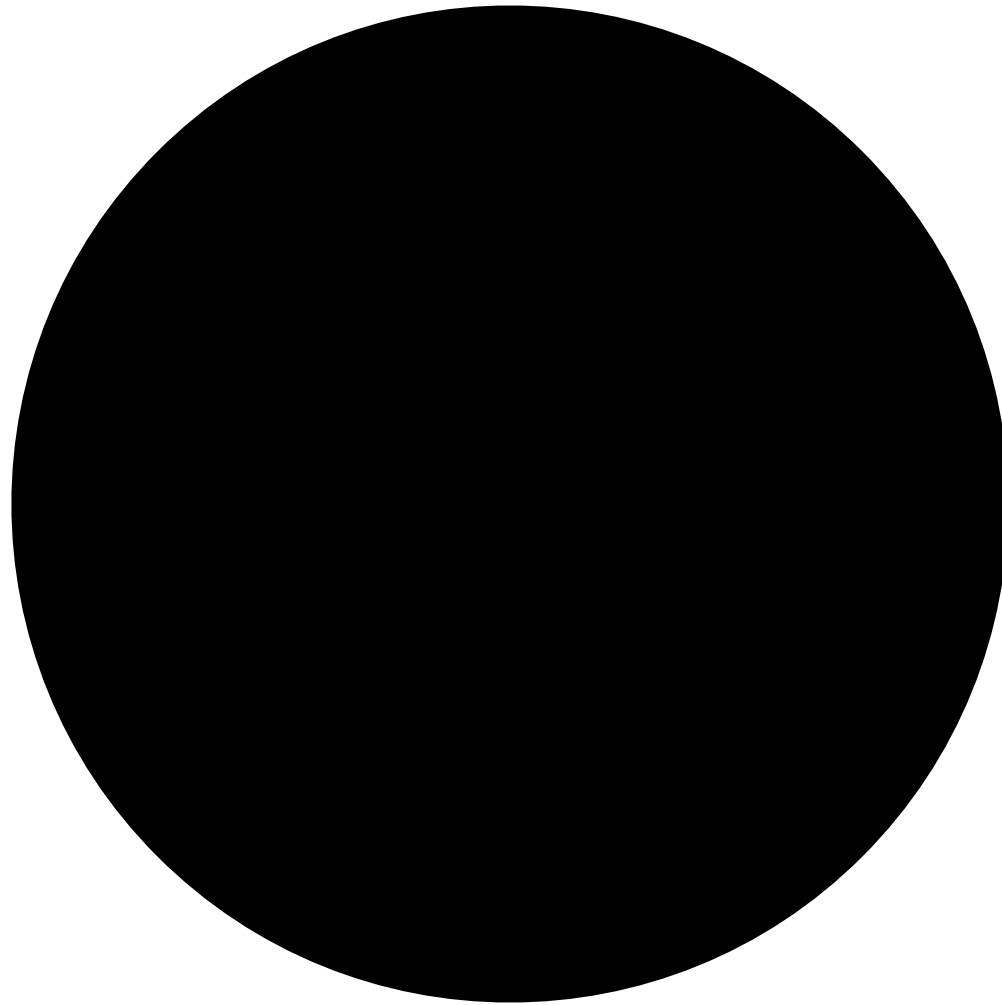


Spacetime wormholes in quantum gravity

Tom Hartman
Cornell University

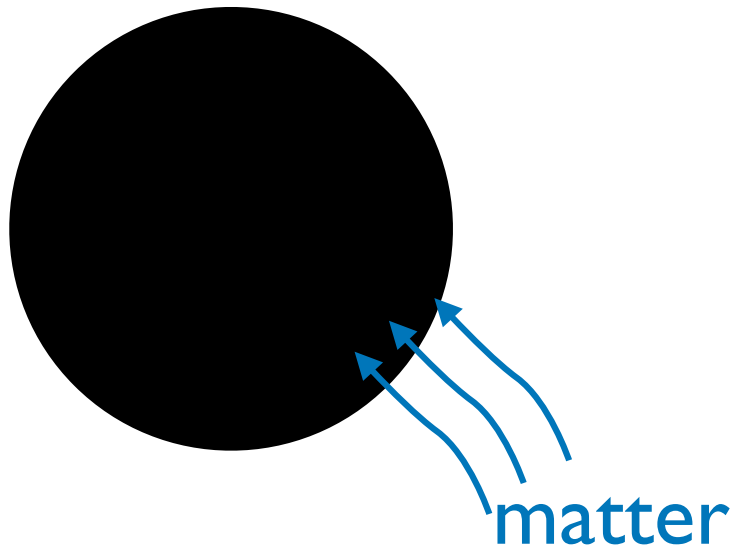
IFT Christmas Workshop ♦ Madrid ♦ December 16, 2022

Classical black holes



***are featureless* objects — pure spacetime curvature**

However, they behave like thermodynamic systems with a very large number of microscopic degrees of freedom:

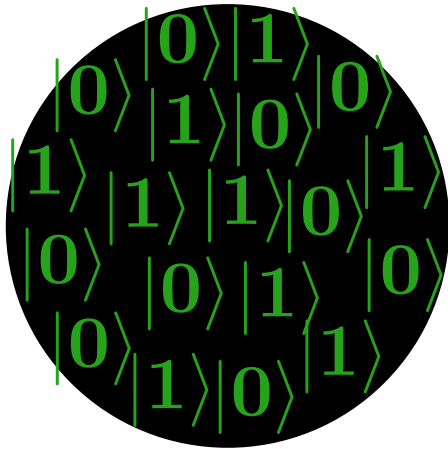


Hawking temperature $T = \frac{\hbar}{4\pi R}$

entropy $S = \frac{\text{Area}}{4\hbar G}$

Sagittarius A^* $S = 2^{88}$

In quantum gravity,



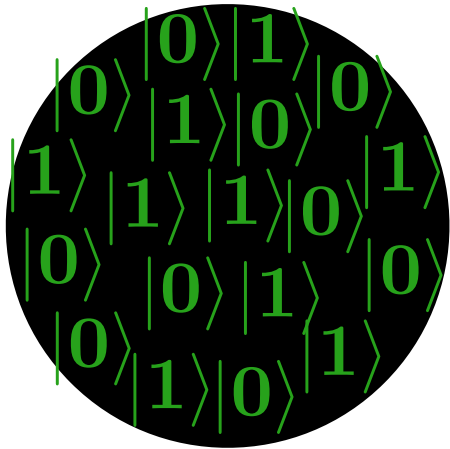
$$\# \text{ states} \approx e^S$$

How is this quantum information encoded?

How does it escape when a black hole evaporates?

What else does this tell us about the UV completion?

In quantum gravity,



$$\# \text{ states} \approx e^S$$

How is this quantum information encoded?

How does it escape when a black hole evaporates?

What else does this tell us about the UV completion?

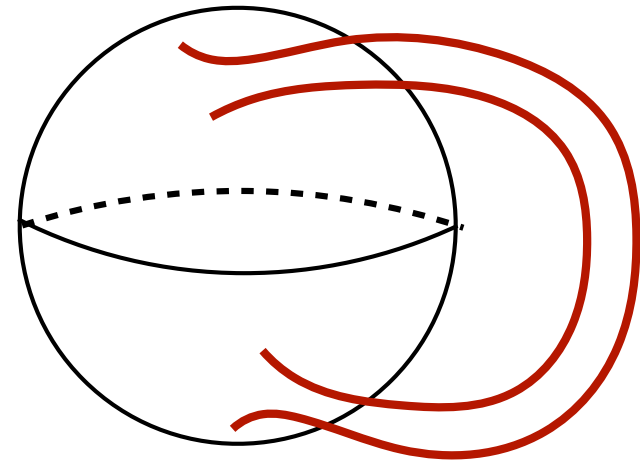
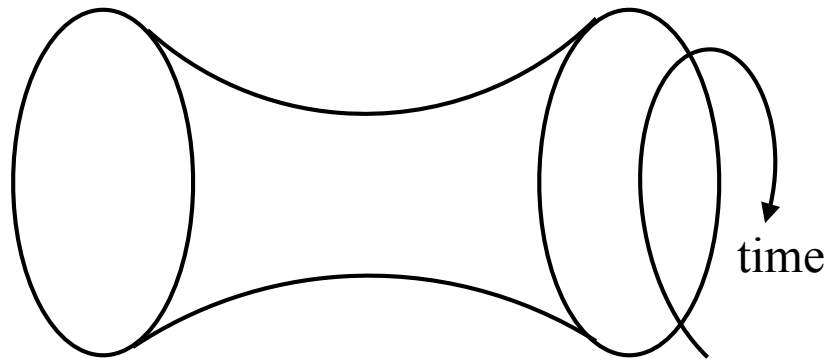
Recent progress

◆ Modern tools of quantum information science

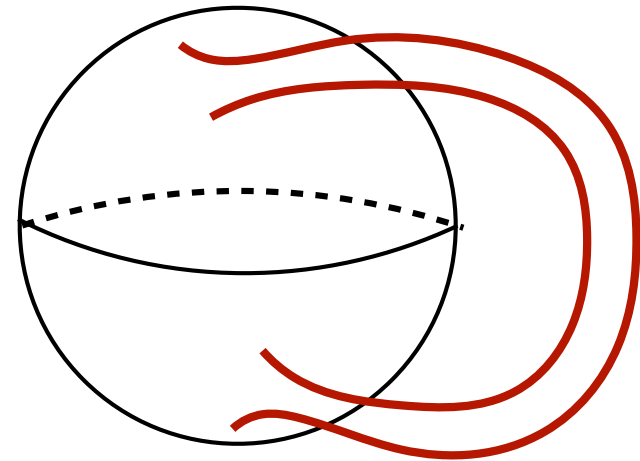
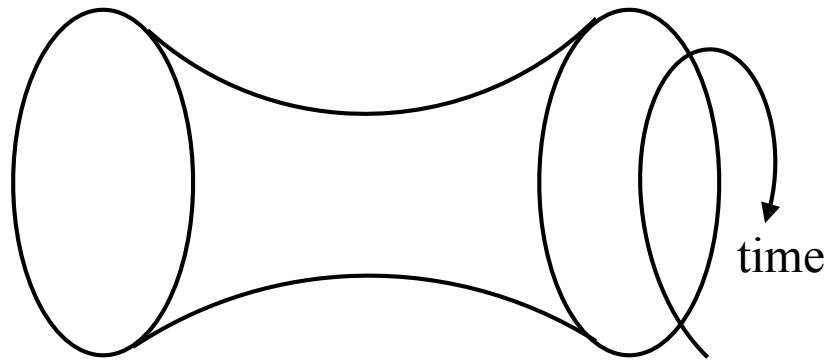
$$S = -\text{tr} \rho \ln \rho , \quad I(A : B) , \quad \mathcal{N}_{A \rightarrow B}$$

◆ Spacetime wormholes

Higher topology spacetime manifolds that contribute to the gravitational path integral

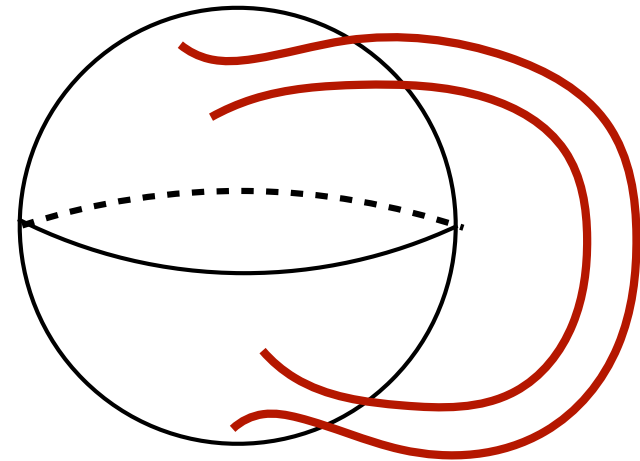
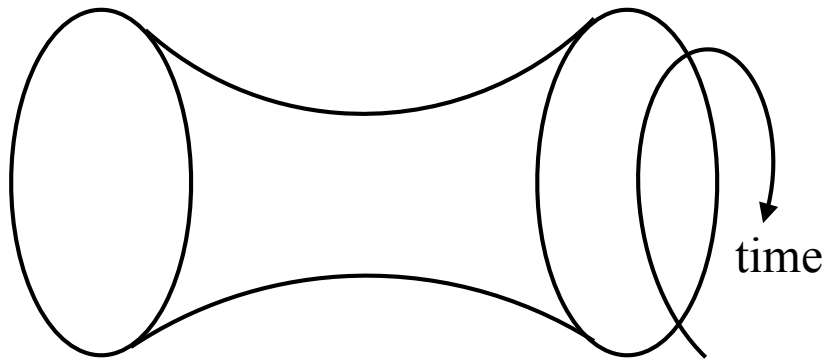


These are not the wormholes of science fiction — those are *spatial* wormholes (and don't exist in theories with positive energy!)



These are not the wormholes of science fiction — those are *spatial* wormholes (and don't exist in theories with positive energy!)

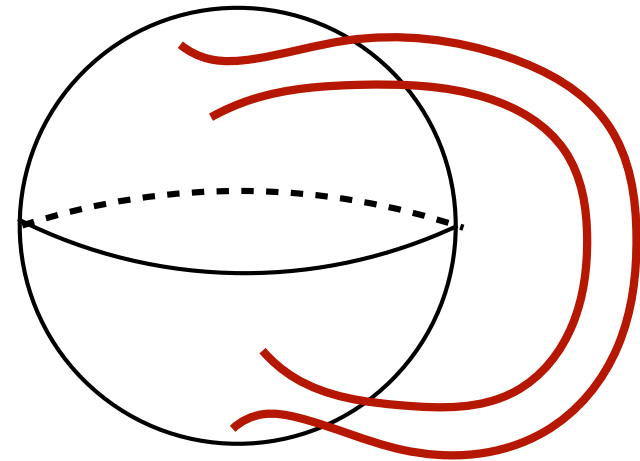
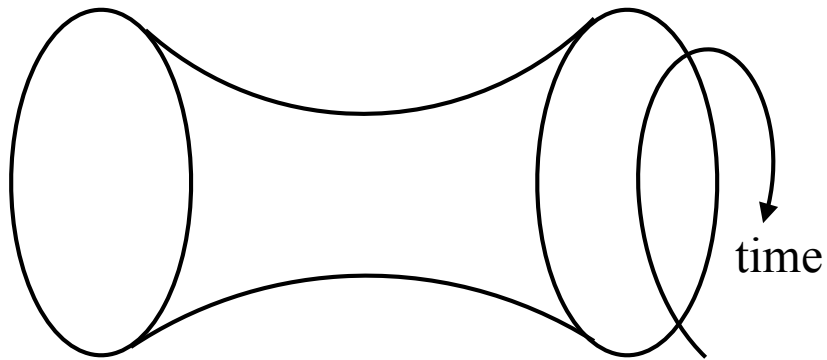
Spacetime wormholes are 4-dimensional solutions of General Relativity that exist whenever there are black holes.



These are not the wormholes of science fiction — those are *spatial* wormholes (and don't exist in theories with positive energy!)

Spacetime wormholes are 4-dimensional solutions of General Relativity that exist whenever there are black holes.

Generic prediction of GR — any Λ , any matter content, no strings required

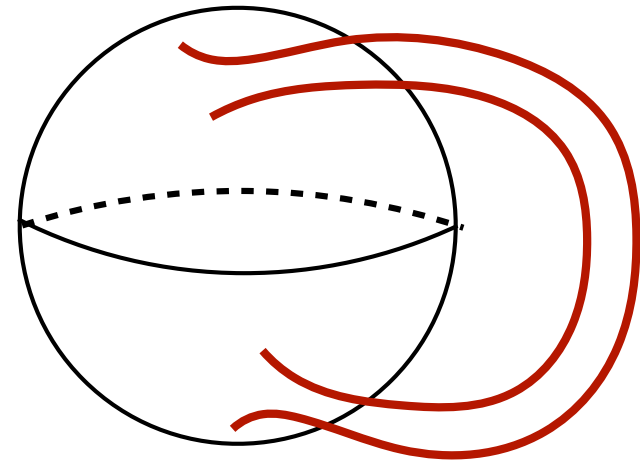
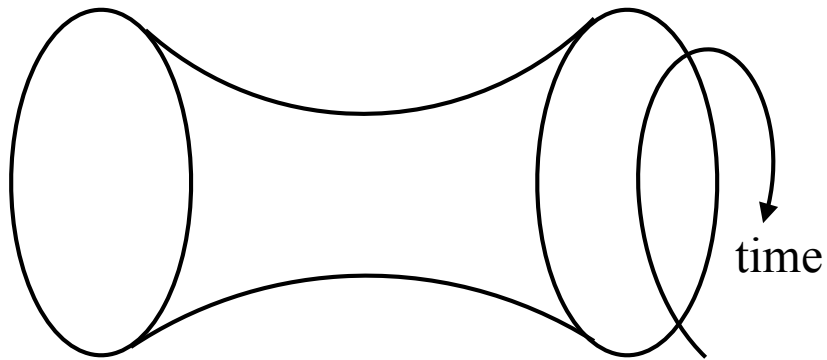


These are not the wormholes of science fiction — those are *spatial* wormholes (and don't exist in theories with positive energy!)

Spacetime wormholes are 4-dimensional solutions of General Relativity that exist whenever there are black holes.

Generic prediction of GR — any Λ , any matter content, no strings required

Complex-valued metric.



These are not the wormholes of science fiction — those are *spatial* wormholes (and don't exist in theories with positive energy!)

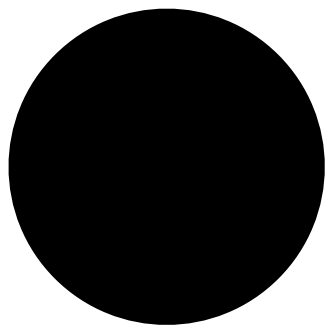
Spacetime wormholes are 4-dimensional solutions of General Relativity that exist whenever there are black holes.

Generic prediction of GR — any Λ , any matter content, no strings required

Complex-valued metric.

Gravitational instantons.

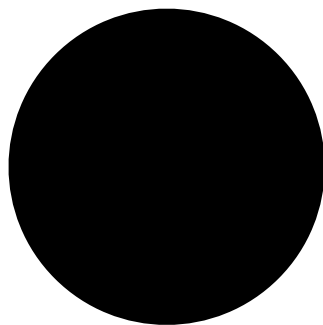
Black hole entropy tells us the average density of states in the UV theory,



$$\Rightarrow \overline{\rho(E)}$$

(Averaged over a small energy window)

Black hole entropy tells us the average density of states in the UV theory,



$$\Rightarrow \overline{\rho(E)}$$

(Averaged over a small energy window)

Spacetime wormholes tell us more detailed statistical information about the UV.

I will discuss two examples:

I. Replica wormholes \Rightarrow entropy of Hawking radiation

II. Thin-shell wormholes \Rightarrow operator statistics and global symmetry violation

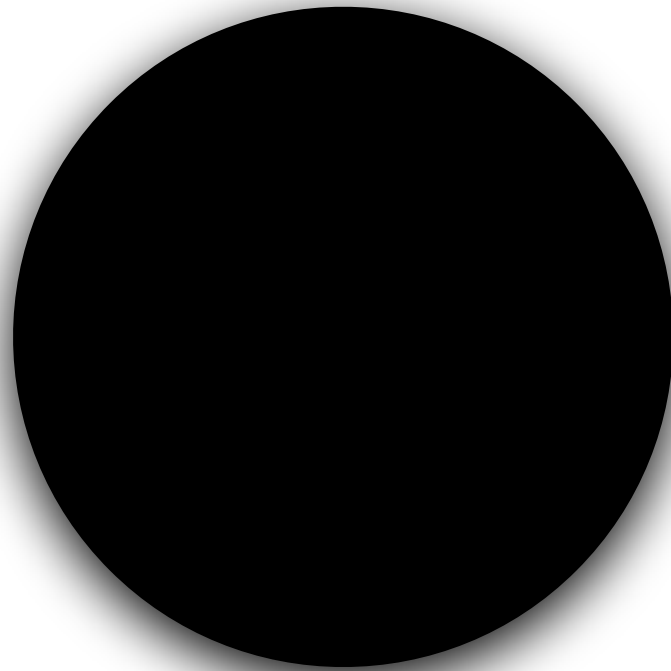
I. Replica Wormholes

Based on:

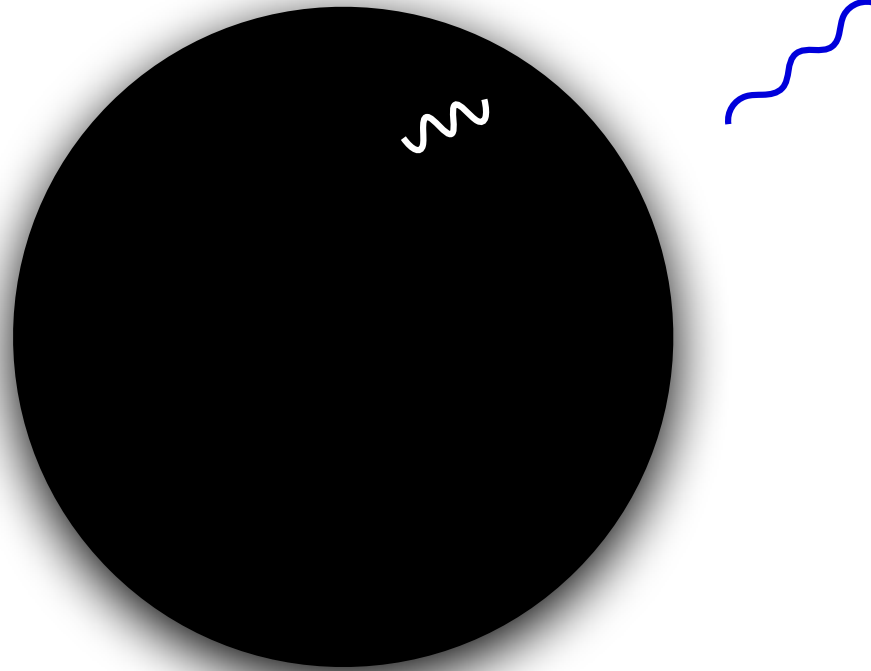
[Almheiri, TH, Maldacena, Shaghoulian, Tajdini '19]

[Penington, Shenker, Stanford, Yang '19]

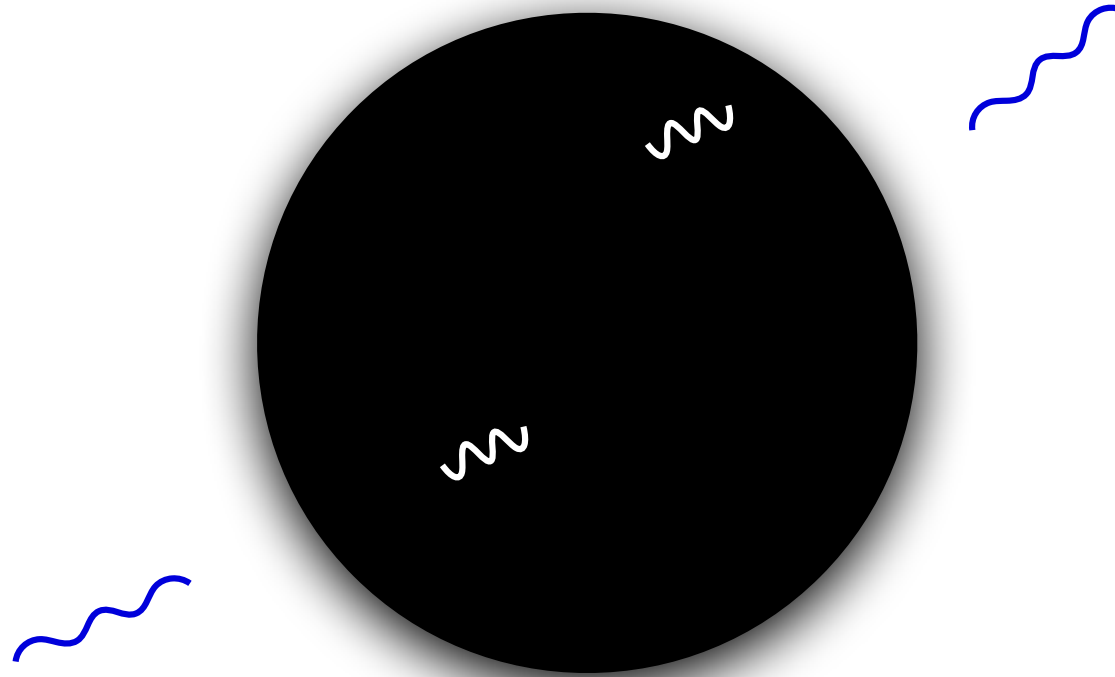
Information loss in black hole evaporation



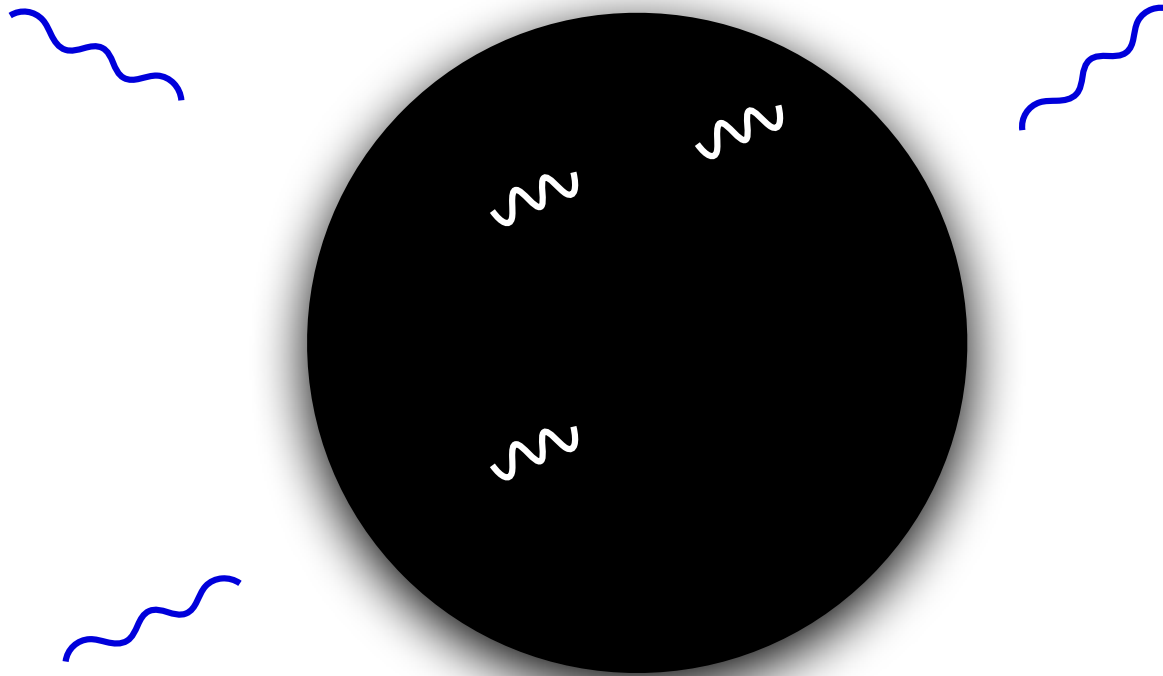
Information loss in black hole evaporation



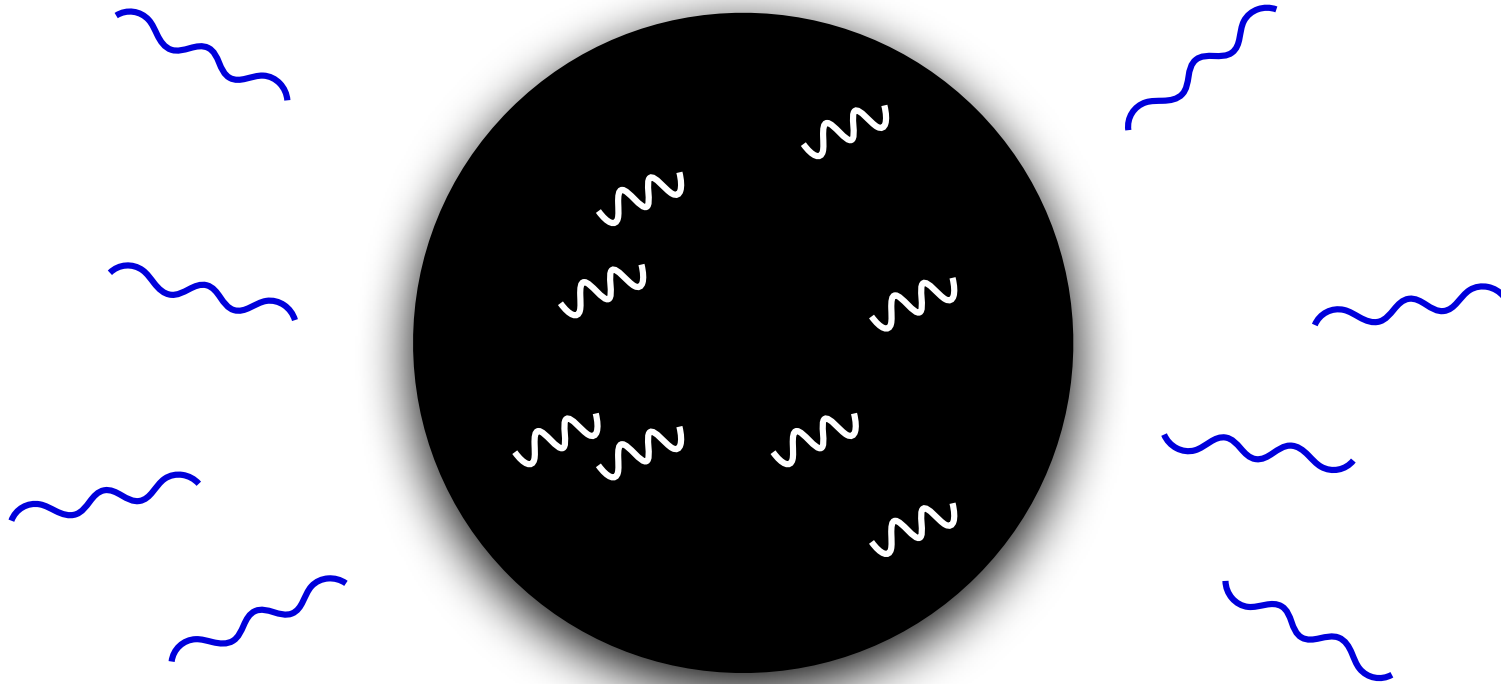
Information loss in black hole evaporation



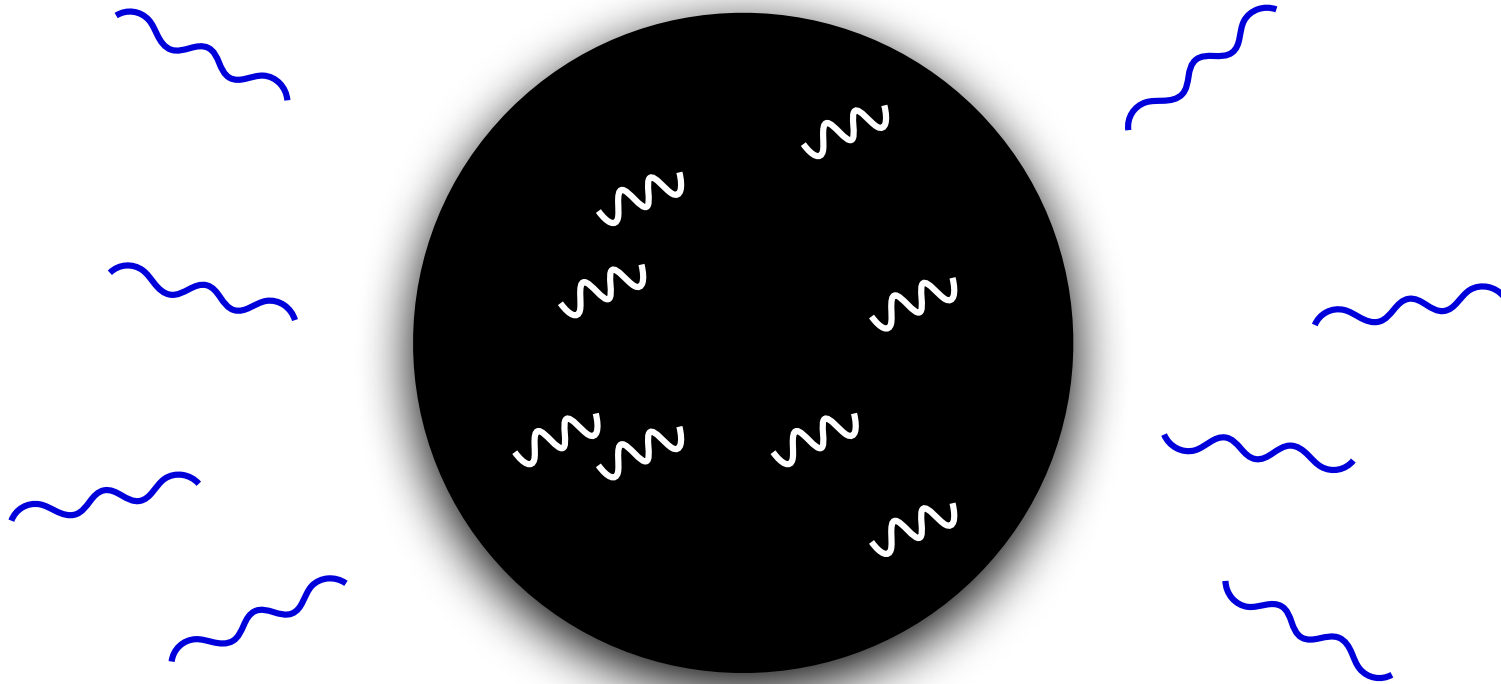
Information loss in black hole evaporation



Information loss in black hole evaporation



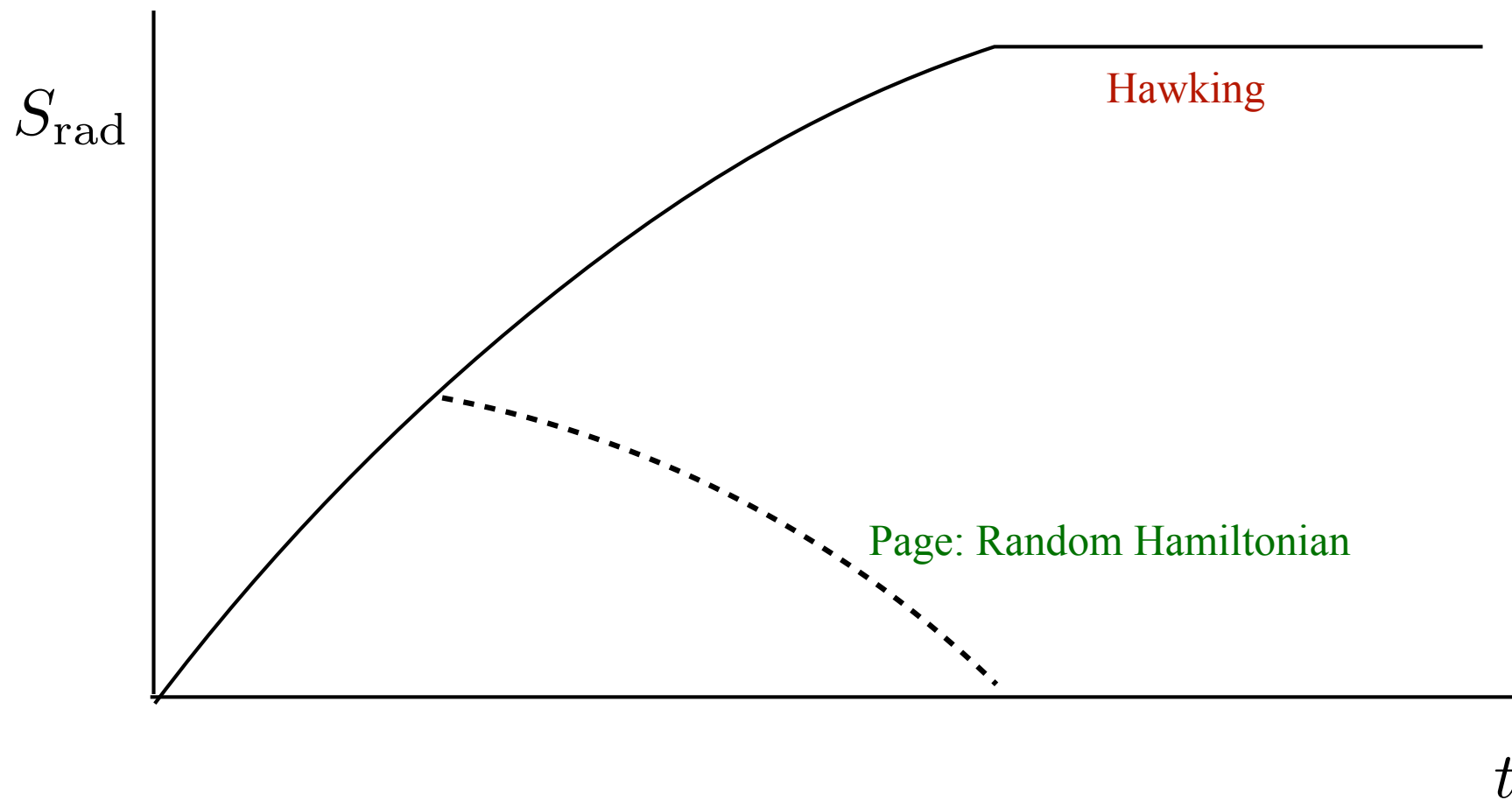
Information loss in black hole evaporation



Hawking radiation is a process of *entanglement production* between the black hole interior and the radiation.

The “paradox” is that at the end, there is nothing for the radiation to be entangled with — QIS.

Entropy of Hawking radiation

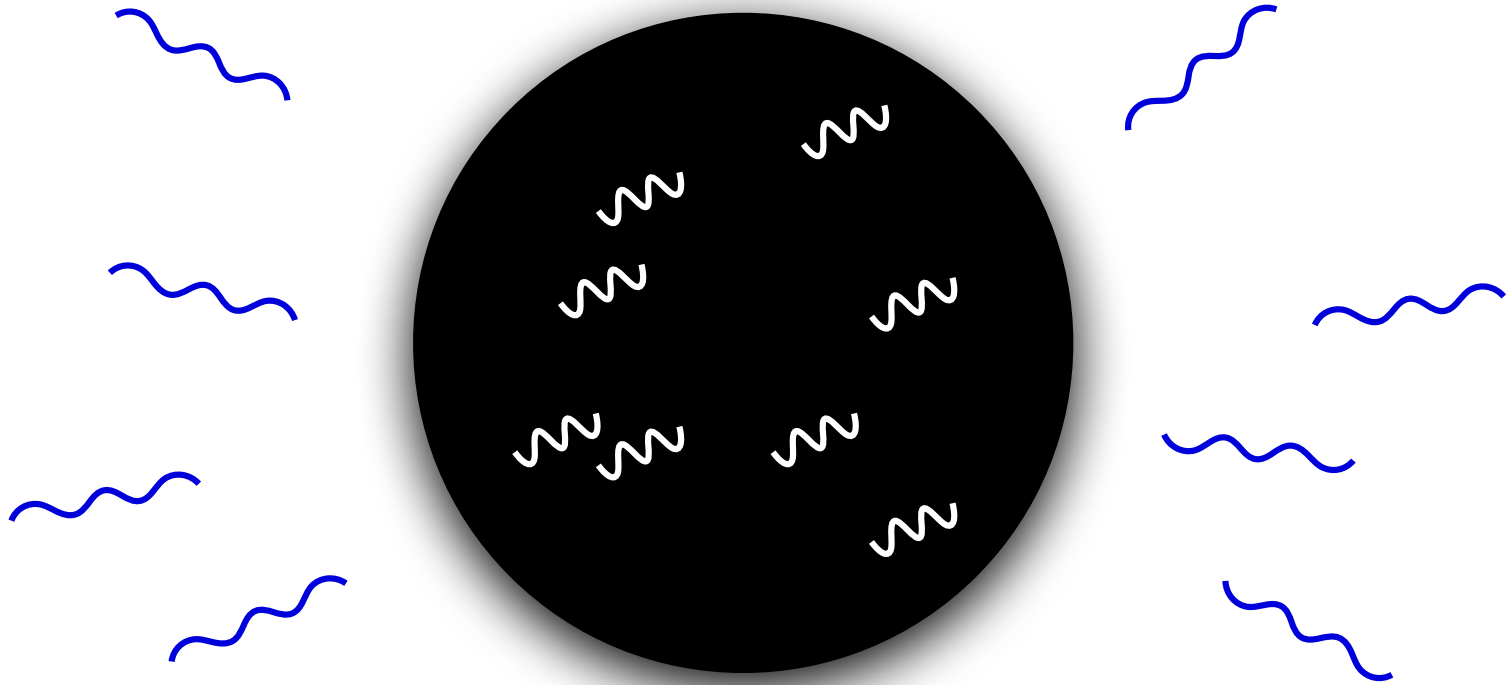


The Island Effect

[Penington, '19]

[Almheiri, Engelhardt, Marolf, Maxfield '19]

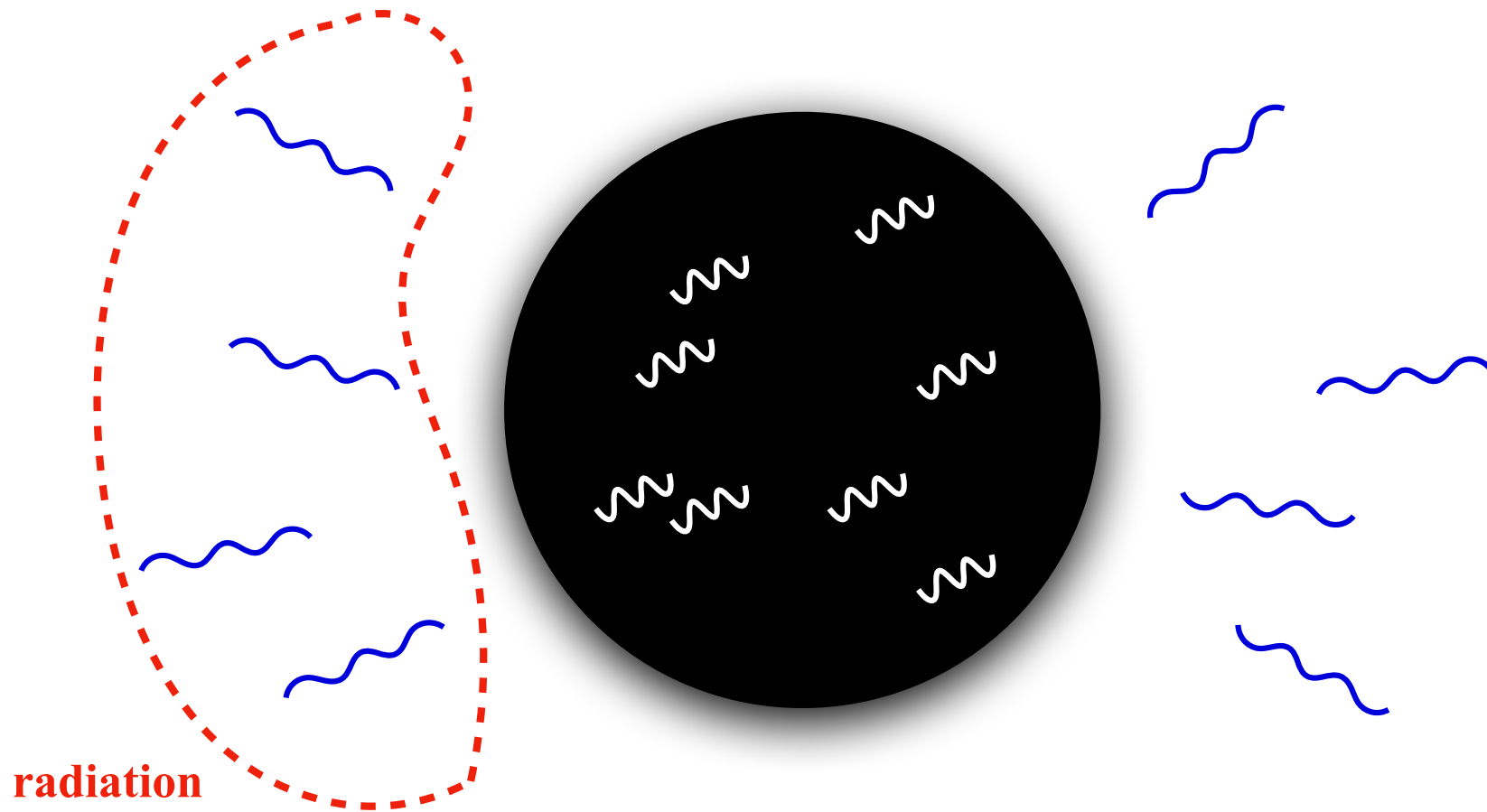
The Island Effect



[Penington, '19]

[Almheiri, Engelhardt, Marolf, Maxfield '19]

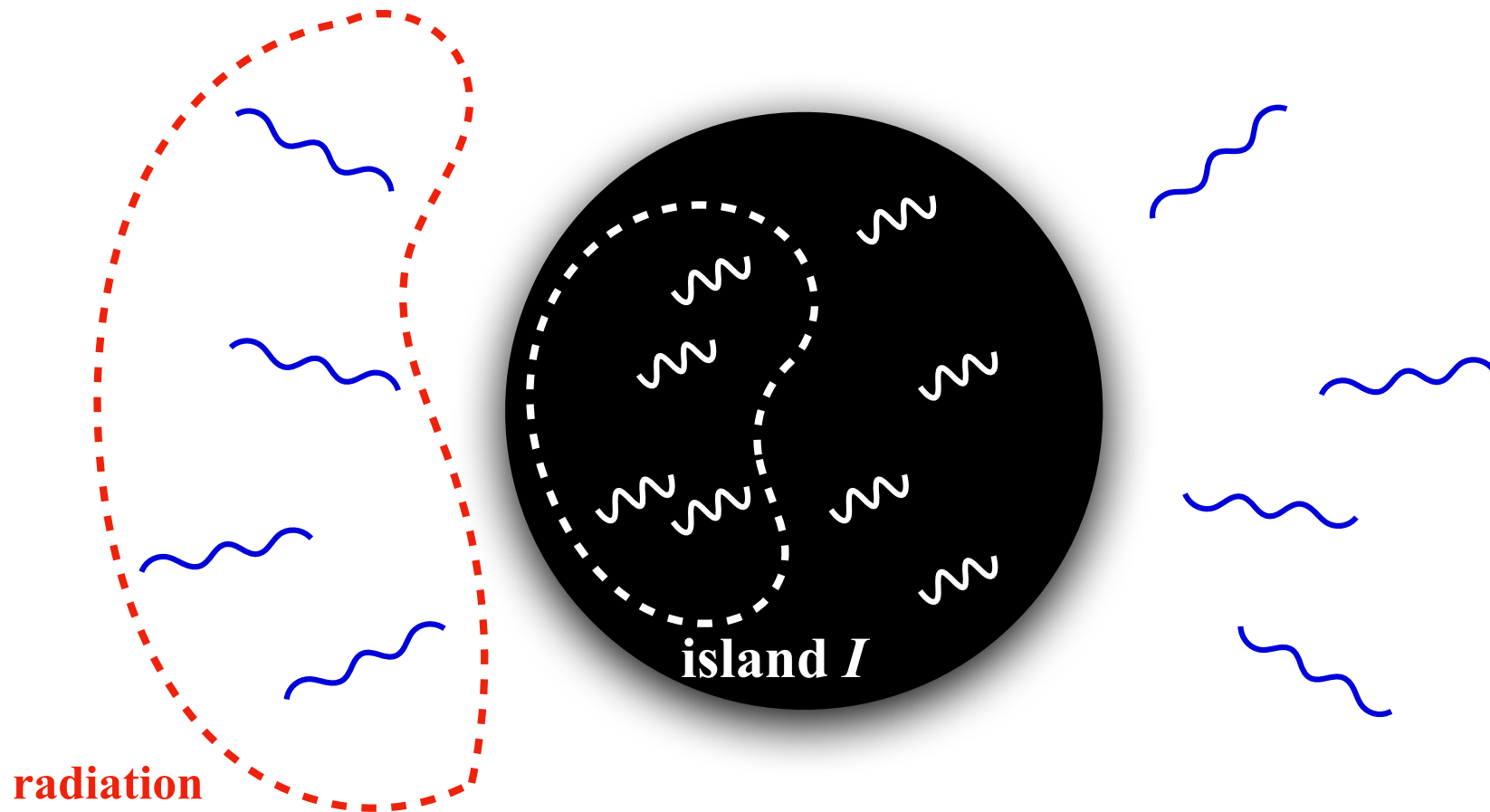
The Island Effect



[Penington, '19]

[Almheiri, Engelhardt, Marolf, Maxfield '19]

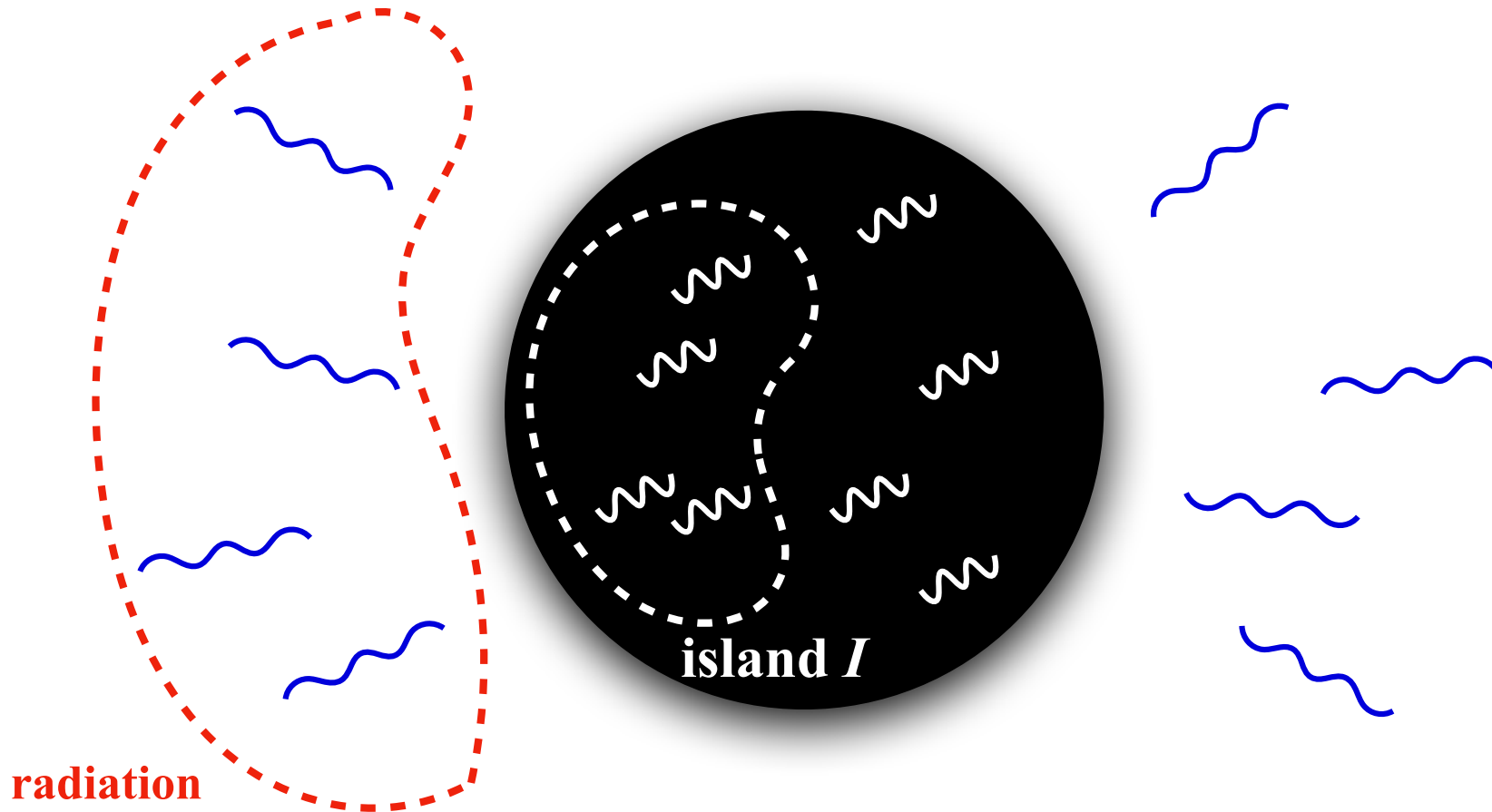
The Island Effect



[Penington, '19]

[Almheiri, Engelhardt, Marolf, Maxfield '19]

The Island Effect



Entropy formula:
$$S(\rho_{\text{rad}}) = \min_I \text{ext}_I \left[\frac{\text{Area}(\partial I)}{4} + S_{\text{QFT}}(I \cup \text{rad}) \right]$$

[Penington, '19]

[Almheiri, Engelhardt, Marolf, Maxfield '19]

Replica wormholes

A “replica wormhole” is a gravitational instanton supported by matter entanglement.

\implies island effect

Replica wormholes

A “replica wormhole” is a gravitational instanton supported by matter entanglement.

\implies island effect

Consider the "purity"

$$Z_2 := \text{tr} (\rho^2)$$

$$\text{pure: } Z_2 = 1$$

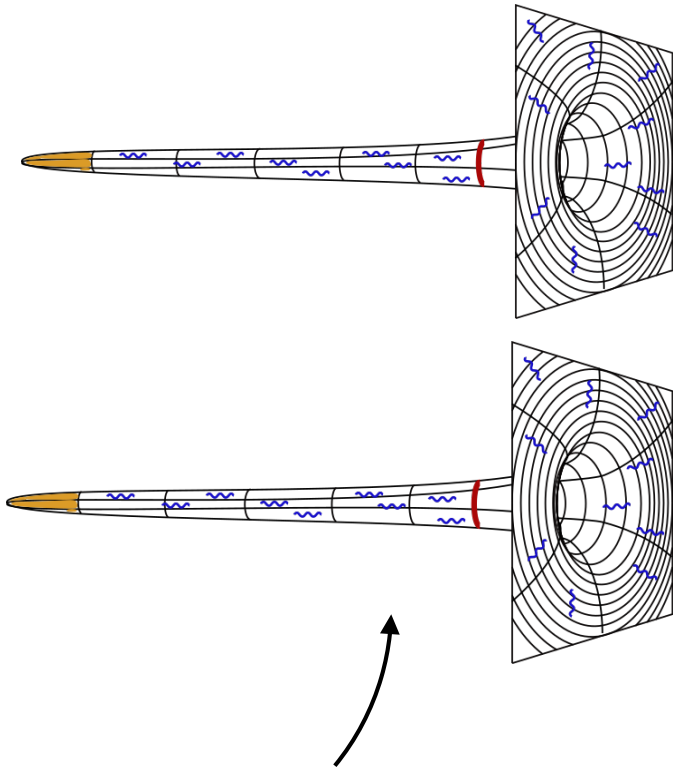
$$\text{mixed: } Z_2 < 1$$

Calculate the purity of the Hawking radiation by a gravitational path integral:

$$\text{tr}(\rho^2) \approx$$

Calculate the purity of the Hawking radiation by a gravitational path integral:

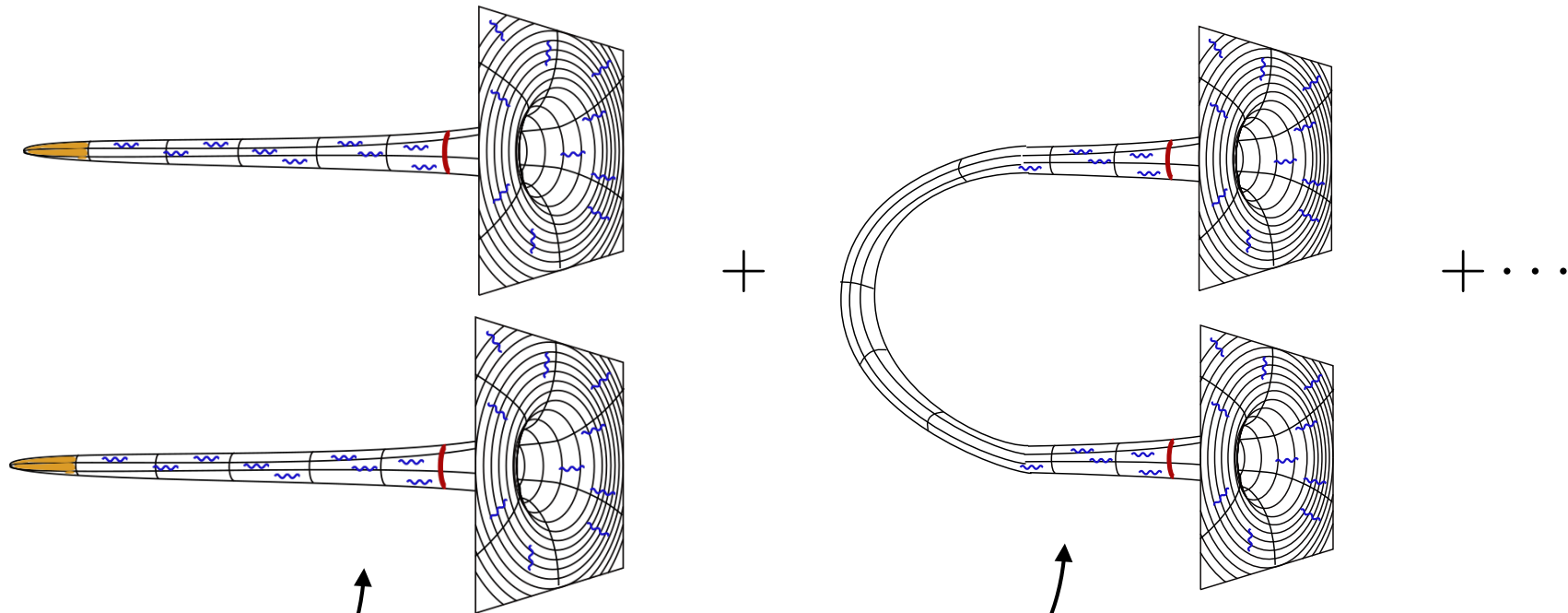
$$\text{tr}(\rho^2) \approx$$



Hawking radiation
is highly mixed

Calculate the purity of the Hawking radiation by a gravitational path integral:

$$\text{tr}(\rho^2) \approx$$

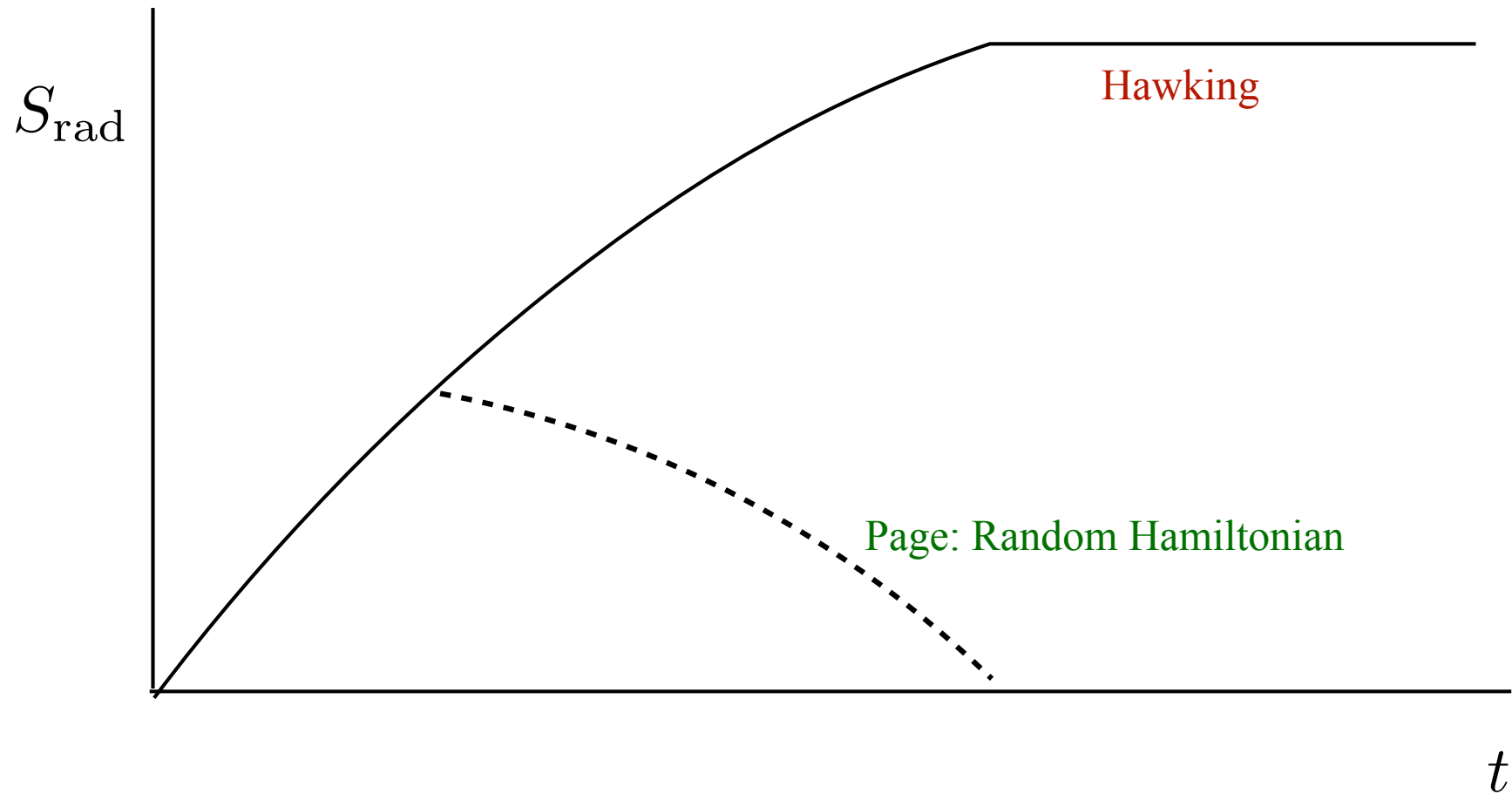


Hawking radiation
is highly mixed

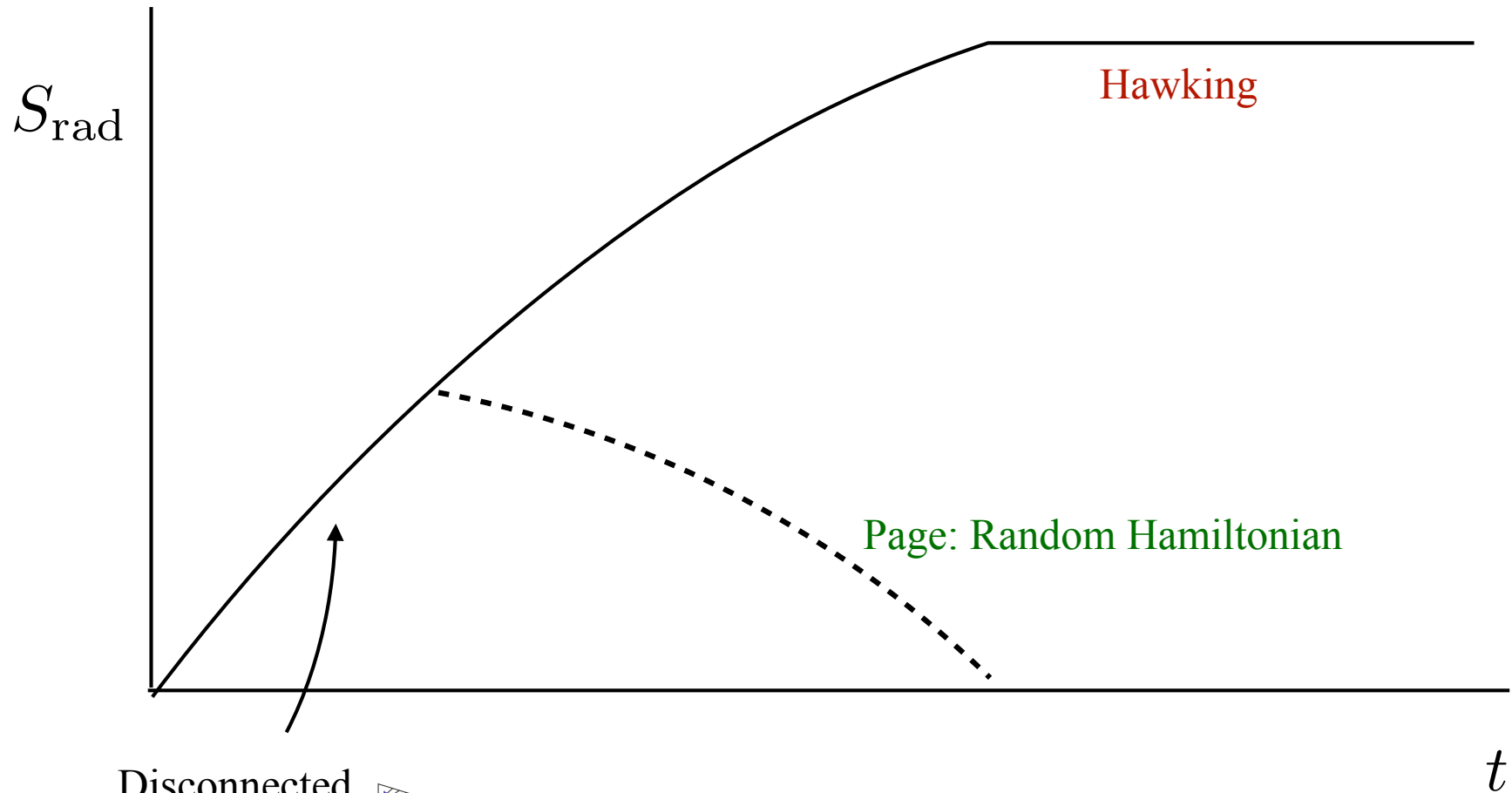
$$G_{\mu\nu} = 8\pi G_N \langle T_{\mu\nu} \rangle$$

Hawking radiation re-purifies

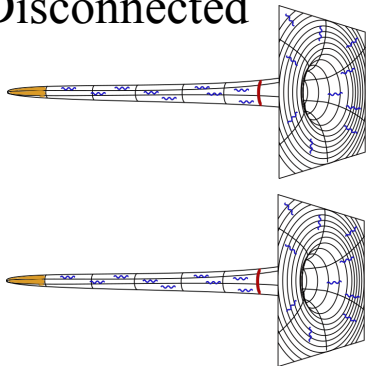
Intriguingly: The resulting entropy agrees precisely with the “Random Hamiltonian” prediction of Page.



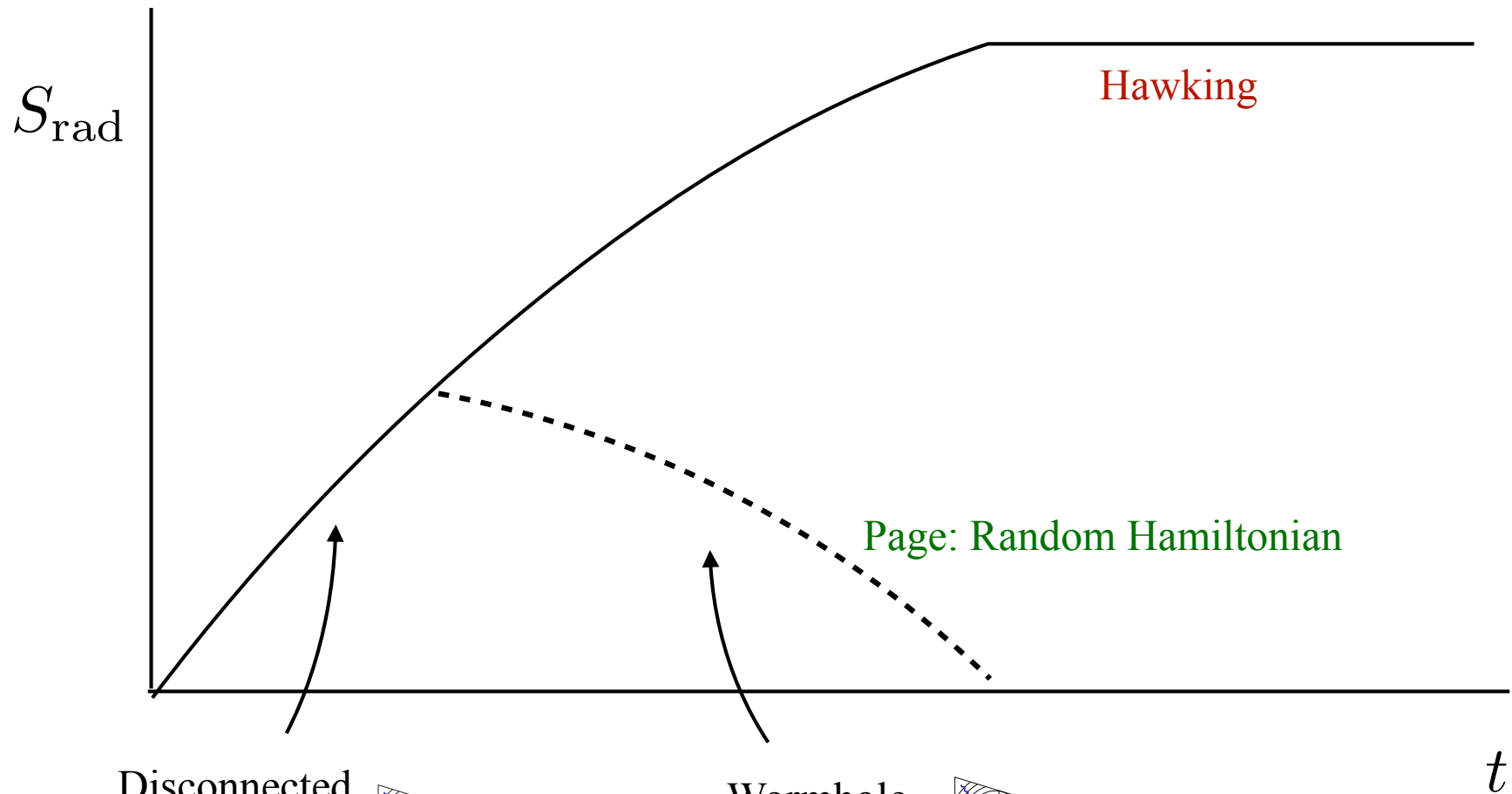
Intriguingly: The resulting entropy agrees precisely with the “Random Hamiltonian” prediction of Page.



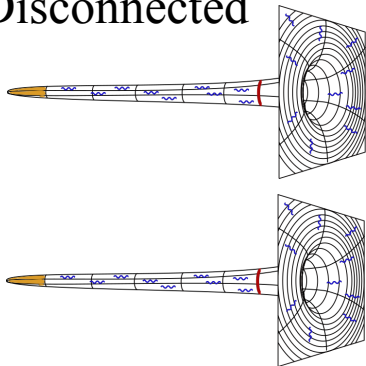
Disconnected



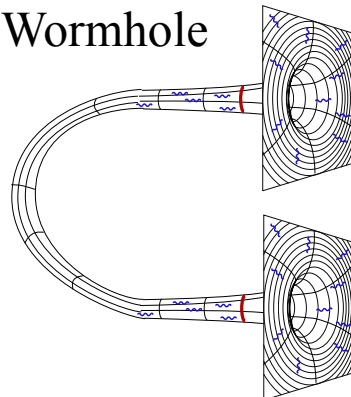
Intriguingly: The resulting entropy agrees precisely with the “Random Hamiltonian” prediction of Page.



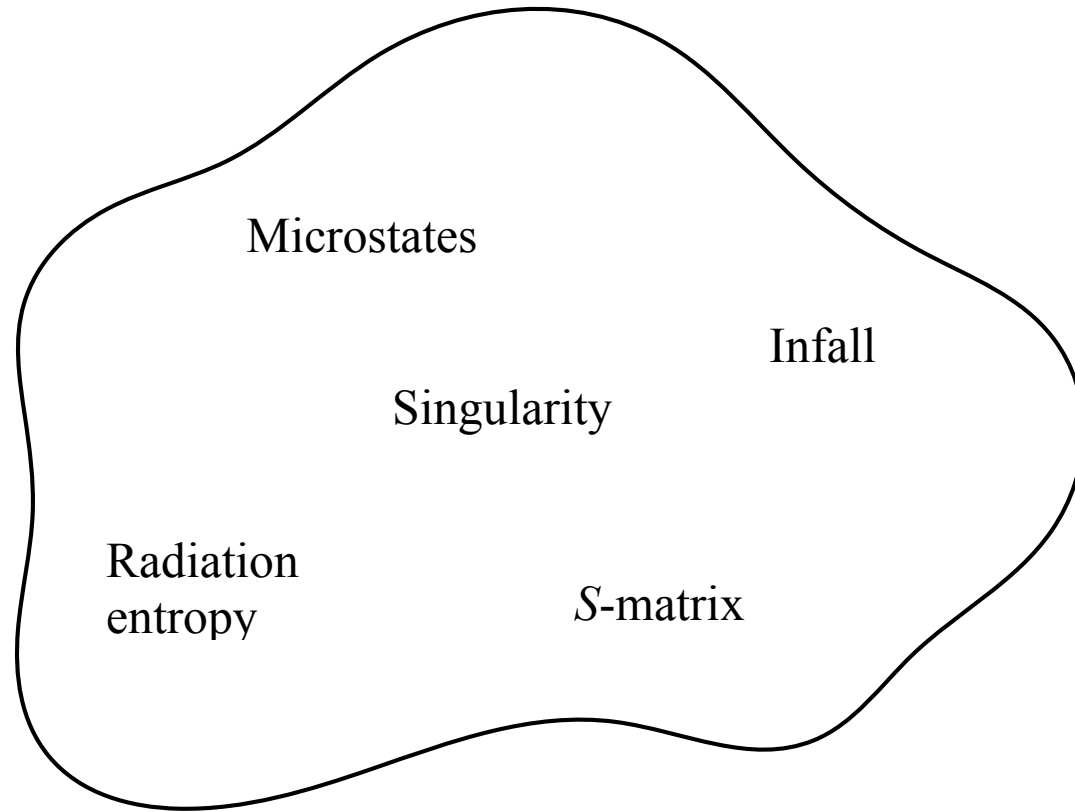
Disconnected



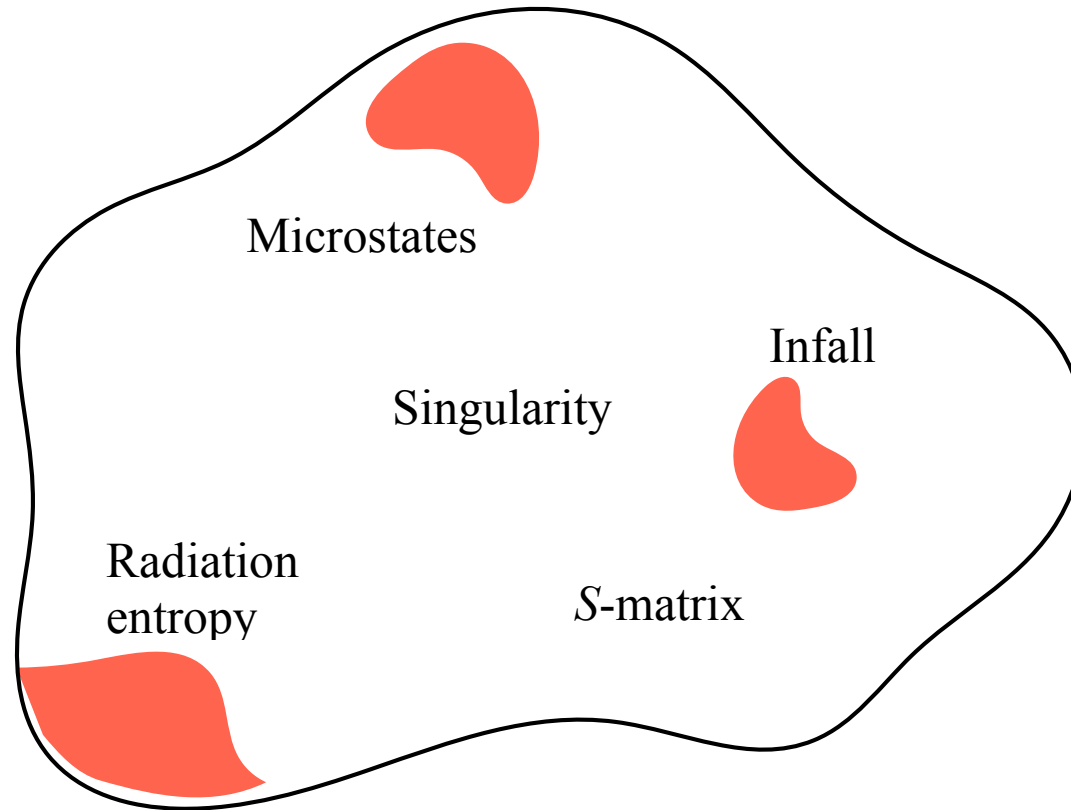
Wormhole



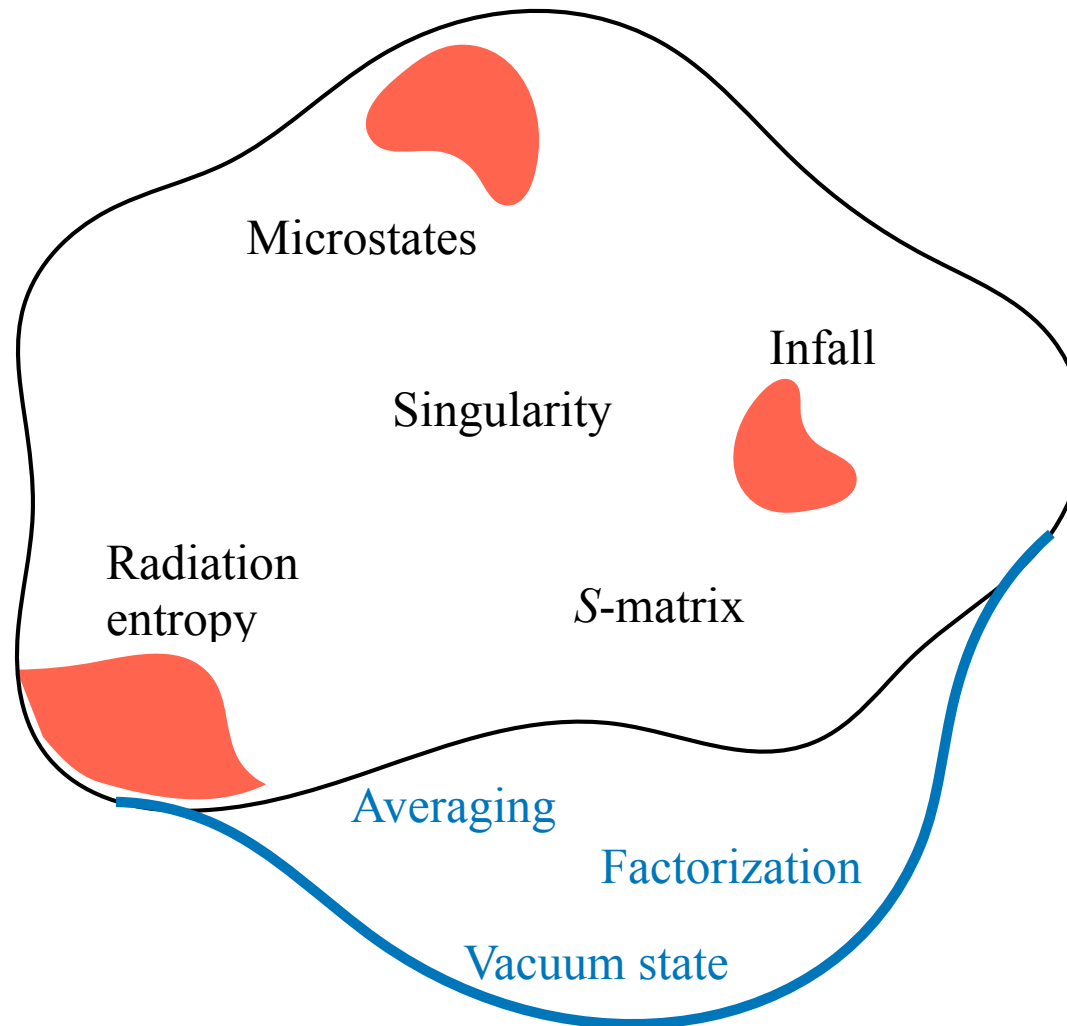
Status of the information problem



Status of the information problem



Status of the information problem



Some corners have been solved; but the scope of the problem is now **BIGGER!**

II. Thin-shell wormholes

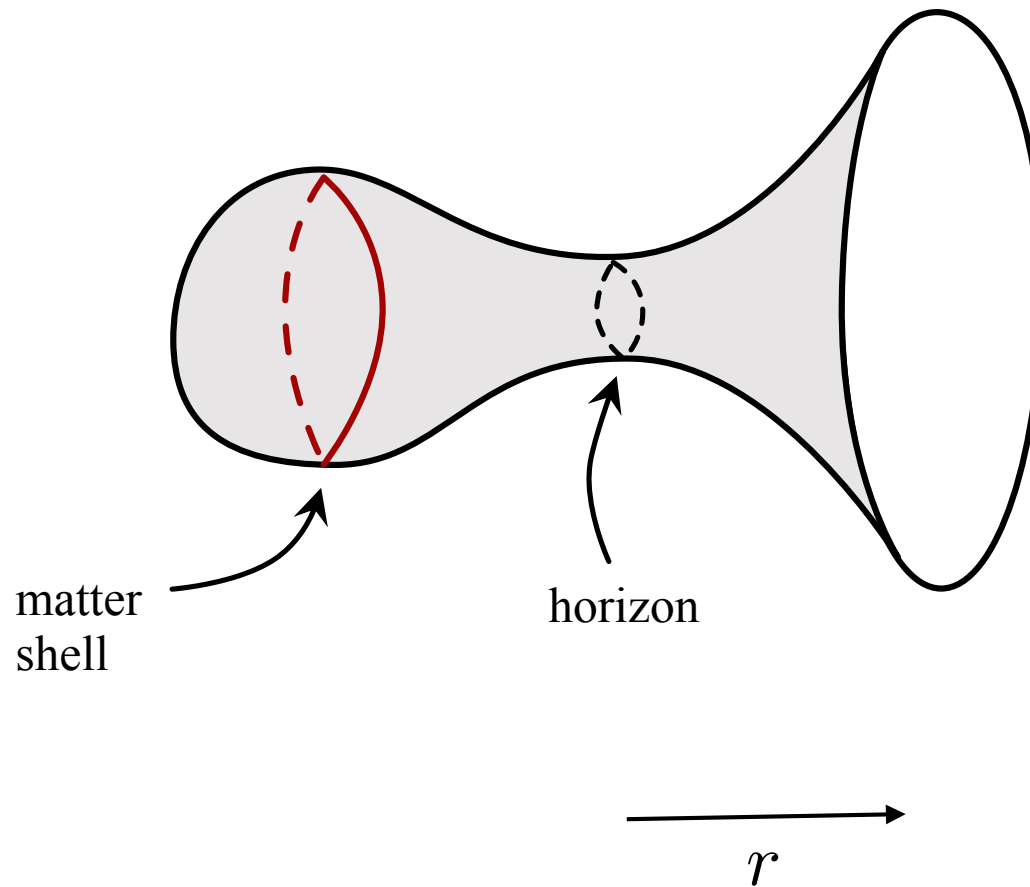
Based on:

[Chandra, Collier, TH, Maloney '22]

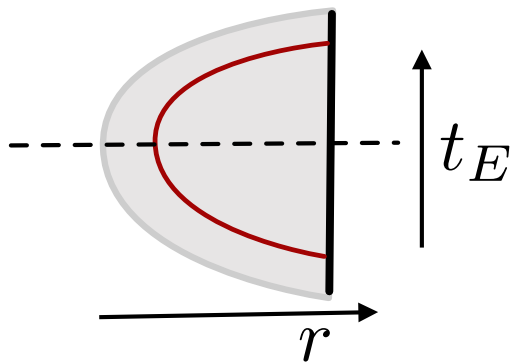
[Chandra, TH '22]

Consider a black hole created by a spherically symmetric thin shell of massive particles in 4 spacetime dimensions.

Spatial geometry



In the Euclidean path integral, this black hole is a solution that looks like this:



Standard black hole thermodynamics from the path integral:

[Gibbons, Hawking]

$$\langle \Psi | \Psi \rangle \approx e^{-I_{cl}} \quad (\text{WKB})$$

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\langle \Psi | \Psi \rangle = \sum_n |\psi_n|^2$$

Therefore: General relativity “knows” the average value of $|\psi_n|^2$

Question: Does General Relativity know higher moments of black hole states?

$$\overline{\psi_{n_1} \psi_{n_2} \psi_{n_3} \psi_{n_4} \cdots}$$

Question: Does General Relativity know higher moments of black hole states?

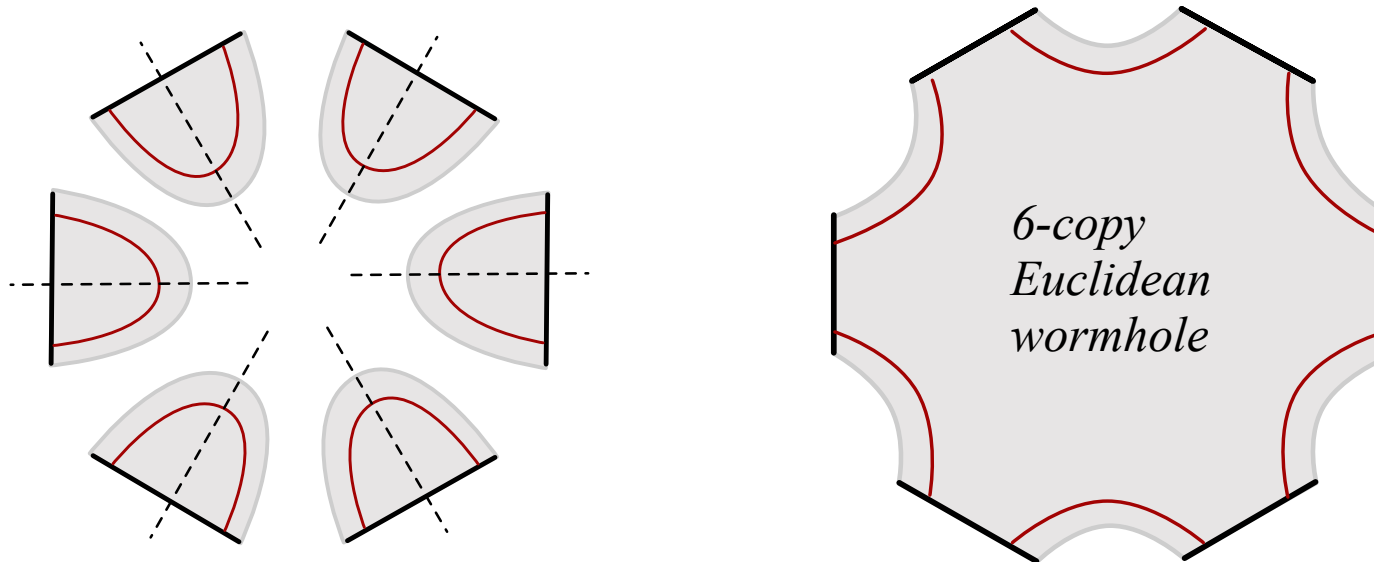
$$\overline{\psi_{n_1} \psi_{n_2} \psi_{n_3} \psi_{n_4} \cdots}$$

A partial answer:

Yes. At least some higher moments are encoded in spacetime wormholes connecting multiple copies of the original black hole.

d = 2 [Penington, Shenker, Stanford, Yang '19]

d > 2 [Chandra, TH '22]



Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:

$$\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots = \sum_{n_1, n_2, \dots} \psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots$$

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:


$$\overline{\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots} = \sum_{n_1, n_2, \dots} \psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots$$

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:


$$\overline{\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots} = \sum_{n_1, n_2, \dots} \psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots$$


Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:

$$\overline{\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots} = \sum_{n_1, n_2, \dots} \psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots$$


The diagram shows the summation over indices n_1, n_2, \dots . Below the summation, there are red and green brackets. The first term $\psi_{n_1}^* \psi_{n_1}$ has a red bracket underneath. The second term $\psi_{n_2}^* \psi_{n_2}$ has a red bracket underneath. A green bracket connects the ψ_{n_1} of the first term to the $\psi_{n_2}^*$ of the second term. Another green bracket connects the ψ_{n_2} of the second term to the $\psi_{n_1}^*$ of the first term. This illustrates the contraction of indices between different copies of the state.

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:

$$\begin{aligned} \overline{\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots} &= \sum_{n_1, n_2, \dots} \overline{\psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots} \\ &= \left(\sum_n |\psi_n|^2 \right)^k + \boxed{\sum_n |\psi_n|^{2k}} + \cdots \end{aligned}$$

Suppose $|\Psi\rangle$ is a random quantum state, drawn from a Gaussian distribution:

$$|\Psi\rangle = \sum_n \psi_n |n\rangle$$

$$\overline{\psi_m^* \psi_n} = f(E_m) \delta_{mn}$$

To study statistics we look at many copies together:

$$\begin{aligned} \overline{\langle \Psi | \Psi \rangle \langle \Psi | \Psi \rangle \cdots} &= \sum_{n_1, n_2, \dots} \overline{\psi_{n_1}^* \psi_{n_1} \psi_{n_2}^* \psi_{n_2} \cdots} \\ &= \left(\sum_n |\psi_n|^2 \right)^k + \sum_n |\psi_n|^{2k} + \cdots \\ &= \text{[Diagram: 4 sectors]} + \text{[Diagram: 8-sided polygon]} + \cdots \end{aligned}$$

The diagram shows the expansion of the average of the product of k copies of the state norm. The first term is the k -th power of the average norm, represented by a red sum. The second term is the sum of the $2k$ -th power of the norm, highlighted in a green box. The diagrams below illustrate the combinatorial structure: the first term is represented by four sectors meeting at a central point, and the second term is represented by an eight-sided polygon with concave sides.

In this case, and many similar examples,



Subject to the constraints of gauge invariance and asymptotic locality.

Another example:

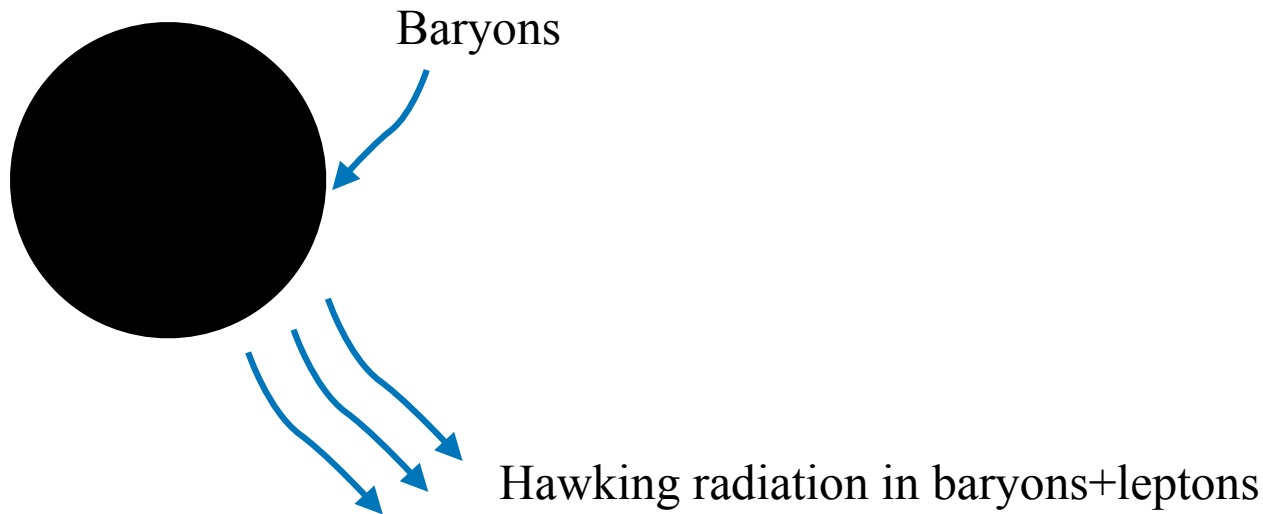
$$\begin{array}{l} \text{Einstein gravity} \\ \text{in } \text{AdS}_3 \end{array} = \begin{array}{l} \text{Average over } \text{CFT}_2\text{'s} \\ \text{subject to conformal bootstrap} \end{array}$$

[Chandra, Collier, TH, Maloney '22]

Violation of global symmetries

There is a general argument that quantum gravity violates all global symmetries.

e.g. [Misner, Wheeler '57]

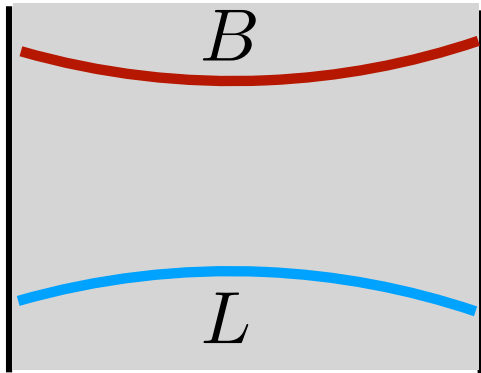


In the Standard Model:

$$U(1)_{B-L} \quad (qqq\ell)$$

Thin-shell wormholes can be used to calculate nonperturbative contributions to this effect.

Two-copy wormhole:



$$e^{-I_{\text{cl}}} \approx \overline{|\langle B|L\rangle|^2}$$

Conclusion

The gravitational path integral encodes the statistics of the UV quantum theory.

Interplay between *quantum information* and *higher topologies*.

However we do not yet know the full story.

- Can all statistics be determined systematically?
- How does this compare to UV-complete examples in string theory?
- Is there a role for higher topology / quantum info in cosmology?

Conclusion

The gravitational path integral encodes the statistics of the UV quantum theory.

Interplay between *quantum information* and *higher topologies*.

However we do not yet know the full story.

- Can all statistics be determined systematically?
- How does this compare to UV-complete examples in string theory?
- Is there a role for higher topology / quantum info in cosmology?

Thank you.