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# DERIVING THE WGC FROM ENTANGLEMENT ENTROPY THEOREMS

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**Utrecht University**

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    - The **WGC**: Three proofs in one workshop  
(see Shiu & Remmen's talks)!
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- I will present a “physics proof” of the WGC, to appear (very?) soon.
  - Like any other proof, it has some **assumptions**:
    - **Holography**: Statements only in AdS, use CFT dual
    - Validity of **large N expansion/semiclassical bulk EFT**
  - **No use** of usual crutches: SUSY or even (directly) String Theory
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Exactly stable extremal black branes lead to contradiction with rigorous quantum info theorems.

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- Argument by contradiction: Take a WGC-violating EFT in AdS, and derive an inconsistency.
- For concreteness, focus on Einstein-Maxwell-AdS (+possibly charged matter).

$$\int d^{d+1}x \left( \frac{R}{2\kappa_{d+1}^2} + \frac{d(d-1)}{\ell^2} \right) - \frac{1}{4g^2} F_{\mu\nu} F^{\mu\nu}$$

- This theory has black brane solutions

$$ds^2 = -U(r)dt^2 + \frac{dr^2}{U(r)} + \frac{r^2}{\ell^2} \left( \sum_{i=1}^{d-1} dx_i^2 \right)$$

$$U(r) \equiv \frac{r^2}{\ell^2} - \frac{m}{r^{d-2}} + \frac{q^2}{r^{2d-4}} \quad A = \mu - \frac{1}{c} \frac{\sqrt{2}g}{\kappa_{d+1}} \frac{q}{r^{d-2}} dt$$

- There is a horizon at a finite radius

$$\frac{d}{d-2} r_+^{2d-2} \geq q^2 \ell^2$$

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- We will play a game with **two players**:
    - **Thermodynamics**
    - Holographic **entanglement entropy**
  - We will look at these two things on states described by the **extremal black brane** mentioned above.
  - Always work on CFT on  $R^d / \text{AdS}_{d+1}$  in the Poincaré patch.
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- Holography 101: In the large  $N$  limit, the bulk EFT tells us about the thermodynamics of the dual CFT!

$$Z_{CFT} = \sum_{\text{Saddles}} e^{-S}$$

- Black brane has temp.  $\beta$  and voltage  $\mu$ . It controls thermo of CFT in  $R^d$  with temp.  $\beta$  and chemical pot  $\mu$  [Chamblin-Empanan-Johnson-Myers '99]

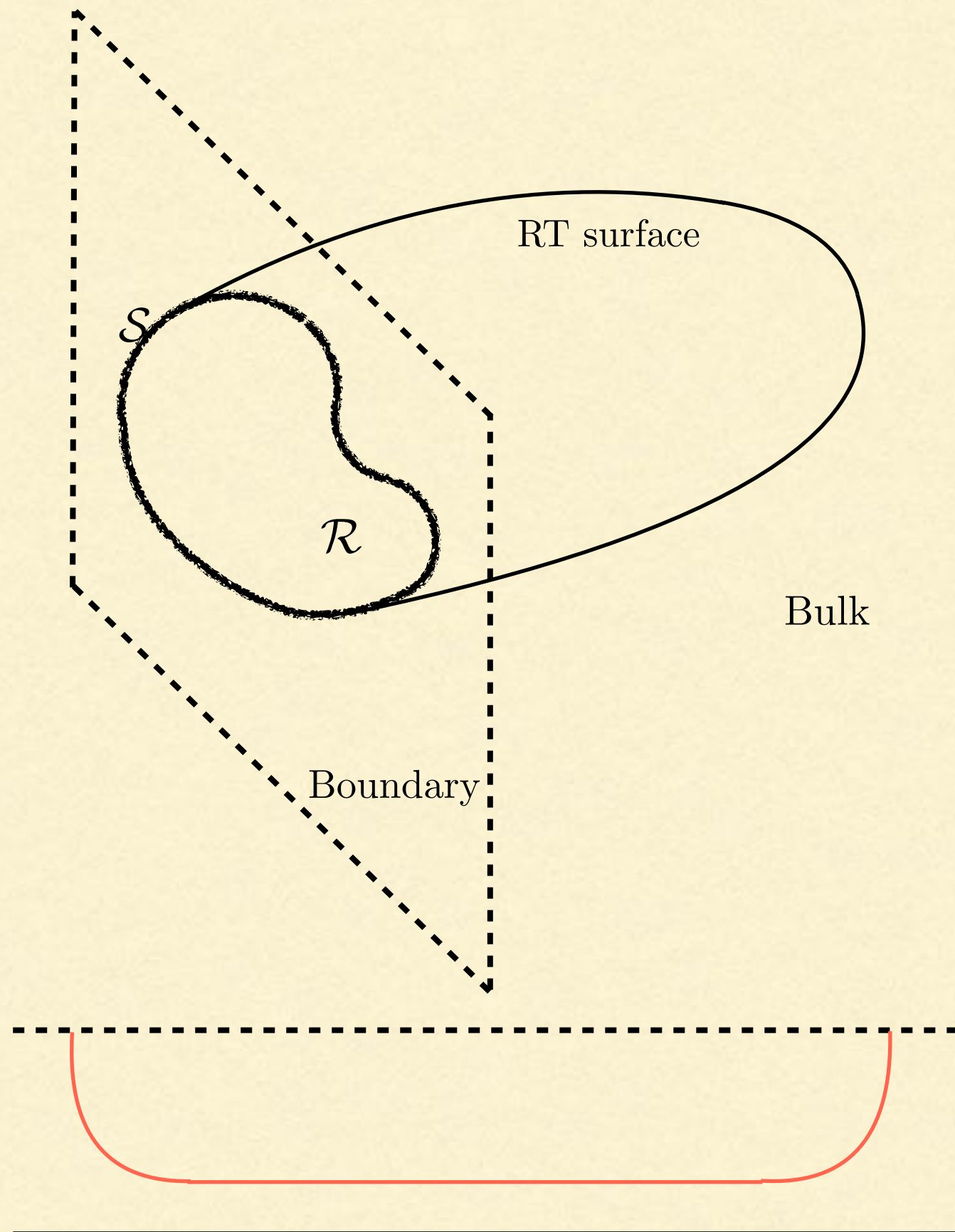
$$H \rightarrow H - \mu Q, \quad Q = \int J^0 d^{d-1}x$$

- At  $\beta \rightarrow \infty$ , ground state of CFT becomes dual to the extremal black brane. This has **finite** entropy density and linear heat capacity at low  $T$ .

$$S = \frac{A}{4G} = \frac{2\pi r_+^{d-1}}{\kappa_{d+1}} \quad \frac{d\langle H - \mu Q \rangle_\beta}{d\beta} \propto T$$

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- Ryu-Takayanagi '06 told us what the large  $N$  limit of EE (on semiclassical states) is:

$$S_{EE} = \frac{A_{\text{Extremal}}}{4G}$$

- We like to study scaling of EE with size of  $S$ . Several possibilities:
  - Volume law
  - Area law (typical of ground states)

- The black brane linear heat capacity leads to an upper **bound** on the entanglement entropy of ground states via the following Theorem: [Brandao-Cramer '14]

Consider a translationally invariant local Hamiltonian on a  $(d-1)$ —dimensional lattice. Let  $S$  be a cubic region of edge length  $L$ , and let  $|\psi\rangle$  be a ground state. If the heat capacity is proportional to  $T^\gamma$  for low  $T$ , then the entanglement entropy of  $S$  satisfies a sub-volume law

$$S_{|\Psi\rangle}(L) \leq C_0 L^{d-2+\frac{1}{\gamma+1}}$$

for  $L$  sufficiently large and some constant  $C_0$  independent of  $L$

- In our case, this means that the ent. entropy of translationally inv. ground states scales at most as **subvolume**

$$S(L) \propto L^{d-3/2} \quad (\text{Area} \propto L^{d-2})$$

$$(\text{Volume} \propto L^{d-1})$$

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Let's check the hypotheses:

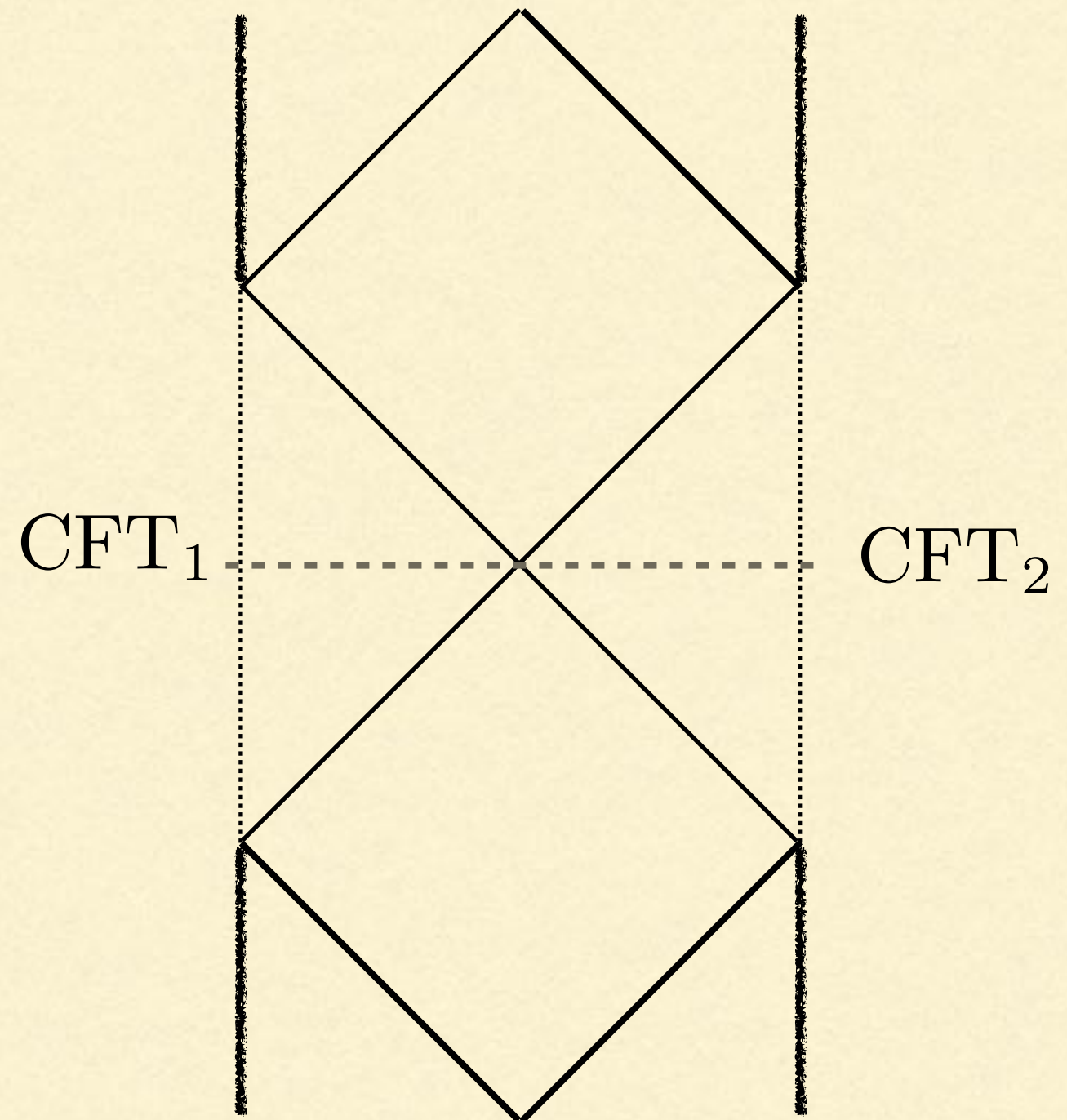
- Lattice? We regularize the CFT to compute EE.
- Local Hamiltonian (nearest-neighbour)? Ensured by locality
- Linear heat capacity,  $\gamma = 1$ .



- We still need a **translationally invariant ground state**.
- We take the **thermofield double state** at zero temp.

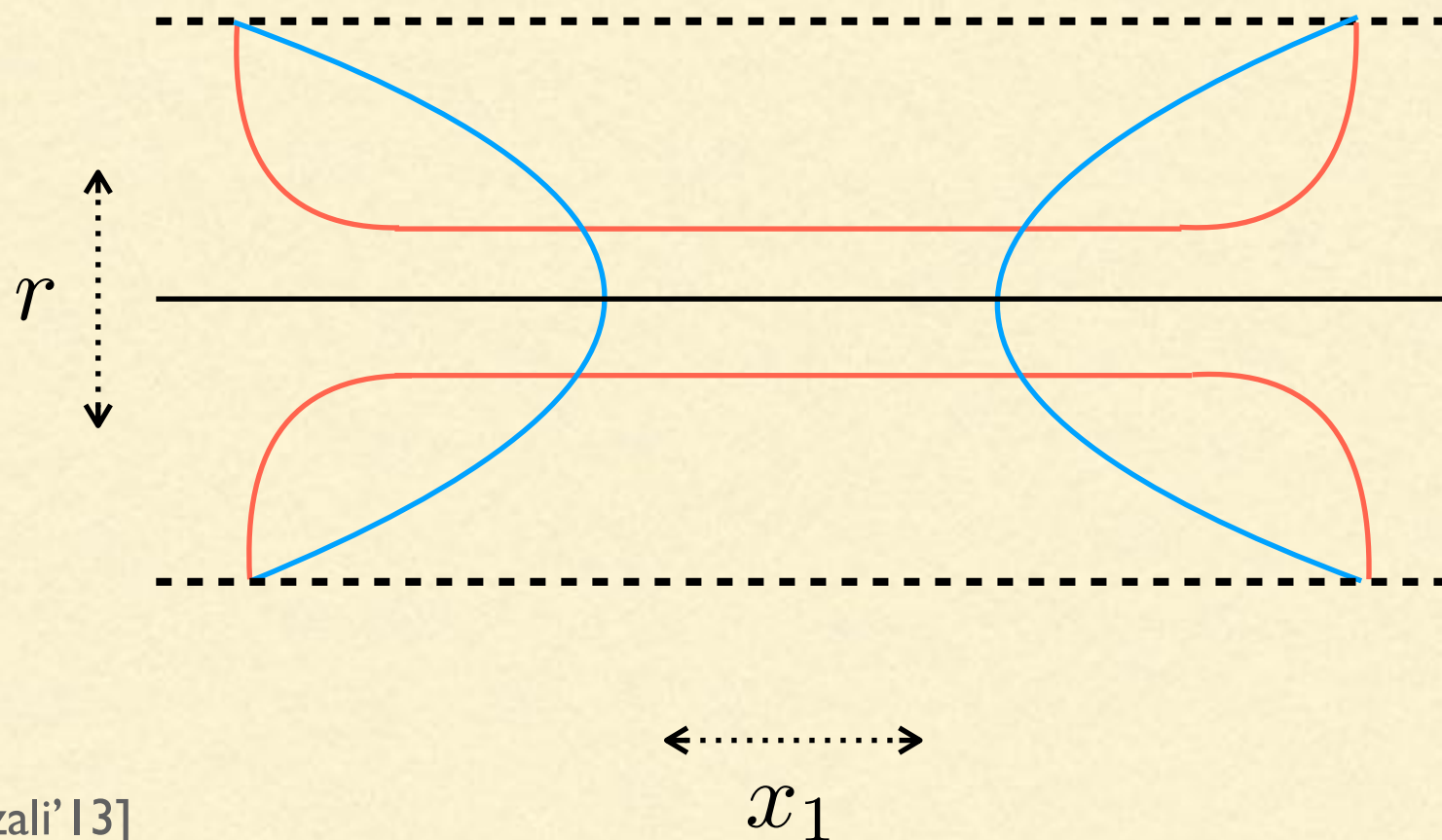
$$|TFD\rangle = \frac{1}{Z} \sum_n e^{-\frac{\beta}{2}(E_n - \mu Q_n)} |n\rangle_L |n\rangle_R$$

- Ground state ✓
- Translationally invariant ✓
- It has a **semiclassical bulk description**: The two-sided black brane [Maldacena '02...]



- There are two competing RT surfaces. We're supposed to take the one with lower  $A$ .

- Blue: Area law
- Red: Volume law



- [Hartman-Maldacena '13,  
Andrade-Fischetti-Marolf-Ross-Rozali'13]

- In the  $T=0$  TFD, the blue surface has **divergent  $A$** .

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- We get a **contradiction**: Volume law when the theorem forbids it. One of our hypothesis must be **wrong**:
    - Validity of EFT/large N framework ✓
    - Don't like  $T=0$ ? Can repeat analysis in the  $T \rightarrow 0$  limit ✓
    - Einstein-Maxwell / Stable black brane ✗
  - More generally:

**Electrically charged extremal black branes must\* be unstable.**

\* Should have polynomial low-T heat capacity. Unstable means there's another solution w. lower free energy.

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- Simplest modification of the EFT one can make to fix the above picture: Introduce **charged fields** with high-enough  $q/m$

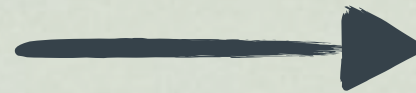
$$m^2 - \frac{g}{\kappa_{d+1}} q_\phi^2 \leq \frac{d(d-1)}{4\ell^2}$$

The (mild) **Weak Gravity Conjecture**.

- Comes with AdS correction; One of the AdS WGC variants advocated in [Nakayama-Nomura '15]
  - BH is unstable; decays to horizonless “electron star”.
  - Only possibility I know of, in the EFT regime. **Higher derivative terms** suppressed by some UV scale should not be very important for huge black branes (with very high  $\mu$ ).
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- Also, connection to another quantum information theorem [Brandao-Horodecki '13]

Exponential decay  
of **all** correlators



Area law for  
entanglement

- Can check exp. decay of correlators in the worldline approximation [Andrade-Fischetti-Marolf-Ross-Rozali'13]
  - Connection to **ER=EPR**? The extremal thermofield double has a wormhole, but its length is infinite.
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# CONCLUSIONS

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- Quantum Information Theorem = No stable electric black branes.
- EFT's that predict these are sick & in the Swampland, leading to the (mild) WGC in AdS.
- Take home message:

## **Entanglement as a Swampland constraint machine**

- Check other conjectures: SDC<sub>[Ooguri-Vafa '06,...]</sub>, Scalar-WGC<sub>[Palti '17]</sub>, Convex-Hull<sub>[Cheung-Remmen '14]</sub>... or prove new ones!
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# Thank you!

¡Gracias!

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## FAQ

- **Q:** What about BPS black holes? They have an horizon & are stable.

**A:** There are **no** examples of asymptotically AdS (only) electrically charged black branes which are BPS (that I know of). The extremality bound is **strictly above** the BPS bound.

- **Q:** Ok, so what about branes with magnetic charges?

**A:** Turning on a magnetic charge = turning on a background magnetic field in the CFT. Breaks translational invariance of hamiltonian!

- **Q:** Bah, who cares. AdS is SUSY anyway. WGC=BPS bound.

**A:** There can be  $U(1)$ 's which are not central charges (flavor  $U(1)$ 's). No BPS bound for these, yet the argument still applies! Also, saturating the BPS bound is not completely trivial...

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