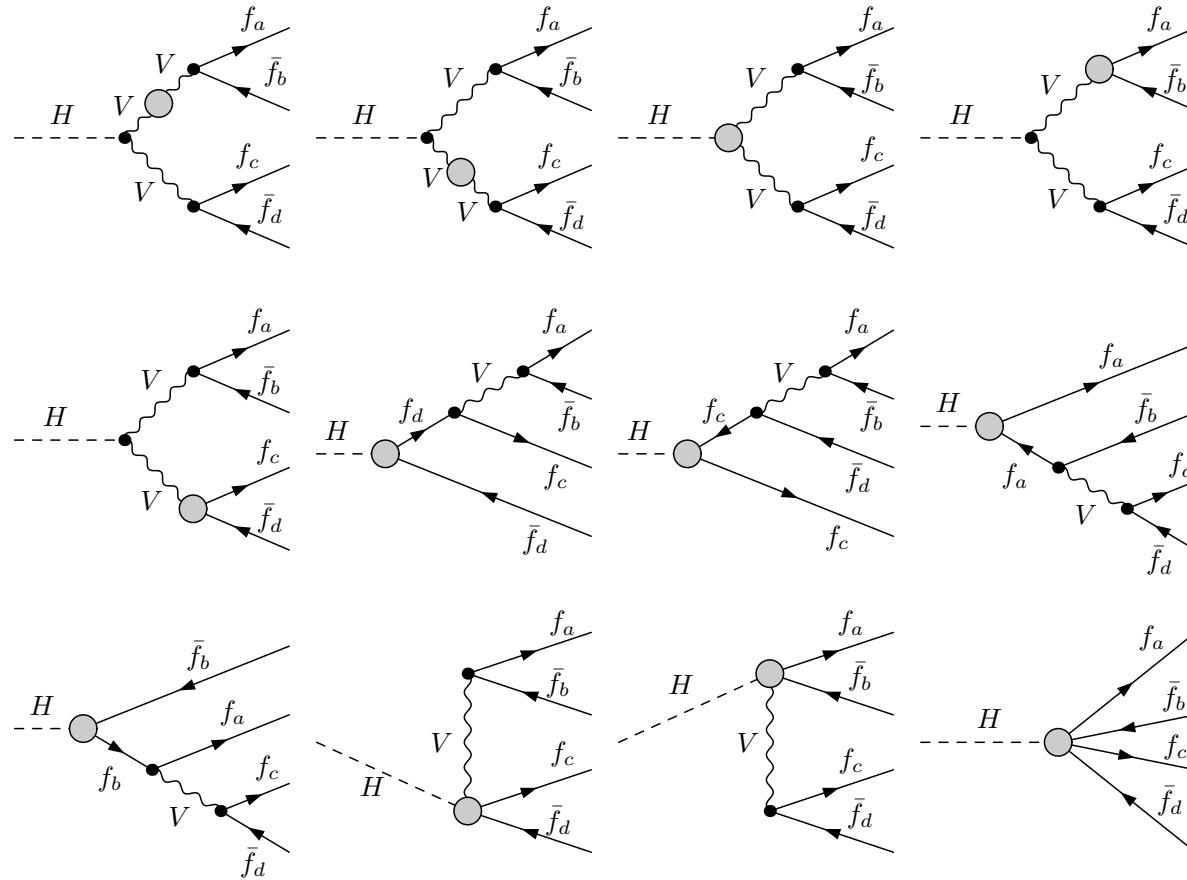
Anastasiou et al. 2015



scale variation
at N3LO ~2%

decays: $H \rightarrow VV \rightarrow 4f$

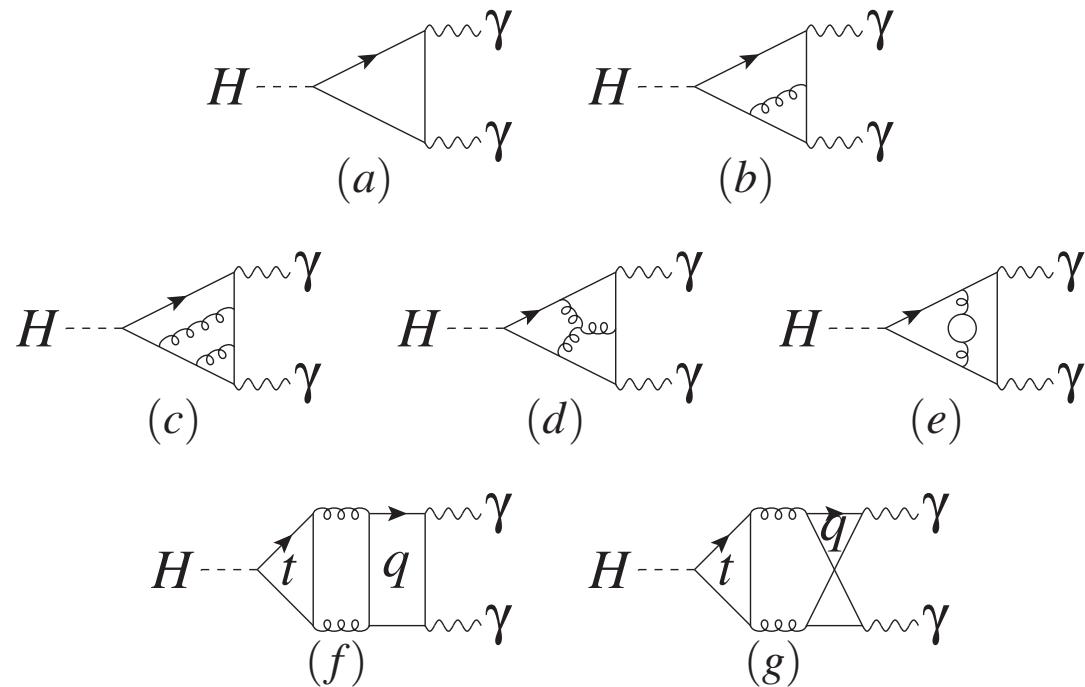
needs also background processes + h.o.



Bredenstein, Denner, Dittmaier, Weber 2007

PROPHECY4f

decays: $H \rightarrow \gamma\gamma$



non-singlet and singlet terms; electroweak corrections (Passarino,...)

$$\Gamma_{H \rightarrow \gamma\gamma} = (9.398 - \frac{0.148}{\text{LO} \times \text{NLO-EW}} + \frac{0.168}{\text{LO} \times \text{NLO-QCD}} + 0.00793) \text{ keV}$$

α_s^2 term dominated by singlet part of prediction,

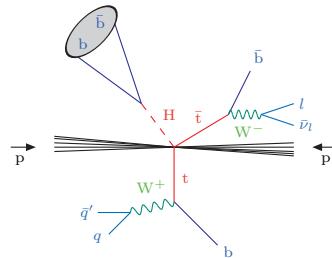
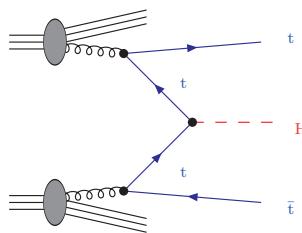
prediction good to 1 permille!

precision physics in the Higgs era

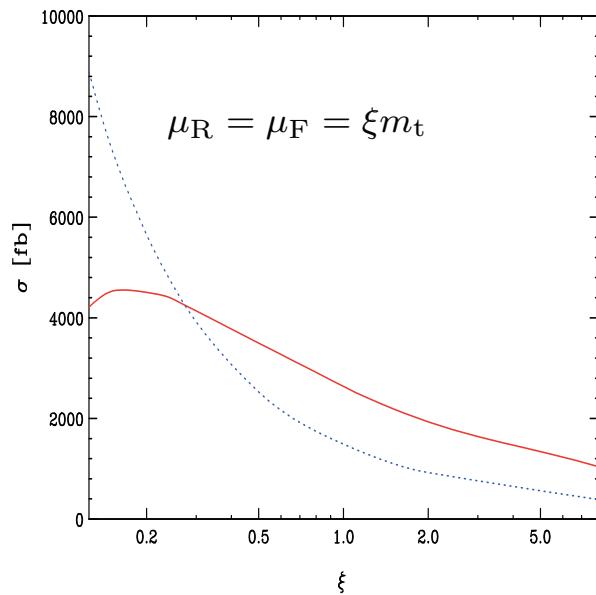
- Higgs boson mass M_H is now an additional precise input parameter
- partial decay widths influenced by quantum effects
 - accordingly, affected also by non-standard physics
 - new set of sensitive precision observables
- M_H itself is a prediction/ precision observable in specific classes of models
- further empirical information available on top of the “classical” set of precision observables

production of $t\bar{t}H$

- measurement of the Higgs–top Yukawa coupling



- requires accurate background calculations, e.g.
for $pp \rightarrow t\bar{t} b\bar{b}$



.... LO

— NLO

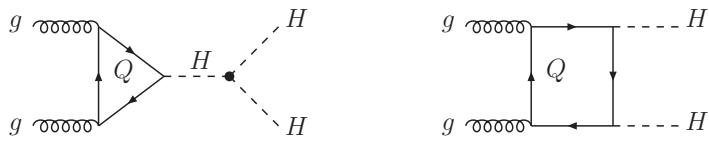
Bredenstein et al.

Bevilacqua et al.

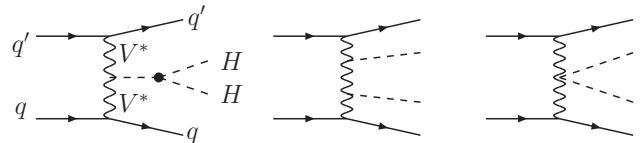
measurement of Higgs self-coupling

HH production channels

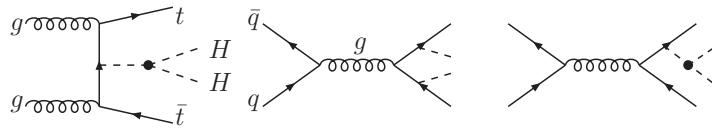
gg fusion



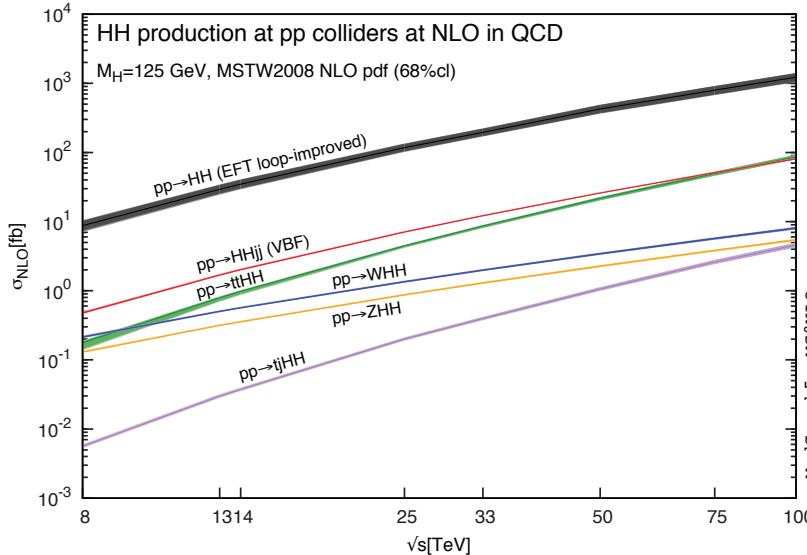
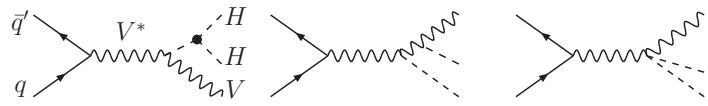
VV fusion



Associated production with top



Higgs-strahlung



\sqrt{s} [TeV]	$\sigma_{gg \rightarrow HH}^{\text{NLO}}$ [fb]	$\sigma_{qq' \rightarrow HHq q'}^{\text{NLO}}$ [fb]	$\sigma_{qq' \rightarrow WHH}^{\text{NNLO}}$ [fb]	$\sigma_{q\bar{q} \rightarrow ZHH}^{\text{NNLO}}$ [fb]	$\sigma_{q\bar{q}/gg \rightarrow t\bar{t}HH}^{\text{LO}}$ [fb]
8	8.16	0.49	0.21	0.14	0.21
14	33.89	2.01	0.57	0.42	1.02
33	207.29	12.05	1.99	1.68	7.91
100	1417.83	79.55	8.00	8.27	77.82

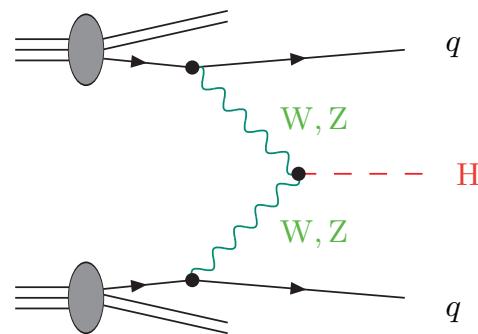
Gluon-gluon fusion dominates
Only some contribute with HHH

production of $VV + 2$ jets

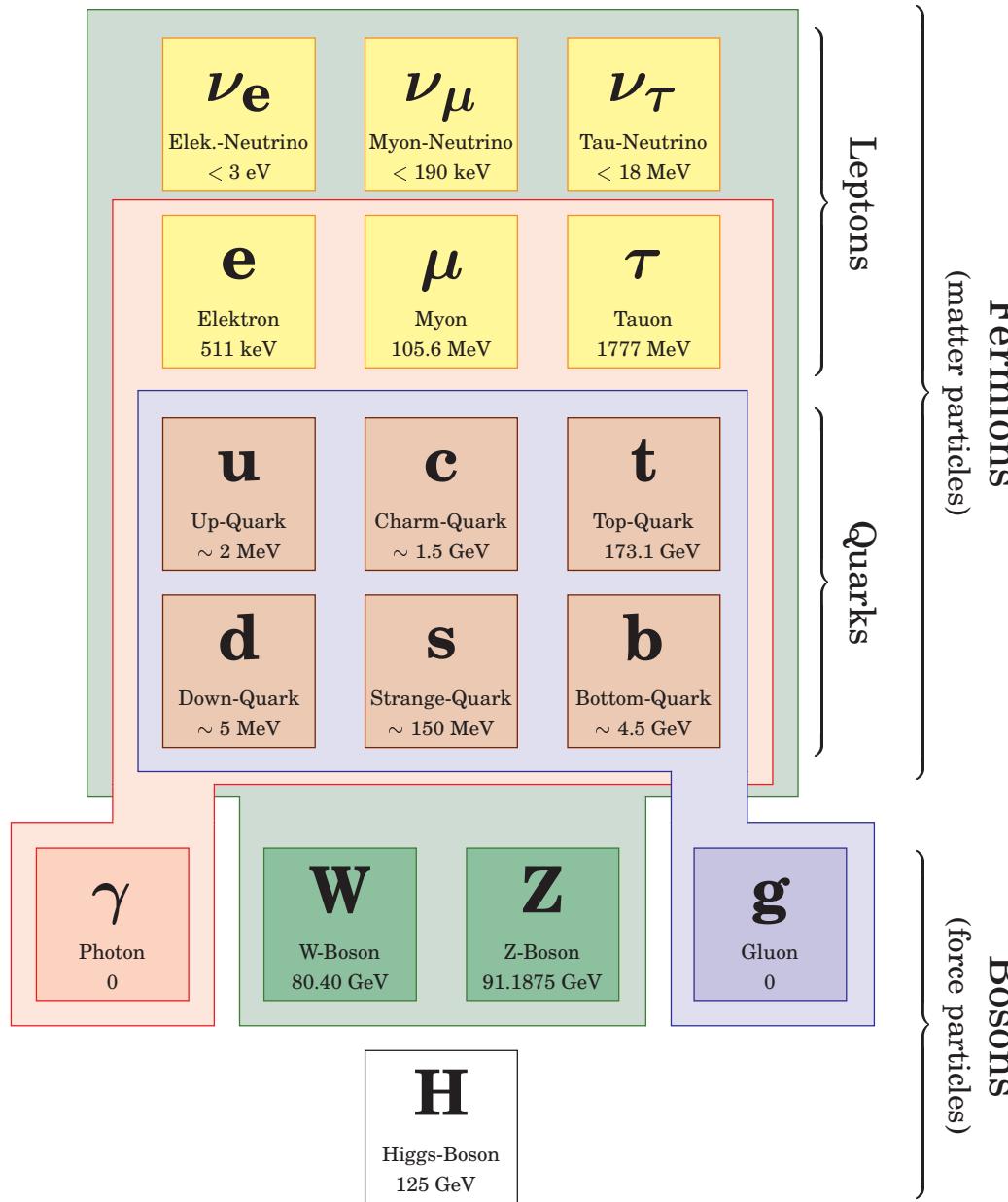
- measure $VV \rightarrow VV$ scattering
sensitive to electroweak symmetry breaking – BSM?



- background for Higgs production via vectorboson fusion with $H \rightarrow VV$

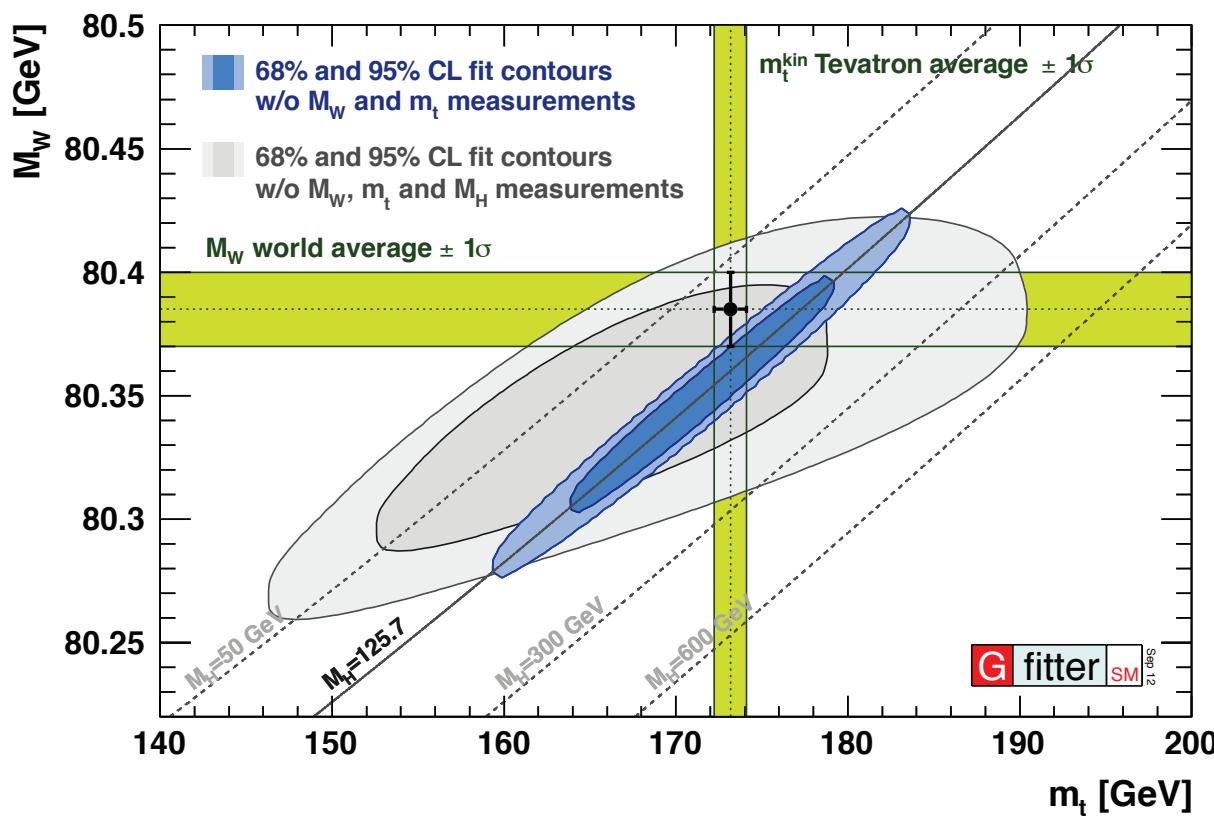


Status of the Standard Model



SM input completely determined \Rightarrow precision observables

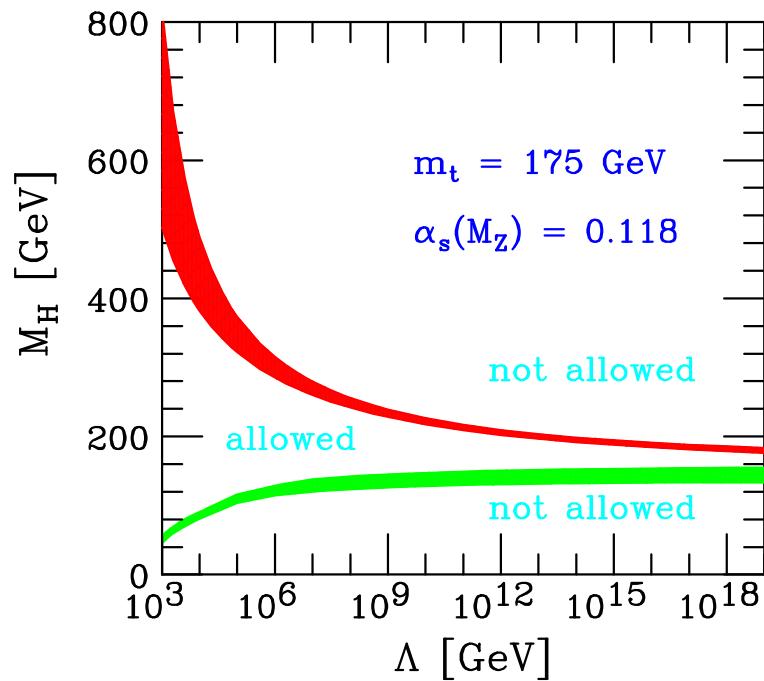
	theo	exp
$\sin^2 \theta_{\text{eff}}$	$0.23152 \pm 0.00005 \pm 0.00005$	0.23153 ± 0.00016
M_W (GeV)	$80.361 \pm 0.006 \pm 0.004$	80.385 ± 0.015



evolution of self-coupling λ with scale Q

$$Q^2 \frac{d\lambda}{dQ^2} = \frac{1}{16\pi^2} (12\lambda^2 - 3g_t^4 + 6\lambda g_t^2 + \dots)$$

vacuum stability: $\lambda(Q) > 0$ up to scale Λ

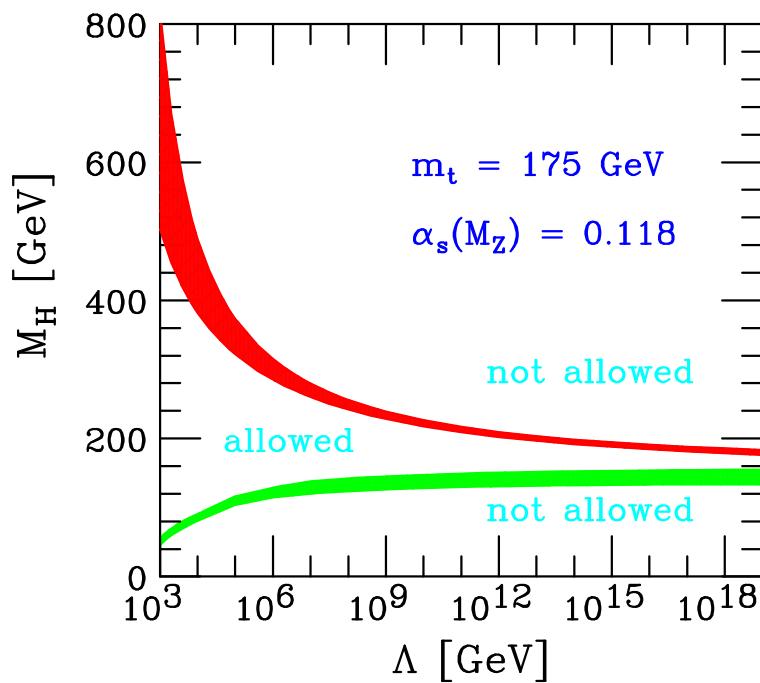


Hambye, Riesselmann 1997

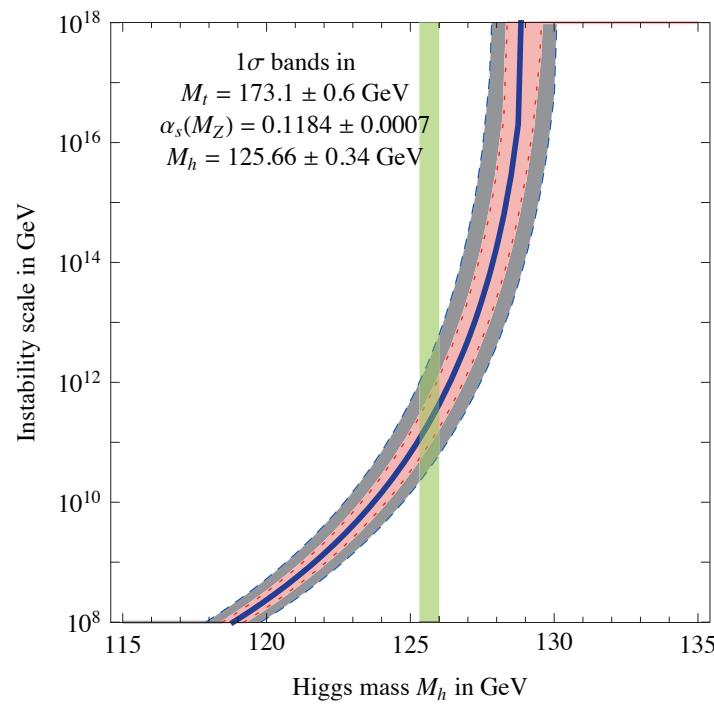
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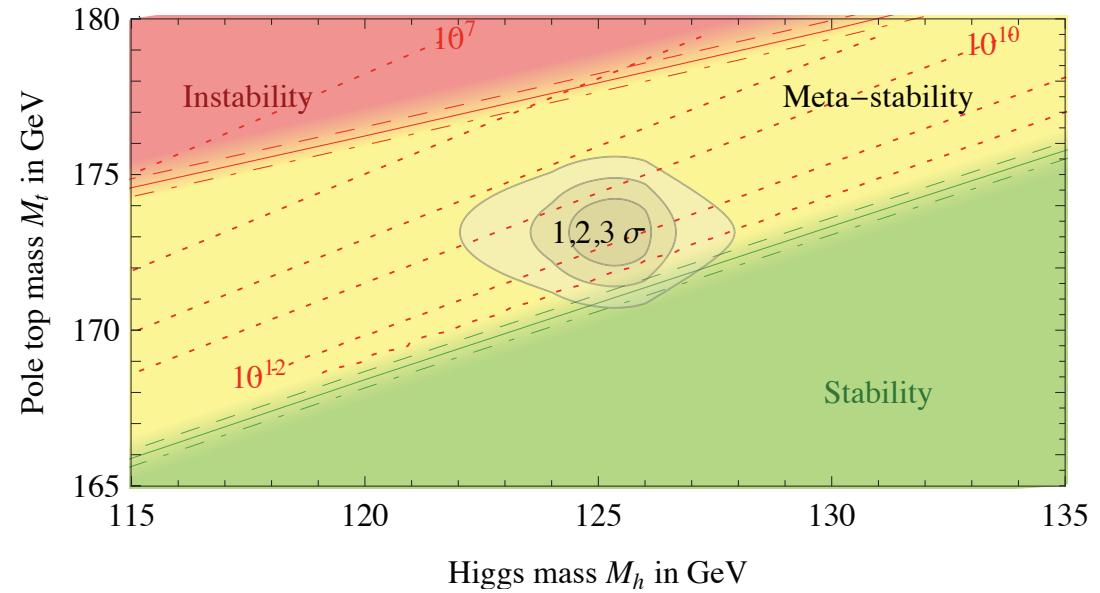
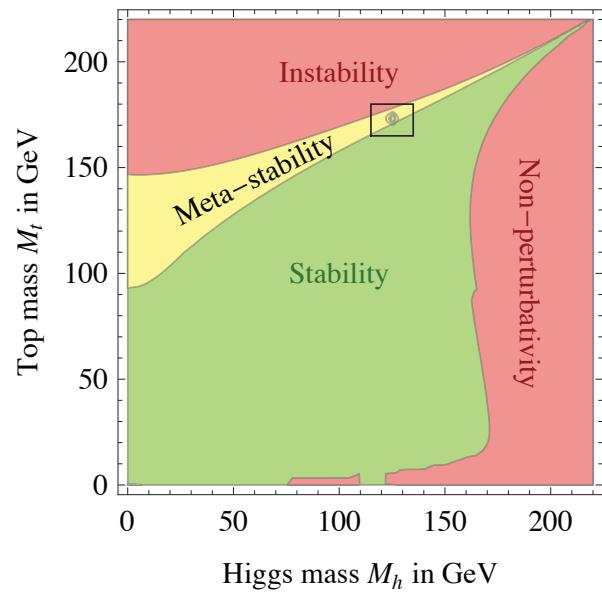
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Hambye, Riesselmann 1997



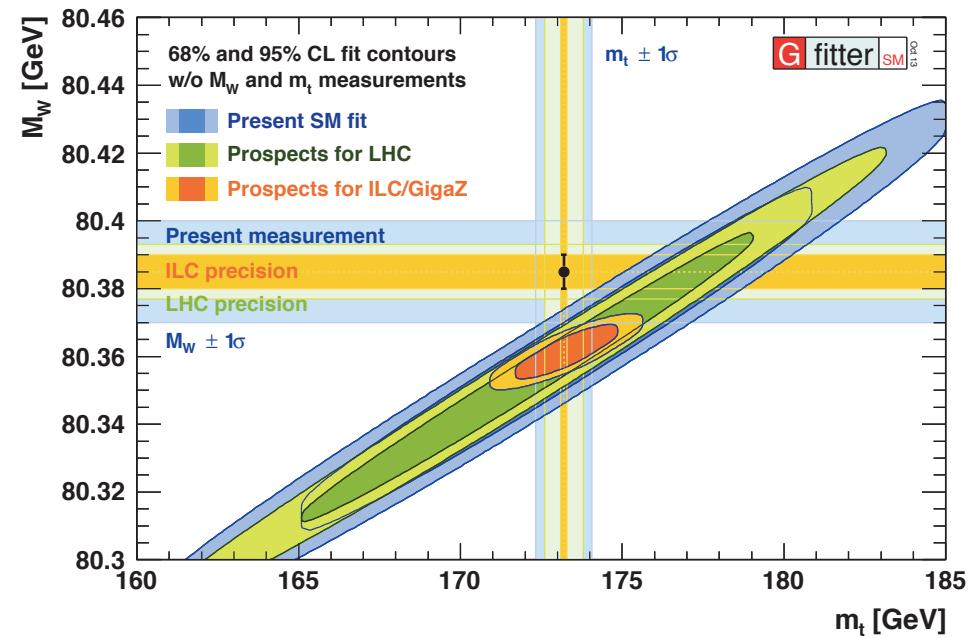
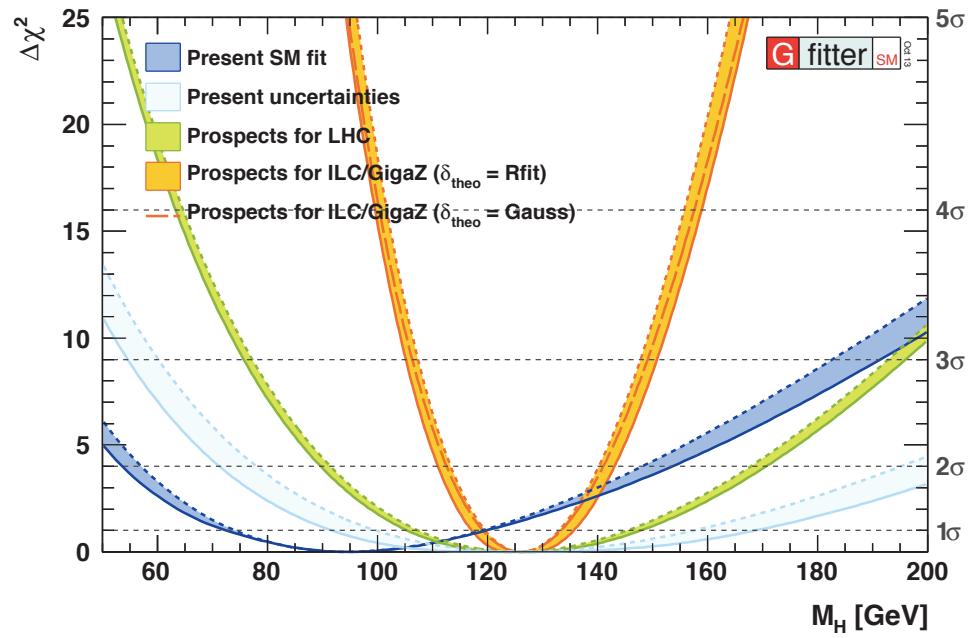
Degrassi et al. 2012



[Degrassi *et al.* 2012]

precise measurement of m_t crucial input

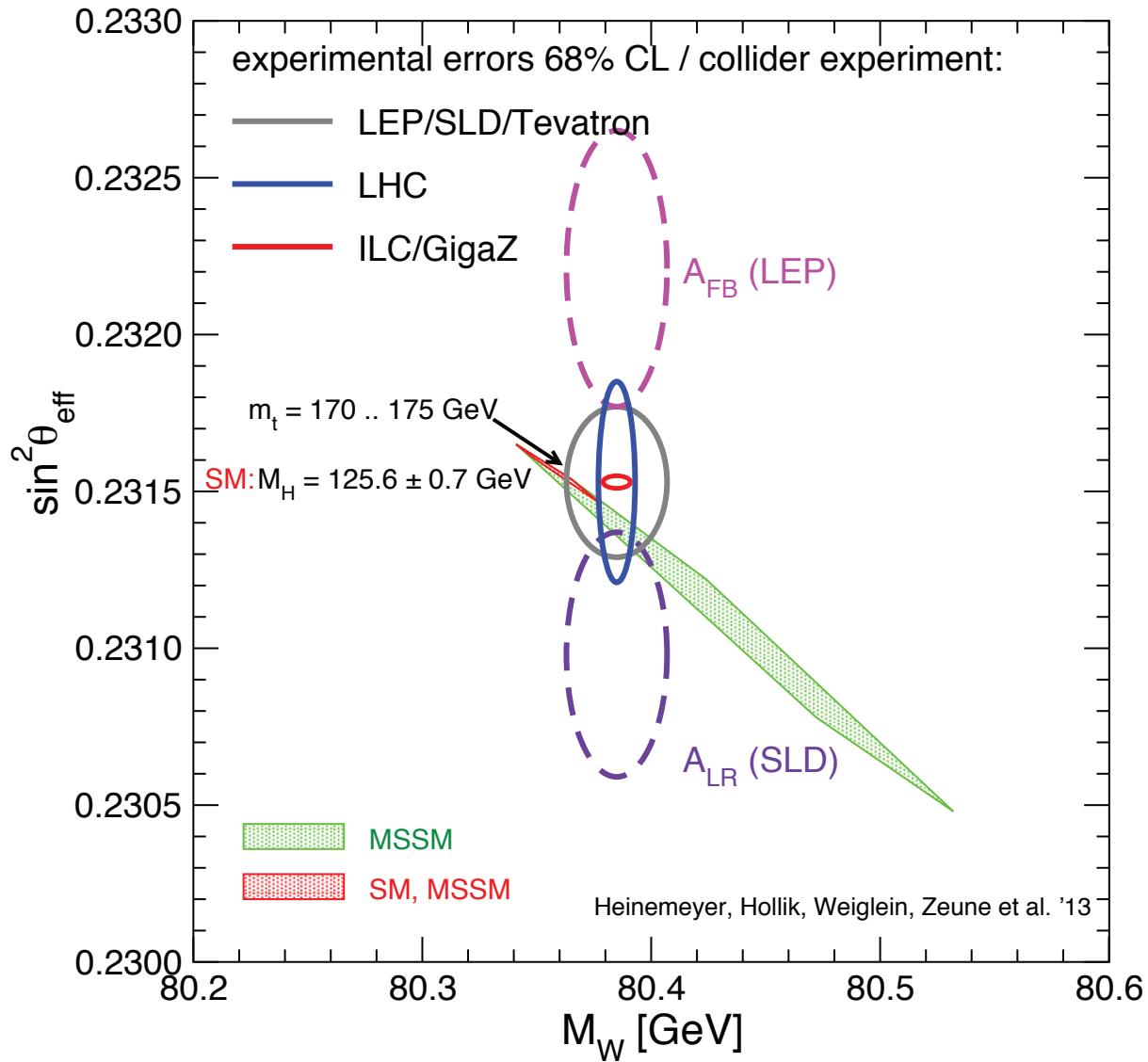
the future of SM tests

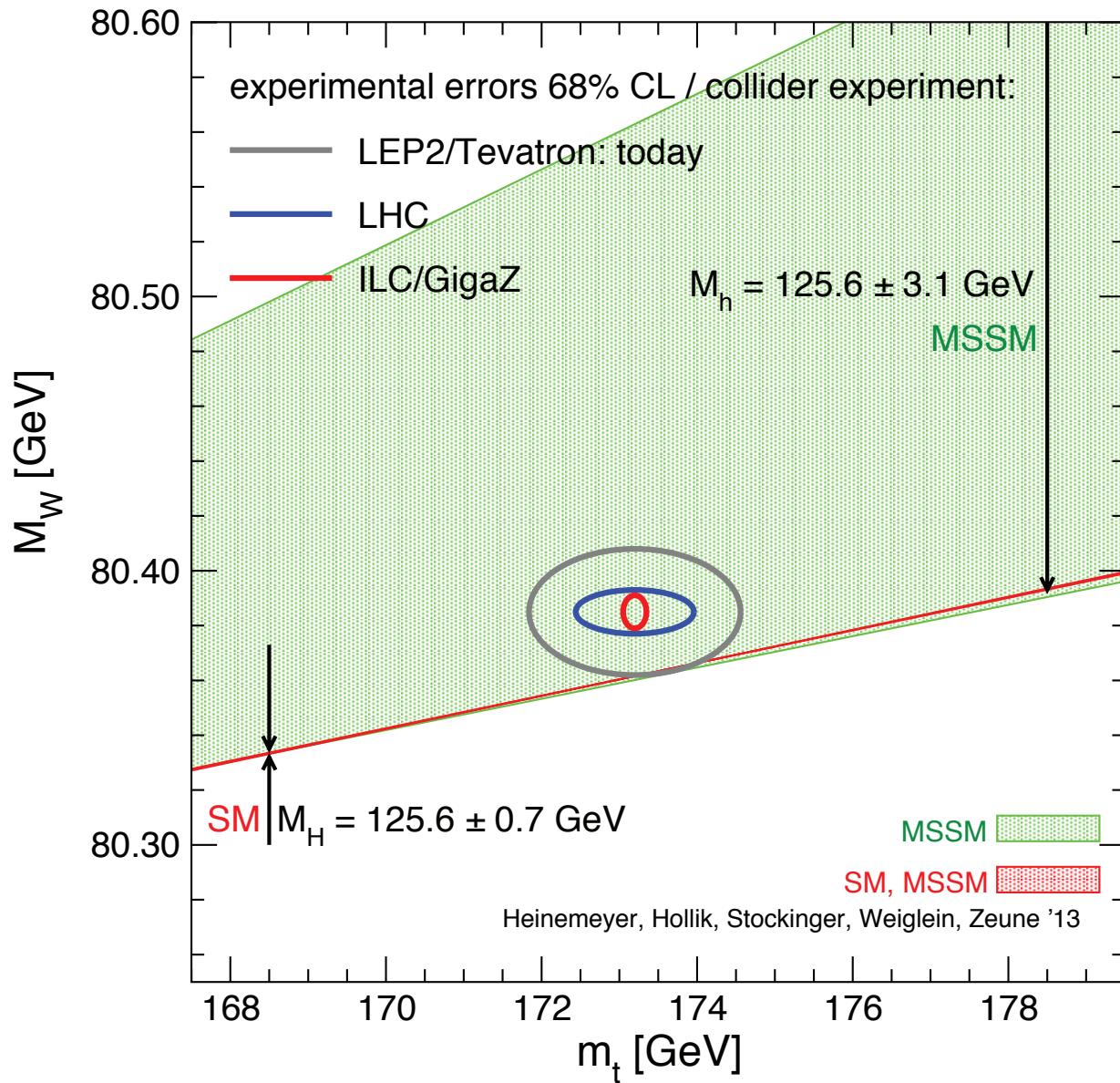


from: *Working Group Report of the 2013 Community Summer Study (Snowmass) on the Energy Frontier Precision Study of Electroweak Interactions*

present and aimed experimental precision

exp. error	LEP/Tev/LHC	LC/GigaZ	FCee
M_Z [MeV]	2.1	2.1	0.1
M_W [MeV]	10	3-5	1
$\sin^2 \theta_{\text{eff}} [10^{-5}]$	17	1.3	0.6
$R_b [10^{-5}]$	66	15	5
m_{top} [GeV]	0.7	0.1	0.05
$\Delta\alpha$	$1 \cdot 10^{-4}$	$5 \cdot 10^{-5}$	$< 5 \cdot 10^{-5}$
α_s	$6 \cdot 10^{-4}$	$< 6 \cdot 10^{-4}$	$1 \cdot 10^{-4}$





Conclusions and perspectives

- impressive confirmation by a huge data sample from low to high energies, no significant deviations
- quantum effects have been established at many σ
- perfect indirect and direct determination of the top quark
- now repeated for the Higgs boson
- scalar particle at 125 GeV: SM Higgs boson or first signal of BSM?
- next steps with upgraded LHC / improved precision
 - confirm the Higgs boson properties
 - check versus electroweak precision measurements
 - or find deviations, new structures

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