How Good is the Standard Model?

Andrea Wulzer



Based on:

D'Agnolo, AW, 2018
D'Agnolo, Grosso, Pierini, AW, Zanetti, 2019
D'Agnolo, Grosso, Pierini, AW, Zanetti, 2021
Letizia, Grosso, AW, et. al., 2022

Statisticians formulate an interesting problem: **g.o.f.***

Be \mathcal{D} a set of data, and R a stat. hyp. for their distribution Does R provide the **right description** of \mathcal{D} ?

^{*}often question emerges after optimising distribution free parameters on the data, as a way to assess fit quality. But the problem is more general

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Answering is more **easy** the more **restrictive** assumptions we make on how the true distribution, if not R, can look like But, more **partial** as well.

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Simple vs Simple hypothesis test



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- Optimal approach provided by Neyman–Pearson Lemma
- Optimal answer to very specific question: **test has no or very limited power if truth** ≠ **H**₁

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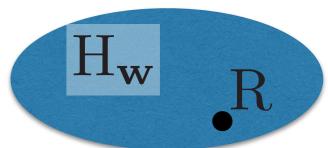
Simple vs Simple hypothesis test





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Simple vs Composite test



- No Optimal solution. But, Likelihood Ratio is Good solution
- Answers a more general question:
 some power if truth is in H_w.
 Generically, larger H_w = less power

The LHC g.o.f. challenge

By analysing the LHC data, we would like to find evidence of failure of the SM theory, suggesting need of BSM.

This is a tremendously hard gof problem!

BSM is tiny departure from SM, or large in tiny prob. region Affecting few (unknown) observables over ∞ many we can measure

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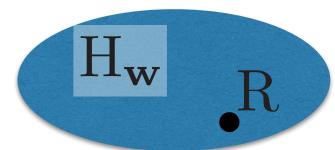
Model-dependent BSM searches

specific BSM model



- Optimise sensitivity to one
- Fail to discover other models. What if the right theoretical model is not yet formulated?

Model-independent searches

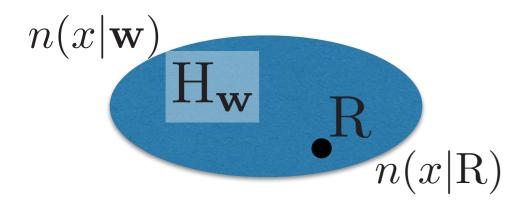


- Could reveal **truly unexpected** new physical laws.
- No hopes to find Optimal strategy. For a Good strategy, we need a **good choice of Hw.**

Data: $\mathcal{D} = \{x_i\}, i = 1, \dots, \mathcal{N}_{\mathcal{D}}$

I.i.d. measurements of, e.g., reconstructed particle momenta in a region of interest

$$n(x) = N P(x)$$
$$N = \int dx \, n(x)$$



$$n(x|\mathbf{w}) = n(x|\mathbf{R}) e^{f(x;\mathbf{w})}$$

 $f(x; \mathbf{w})$ is a **neural network**, or other flexible functional approximant with good properties in many dimensions, like **kernels**

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Strategy is to evaluate the classical Likelihood Ratio test statistic

$$t(\mathcal{D}) = 2 \log \frac{\max_{\mathbf{w}} [\mathcal{L}(\mathbf{H}_{\mathbf{w}}|\mathcal{D})]}{\mathcal{L}(\mathbf{R}|\mathcal{D})} = 2 \max_{\mathbf{w}} \left\{ \log \left[\frac{e^{-N(\mathbf{w})}}{e^{-N(\mathbf{R})}} \prod_{i=1}^{N_{\mathcal{D}}} \frac{n(x_i|\mathbf{w})}{n(x_i|\mathbf{R})} \right] \right\}$$

by supervised training Data vs Reference (background) sample.

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By using a special loss function:

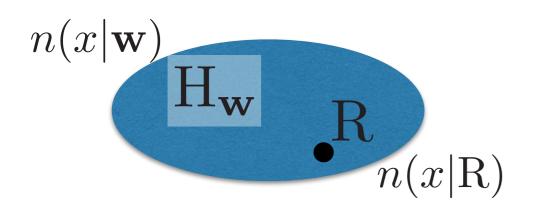
$$L[f] = \sum_{(x,y)} \left[(1-y) \frac{N(\mathbf{R})}{\mathcal{N}_{\mathcal{R}}} (e^{f(x)} - 1) - y f(x) \right] \longrightarrow t(\mathcal{D}) = -2 \min_{\{\mathbf{w}\}} L[f(\cdot, \mathbf{w})]$$

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Three-lines derivation:

$$t(\mathcal{D}) = 2 \max_{\mathbf{w}} \left\{ \log \left[\frac{e^{-N(\mathbf{w})}}{e^{-N(\mathbf{R})}} \prod_{i=1}^{\mathcal{N}_{\mathcal{D}}} \frac{n(x_i | \mathbf{w})}{n(x_i | \mathbf{R})} \right] \right\} = -2 \min_{\mathbf{w}} \left[N(\mathbf{w}) - N(\mathbf{R}) - \sum_{i=1}^{\mathcal{N}_{\mathcal{D}}} f(x_i; \mathbf{w}) \right]$$

Approximate integral as Monte Carlo sum:

$$N(\mathbf{w}) = \int dx \, n(x|\mathbf{R}) \, e^{f(x;\mathbf{w})} = \frac{N(\mathbf{R})}{\mathcal{N}_{\mathcal{R}}} \sum_{x \in \mathcal{R}} e^{f(x;\mathbf{w})}$$

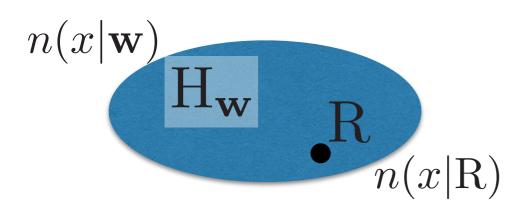
$$t(\mathcal{D}) = -2 \min_{\{\mathbf{w}\}} \left[\frac{N(\mathbf{R})}{\mathcal{N}_{\mathcal{R}}} \sum_{x \in \mathcal{R}} (e^{f(x;\mathbf{w})} - 1) - \sum_{x \in \mathcal{D}} f(x;\mathbf{w}) \right] \equiv -2 \min_{\{\mathbf{w}\}} L[f(\cdot, \mathbf{w})]$$

 $\mathcal{D} = \{x_i\}, i = 1, \dots, \mathcal{N}_{\mathcal{D}}$ Data:

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 Like saying that $n(x | \mathbf{R})$ is "known", as it is infinitely samplable. Factor few enough.

Approximate integral as Monte Carlo sum:

$$n(x|\mathbf{w}) = n(x|\mathbf{R}) e^{f(x;\mathbf{w})}$$

In order to read this as "equal", we need

$$\mathcal{N}_{\mathcal{R}} \gg N(\mathbf{R})$$

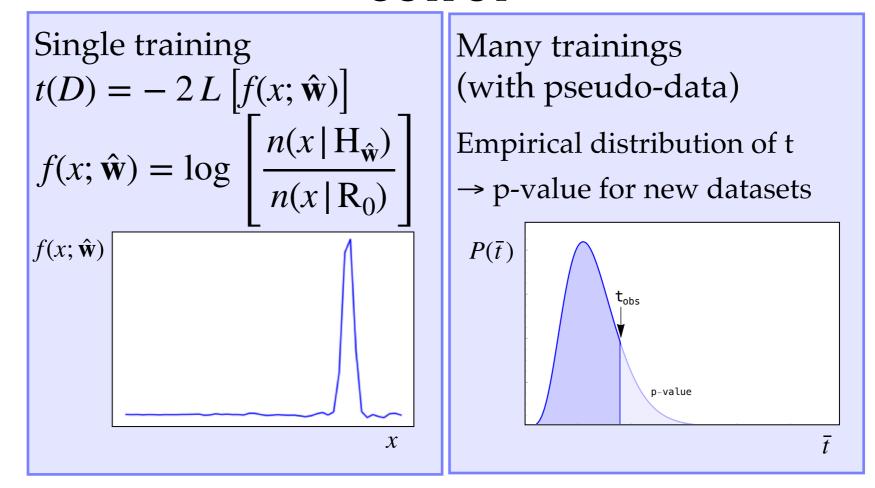
e Carlo sum:
$$N(\mathbf{w}) = \int \! dx \, n(x|\mathbf{R}) \, e^{f(x;\mathbf{w})} = \frac{N(\mathbf{R})}{\mathcal{N}_{\mathcal{R}}} \sum_{x \in \mathcal{R}} \, e^{f(x;\mathbf{w})}$$

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Reference sample (R) label=0 Data sample (D) label=1 NN training $\hat{\mathbf{w}}$

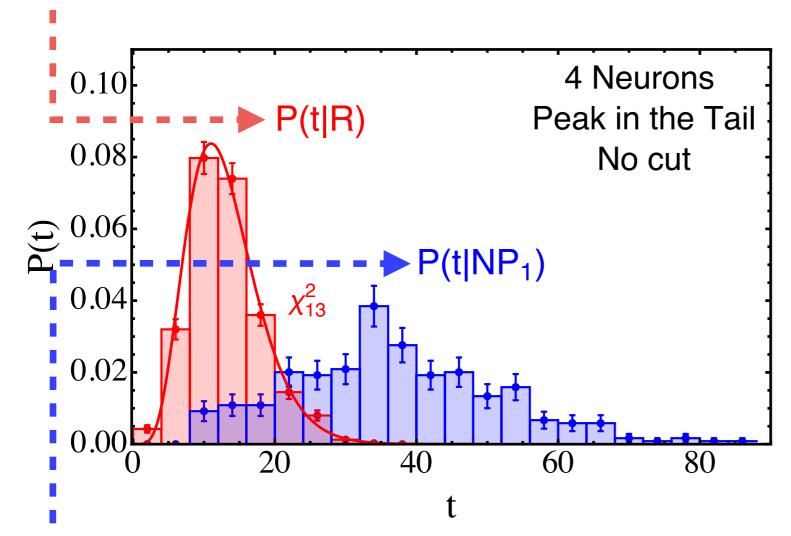
<u>Unbinned</u> training samples!

OUTPUT



(Simple 1d example with exponential Reference)

Distribution of the test statistic "t" in Reference Hypothesis

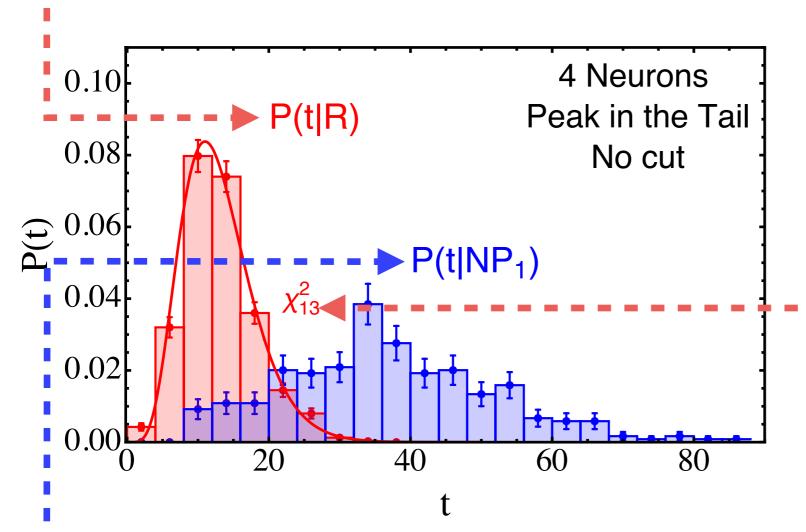


Distribution of "t" in one New Physics Model Hypothesis

$$t \rightarrow p \rightarrow Z$$
-score (we use $Z = \Phi^{-1}(1 - p)$)

(Simple 1d example with exponential Reference)

Distribution of the test statistic "t" in Reference Hypothesis



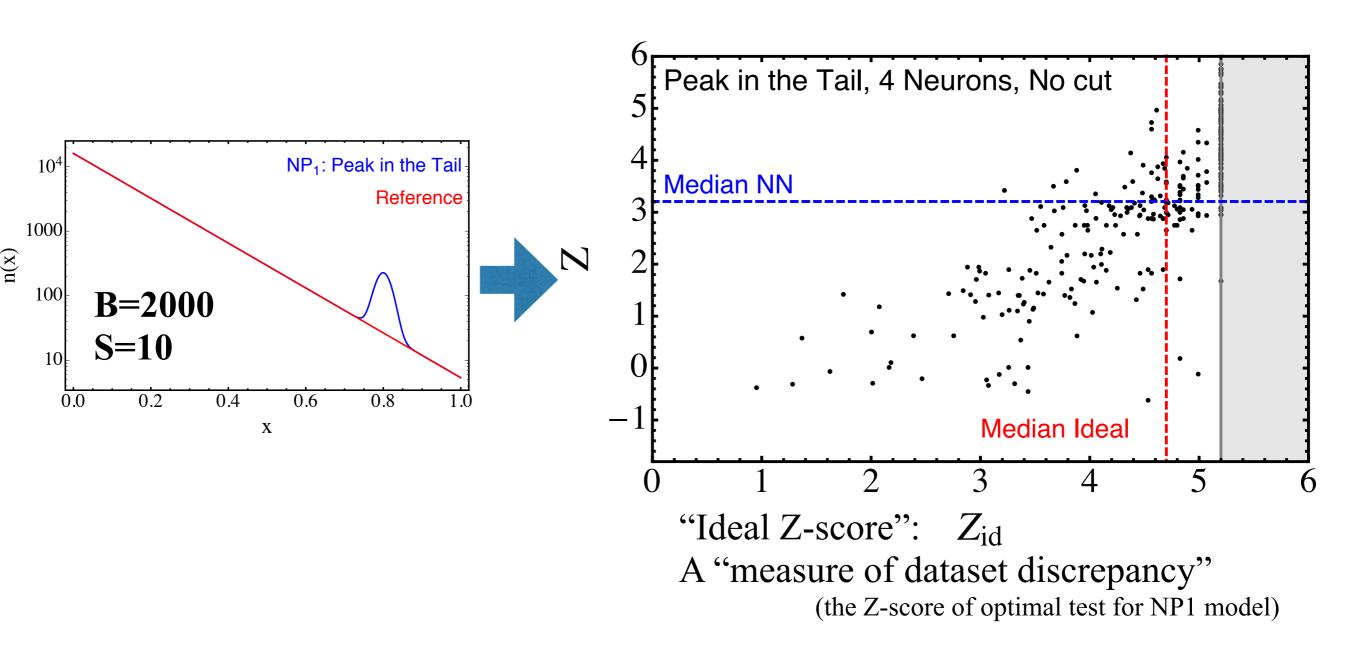
Notice agreement with Wilks' Formula:

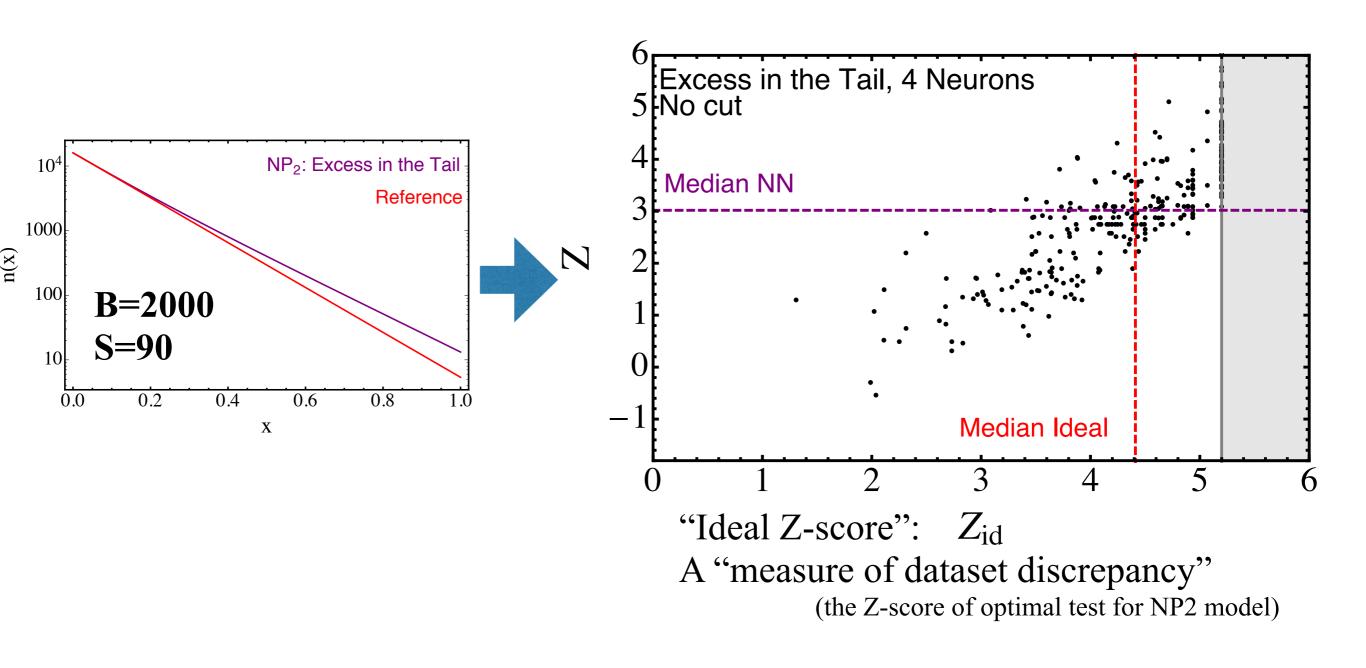
Sufficiently regularised networks found to behave as if their number of d.o.f. was equal to number of parameters.

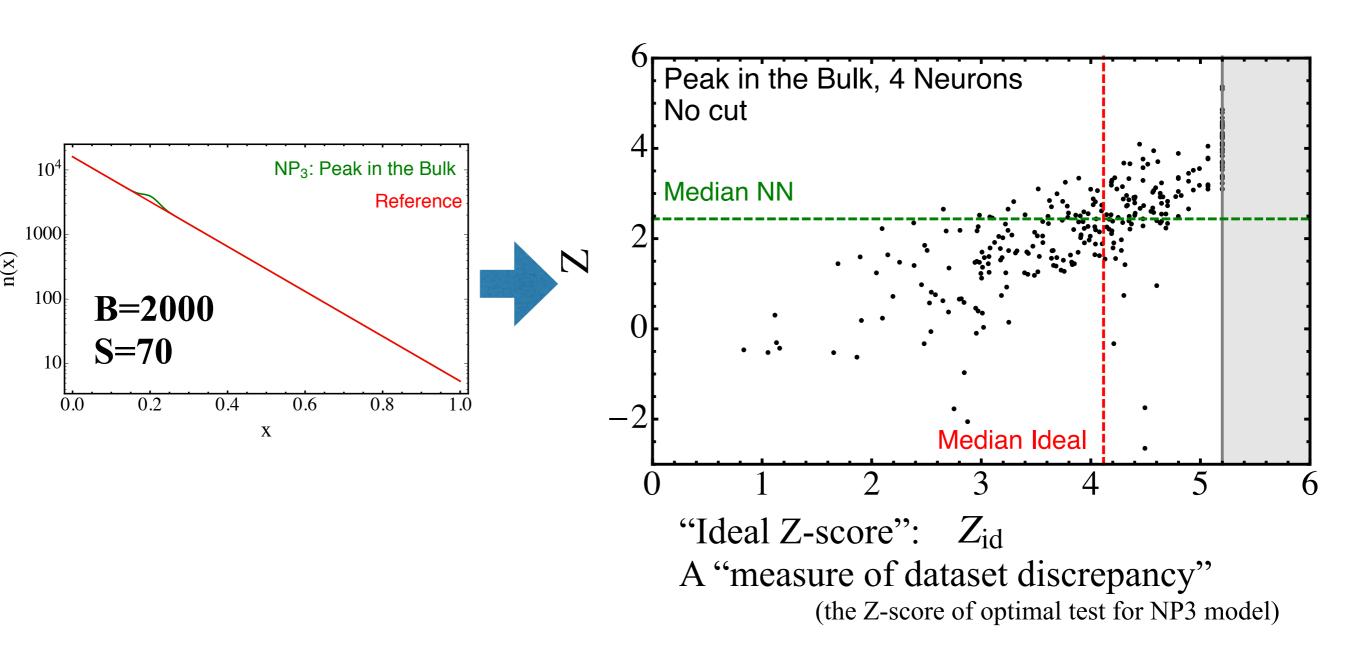
Theoretical reason mysterious

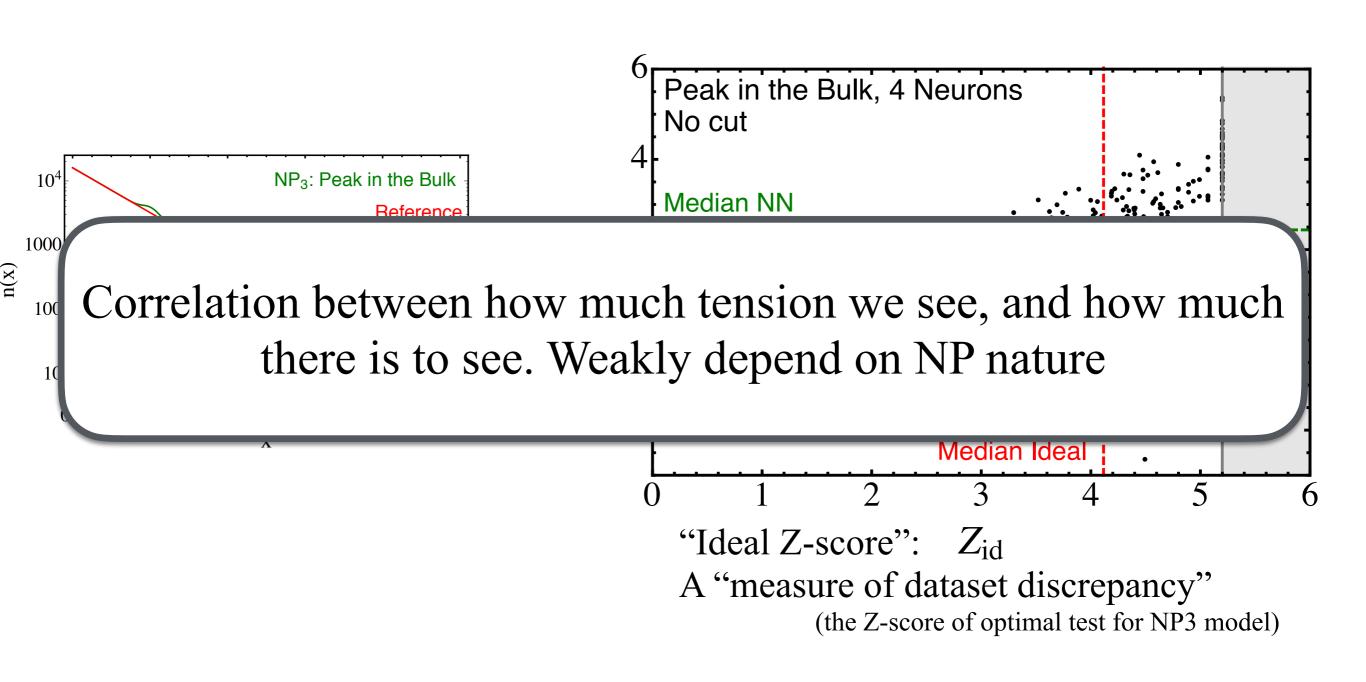
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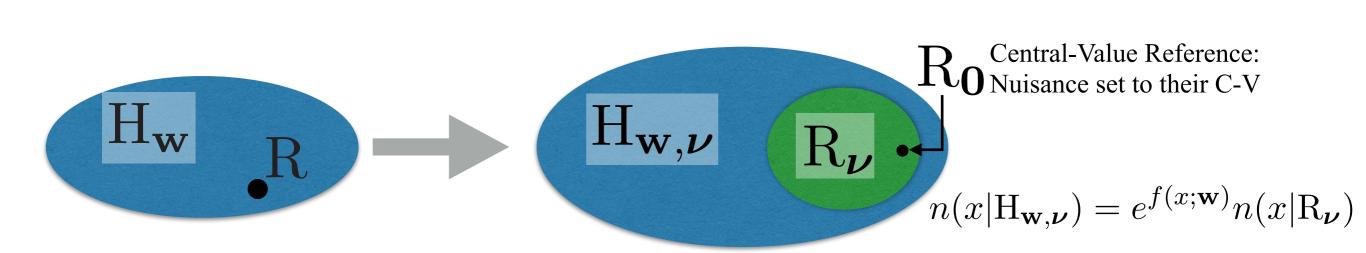


Imperfect Machine

Reference Sample is an imperfect representation of SM e.g., PDF/Lumi/Detector Modeling ...

Imperfections are Nuisance Parameters

Constrained by Auxiliary Measurements Define a composite Reference hypothesis



$$t(\mathcal{D}, \mathcal{A}) = 2 \log \frac{\max_{\mathbf{w}, \boldsymbol{\nu}} \left[\mathcal{L}(\mathbf{H}_{\mathbf{w}, \boldsymbol{\nu}} | \mathcal{D}) \cdot \mathcal{L}(\boldsymbol{\nu} | \mathcal{A}) \right]}{\max_{\boldsymbol{\nu}} \left[\mathcal{L}(\mathbf{R}_{\boldsymbol{\nu}} | \mathcal{D}) \cdot \mathcal{L}(\boldsymbol{\nu} | \mathcal{A}) \right]}$$

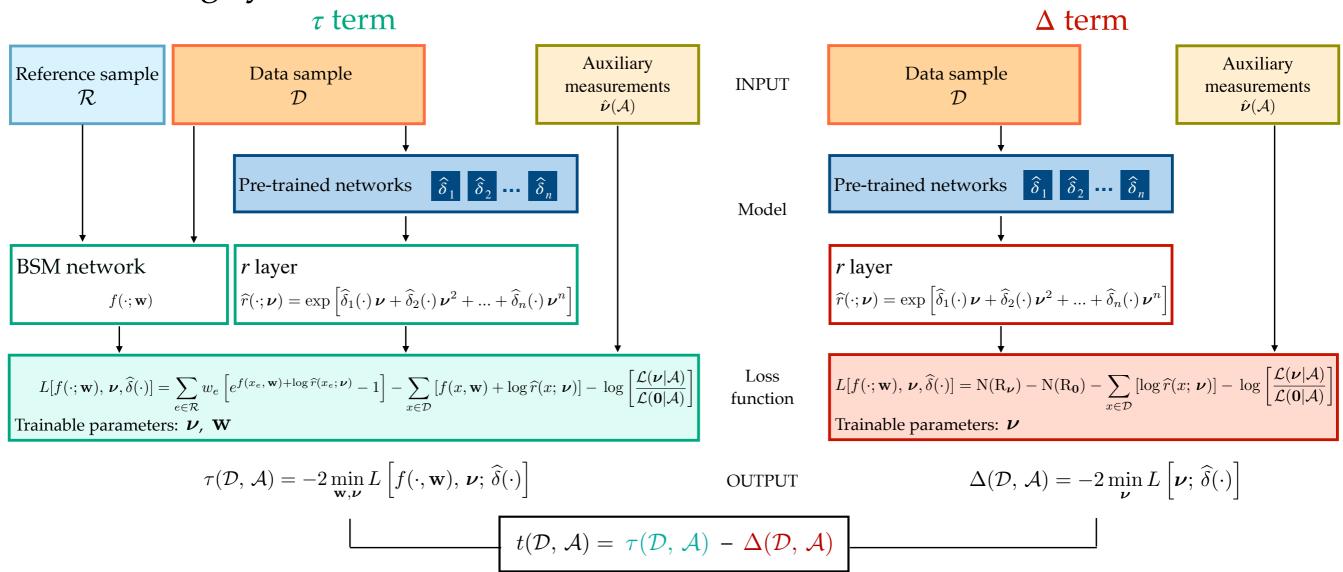
Strategy conceptually unchanged.
$$t(\mathcal{D}, \mathcal{A}) = 2 \log \frac{\max\limits_{\mathbf{w}, \boldsymbol{\nu}} \left[\mathcal{L}(\mathbf{H}_{\mathbf{w}, \boldsymbol{\nu}} | \mathcal{D}) \cdot \mathcal{L}(\boldsymbol{\nu} | \mathcal{A}) \right]}{\max\limits_{\boldsymbol{\nu}} \left[\mathcal{L}(\mathbf{R}_{\boldsymbol{\nu}} | \mathcal{D}) \cdot \mathcal{L}(\boldsymbol{\nu} | \mathcal{A}) \right]}$$
$$= 2 \max\limits_{\mathbf{w}, \boldsymbol{\nu}} \log \left[\frac{\mathcal{L}(\mathbf{H}_{\mathbf{w}, \boldsymbol{\nu}} | \mathcal{D})}{\mathcal{L}(\mathbf{R}_{\mathbf{0}} | \mathcal{D})} \cdot \frac{\mathcal{L}(\boldsymbol{\nu} | \mathcal{A})}{\mathcal{L}(\mathbf{0} | \mathcal{A})} \right] - 2 \max\limits_{\boldsymbol{\nu}} \log \left[\frac{\mathcal{L}(\mathbf{R}_{\boldsymbol{\nu}} | \mathcal{D})}{\mathcal{L}(\mathbf{R}_{\mathbf{0}} | \mathcal{D})} \cdot \frac{\mathcal{L}(\boldsymbol{\nu} | \mathcal{A})}{\mathcal{L}(\mathbf{0} | \mathcal{A})} \right] = \tau(\mathcal{D}, \mathcal{A}) - \Delta(\mathcal{D}, \mathcal{A})$$

Implementation slightly more complex

Imperfect Machine

New Physics Learning Machine (NPLM)

Including systematic uncertainties

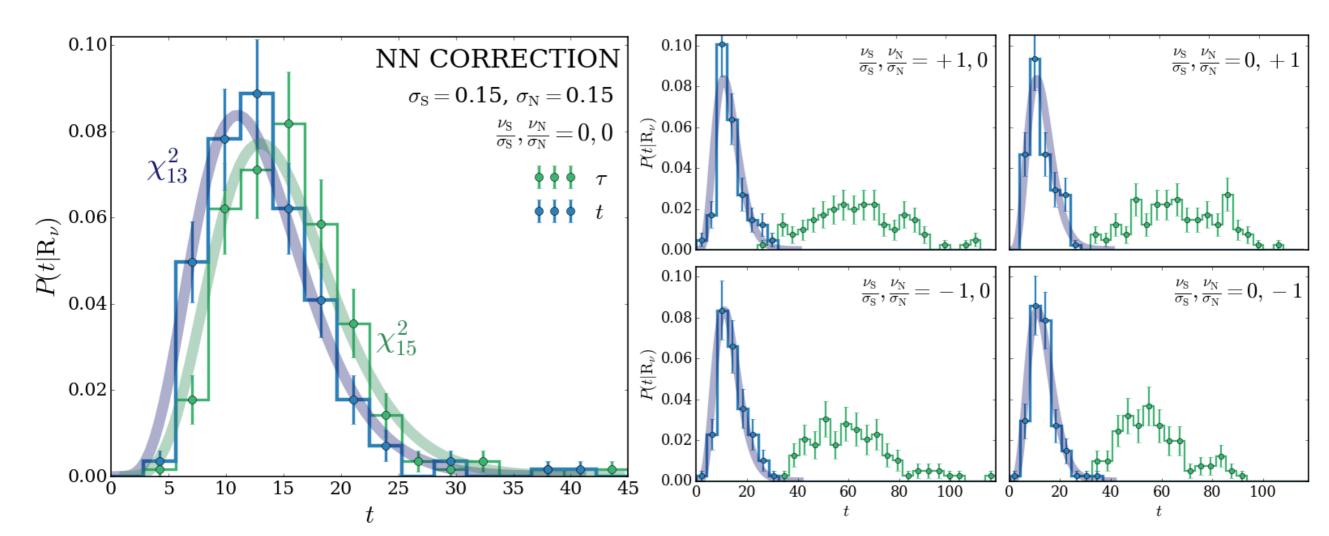


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An Imperfect Machine at Work

(Simple 1d example with exponential Reference)

Tau distribution distorted by non-central value nuisance if not corrected, produces false positives



t = Tau-Delta independent of true nuisance value this is essential for a feasible test

Towards LHC

Our proposed strategy is fully defined, including:

- Hyperparameters and regularisation selection
- Systematic approach to Reference mis-modelling

Validated on problems of realistic scale of complexity:

- 2-body final state with uncertainties (5D)
- 11+MET "SUSY" (8D)
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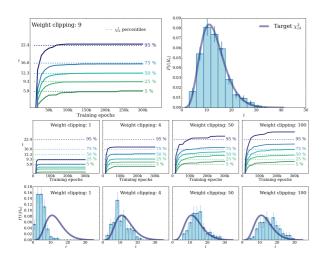
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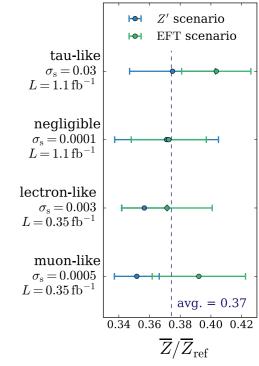
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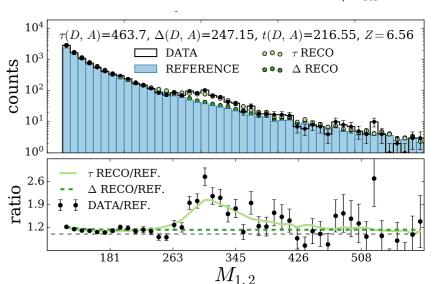
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Results in summary:

- model-selection strategy converges
- sensitivity to resonant or non-resonant NP
- "uniform" response to NP of different nature
- trained network reconstruct NP







Outlook

Next step is implementation with true LHC data.

Open theoretical questions

- Why exactly we get chi-squared distributed "t"?
- Regularisation selects space of alternatives, where we are looking for NP A principled approach to regularisation and "reasonable" alternatives?

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- Data Validation/DQM
- Other GoF problems

First Real-Life Application?

[Grosso, Lai, Letizia, Pazzini, Rando, Wulzer, Zanetti, to appear]

nD DQM

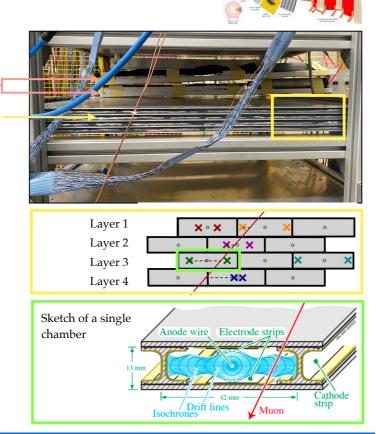
Online monitoring of a DT chamber:

Setup (Legnaro INFN national laboratory):

- 2 scintillators as signal trigger
- 1 drift tube chamber: 4 layers 16 wires each (16x4=64 wires)
- Source of signals: cosmic muons (triggered rate ~3 MHz)
- **Event**: muon track reconstructed interpolating 3/4 hits (one per layer)

Observables (6D problem):

- 4 drift times [$t_{\text{drift}, 1}$, $t_{\text{drift}, 2}$, $t_{\text{drift}, 3}$, $t_{\text{drift}, 4}$]: time for the ionised electrons to reach the wire from the interaction point ($v_{\text{drift}} = \text{cm/s}$).
- θ : reconstructed track angle
- N_{hits}: average number of hits per time window ("orbit")



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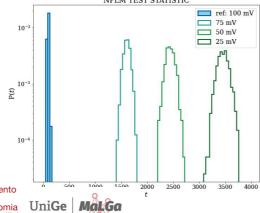


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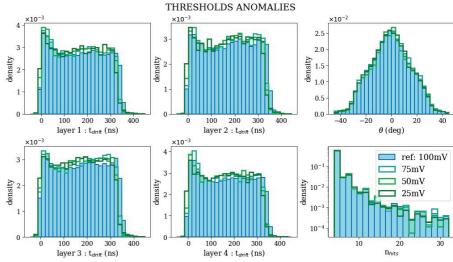
Online monitoring of a DT chamber:

- Reference sample: long run in optimal conditions
- Anomalous samples: short runs acquired in presence of a controlled anomaly in the value of the threshold tension of the DT chamber
- Result of the test statistics
 Complete separation of the distributions!



NPLM with Falkon $M = 50, \sigma = 4.84, \lambda = 10^{-7}$ N(D) = 5000 $N_{\rm ref} = 200\,000$

Execution time: $\sim 1.5 \, \text{s}$



Distribution of the observables at different values of the threshold tension

→ more about this in Marco's talk tomorrow!

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But, maybe, New Physics as well!!

Thank You

Backup

The LHC g.o.f. challenge

From a theorists' perspective:

Non-discovering **model-dependent** searches can be turned into **exclusions** of the targeted BSM. They are still informative as tell us what has **not been** discovered.

Notice however that they would **not tell us what has been discovered** any better than model-independent search, in general. Jet plus MET could have been anything.

How probable that reality is so much different from theory that we cannot envisage it before experiments? This would be great! (... right?)

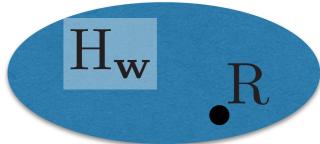
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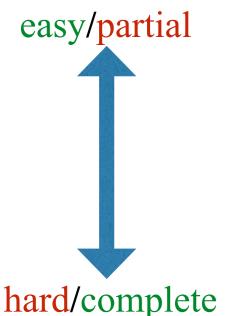
The major concern of any scientist:

Am I doing everything right?

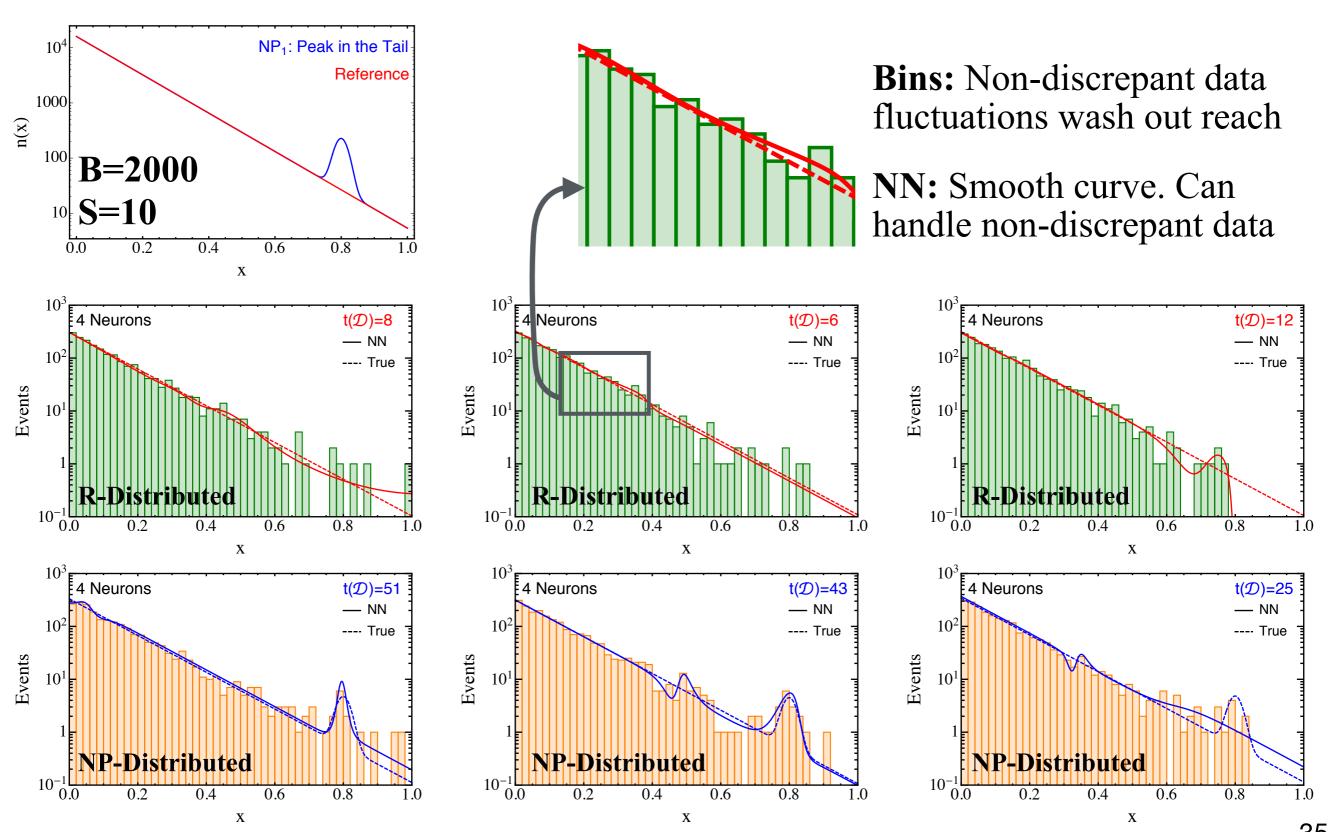
Being unable to answer, we turn to an easier question:

What could be wrong? and we check that

Cross-checks are more easy the more specifically we characterise the possible failure. But also less powerful

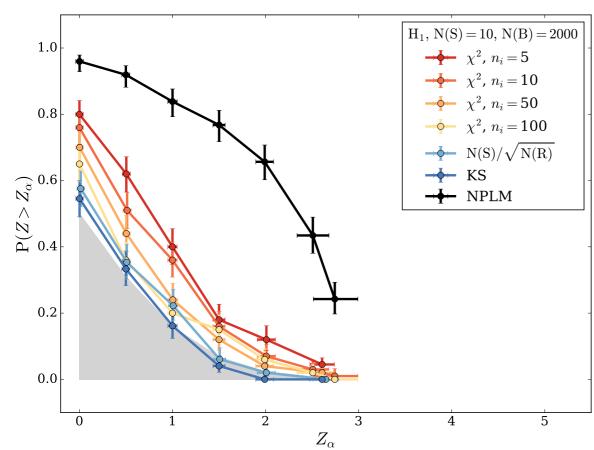


- did I turn QED showering on, in my PYTHIA simulation?
- is the power plug of my detector connected?
- •
- is my detector system working "normally"?
- •
- is my state-of-the-art knowledge of fundamental interactions (the **SM**) **correct**, or it **fails** to describe the LHC data?



(Simple 1d example with exponential Reference)

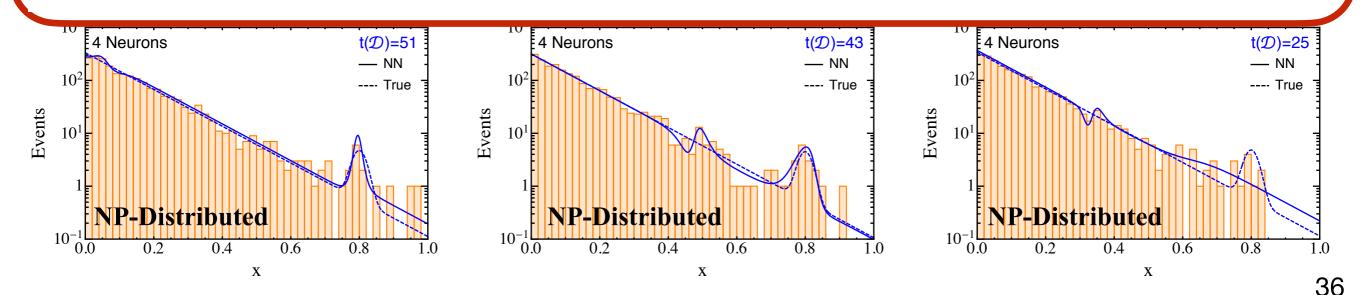
Probability to find evidence of R being wrong at some level of confidence.



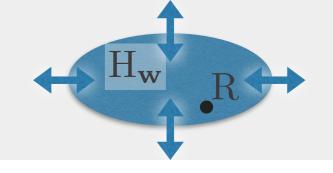
We are better than binned χ^2 because our model has less parameters but same effective expressive power.

Same reason why bins are outdated as statistical models.

Gap to bins grows (exponentially) with (the curse of) dimensionality.

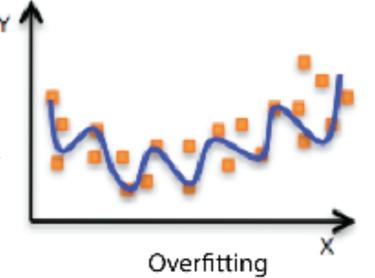


Model Selection



Which hypotheses (distributions) our (statistical) model contains?

- •Not "all of them", otherwise it would fail (overfitting)
- •It should contain approximations of all the reasonable ones
- •No Statistical Learning notion of model capacity seems reasonable physics measure of volume or boundaries of Hw
- •Minimal allowed variation scale would sound reasonable, but no theory developed



Waiting for principled approach, solution is χ^2 -compatibility:

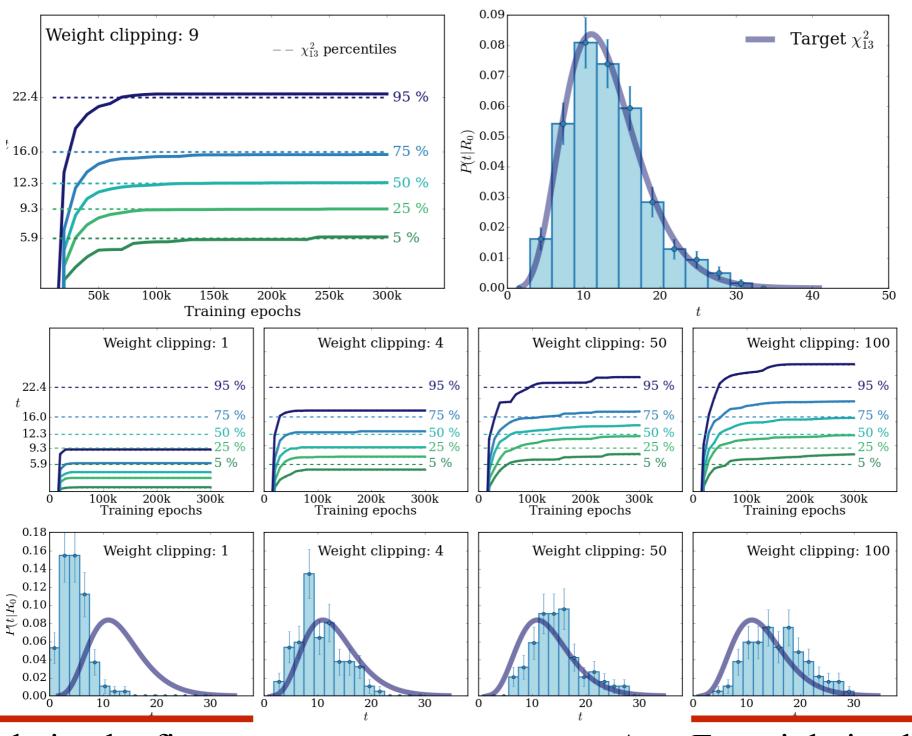
•Naive Wilks Theorem application:

P(t|R) is χ^2 , with as many d.o.f. as fit parameters (for us, num. of NN par.s) Provided statistics is large relative to fitted model "complexity" ... or, which is the same ...

Provided model is "simple enough", for given data statistics

- •Asy. For. violation = sensitivity to low-statistics portion of dataset = overfitting
- •Regularisation by Weight Clipping, that forbids sharp variations
- $\bullet NN$ with too many parameters cannot be made χ^2 -compatible. Take largest allowed

Weight Clipping Selection



Asy. For. violation by fit parameters boundary

Asy. For. violation by sensitivity to sparse data points