## Disconnected contributions from GPU's

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## Outline

- Brief introduction to disconnected contributions
- Stochastic procedures
- Truncated Solver Method (TSM)
- The one-end trick and other improvements
- GPU performance and scaling
- The summation method
- Results
- Conclusions and future plans


## Motivation



- Determination of flavour singlet quantities $\eta$ mass, nucleon form factors...
- Computations of non-perturbative nature
- We must rely on lattice methods

$$
L(x)=\operatorname{Tr}[\Gamma G(x ; x)]
$$

## Disconnected contributions



- For the evaluation of disconnected diagrams we need to compute all-to-all propagators
- Very expensive from the computational point of view
- Neglected in most hadron structure studies
$L(x)=\operatorname{Tr}[\Gamma G(x ; x)]$


## Stochastic procedures

- Exact computation of the all-to-all unfeasible nowadays
- We can use stochastic techniques
- Invert a random set of sources $\left|\eta_{j}\right\rangle$ that form a basis up to stochastic errors
- Properties $\left\{\begin{aligned} \frac{1}{N} \sum_{j=1}^{N}\left|\eta_{j}\right\rangle & =O\left(\frac{1}{\sqrt{N}}\right) \\ \frac{1}{N} \sum_{j=1}^{N}\left|\eta_{j}\right\rangle\left\langle\eta_{j}\right| & =I+O\left(\frac{1}{\sqrt{N}}\right)\end{aligned}\right.$
- In this work we use $\mathbf{Z}_{2}$ and $\mathbf{Z}_{4}$ noise sources
" So we get an unbiased estimation of the all-to-all propagator

$$
M\left|s_{j}\right\rangle=\left|\eta_{j}\right\rangle \longrightarrow M_{E}^{-1}:=\frac{1}{N} \sum_{j=1}^{N}\left|s_{j}\right\rangle\left\langle\eta_{j}\right| \approx M^{-1}
$$

## Stochastic procedures

- Error decresases as $1 / \sqrt{N}$, we usually need a large number of stochastic sources $N$
- Each stochastic source requires an inversion of the fermionic matrix

$$
M\left|s_{j}\right\rangle=\left|\eta_{j}\right\rangle
$$

- To reduce the stochastic noise, we want to increase $N$ at a reduced cost
- The Truncated Solver Method (TSM) allows us to do this


## The Truncated Solver Method

- First published in PoSLaT2007, 141 (G. Bali, S. Collins and A. Schäffer)
- Instead of solving $M\left|s_{j}\right\rangle=\left|\eta_{j}\right\rangle$ exactly, we aim at a low precision estimation
- Cut the inverter (CG) at a certain number of iterations OR at a given precision $\rho^{2} \sim 10^{-4}$
- This is cheap (fast inversions) but inaccurate $\longrightarrow$ We introduce a bias
- We compute the correction of the bias stochastically

$$
M_{E}^{-1}:=\frac{1}{N_{H P}} \sum_{j=1}^{N_{H P}}\left(\left|s_{j}\right\rangle\left\langle\left.\eta_{j}\right|_{H P}-\mid s_{j}\right\rangle\left\langle\left.\eta_{j}\right|_{L P}\right)+\frac{1}{N_{L P}} \sum_{j=N_{H P}+1}^{N_{H P}+N_{L P}}\left|s_{j}\right\rangle\left\langle\left.\eta_{j}\right|_{L P}\right.\right.
$$

## The Truncated Solver Method

- If the convergence in the inversions is fast, we can get away with a low $N_{H P}$
- Error should decrease essentially as $1 / \sqrt{N_{L P}}$
- Since the LP sources don't require an accurate inversion, we can take advantage of the half precision algorithms for GPU's

Mixed double/single $\sim 100$ GFlops Mixed double/half $\sim 170$ GFlops

- Two parameters to tune: precision of LP and $N_{H P} / N_{L P}$ ratio
- Fine-tuning depends on the loop to be computed


## Determination of the TSM parameters



- Data for $\bar{\psi} \gamma_{3} D_{3} \psi$, a piece of $\langle x\rangle_{u+d}$
- After 300 LP, hard to improve
- $24 \mathrm{HP} / 300 \mathrm{LP} \approx$ 48HP are enough for all loops with local and one-derivative insertion


## The one-end trick

- For twisted mass fermions, the difference of propagators in the twisted basis is

$$
M_{u}^{-1}-M_{d}^{-1}=-2 i \mu M_{d}^{-1} \gamma_{5} M_{u}^{-1}
$$

- So, instead of computing the I.h.s. we do the r.h.s. as

$$
\sum X\left(M_{u}^{-1}-M_{d}^{-1}\right)=-2 i \mu \sum_{r}\left\langle s^{\dagger} X \gamma_{5} s\right\rangle_{r}
$$

- In principle, the trick only works for the difference, but an alternative version can be developed for the sum

$$
\sum X\left(M_{u}^{-1}+M_{d}^{-1}\right)=2 \sum_{r}\left\langle s^{\dagger} \gamma_{5} X \gamma_{5} D_{W} s\right\rangle_{r}
$$

## The one-end trick

- Below a list of bilinears with its appropiate version of the one-end trick

| Bilinear | Twisted Basis | Standard vv | Generalized vv |
| :---: | :---: | :---: | :---: |
| $\bar{\psi} \psi$ | $i \bar{\psi} \gamma_{5} \tau_{3} \psi$ | $\checkmark$ | $\mathbf{x}$ |
| $\bar{\psi} \tau_{3} \psi$ | $i \bar{\psi} \gamma_{5} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $i \bar{\psi} \gamma_{5} \psi$ | $-\bar{\psi} \tau_{3} \psi$ | $\checkmark$ | $\mathbf{x}$ |
| $i \bar{\psi} \gamma_{5} \tau_{3} \psi$ | $-\bar{\psi} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $\bar{\psi} \gamma_{\mu} \psi$ | $\bar{\psi} \gamma_{\mu} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $\bar{\psi} \gamma_{5} \gamma_{\mu} \psi$ | $\bar{\psi} \gamma_{5} \gamma_{\mu} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $\bar{\psi} \gamma_{\mu} D_{\nu} \psi$ | $\bar{\psi} \gamma_{\mu} D_{\nu} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $\bar{\psi} \gamma_{\mu} D_{\nu} \tau_{3} \psi$ | $\bar{\psi} \gamma_{\mu} D_{\nu} \tau_{3} \psi$ | $\checkmark$ | $\mathbf{x}$ |
| $\bar{\psi} \gamma_{5} \gamma_{\mu} D_{\nu} \psi$ | $\bar{\psi} \gamma_{5} \gamma_{\mu} D_{\nu} \psi$ | $\mathbf{x}$ | $\checkmark$ |
| $\bar{\psi} \gamma_{5} \gamma_{\mu} D_{\nu} \tau_{3} \psi$ | $\bar{\psi} \gamma_{5} \gamma_{\mu} D_{\nu} \tau_{3} \psi$ | $\checkmark$ | $\mathbf{x}$ |

## GPU performance and scaling




- Strong scaling competitive up to 8 gpu
- Strongly depends on local volume


## GPU performance and scaling



- Almost perfect weak scaling up to 16 gpu


## Other improvements

- Generation of noise sources on-the-fly
- Don't store propagators/sources, saving I/O - storage (very important for LP sources)
- Implementation of contractions directly on GPU's
- We take advantage of a massively parallel architecture
- Usage of cudaFFT library to generate all momenta


## The summation method

- Alternative to the plateau method for computing ratios
- Excited states suppressed by $e^{-m t_{s}}$, instead of $e^{-m t_{i}}$
- Requires the 3pt at several $t_{s} \rightarrow$ more expensive for the connected
- L. Maiani, G. Martinelli, M. L. Paciello, B. Taglienti, Nucl. Phys. B293, 420 (1987)
- S. Güsken, arXiv:hep-lat/9906034v1
- S. Capitani, B. Knippschild, M. Della Morte, H. Wittig, PoSLaT2010 147
- The 3pt is summed over $t_{i}$ up to $t_{\text {Sink }}$

$$
R_{\text {SUM }}\left(t_{s}\right)=\sum_{t_{i}=0}^{t_{s}} R_{P L A}\left(t_{s}, t_{i}\right)
$$

## The summation method

- We usually require $t_{i} \gg 1$ and $t_{s}-t_{i} \gg 1$ in order to remove undesirable contributions from excited states

$$
R_{P L A}\left(t_{s}, t_{i}\right)=R_{G S}+O\left(e^{-K t_{i}}\right)+O\left(e^{-K^{\prime}\left(t_{s}-t_{i}\right)}\right)
$$

- However, if we sum up to $t_{s}$, the unwanted contributions form a geometric series and become

$$
R_{S U M}\left(t_{s}\right)=t_{s} R_{G S}+c\left(K, K^{\prime}\right)+O\left(e^{-K t_{s}}\right)+O\left(e^{-K^{\prime} t_{s}}\right)
$$

- This way the contribution of excited states is always supressed by an $e^{-K t_{s}}$ factor


## Results: $\sigma_{\pi N}$ disconnected




- The one-end trick for the difference works very well suppressing the noise
- Clear rise of the plateau as $t_{S}$ grows
- Disconnected piece around $\sim 10-15 \%$ of the connected piece


## Results: Strange content of the nucleon




- Again, one-end trick for the difference
- Preliminary result, we will increase statistics
- The plateau seems to perfom better in this case


## Results: $\sigma_{K \Delta}$ disconnected



- Preliminary result with small statistics, we need to investigate with more configurations if we have contamination here
(2) The Cyprus


## Results: $g_{A}$ disconnected




- One-end trick for the sum doesn't provide $\mu$ noise suppression
- Consistent with former determinations of disconnected $g_{A}$, G. Bali et al. Phys.Rev.Lett. 108 (2012), 222001
- Disconnected piece negative and $\sim 5-10 \%$ of the connected piece


## Conclusions

- GPU's suitable for computation of disconnected diagrams
- TSM highly reduces the variance while keeping the same computer cost
- The one-end trick for $u-d$ gives great results at low cost, but excludes flavour singlets
- The version for $u+d$ doesn't perform so well
- Noisy due to the presence of $D_{W}$ or the lack of the noise suppression factor $\mu$
- Our current (partial) results for $A_{20}$ and $\tilde{A}_{20}$ are inconclusive
- The trick computes all time slices in a single inversion
- We can use both the plateau and the summation method
- Plateau and summation methods give consistent results
- It seems that some observables are affected by contamination coming from higher excitations


## Future plans

- Improve the signal for all the successful observables using all the available momenta and statistics
- Compare the one-end trick for the sum with time dilution, taking into account other improvements
- Always in GPU's, to achieve maximum performance
- Focus on flavour singlets

