Heterotic Cosmic String Vacua and Generalized Half-flat Compactifications

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in collaboration with

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 - Cosmic strings
 - Spin(7) and generalized half-flat manifolds

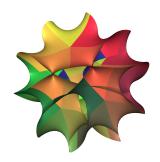
Why the heterotic string?

Some reasons to study heterotic string theory:

- String theory: consistent theory of quantum gravity!
- Comes equipped with an E₈×E₈ (or SO(32)) gauge group
 ⇒ good framework for grand unified models.
- Many candidate Standard Model compactifications known.
- Calabi–Yau compactification gives N = 1 SUSY in d = 4
 → suitable for MSSM-style models.
- Appealing mathematical framework, reasonably well-studied.

Calabi-Yau compactification

- Superstring theory is self-consistent only in 10 spacetime dimensions.
- Assume the extra 6 spatial dimensions are compactified.
- Lots of supersymmetry in d = 10
 → want to break most of it.
- Amount of broken SUSY ⇒ holonomy group of compactification manifold.



- Maximum holonomy is $SO(6) \cong SU(4) \Rightarrow \text{no SUSY preserved}$.
- Calabi–Yau manifold: SU(3) holonomy ⇒ 1/4 SUSY preserved
 - e.g. heterotic Calabi-Yau: 4 of 16 supercharges unbroken.

Moduli stabilization: Type IIB example

Compactification gives rise to $\frac{\text{moduli}}{\text{moduli}}$: flat directions in the potential \rightarrow need to stabilize.

- Example: moduli stabilization in type IIB string theory.
- Theory contains R-R 3-form flux F₃ and NS-NS 3-form flux H₃.
- Compactify such that on the manifold, F_3 and H_3 are non-zero \rightarrow flux compactification.
- In the right combination, dilaton and all complex structure moduli can be stabilized by these fluxes
- Kähler moduli remain unstabilized, can fix with eg.
 - non-perturbative effects (KKLT),
 - non-perturbative effects and perturbative α' corrections (LVS).
- All moduli stabilized!



Problems with heterotic moduli stabilization

- In heterotic string theory, only have NS-NS flux H_3 .
- Can stabilize complex structure moduli... then what?
- Dilaton can be stabilized by gaugino condensation.
- No other non-perturbative effects, no other options for flux quantisation.
- In fact, problem is even worse:

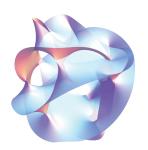
Strominger, 1986

If a heterotic compactification on a manifold Y has a maximally symmetric (e.g. Poincaré) vacuum and non-vanishing H_3 , Y is non-Calabi–Yau.

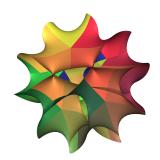
• Hence for a Calabi–Yau compactification, $H_3 = 0!$



Mirror symmetry







- A related issue is mirror symmetry.
- Type IIA compactified on $Y \leftrightarrow \text{type IIB compactified on } \widetilde{Y}$.
- Flux compactifications: R-R flux $F_3 \leftrightarrow F_0$, F_2 , F_4 , F_6 .
- Problem: no obvious mirror dual for NS-NS flux H₃!

What is an SU(3) structure manifold?

- Mirror dual: manifold with SU(3) structure, but not Calabi-Yau hep-th/0008142 (Vafa), hep-th/0211102 (Gurrieri, Louis, Micu, Waldram).
- SU(3) structure: there is a globally-defined spinor ζ that leaves 1/4
 of the SUSY unbroken.
- Calabi–Yau case: ζ is covariantly constant with respect to the Levi-Civita connection ∇ .
- Non-CY case: $\nabla \zeta \sim T^0 \zeta$ (note: Γ matrices/indices suppressed).
- T⁰ is the intrinsic torsion of the manifold.
- SU(3) decomposition: torsion splits into 5 torsion classes,

$$T^0 \in \mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3 \oplus \mathcal{W}_4 \oplus \mathcal{W}_5$$
.



Half-flat manifolds

Two (not mutually exclusive) ways to satisfy Strominger's theorem:

Option 1:

Study compactifications on SU(3) structure manifolds with torsion.

- Has been studied in eg. hep-th/0408121 (Gurrieri, Lukas, Micu), hep-th/0507173 (de Carlos, Gurrieri, Lukas, Micu).
- Torsion quantization understood for half-flat manifolds.
- Expanding the SU(3) invariant forms on appropriate bases, the only non-closed basis forms in the half-flat case satisfy

$$d\omega_i = e_i \beta^0$$
, $d\alpha_0 = e_i \tilde{\omega}^i$.

For half-flat manifolds, torsion falls into the SU(3) classes

$$T^0 \in \mathcal{W}_1^+ \oplus \mathcal{W}_2^+ \oplus \mathcal{W}_3$$
,

where + denotes the real part of W.



Domain wall vacuum

Option 2:

Break maximal symmetry of d = 4 spacetime.

- Compactification with H-flux on a half-flat or Calabi-Yau manifold.
- There exist 1/2-BPS domain wall solutions 1305.0594 (Klaput, Lukas, Svanes).
- 1/2-BPS: 2 of the 4 SUSY generators in d = 4, $\mathcal{N} = 1$ unbroken.
- d = (2 + 1) Poincaré symmetry preserved; DW breaks symmetry in transverse y direction.
- Moduli satisfy flow equations in the y coordinate.
- 10d perspective: SU(3) fibred over y → G₂ structure arXiv:1005.5302 (Lukas, Matti).



Extension: cosmic strings

The following is based on research in progress, in collaboration with Cyril Matti (City University, London) and Eirik Svanes (LPTHE, Paris).

 The 1/2-BPS domain wall breaks maximal symmetry along one noncompact direction.

Proposition:

Fibre the SU(3) structure over an interval that is not a Cartesian coordinate direction.

- Cylindrical polar coordinates (ρ, ϕ, z) : fibre along $\rho \Rightarrow$ cosmic string.
- 4d metric:

$$extit{d} s_4^2 = e^{-2B(
ho)} \left(\eta_{lphaeta} extit{d} x^lpha extit{d} x^eta + extit{d}
ho^2 +
ho^2 extit{d} \phi^2
ight) \,.$$



Cosmic string flow equations

The Killing spinor equations reduce to

$$egin{aligned} \partial_{
ho}A^I &= -ie^{-B}e^{K/2}K^{IJ^*}D_{J_*}W^* \;, \ \partial_{
ho}B &= ie^{-B}e^{K/2}W \;, \ 0 &= \operatorname{Im}(K_I\partial_{
ho}A^I) \;, \ 2\partial_{
ho}\zeta &= -\partial_{
ho}B\zeta \;, \end{aligned}$$

with the 1/2-BPS constraint $\overline{\zeta} = \sigma^{\underline{r}} \zeta$, where

$$\sigma^{\underline{\rho}} = \frac{\widehat{\mathbf{x}}}{\rho} \sigma^{\underline{1}} + \frac{\widehat{\mathbf{y}}}{\rho} \sigma^{\underline{2}} = \left(\begin{array}{cc} \mathbf{0} & \mathbf{e}^{-i\phi} \\ \mathbf{e}^{i\phi} & \mathbf{0} \end{array} \right) \ .$$

- Note: same structure as in the domain wall case (with $\rho \leftrightarrow y$).
- Related to DW by conformal transformation: $\rho = e^y$; $x = -\phi$.

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Spin(7): two transverse coordinates

The cosmic string solution is a special case of the metric

$$ds_4^2 = e^{-2B(x_a)} \left(\eta_{lphaeta} d ilde{x}^lpha d ilde{x}^eta + g_{ab} dx^a dx^b
ight) \; .$$

- We can consider more general codimension-2 topological defects.
- 10d perspective: looks like an 8-dimensional Spin(7) structure.
- For the corresponding 6d compact SU(3) structure manifold, consider a generalized half-flat manifold, which satisfies

$$d\omega_i = p_{Ai}\beta^A - q_i^A\alpha_A$$
, $d\alpha_A = p_{Ai}\tilde{\omega}^i$, $d\beta^A = q_i^A\tilde{\omega}^i$, $d\tilde{\omega}^i = 0$,

where ω_i and (α_A, β^B) are basis 2- and 3-forms, respectively.

Relevant SU(3) torsion classes are now

$$T^0 \in \mathcal{W}_1 \oplus \mathcal{W}_2 \oplus \mathcal{W}_3$$
.

Simple example

• 1/4-BPS ansatz: $\overline{\zeta} = \sigma^{1}\zeta = -i\sigma^{2}\zeta$ gives Killing spinor equations

$$\begin{split} \left(\partial_{x}-i\partial_{y}\right)A^{I}&=-ie^{-B}e^{K/2}K^{IJ^{*}}D_{J*}W^{*}\;,\\ \left(\partial_{x}+i\partial_{y}\right)B&=ie^{-B}e^{K/2}W\;,\\ 0&=\text{Im}(K_{I}\partial_{a}A^{I})\;,\\ 2\partial_{a}\,\zeta&=-\partial_{a}B\zeta\;. \end{split}$$

- Simple case: $p_{0i} = q_i^A = 0$, H = 0, vanishing axions.
- In this simple example, 4d and 10d equations match when x-dependence vanishes
 - \Rightarrow alternative 1/2-BPS DW solution, G_2 sublocus of Spin(7).
- KSEs satisfy $\frac{\mathcal{K}'}{\mathcal{K}} = \frac{\tilde{\mathcal{K}}'}{\tilde{\mathcal{K}}} = -2\phi' = \frac{3}{\sqrt{\mathcal{K}\tilde{\mathcal{K}}}} p_{ai} v^i w^a \; ,$

where
$$(v^i, w^a) = \operatorname{Im}(T^i, Z^a)$$
, $\mathcal{K} = \mathcal{K}_{ijk} v^i v^j v^k$, $\tilde{\mathcal{K}} = \tilde{\mathcal{K}}_{abc} w^a w^b w^c$.

Stephen Angus

Summary

- String compactifications generate moduli, which must be stabilized. This can be done using fluxes.
- For the heterotic string, only H_3 present. One can compactify on SU(3) structure manifolds which are not Calabi–Yau, and/or sacrifice maximal symmetry in d=4. Domain wall solutions have been considered.
- We found a cosmic string solution, which is a 1/2-BPS G_2 fibration. It is related to the domain wall case by a conformal transformation.
- In the more general codimension-2 case, for vanishing H-flux we found a new type of domain wall solution at a G₂ locus of Spin(7).
- Goal: include flux; 1/4-BPS solutions (e.g. intersecting domain walls → extra warp factor?); more general Spin(7) structure ... work in progress!

